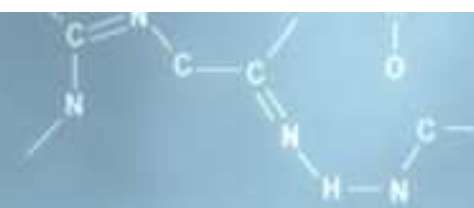


HQL 2018



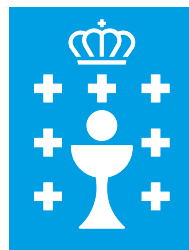
Charmless b Decays

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on behalf of the LHCb collaboration

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28 May 2018



**XUNTA
DE GALICIA**



Outline

1. Types of CP violation

- Direct, mixing-induced

2. 2-body

- $B^0 \rightarrow K^+ \pi^-, \pi^+ \pi^-, B_s^0 \rightarrow K^+ K^-$

3. 3-body

- $B^0 \rightarrow K_S^0 \pi^+ \pi^-$

4. 4-body

- $\Lambda_b^0 \rightarrow p \pi^- \pi^+ \pi^-, \Lambda_b^0 \rightarrow p K^- h^+ h^-$ ^{New!}, $\Xi_b^0 \rightarrow p K^- \pi^+ K^-$ ^{New!}
- $B_s^0 \rightarrow \phi \phi, K^* \bar{K}^*$

Conditions for Direct CP Violation

In charged B decays, presence of multiple amplitudes may lead to direct CP violation

$$A(B \rightarrow f) = \sum_i |A_i| e^{i(\delta_i + \phi_i)}$$

$$\bar{A}(\bar{B} \rightarrow \bar{f}) = \sum_i |A_i| e^{i(\delta_i - \phi_i)}$$

Strong phase (δ) invariant under CP , while weak phase (ϕ) changes sign under CP

$$\mathcal{A}_{CP}(B \rightarrow f) \equiv \frac{|\bar{A}|^2 - |A|^2}{|\bar{A}|^2 + |A|^2} \propto \sum_{i,j} |A_i| |A_j| \sin(\delta_i - \delta_j) \sin(\phi_i - \phi_j)$$

3 conditions required for direct CP violation

At least 2 amplitudes

Non-zero strong phase difference, $\delta_i - \delta_j \neq 0$

Non-zero weak phase difference, $\phi_i - \phi_j \neq 0$

Source of weak phase differences come from different CKM phases of each amplitude

Short-Distance Contributions

Multiple sources of strong phase

1. Short-distance contributions (quark level)

BSS mechanism, PRL **43** 242 (1979)

Tree contribution (a)

Penguin diagram (b) contains 3 quark generations in loop

S -matrix unitarity, CPT require absorptive amplitude

If gluon in penguin is timelike (on-shell)

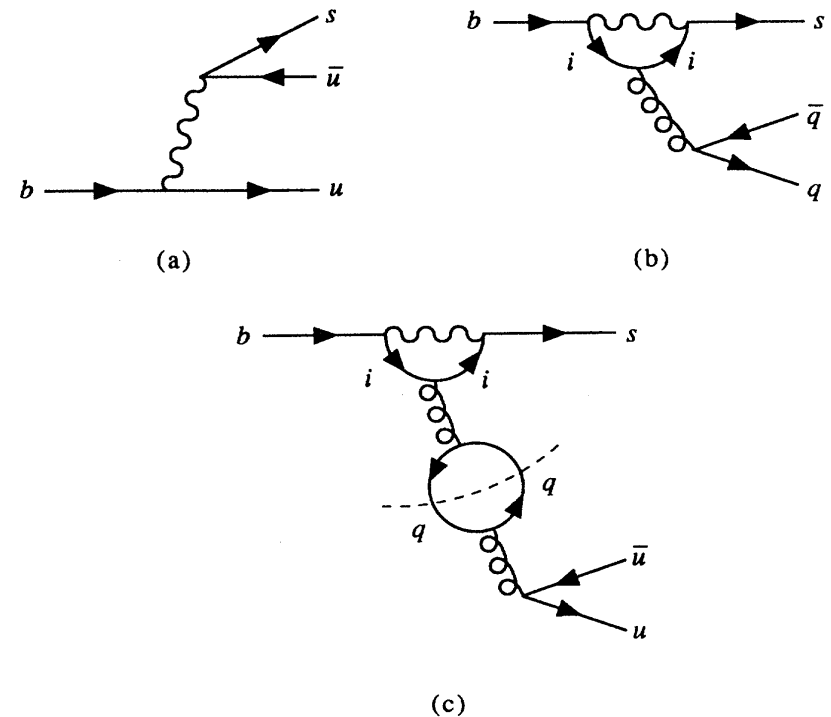
Momentum transfer $q^2 > 4m_i^2$ where $i = u, c$

Imaginary part depends on quark masses

Particle rescattering (c) generates a phase difference

CP violation in 2-body processes caused by this effect

eg. $B^0 \rightarrow K^+ \pi^-$



Long-Distance Contributions

Remaining sources endemic to multibody decays

Long-distance contributions ($q\bar{q}$ level)

2. Breit-Wigner phase

Propagator represents intermediate resonance states

$$F_R^{\text{BW}}(s) = \frac{1}{m_R^2 - s - im_R\Gamma_R(s)}$$

Phase varies across the Dalitz plot

3. Relative CP -even phase in the isobar model

$$A(B \rightarrow f) = \sum_i |A_i| e^{i(\delta_i + \phi_i)}$$

$$\bar{A}(\bar{B} \rightarrow \bar{f}) = \sum_i |\bar{A}_i| e^{i(\delta_i - \phi_i)}$$

Related to final state interactions between different resonances

Neutral Meson Mixing

Mixing arises from a difference between the mass and flavour eigenstates

$$|P_H\rangle = p|P^0\rangle + q|\bar{P}^0\rangle, \quad |P_L\rangle = p|P^0\rangle - q|\bar{P}^0\rangle$$

p, q are complex mixing parameters

Mixing can be described by the effective 2x2 Hamiltonian

$$H_{ij} = M_{ij} - i\Gamma_{ij}/2$$

M is the mass term

Γ provides the decay term due to the $-i$

Solving the Schrödinger Equation

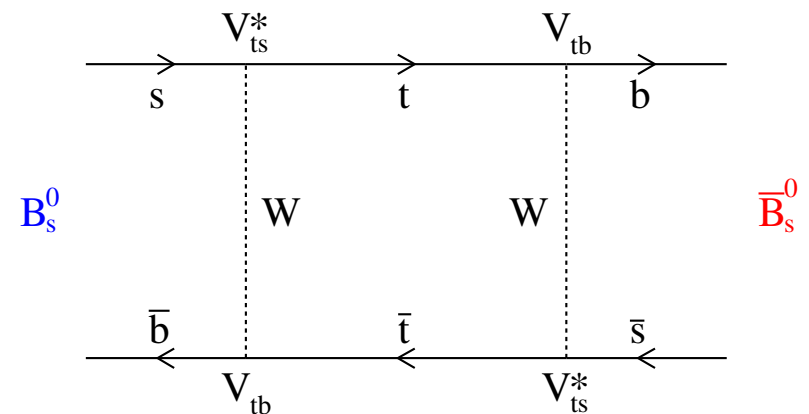
3 mixing physical observables

$\Delta m \equiv m_H - m_L$: mixing frequency in time evolution

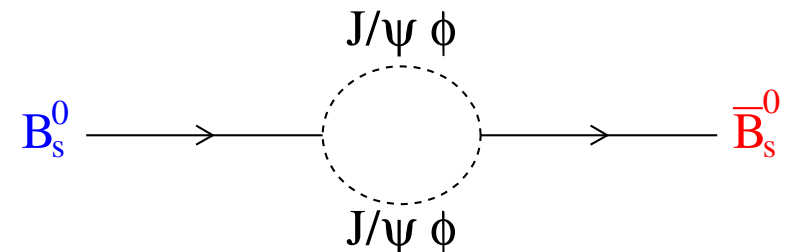
$\Delta\Gamma \equiv \Gamma_H - \Gamma_L$: lifetime difference

$\phi_{\text{mix}} = -\arg(M_{12}/\Gamma_{12})$: CP -violating mixing phase

M_{12} : short-distance (off-shell)



$-i\Gamma_{12}/2$: long-distance (on-shell)



CP Violation in Neutral Mesons

CP violation in neutral meson system governed by complex parameter

$$\lambda_{CP} \equiv \frac{q}{p} \frac{\bar{A}(\bar{P}^0 \rightarrow f_{CP})}{A(P^0 \rightarrow f_{CP})}$$

Access experimentally through time-dependent rate asymmetry in neutral mesons

$$a_{CP}(t) \equiv \frac{\Gamma(\bar{P}^0 \rightarrow f_{CP}) - \Gamma(P^0 \rightarrow f_{CP})}{\Gamma(\bar{P}^0 \rightarrow f_{CP}) + \Gamma(P^0 \rightarrow f_{CP})} = \frac{-\mathcal{C}_{CP} \cos(\Delta mt) + \mathcal{S}_{CP} \sin(\Delta mt)}{\cosh(\Delta\Gamma t/2) + \mathcal{A}_{\Delta\Gamma} \sinh(\Delta\Gamma t/2)}$$

Sensitive to 3 physical observables

\mathcal{C}_{CP} : CP violation in the decay, $|\bar{A}| \neq |A|$

$$\mathcal{C}_{CP} \equiv \frac{|\lambda_{CP}|^2 - 1}{|\lambda_{CP}|^2 + 1}$$

\mathcal{S}_{CP} : Mixing-induced CP violation, $\arg(\lambda_{CP}) \neq 0$

$$\mathcal{S}_{CP} \equiv -\eta_{CP} \frac{2\Im(\lambda_{CP})}{|\lambda_{CP}|^2 + 1}$$

$\mathcal{A}_{\Delta\Gamma}$: Admixture of P_H and P_L that decay to final state

$$\mathcal{A}_{\Delta\Gamma} \equiv -\frac{2\Re(\lambda_{CP})}{|\lambda_{CP}|^2 + 1}$$

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- $B^0 \rightarrow K_S^0 \pi^+ \pi^-$

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- $B_s^0 \rightarrow \phi \phi, K^* \bar{K}^*$

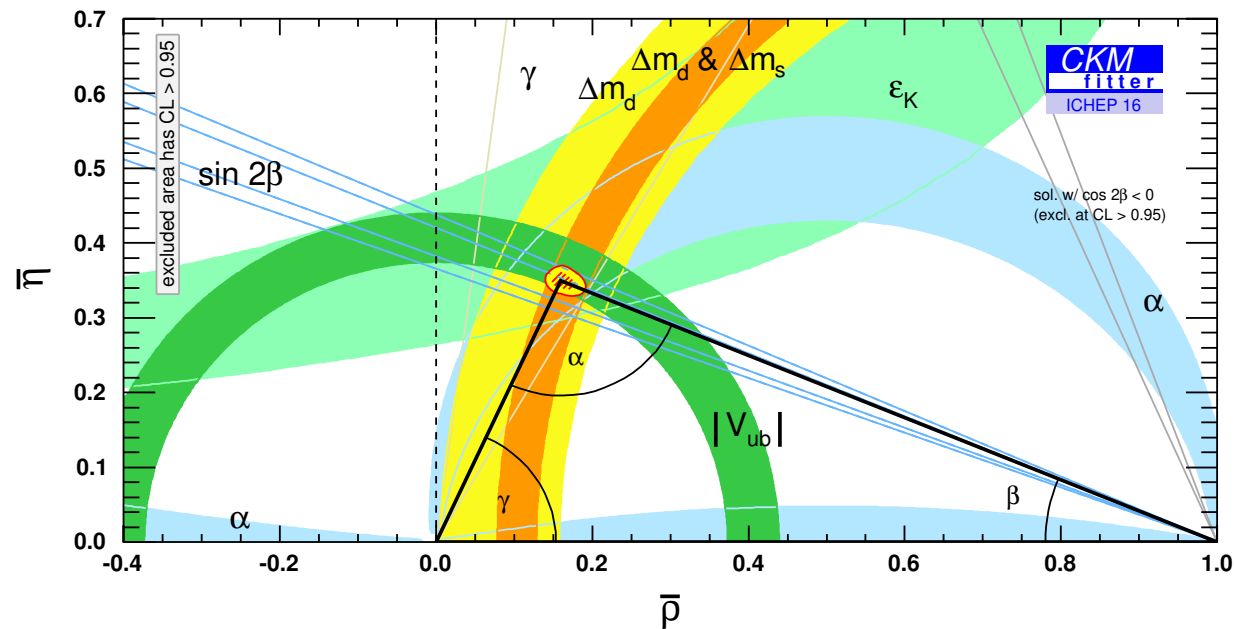
$$B \rightarrow h^+ h^-'$$

Simultaneous analysis includes $B^0 \rightarrow K^- \pi^+, \pi^+ \pi^-$ and $B_s^0 \rightarrow K^+ K^-$

Based on 2011+2012 data (3.0 fb^{-1})

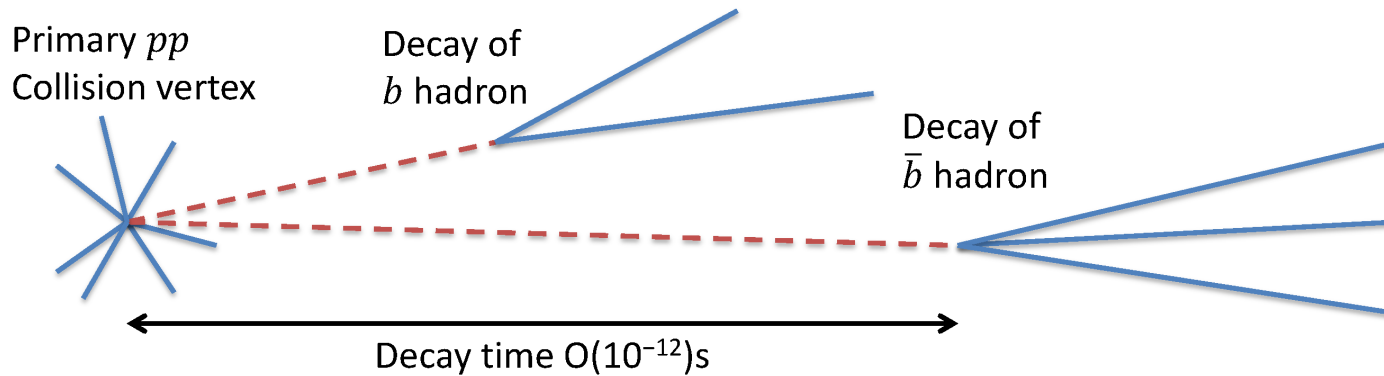
Sensitive to direct and mixing-induced CP violation

Constrain α, γ and $-2\beta_s$



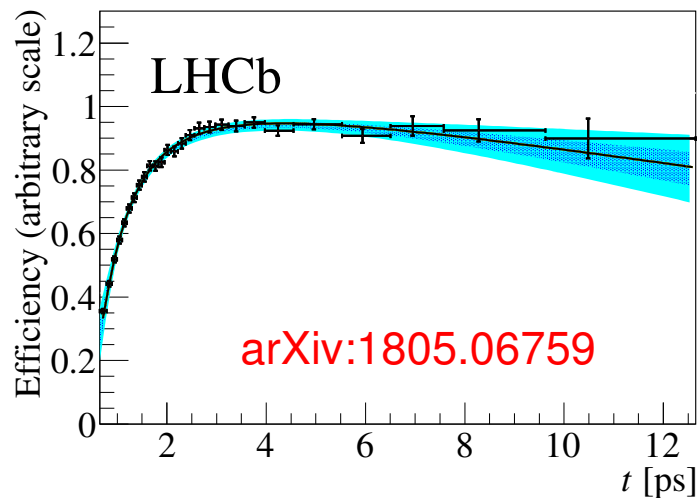
Requires time-dependent and flavour-tagged analysis

Decay Time Distribution



Decay times precisely measured due LHCb VELO vertex measurements

Time distribution affected by acceptance effects due to trigger and selection criteria



Shape determined from $B^0 \rightarrow K^+ \pi^-$ data

Transformation to other final states from simulation

Decay Time Resolution

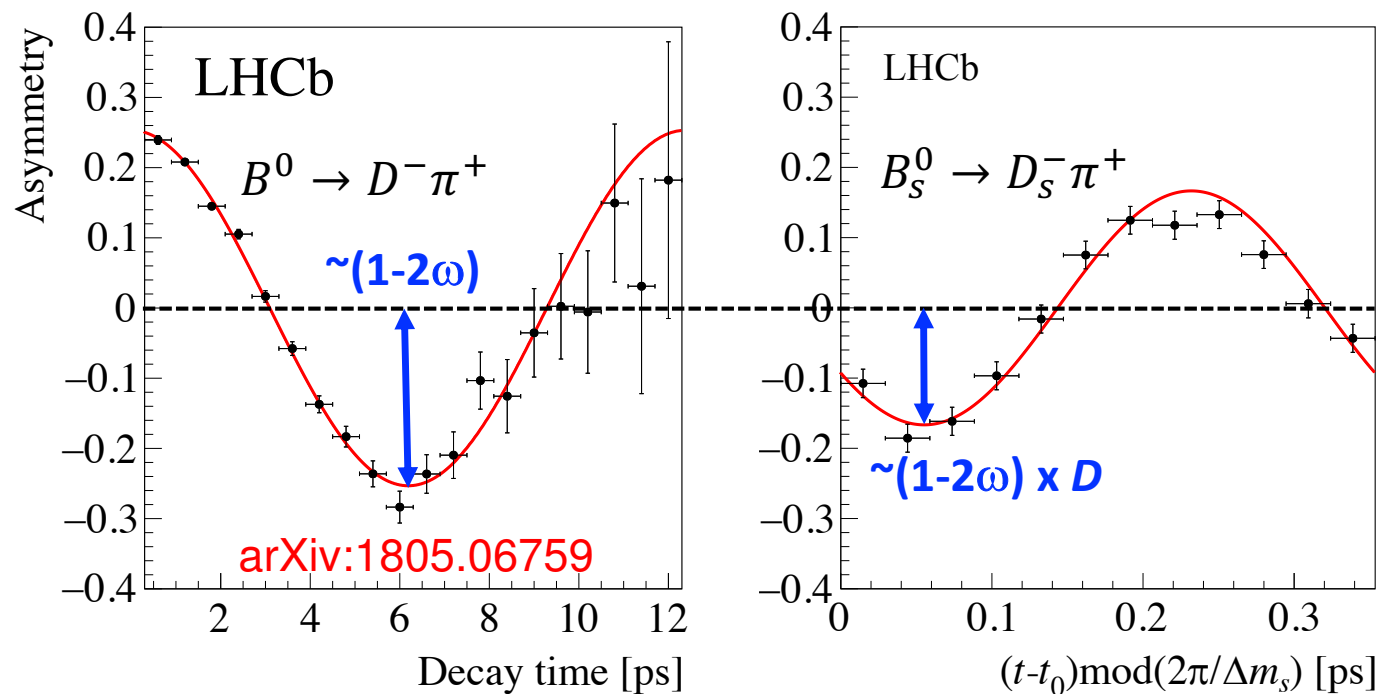
Event-dependent decay time resolution σ_t

Dilutes oscillation amplitudes $D = \exp(\frac{1}{2}\Delta m^2 \sigma_t^2)$

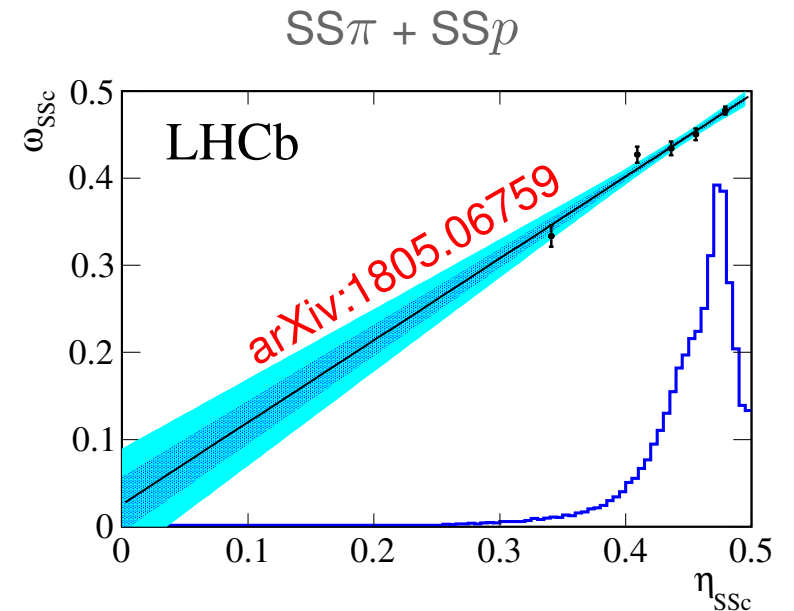
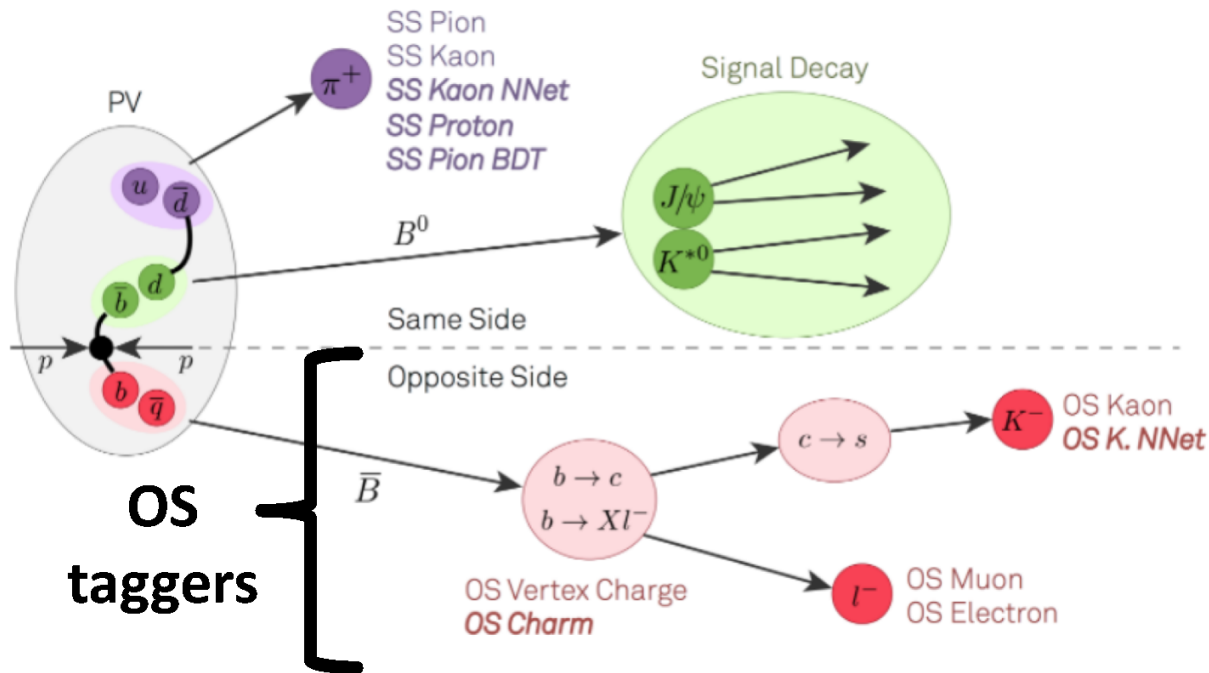
Negligible in B^0 decays due to small Δm_d

Linearly dependent on per-event decay time error

Calibrated from time-dependent asymmetry of $B \rightarrow D\pi$ control samples



Flavour Tagging



Employs Opposite Side (OS) and Same Side (SS) taggers

Calibrated vs Uncalibrated mistag

Algorithm produces per-event tagging decision and associated wrong tag probability

Wrong tag probability linearly calibrated with various control samples

B^0 tagging power: $(4.08 \pm 0.20)\%$, B_s^0 tagging power: $(3.65 \pm 0.21)\%$

$B \rightarrow h^+ h^-$ Results

First error statistical, second systematic

$$\mathcal{C}_{\pi^+\pi^-} = -0.34 \pm 0.06 \pm 0.01,$$

OS: $\pi^+\pi^-$

$$\mathcal{S}_{\pi^+\pi^-} = -0.63 \pm 0.05 \pm 0.01,$$

$$\mathcal{C}_{K^+K^-} = +0.20 \pm 0.06 \pm 0.02,$$

$$\mathcal{S}_{K^+K^-} = +0.18 \pm 0.06 \pm 0.02,$$

$$A_{K^+K^-}^{\Delta\Gamma} = -0.79 \pm 0.07 \pm 0.10,$$

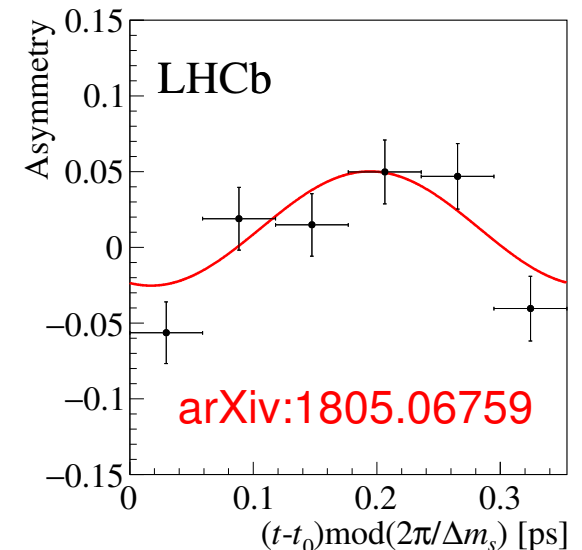
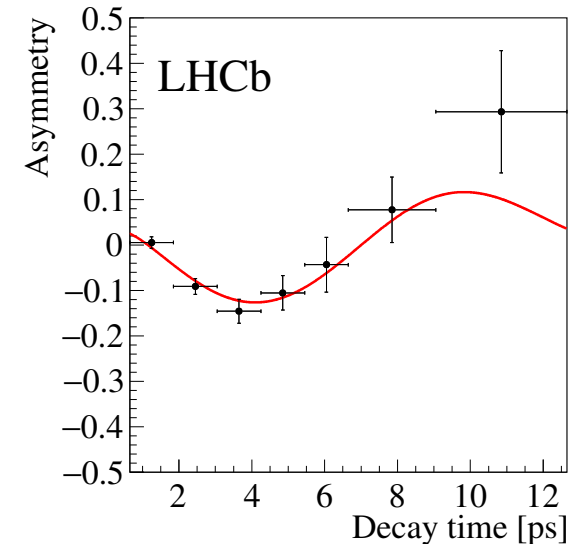
$$\mathcal{A}_{CP}^{B^0 \rightarrow K^+\pi^-} = -0.084 \pm 0.004 \pm 0.003,$$

$$\mathcal{A}_{CP}^{B_s^0 \rightarrow K^-\pi^+} = +0.213 \pm 0.015 \pm 0.007$$

Most precise single measurement

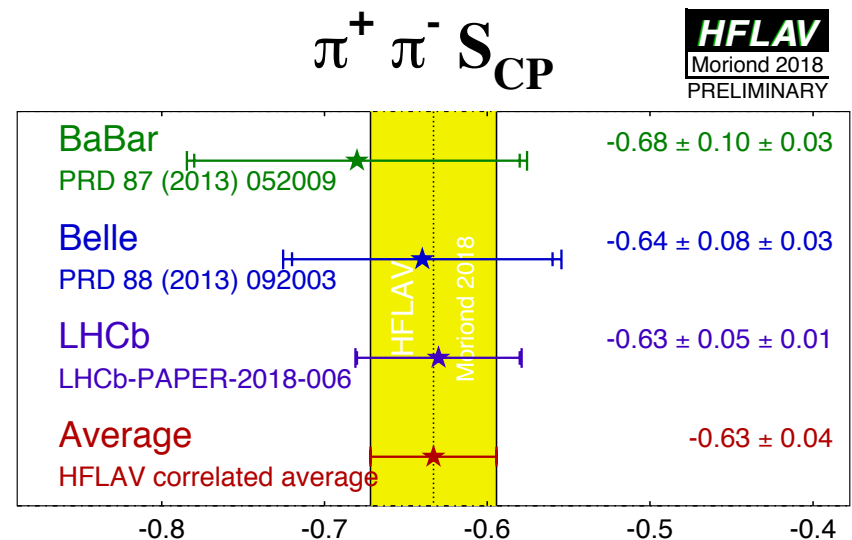
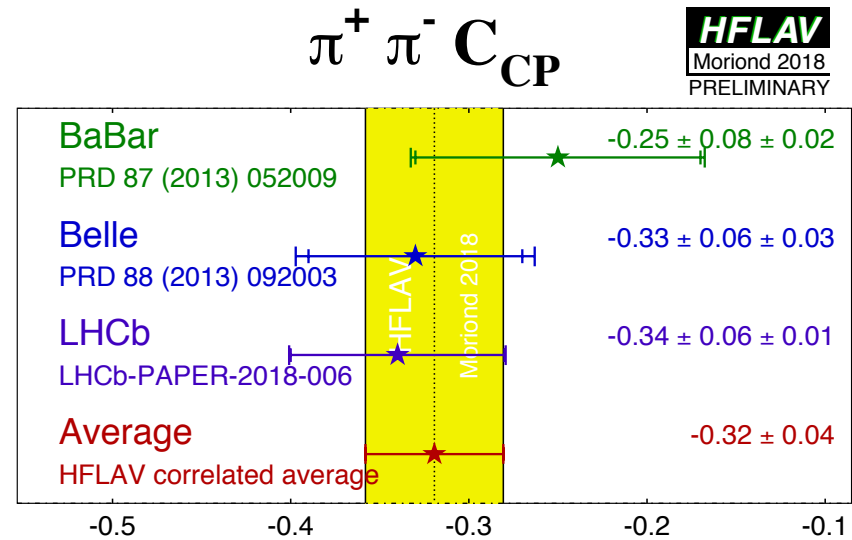
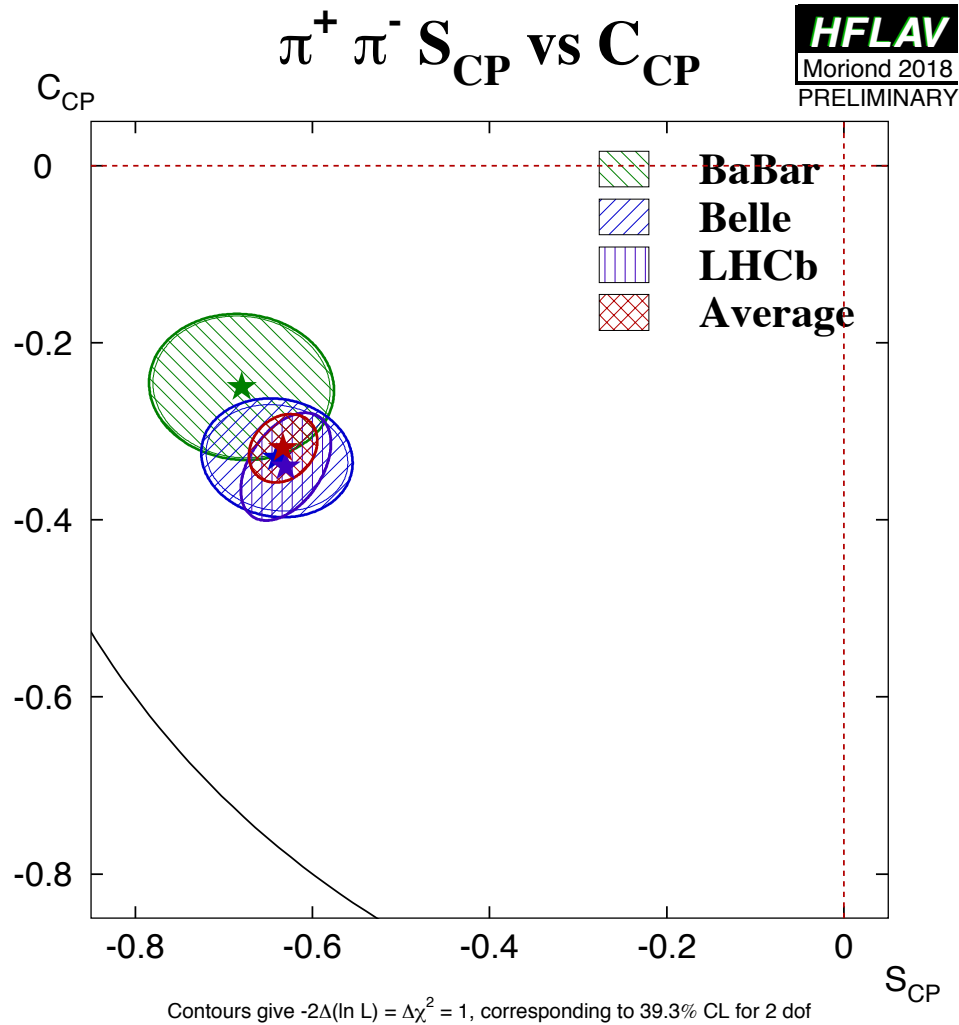
First determination of $A_{K^+K^-}^{\Delta\Gamma}$

4σ evidence for CP violation in $B_s^0 \rightarrow K^+K^-$



OS: K^+K^-

$B \rightarrow h^+ h^-'$ World Average



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3. 3-body

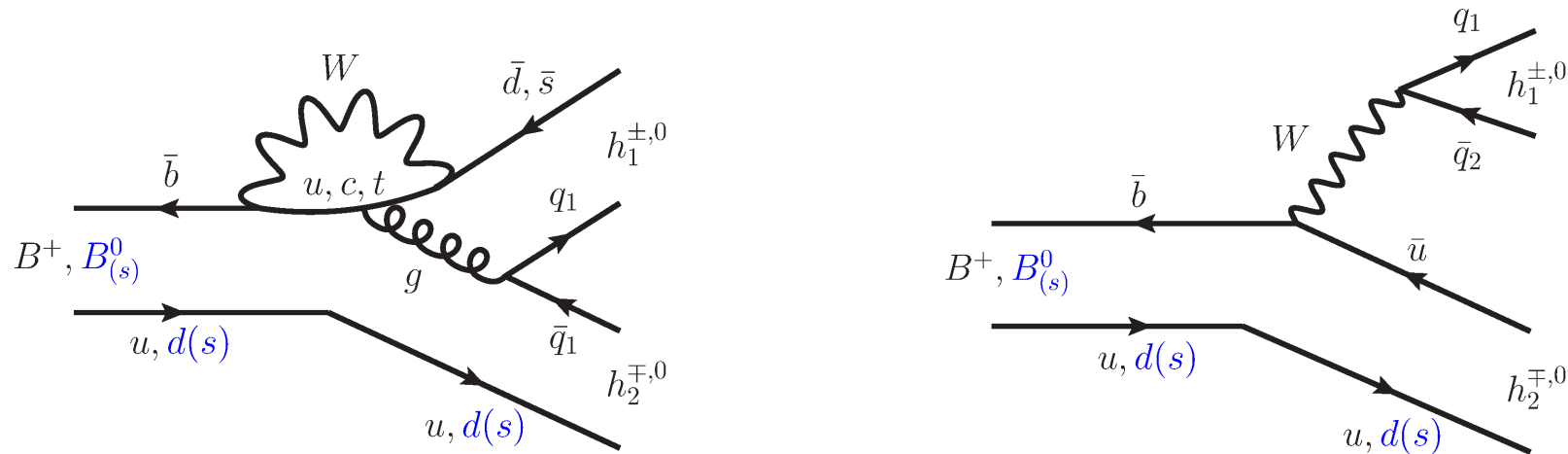
- $B^0 \rightarrow K_S^0 \pi^+ \pi^-$

4. 4-body

- $\Lambda_b^0 \rightarrow p \pi^- \pi^+ \pi^-, \Lambda_b^0 \rightarrow p K^- h^+ h^-$ ^{New!}, $\Xi_b^0 \rightarrow p K^- \pi^+ K^-$ ^{New!}
- $B_s^0 \rightarrow \phi \phi, K^* \bar{K}^*$

$$B^0 \rightarrow K_S^0 \pi^+ \pi^-$$

Mediated by penguin and tree processes



Time-independent amplitude analysis (today)

Sensitive to direct CP violation for intermediate states

Time-dependent amplitude analysis (long-term plan)

Direct measurement of CP violating phase β from CP eigenstate intermediate states

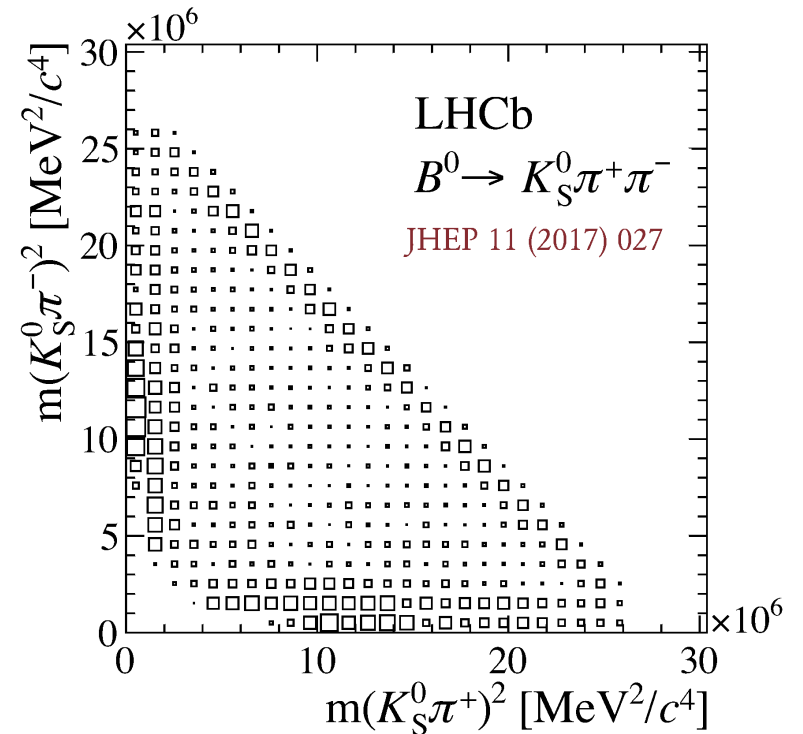
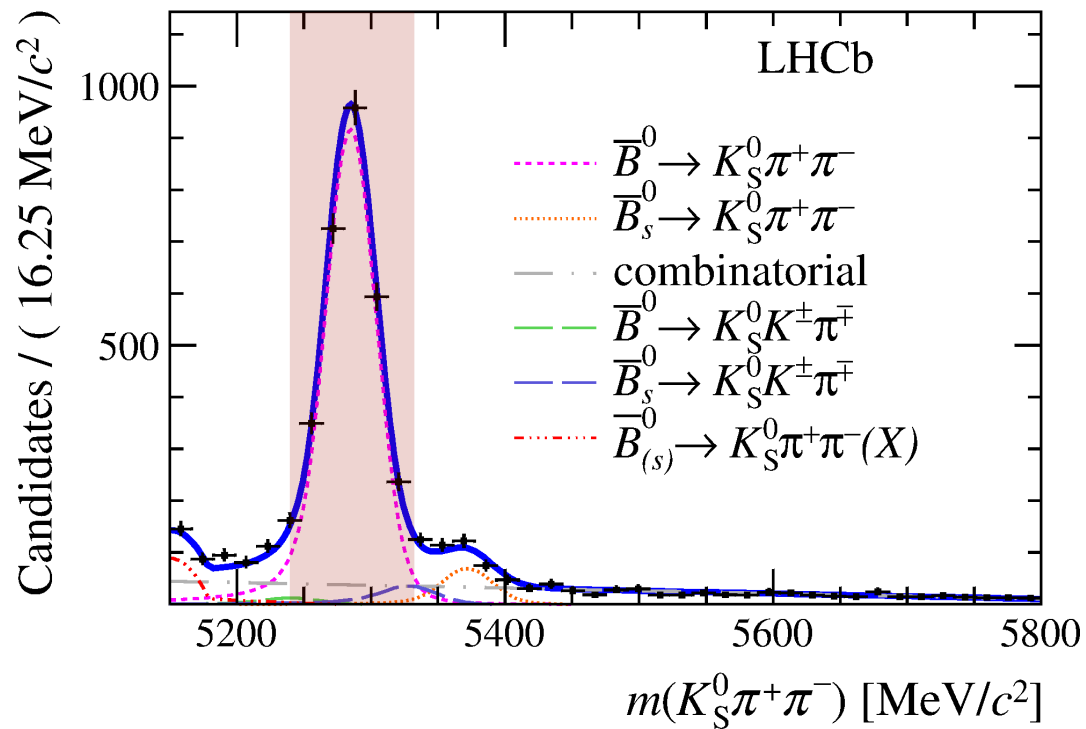
Flavour-specific intermediate states contribute information towards γ measurement

$B^0 \rightarrow K_S^0 \pi^+ \pi^-$ Yield

Analysis performed with 2011+2012 data (3.0 fb^{-1})

Around 3200 signal events in signal region with $\sim 90\%$ purity

[arXiv:1712.09320](https://arxiv.org/abs/1712.09320)



$B^0 \rightarrow K_S^0 \pi^+ \pi^-$ Amplitude

Isobar approach

$$A = \sum_i c_i F_i(m_{12}^2, m_{23}^2)$$

$F_i(m_{12}^2, m_{23}^2)$: strong dynamics form factor

Contains lineshape and spin density

c_i : CP -violating complex fit coefficients

$$\mathcal{A}_{CP}^{\text{Raw},i} = \frac{|\bar{c}_i|^2 - |c_i|^2}{|\bar{c}_i|^2 + |c_i|^2}$$

Raw \mathcal{A}_{CP} corrections

B^0 / \bar{B}^0 production asymmetry

$$(-0.35 \pm 0.81)\%$$

π^+ / π^- detection asymmetry

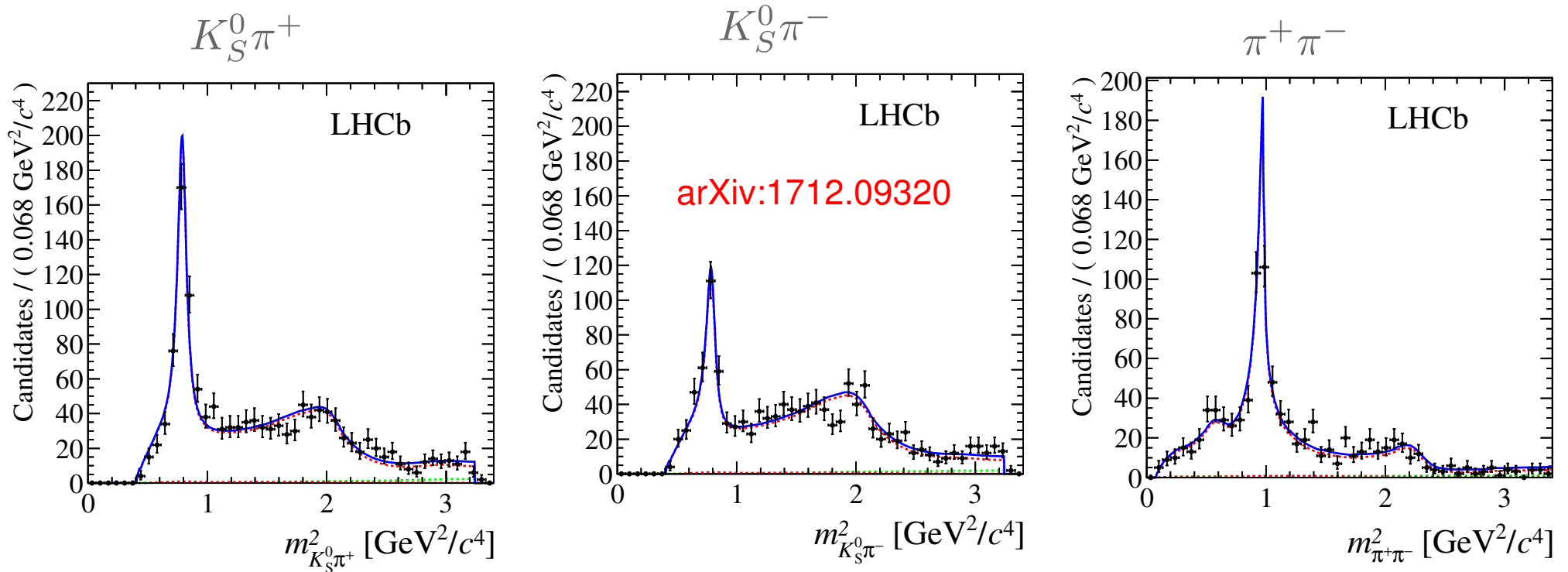
$$(0.00 \pm 0.25)\%$$

Resonance	Parameters	Lineshape
$K^*(892)^-$	$m_0 = 891.66 \pm 0.26$ $\Gamma_0 = 50.8 \pm 0.9$	RBW
$(K\pi)_0^-$	$\mathcal{R}e(\lambda_0) = 0.204 \pm 0.103$ $\mathcal{I}m(\lambda_0) = 0$ $\mathcal{R}e(\lambda_1) = 1$ $\mathcal{I}m(\lambda_1) = 0$	EFKLLM
$K_2^*(1430)^-$	$m_0 = 1425.6 \pm 1.5$ $\Gamma_0 = 98.5 \pm 2.7$	RBW
$K^*(1680)^-$	$m_0 = 1717 \pm 27$ $\Gamma_0 = 332 \pm 110$	Flatté
$f_0(500)$	$m_0 = 513 \pm 32$ $\Gamma_0 = 335 \pm 67$	RBW
$\rho(770)^0$	$m_0 = 775.26 \pm 0.25$ $\Gamma_0 = 149.8 \pm 0.8$	GS
$f_0(980)$	$m_0 = 965 \pm 10$ $g_\pi = 0.165 \pm 0.025 \text{ GeV}$ $g_K = 0.695 \pm 0.119 \text{ GeV}$	Flatté
$f_0(1500)$	$m_0 = 1505 \pm 6$ $\Gamma_0 = 109 \pm 7$	RBW
χ_{c0}	$m_0 = 3414.75 \pm 0.31$ $\Gamma_0 = 10.5 \pm 0.6$	RBW
Nonresonant (NR)		Phase space

EFKLLM: $(K\pi)^0$ form factor from QCDF

Phys. Rev. D **79**, 094005 (2009)

$B^0 \rightarrow K_S^0 \pi^+ \pi^-$ Amplitude



$$\begin{aligned}
 \mathcal{A}_{CP}(K^*(892)^- \pi^+) &= -0.308 \pm 0.060 \pm 0.011 \pm 0.012 \\
 \mathcal{A}_{CP}((K\pi)_0^- \pi^+) &= -0.032 \pm 0.047 \pm 0.016 \pm 0.027 \\
 \mathcal{A}_{CP}(K_2^*(1430)^- \pi^+) &= -0.29 \pm 0.22 \pm 0.09 \pm 0.03 \\
 \mathcal{A}_{CP}(K^*(1680)^- \pi^+) &= -0.07 \pm 0.13 \pm 0.02 \pm 0.03 \\
 \mathcal{A}_{CP}(f_0(980)K_S^0) &= 0.28 \pm 0.27 \pm 0.05 \pm 0.14
 \end{aligned}$$

First observation of CP violation in $B^0 \rightarrow K^{*+}(892)\pi^-$ (6σ significance)

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- $B^0 \rightarrow K_S^0 \pi^+ \pi^-$

4. 4-body

- $\Lambda_b^0 \rightarrow p \pi^- \pi^+ \pi^-, \Lambda_b^0 \rightarrow p K^- h^+ h^-$ ^{New!}, $\Xi_b^0 \rightarrow p K^- \pi^+ K^-$ ^{New!}
- $B_s^0 \rightarrow \phi \phi, K^* \bar{K}^*$

4-body Baryonic Decays

Rich underlying resonant structure

Probe CP violation with integrated and scalar triple-product asymmetry measurements

P -odd triple products

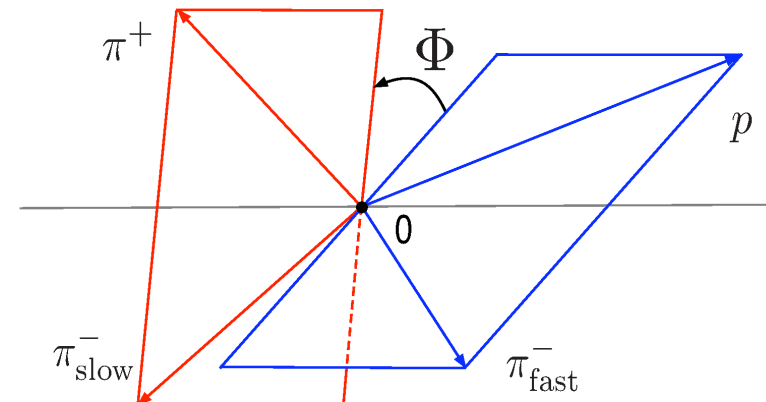
$$\Lambda_b^0: C_{\hat{T}} = \vec{p}_p \cdot (\vec{p}_{h_1^-} \times \vec{p}_{h_2^+}) \propto \sin \Phi$$

$$\bar{\Lambda}_b^0: \bar{C}_{\hat{T}} = \vec{p}_{\bar{p}} \cdot (\vec{p}_{h_1^+} \times \vec{p}_{h_2^-}) \propto \sin \bar{\Phi}$$

P -odd asymmetries of \hat{T} operator

$$A_{\hat{T}} = \frac{N(C_{\hat{T}} > 0) - N(C_{\hat{T}} < 0)}{N(C_{\hat{T}} > 0) + N(C_{\hat{T}} < 0)}$$

$$\bar{A}_{\hat{T}} = \frac{\bar{N}(-\bar{C}_{\hat{T}} > 0) - \bar{N}(-\bar{C}_{\hat{T}} < 0)}{\bar{N}(-\bar{C}_{\hat{T}} > 0) + \bar{N}(-\bar{C}_{\hat{T}} < 0)}$$



P -odd observable

$$a_P^{\hat{T}\text{-odd}} = \frac{1}{2}(A_{\hat{T}} + \bar{A}_{\hat{T}})$$

CP -odd observable

$$a_{CP}^{\hat{T}\text{-odd}} = \frac{1}{2}(A_{\hat{T}} - \bar{A}_{\hat{T}})$$

Sensitive to interference between P -even and P -odd amplitudes

$\Lambda_b^0 \rightarrow p\pi^- h^+ h^-$ Results

Based on 2011-12 data (3.0 fb^{-1})

Nature Physics 13 (2017) 391

No CP violation in integrated phase space

Divide into bins

Scheme A:

Based on dominant resonant structure

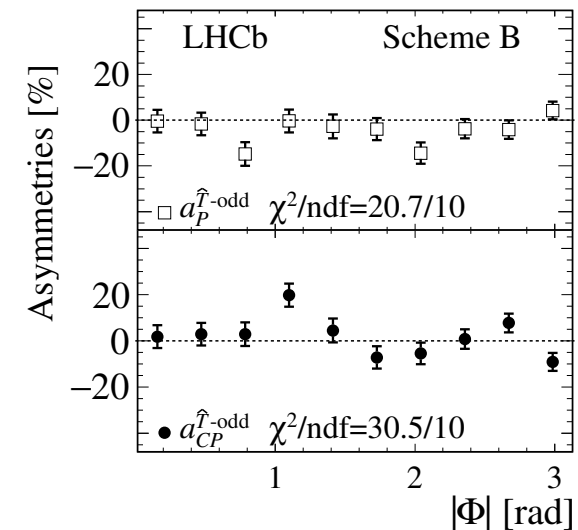
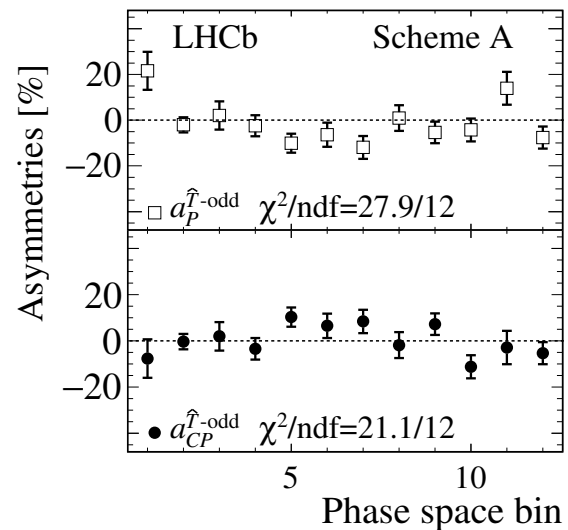
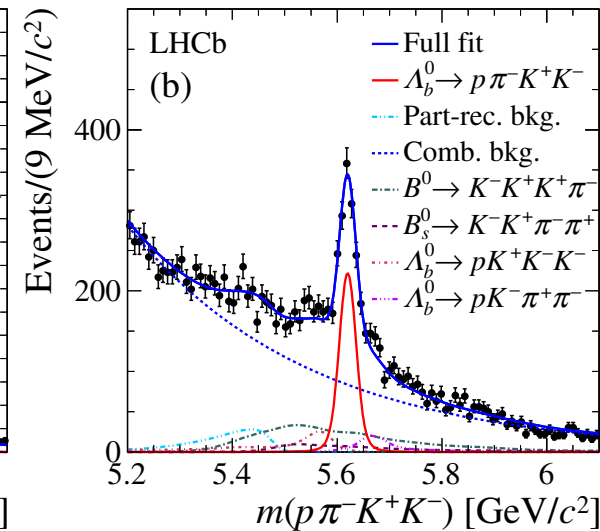
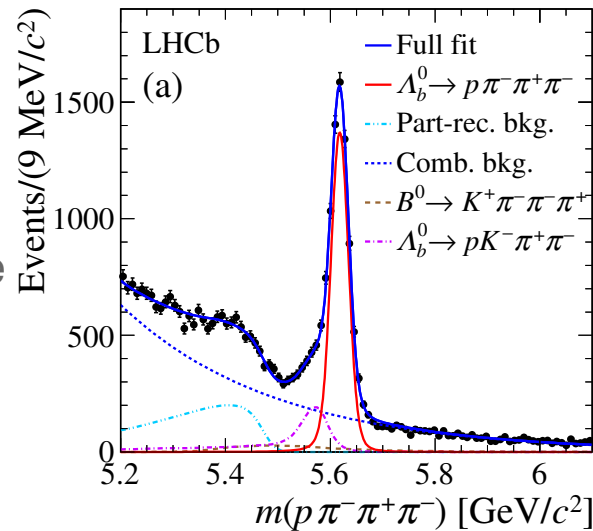
eg. Δ^{++} , N^* , $\rho(770)$

Scheme B:

Function of angle between decay planes

First evidence for CP violation (3.3σ)

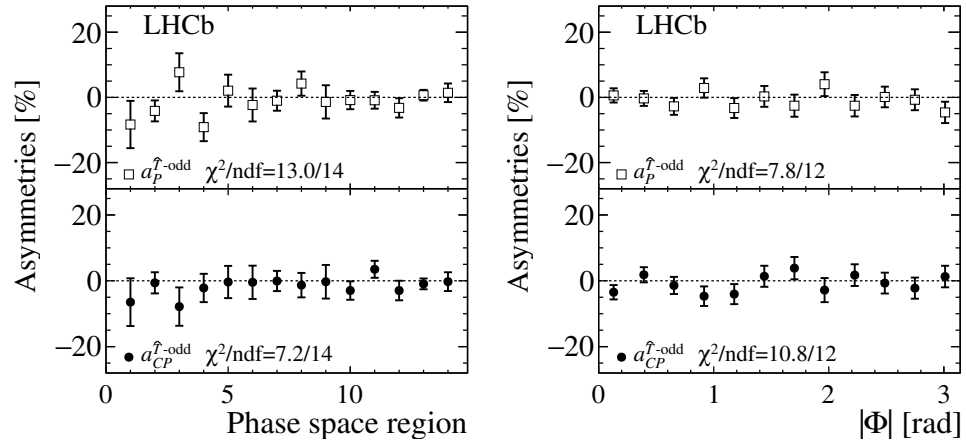
$$\Lambda_b^0 \rightarrow p\pi^- \pi^+ \pi^- \quad \Lambda_b^0 \rightarrow p\pi^- K^+ K^-$$



$\Lambda_b^0 \rightarrow pK^- h^+ h^-$, $\Xi_b^0 \rightarrow pK^- \pi^+ K^-$

New preliminary result based on 3.0 fb^{-1} ($p\pi^- h^+ h^- \rightarrow pK^- h^+ h^-$)

$\Lambda_b^0 \rightarrow pK^- \pi^+ \pi^-$



$\Lambda_b^0 \rightarrow pK^- \pi^+ \pi^-$ Yield: 19877 ± 195

$\Lambda_b^0 \rightarrow pK^- K^+ K^-$ Yield: 5297 ± 83

$\Xi_b^0 \rightarrow pK^- \pi^+ K^-$ Yield: 709 ± 45

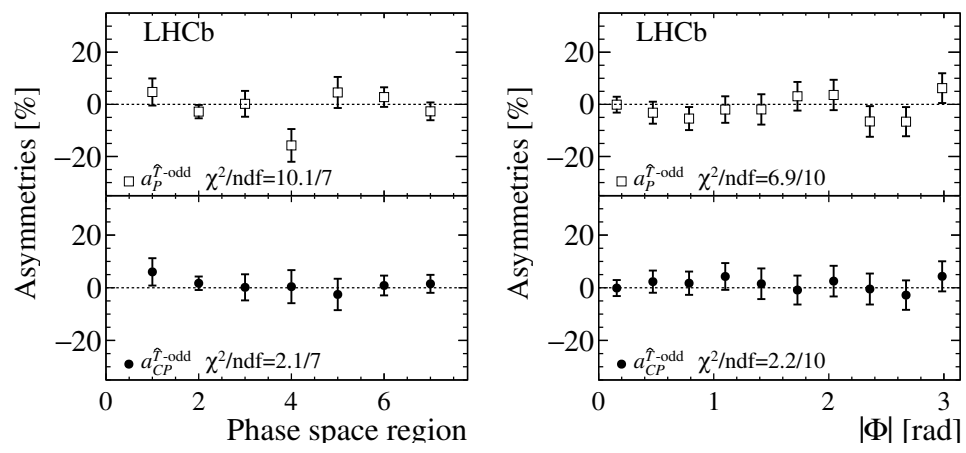
Left: Scheme A, Right: Scheme B

Scheme A: Binned in dominant resonances

Scheme B: Binned in Φ

Additional binned search in mass combinations

$\Lambda_b^0 \rightarrow pK^- K^+ K^-$



No significant asymmetries found

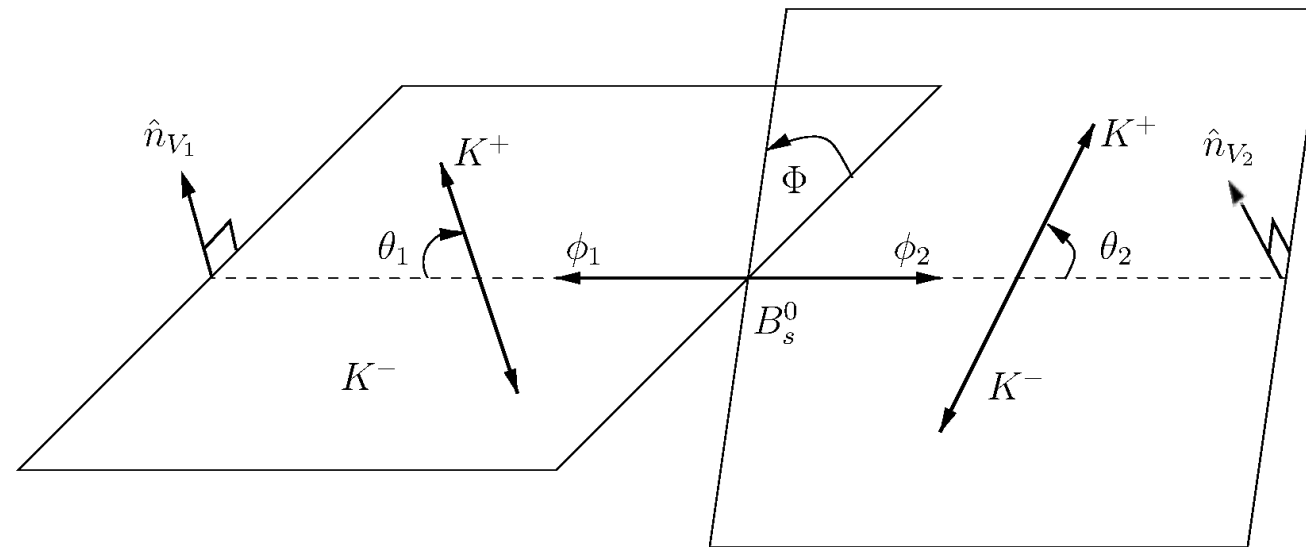
[arXiv:1805.03941](https://arxiv.org/abs/1805.03941)

Time-dependent VV Final States

$B_s^0 \rightarrow \phi\phi$ ($b \rightarrow s\bar{s}s$), $K^*\bar{K}^*$ ($b \rightarrow s\bar{d}d$) penguin dominated final states

Highly sensitive to New Physics amplitudes in the mixing and decay processes

Final state is CP admixture, time-dependent angular analysis to disentangle

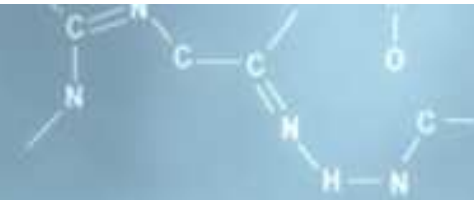


Measure CP -violating mixing phase $\phi_s^{s\bar{s}s}$, $\phi_s^{s\bar{d}d}$

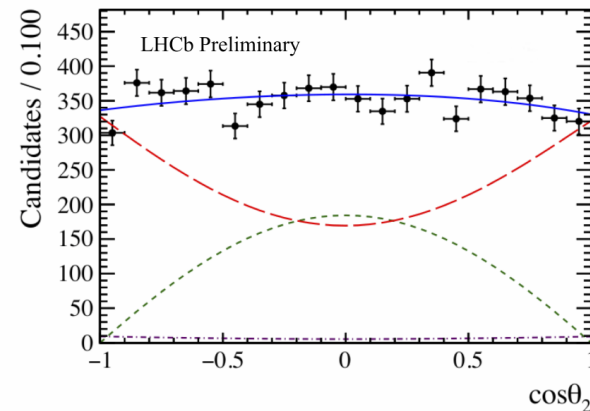
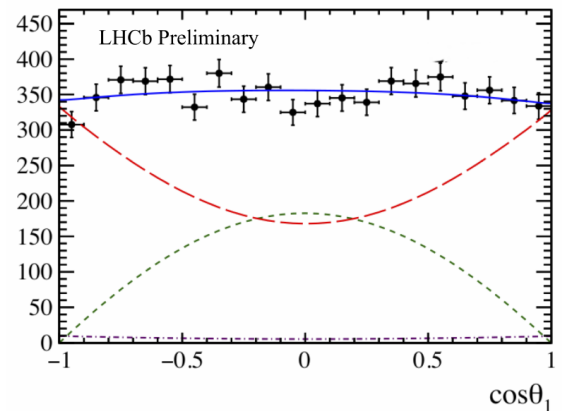
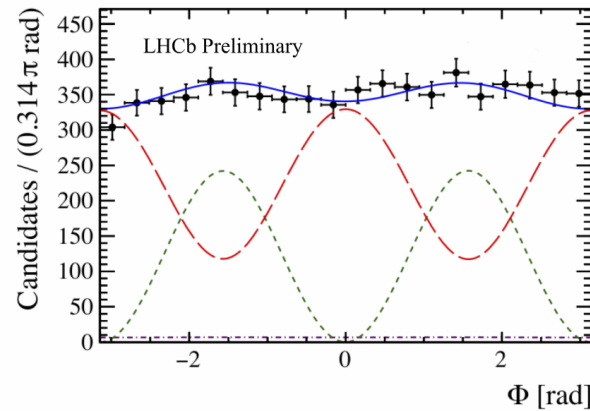
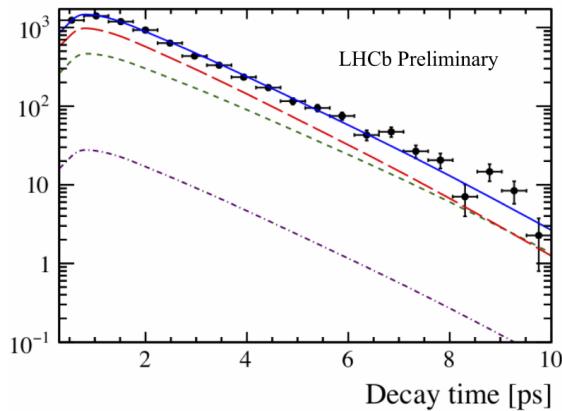
Theory: $|\phi_s^{s\bar{s}s}| < 0.02$ rad

arXiv:0810.0249, Nucl. Phys. B **774**, 64 (2007), Phys. Rev. D **80**, 114026 (2009)

$B_S^0 \rightarrow \phi\phi$



Analysis based on Run 1 and 2015+16 data ($5 fb^{-1}$), LHCb-CONF-2018-001



Effective tagging efficiency
 $(5.74 \pm 0.43)\%$

Red: CP -even VV

Green: CP -odd VV

Purple: $SV + SS$

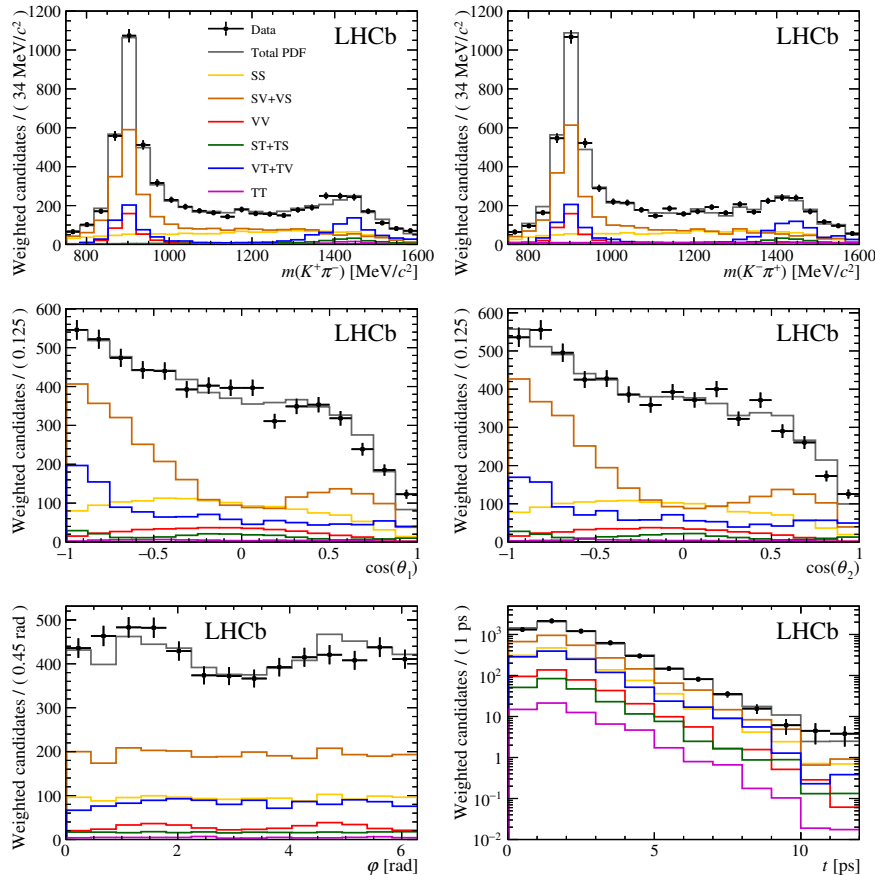
$$\phi_s^{S\bar{S}S} = -0.07 \pm 0.13 \text{ (stat)} \pm 0.03 \text{ (syst)} \text{ rad}$$

$$|\lambda_{CP}| = 1.02 \pm 0.05 \text{ (stat)} \pm 0.03 \text{ (syst)}$$

Additional search with triple product asymmetries shows no CP violation

$$B_s^0 \rightarrow K^* \bar{K}^*$$

Analysis based on 2011+12 data (3 fb^{-1})



World's first measurement

JHEP 03 (2018) 140

$K\pi$ mass distribution modelled

Effective tagging efficiency: $(5.17 \pm 0.17)\%$

Systematics dominated by multi-dimensional acceptance

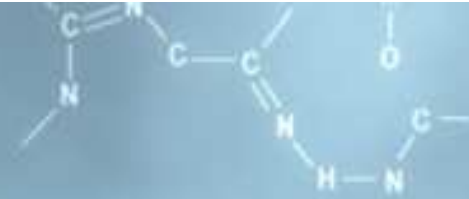
No evidence for CP violation

Results consistent with $B_s^0 \rightarrow \phi\phi$

$$\phi_s^{s\bar{d}d} = -0.10 \pm 0.13 \text{ (stat)} \pm 0.14 \text{ (syst)} \text{ rad}$$

$$|\lambda_{CP}| = 1.035 \pm 0.034 \text{ (stat)} \pm 0.089 \text{ (syst)}$$

Summary



LHCb provides a rich environment to search for various manifestations of CP violation

Time-dependent measurement of CP violation in $B \rightarrow h^+ h^-$

Most precise single measurement

Amplitude analysis of $B^0 \rightarrow K_S^0 \pi^+ \pi^-$

First observation of CP violation in $B^0 \rightarrow K^{*+} \pi^-$

Search for CP violation in 4-body baryonic b decays

First evidence of CP violation in $\Lambda_b^0 \rightarrow p \pi^- \pi^+ \pi^-$ with triple product constructs

Time-dependent measurements of ϕ_s with $B^0 \rightarrow VV$ channels

$B_s^0 \rightarrow \phi\phi$ consistent with SM predictions

First measurement with $B_s^0 \rightarrow K^* \bar{K}^*$