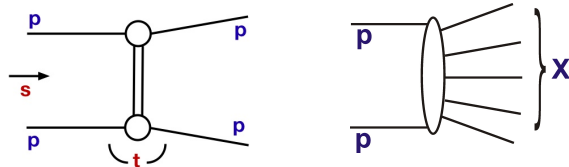


Measurements of the Total Proton-Proton Cross Section with the ATLAS Detector



Paul Newman
(University of Birmingham)
for the ATLAS Collaboration

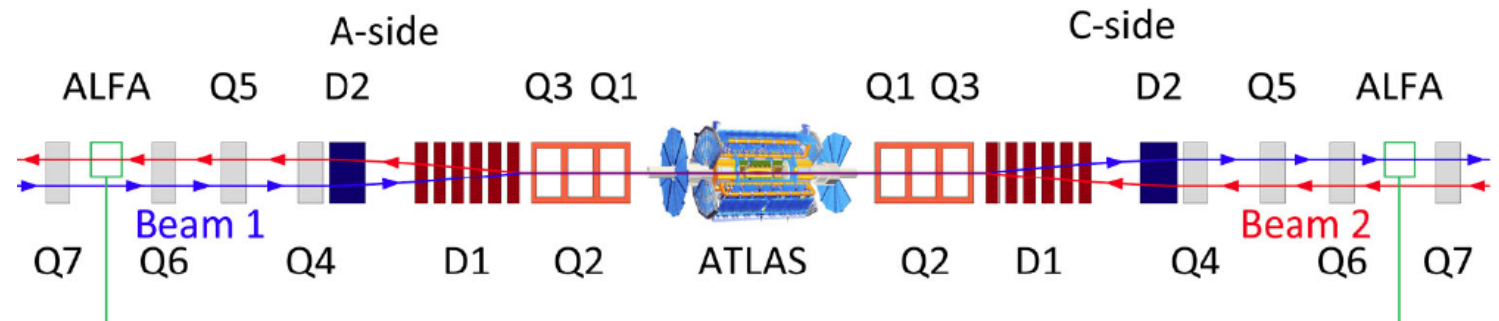


EDS Blois 2017, Prague
26 June 2016



- Elastic & total cross sections with ALFA Roman pots
- Total inelastic cross section from minimum bias data
- Including a bit on soft low mass diffraction

Methods



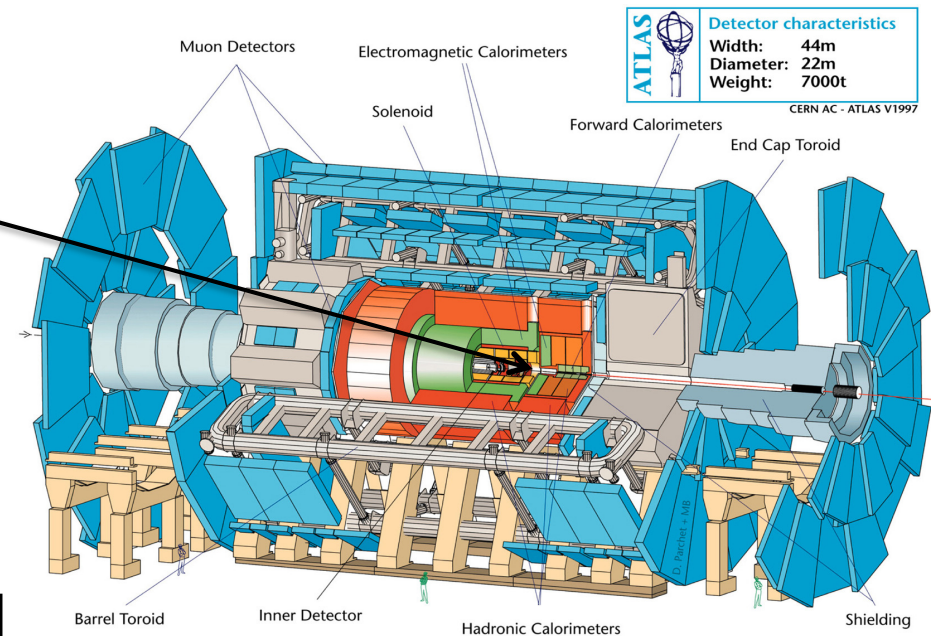
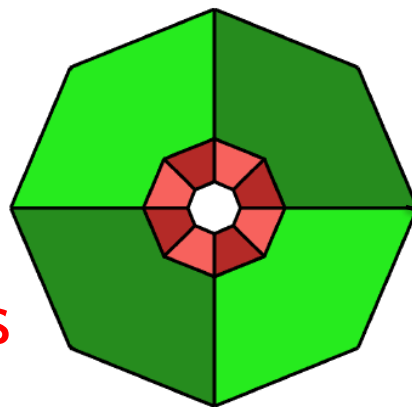
1) Directly tag outgoing protons in ALFA Roman pot spectrometer [4 stations at ~240m from interaction point].

... obtain σ_{el} directly and apply optical theorem for σ_{tot} .

2) Obtain minimally biased samples using the Minimum Bias Trigger Scintillator (MBTS) detectors

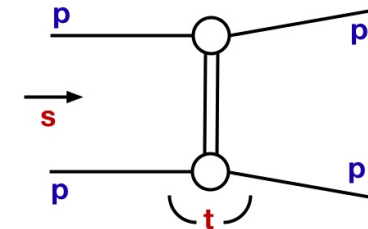
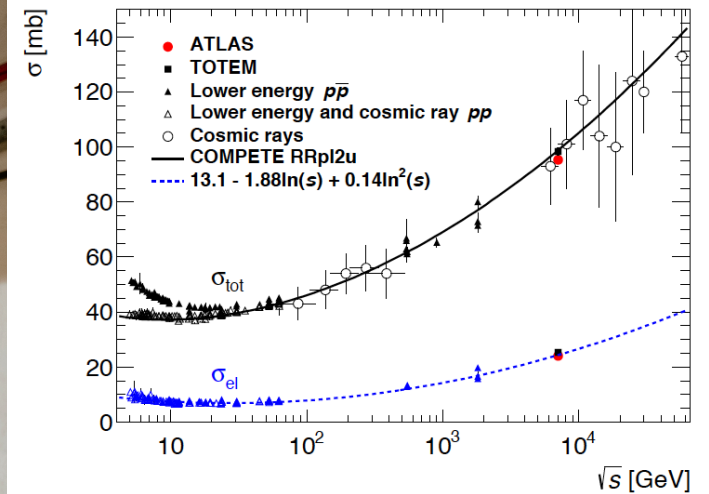
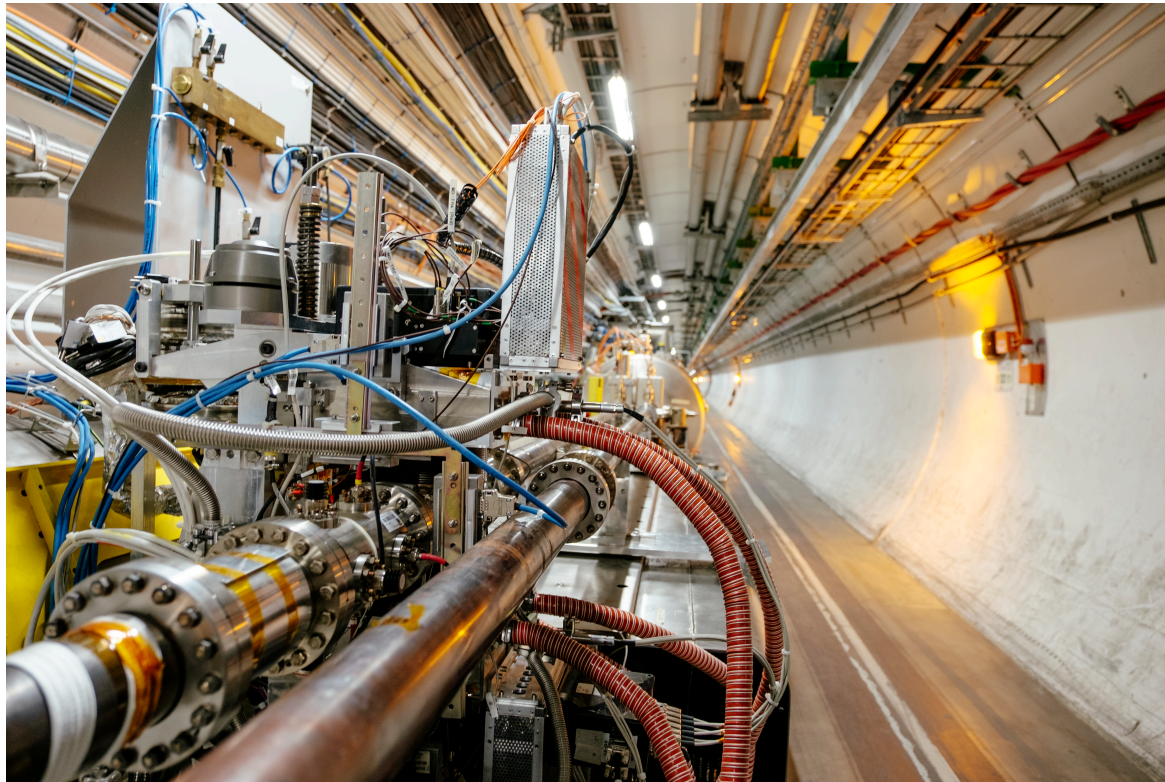
[$z = \pm 3.6\text{m}$, $2.1 < |\eta| < 3.8$, 8-fold Segmentation nearest beam-pipe, 4-fold further out].

... obtain almost all of σ_{inel} directly.



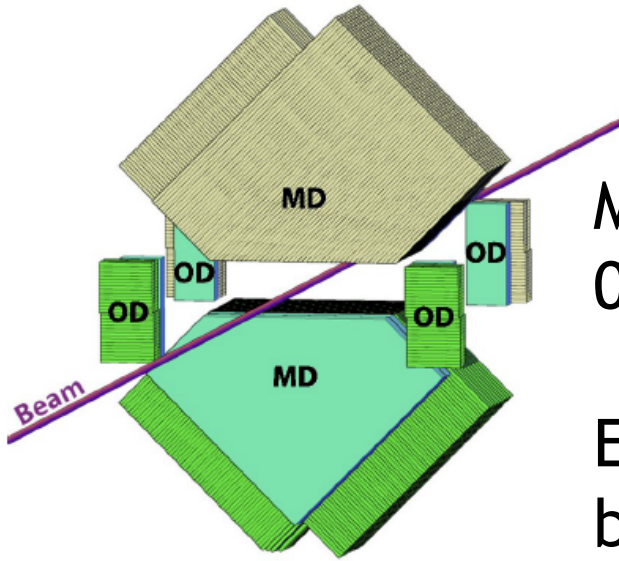
1) Proton Tagging Method

- Earlier result at $\sqrt{s}=7$ TeV: Nucl Phys B889 (2014) 486
- Presented here, $\sqrt{s}=8$ TeV result: Phys Lett B761 (2016) 158



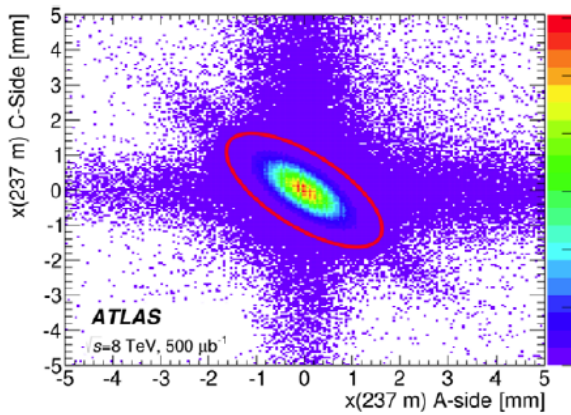
2012 $\beta^* = 90$ m run, allowing access to small t
 \rightarrow 3.8 million events after selection, background $\sim 0.1\%$

Measurement Principle

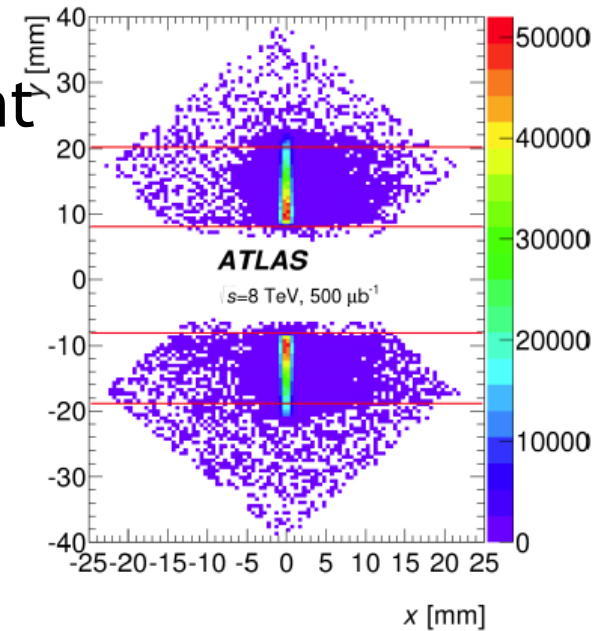
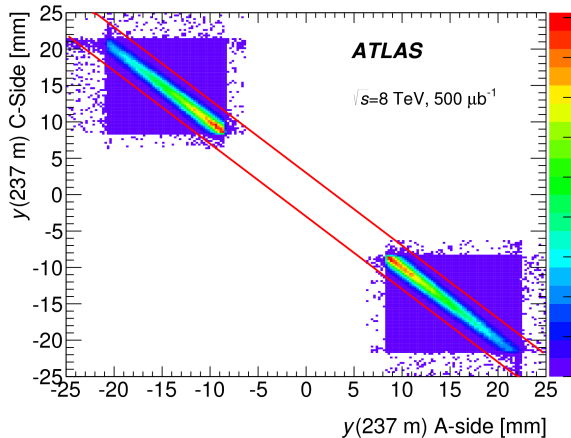


Main detectors (MD) are 2x10 layers of 0.5mm² square fibres per pot.

Elastically scattered protons characterised by back-to-back topology (anti-correlation in x and y between A and C sides of ATLAS)



At high β^* , parallel to point focusing such that y coordinate maps linearly onto proton p_T ; hence quadratically onto t:

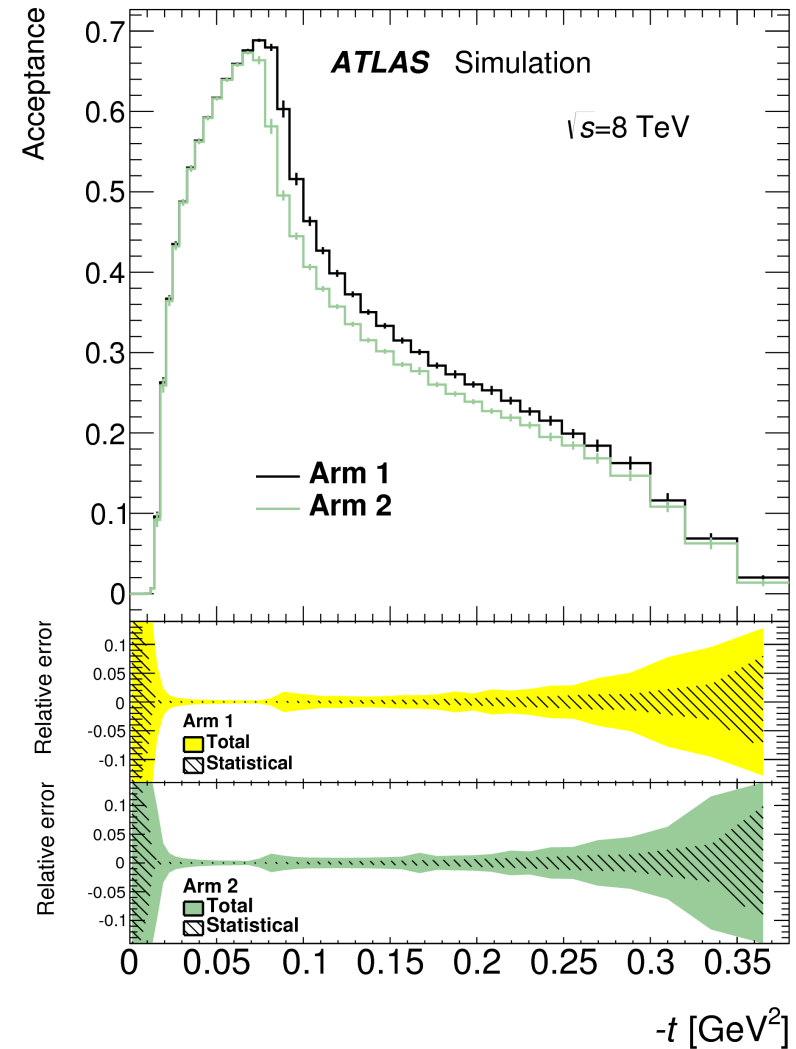
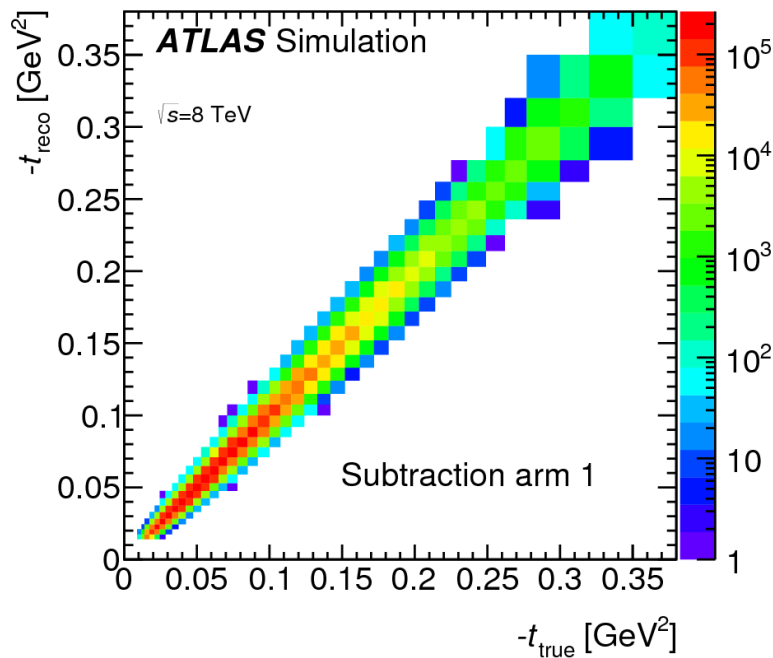


$$\begin{pmatrix} y \\ \theta_y \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} y^* \\ \theta_y^* \end{pmatrix}$$

$$t = -(p\theta^*)^2 \quad 4$$

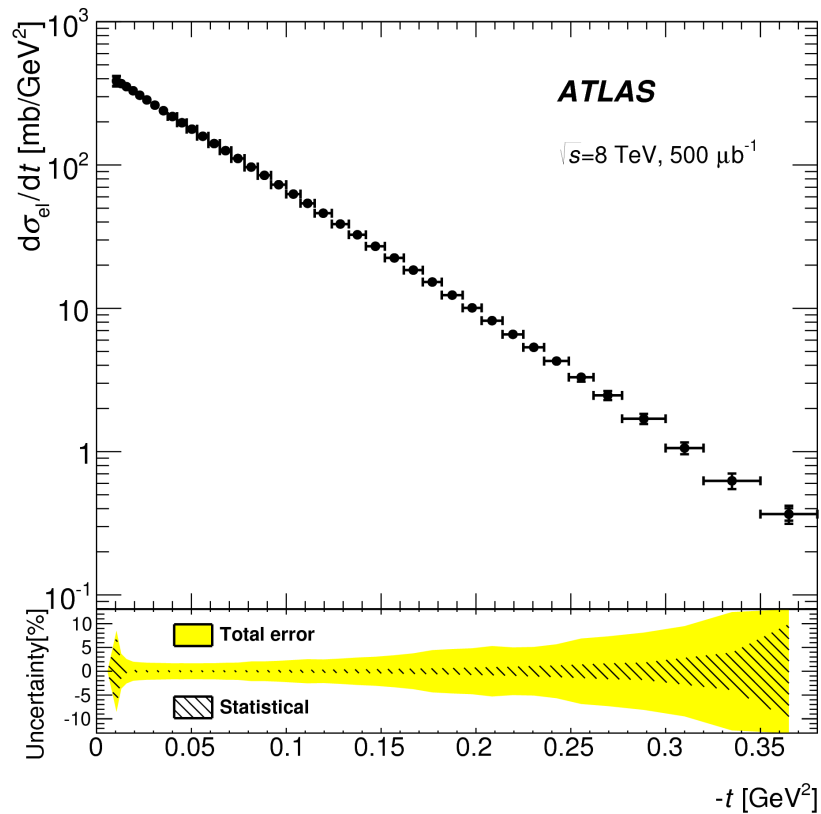
Correcting for Instrumental Effects

- Reconstruction ($\sim 90\%$) and trigger ($>99.9\%$) efficiencies determined from data.
- Beam optics model tuned with ALFA constraints and applied to resolution and acceptance considerations via PYTHIA8



Dedicated lumi determination \rightarrow 1.5% uncertainty

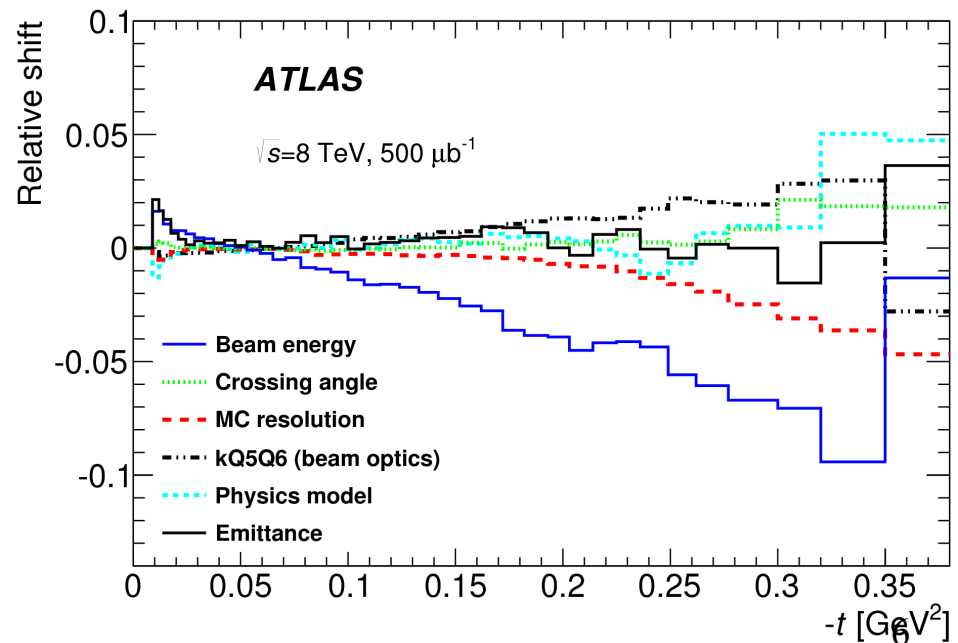
Unfolded Data



- Data cover the region $0.01 < |t| < 0.36\ \text{GeV}^2$, systematic limited throughout

- Dominant systematic at small $|t|$ is luminosity (1.5% normalisation)

- 0.65% uncertainty on beam energy generates largest systematic at high $|t|$



(a)

Fitting the Data

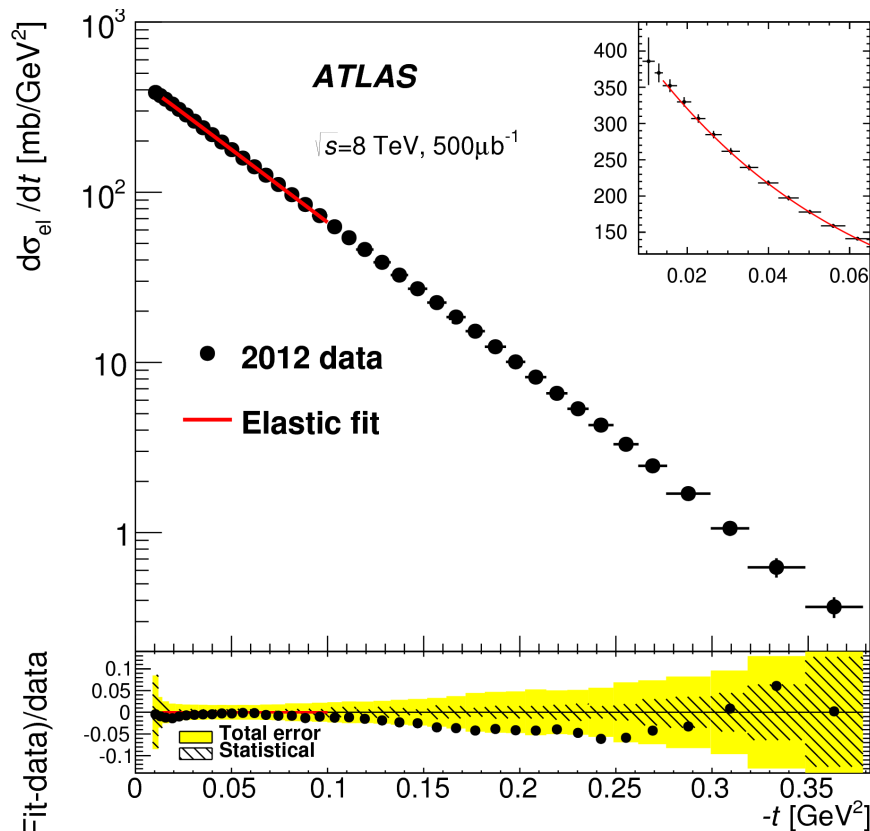
$$\frac{d\sigma}{dt} = \frac{1}{16\pi} |f_N(t) + f_C(t)e^{i\alpha\phi(t)}|^2,$$

$$f_C(t) = -8\pi\alpha\hbar c \frac{G^2(t)}{|t|},$$

$$f_N(t) = (\rho + i) \frac{\sigma_{\text{tot}}}{\hbar c} e^{-B|t|/2},$$

Fixed parameters:

- ρ (~ 0.14) = ratio of Re/Im amplitudes at $t=0$
- G = proton electric form factor
- ϕ = phase of Coulomb-nuclear interference at $t=0$



(b)

Influence of interference term small.

Data in hadronic region compatible with pure exponential: $\sim e^{Bt}$

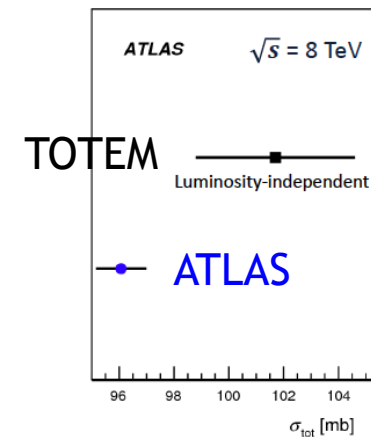
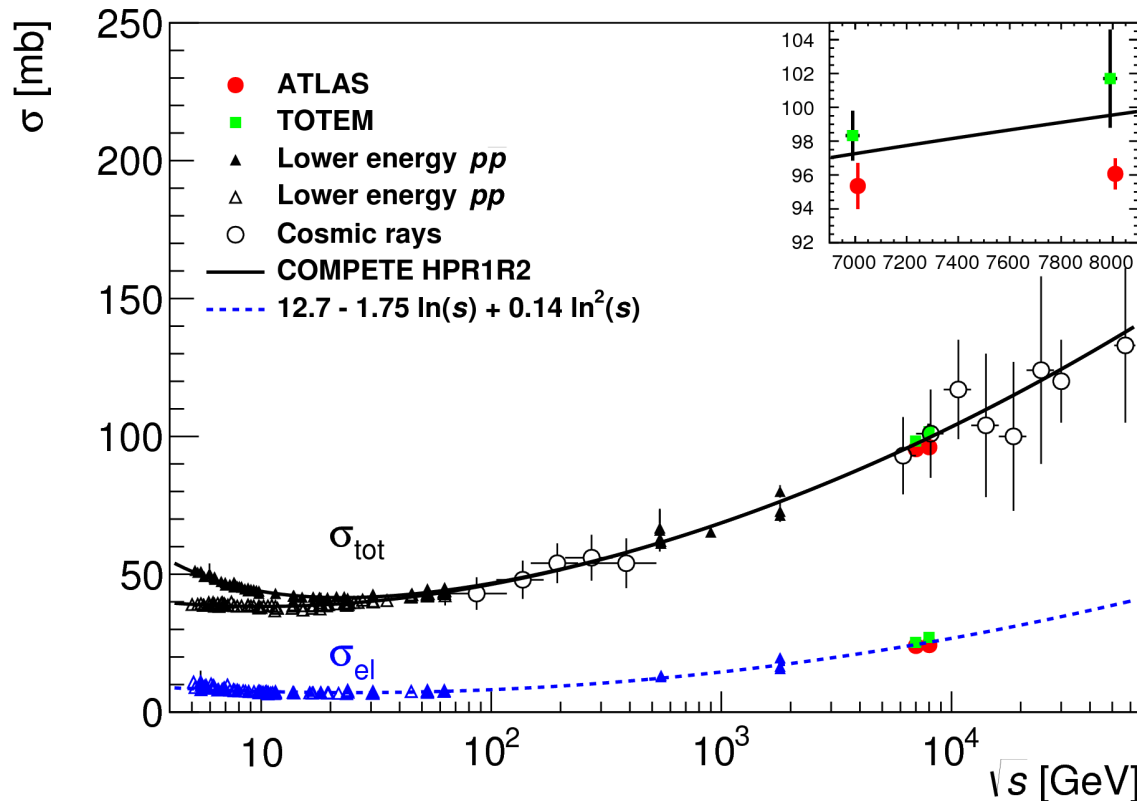
Total hadronic cross section emerges via fit normalisation

Results and Energy Dependence: σ_{tot}

$$\sigma_{tot}(8 \text{ TeV}) = 96.07 \pm 0.18(\text{stat.}) \pm 0.85(\text{exp.}) \pm 0.31(\text{extr.}) \text{ mb}$$

$$\sigma_{el}(8 \text{ TeV}) = 24.33 \pm 0.04(\text{stat}) \pm 0.39(\text{syst}) \text{ mb}$$

$$\sigma_{inel}(8 \text{ TeV}) = 71.73 \pm 0.15(\text{stat}) \pm 0.69(\text{syst}) \text{ mb}$$

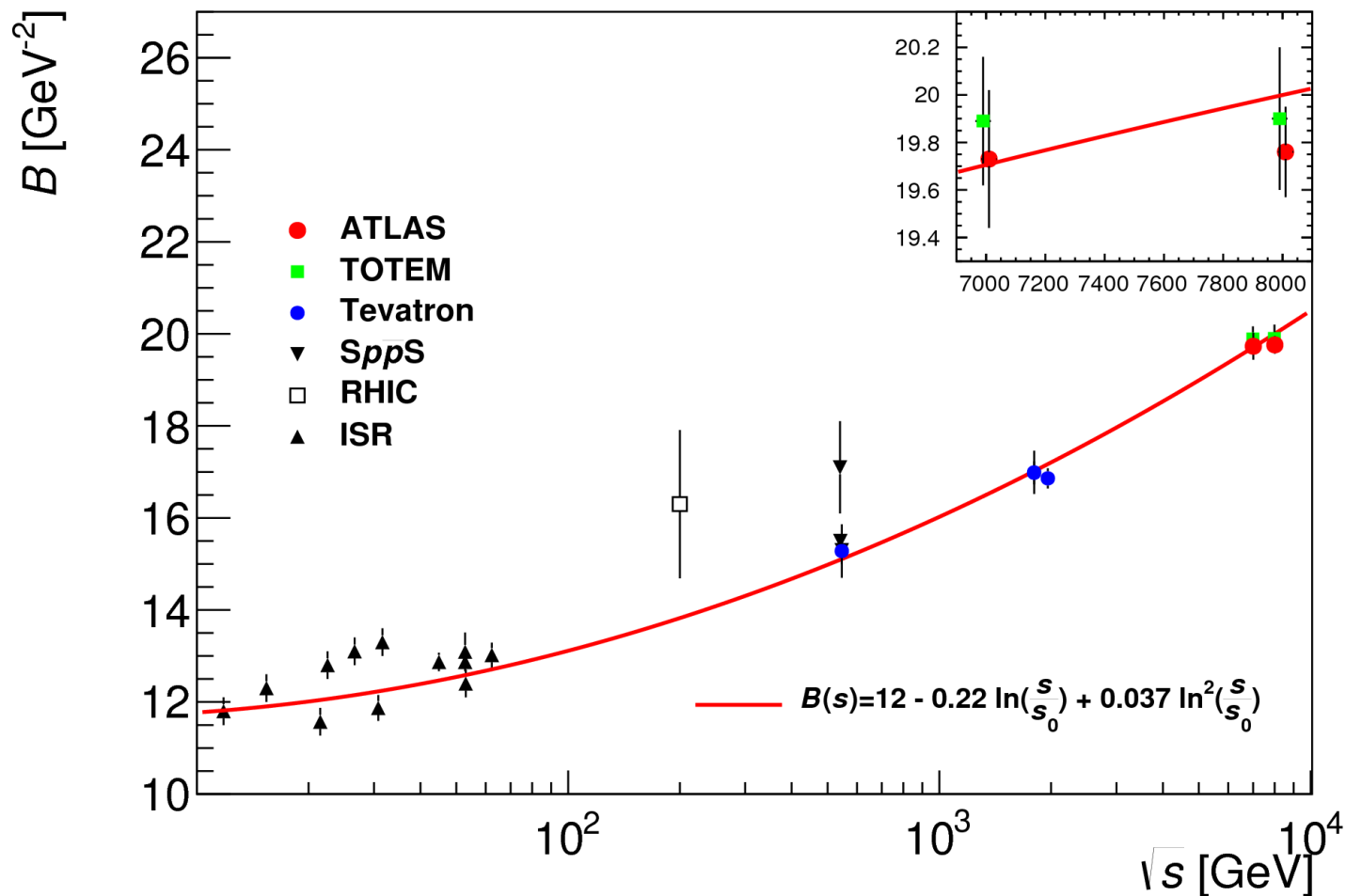


- ATLAS 1.9σ lower than TOTEM lumi-independent result (cf 1.3σ at 7TeV)
 - Both uncertainties dominated by luminosity

Data remain compatible with slow (log / power) growth with \sqrt{s}

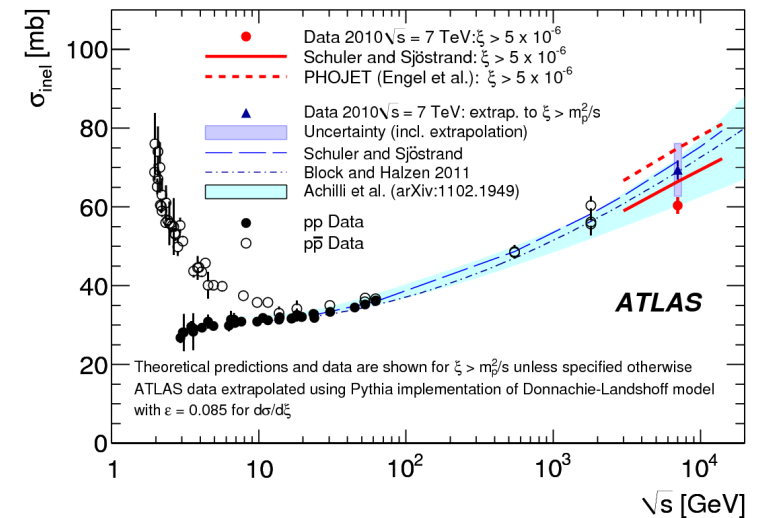
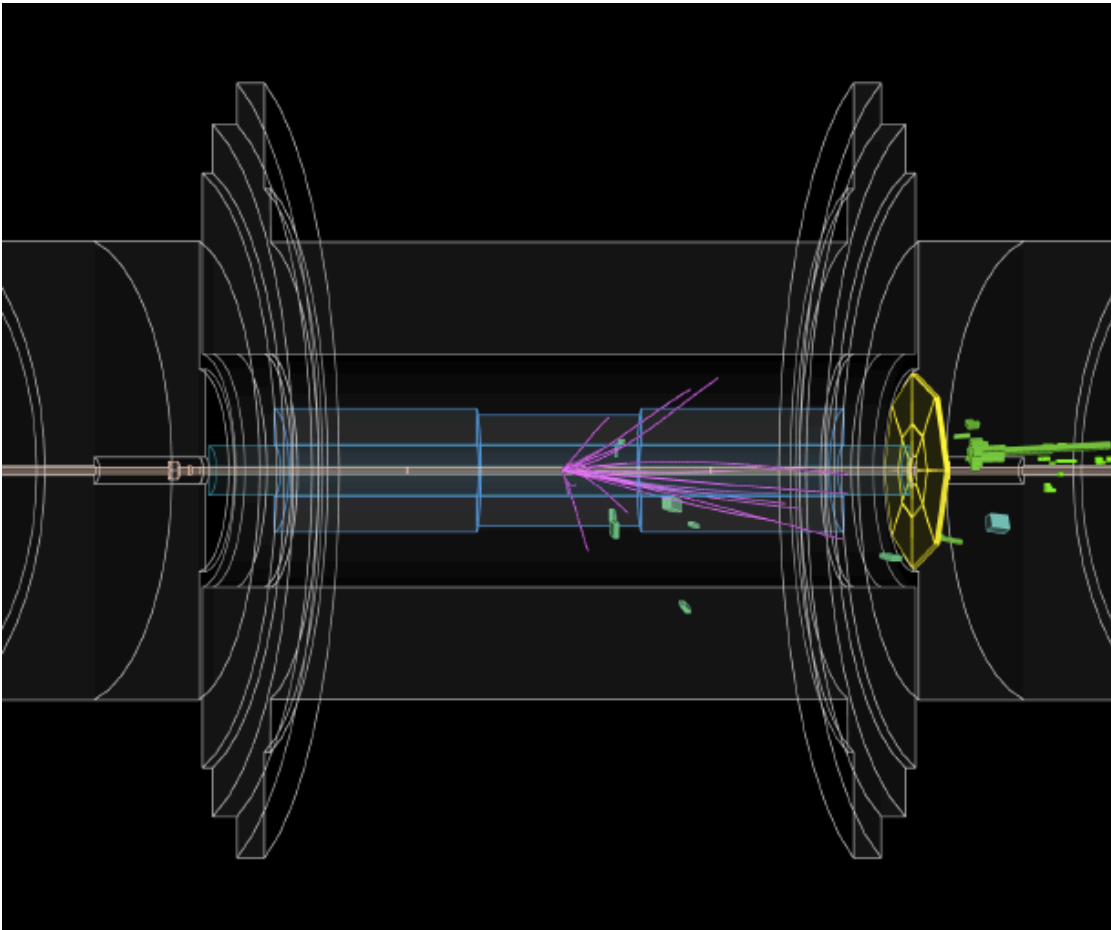
Results and Energy Dependence: B

$$B(8 \text{ TeV}) = 19.74 \pm 0.05(\text{stat.}) \pm 0.16(\text{exp.}) \pm 0.15(\text{extr.}) \text{ GeV}^{-2}$$



Data remain compatible with shrinkage of forward elastic peak with \sqrt{s} . Logarithmic model shown is Schegelsky & Ryskin

2) Minimum Bias Method



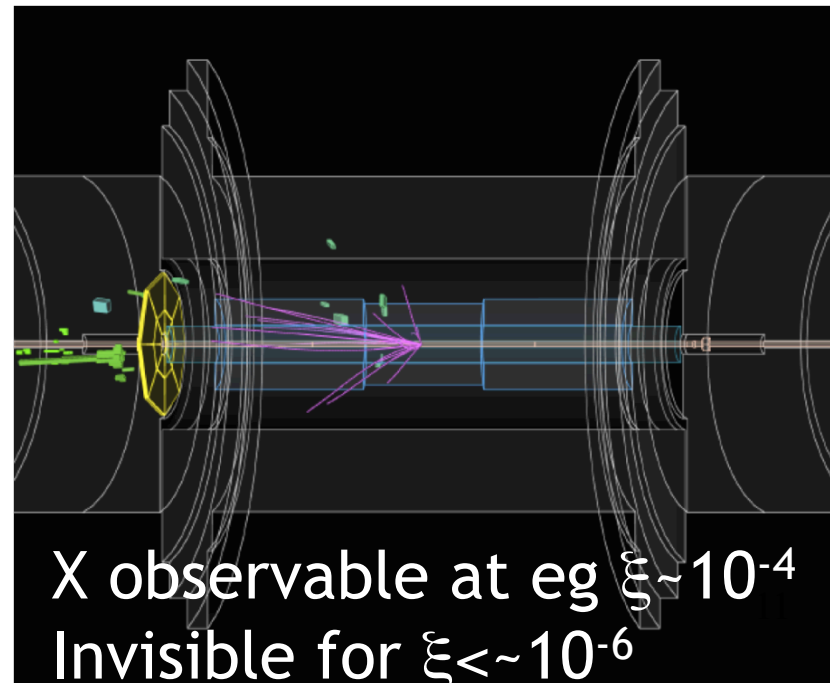
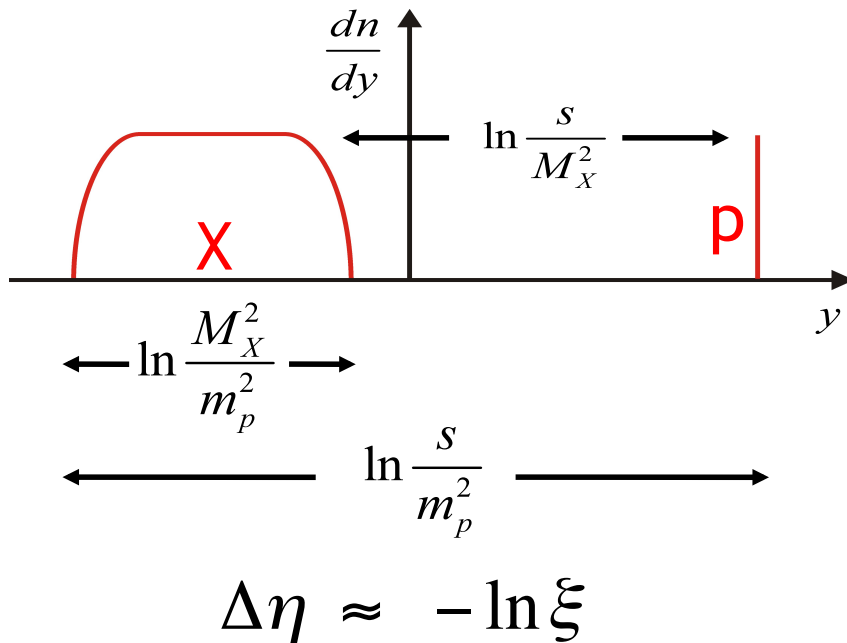
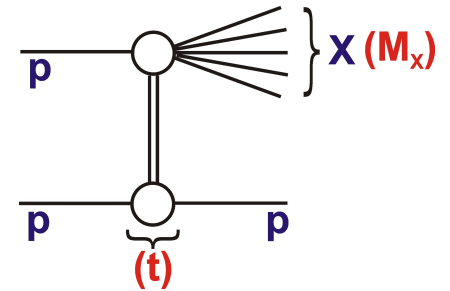
- Earlier result at $\sqrt{s}=7$ TeV: Nature Commun 2 (2011) 463
- Presented here, 13TeV result: PRL 117 (2016) 182002
(from short low pile-up run taken in June 2015)

Total Inelastic Cross Section

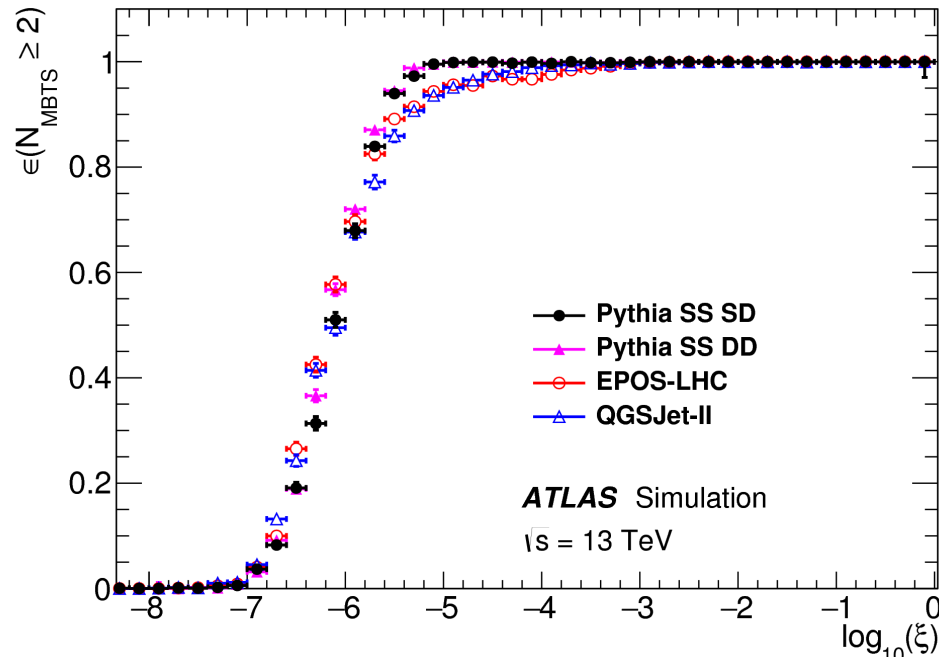
- MBTS sees 90-95% of all inelastic events \rightarrow “simple” counting experiment.

- Complication: controlling low mass diffractive dissociation that leaves no signal in MBTS

$$\xi = \frac{M_X^2}{s}$$



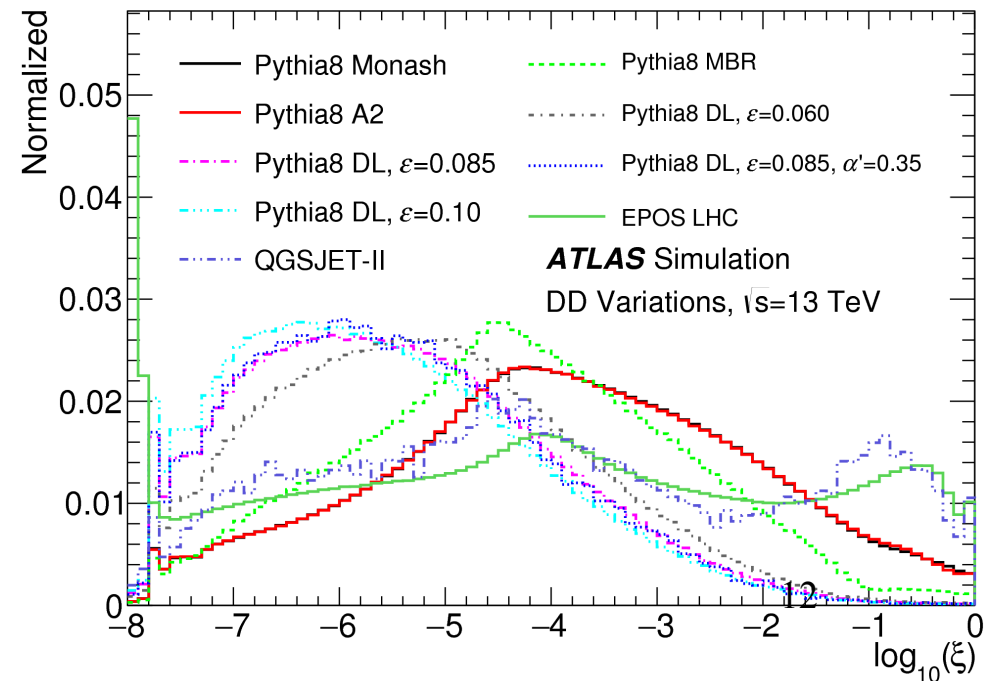
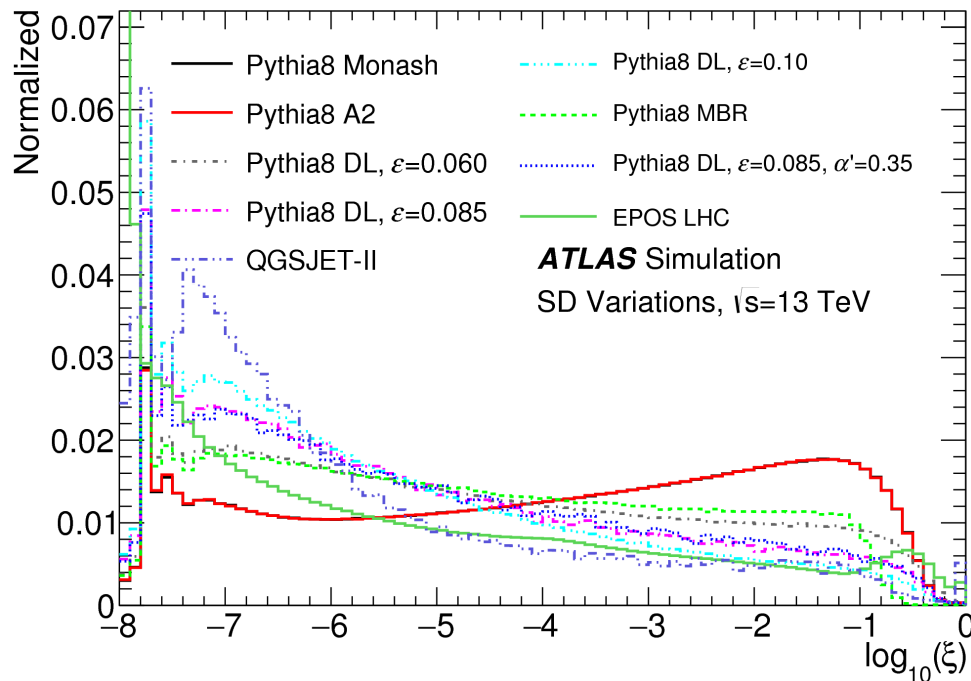
Why low mass diffraction is a problem



$$\Delta\eta \approx -\ln \xi$$

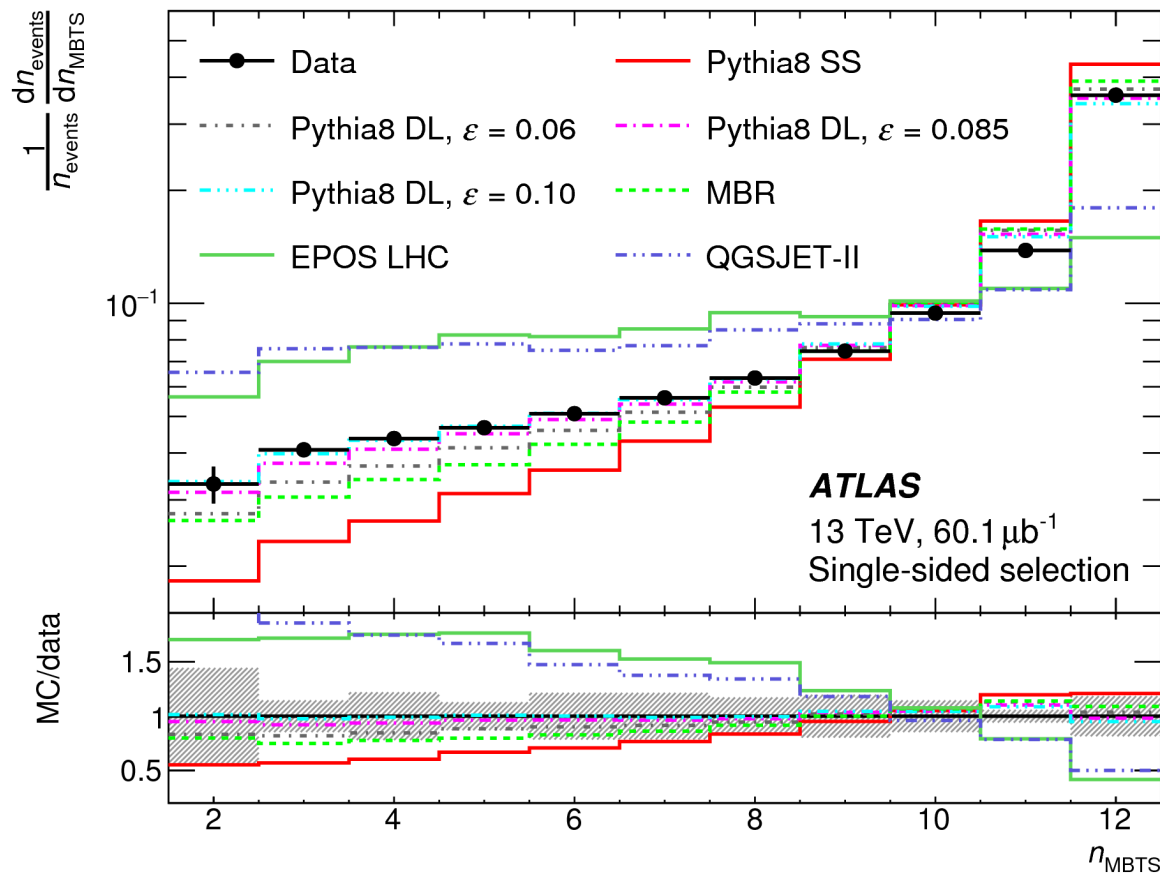
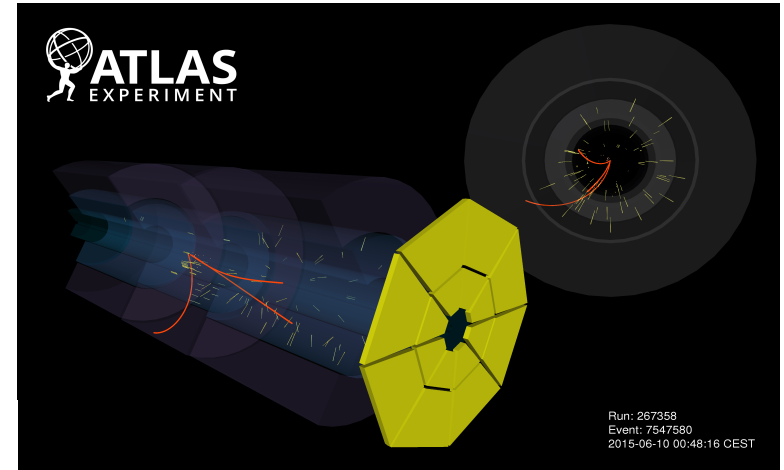
Acceptance limit of MBTS ($|\eta| < \sim 4$) corresponds to $\xi \sim 10^{-6}$

Cross section for $\xi < 10^{-6}$ poorly constrained \rightarrow fiducial region defined as $\xi > 10^{-6}$



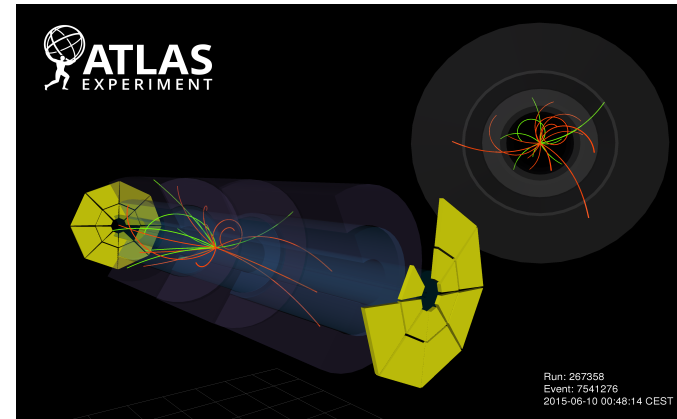
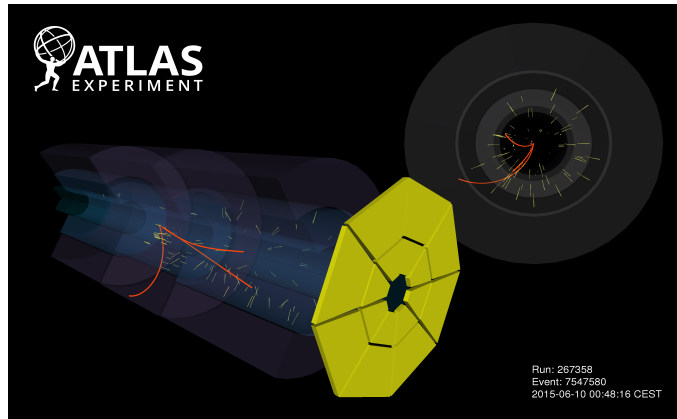
Benchmarking Diffractive MC models

“Single Sided” sample:
 ... activity on one side of MBTS,
 empty on other: enriched in SD events



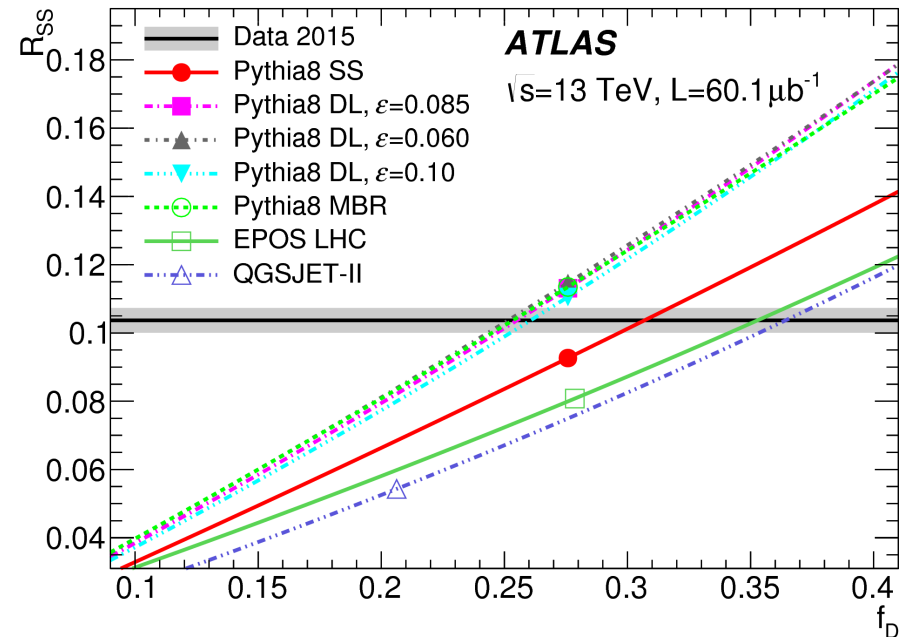
→ MBTS multiplicity in single sided sample distinguishes between MC diffraction models

Tuning Diffractive MC models



R_{SS} = Ratio of single sided to double sided MBTS samples ... used to tune fractions of events considered diffractive in each MC model

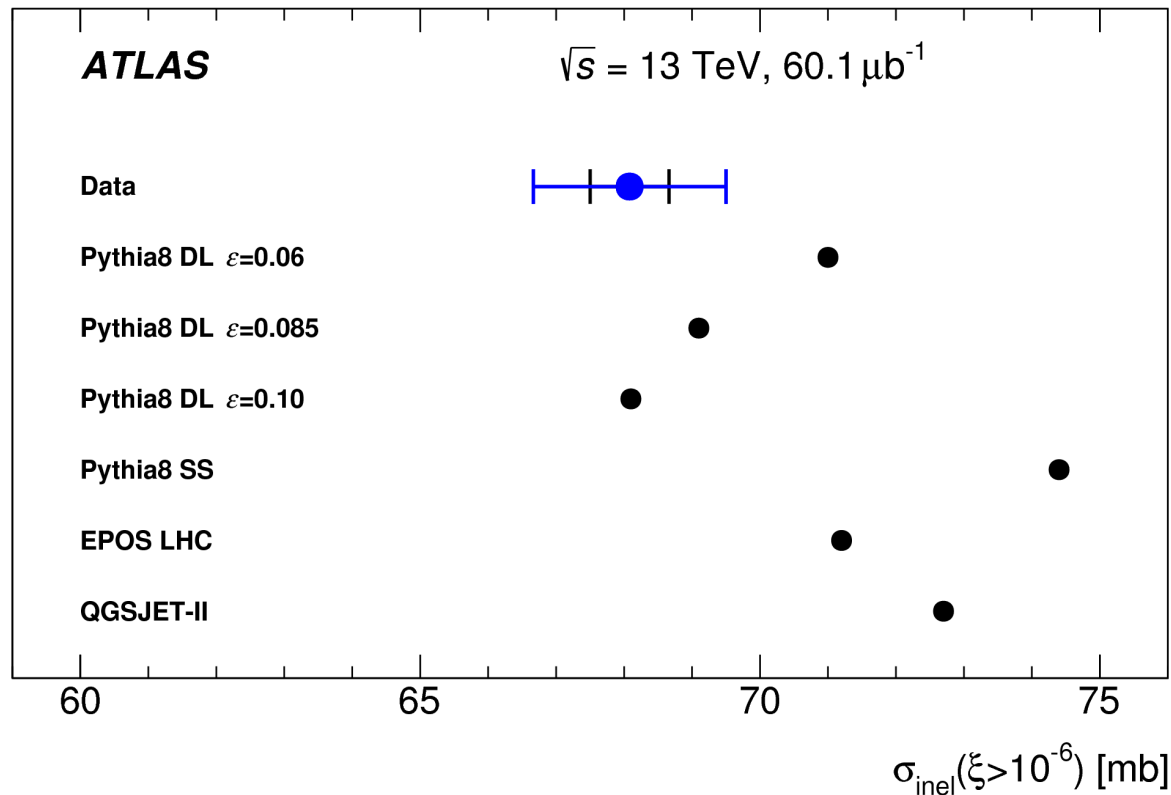
Baseline MC is PYTHIA8 with DL pomeron flux and $\alpha_{IP}(0) = 1.085$



- Uncertainty in $\xi > 10^{-6}$ fiducial region dominated by luminosity
- After extrapolation to full σ_{inel} , model uncertainty dominates

Cross Section in Fiducial Range

$$\sigma_{\text{inel}}^{\text{fid}}(13 \text{ TeV}) = 68.1 \pm 0.6 (\text{exp.}) \pm 1.3 (\text{lumi}) \text{ mb}$$



- Donnachie-Landshoff implementation in PYTHIA8 consistent with data within $\sim 2\sigma$ for $\alpha_{\text{IP}}(0) = 1.06 \dots 1.14$

- EPOS, QGSJET, PYTHIA8 S&S ($\alpha_{\text{IP}}(0) = 1$) exceed result by $> 2\sigma$

Extrapolation to Full Inelastic Cross Sec

Data-driven extrapolation into region with $\xi < 10^{-6}$, with minimal dependence on MC models:

$$\sigma_{\text{inel}} = \sigma_{\text{inel}}^{\text{fid}} + \sigma^{7 \text{ TeV}}(\xi < 5 \times 10^{-6}) \times \frac{\sigma^{\text{MC}}(\xi < 10^{-6})}{\sigma^{7 \text{ TeV, MC}}(\xi < 5 \times 10^{-6})}$$

$$\sigma^{7 \text{ TeV}}(\xi < 5 \times 10^{-6}) = \sigma_{\text{inel}}^{7 \text{ TeV}} - \sigma^{7 \text{ TeV}}(\xi > 5 \times 10^{-6}) = 11.0 \pm 2.3 \text{ mb}$$

from ALFA result

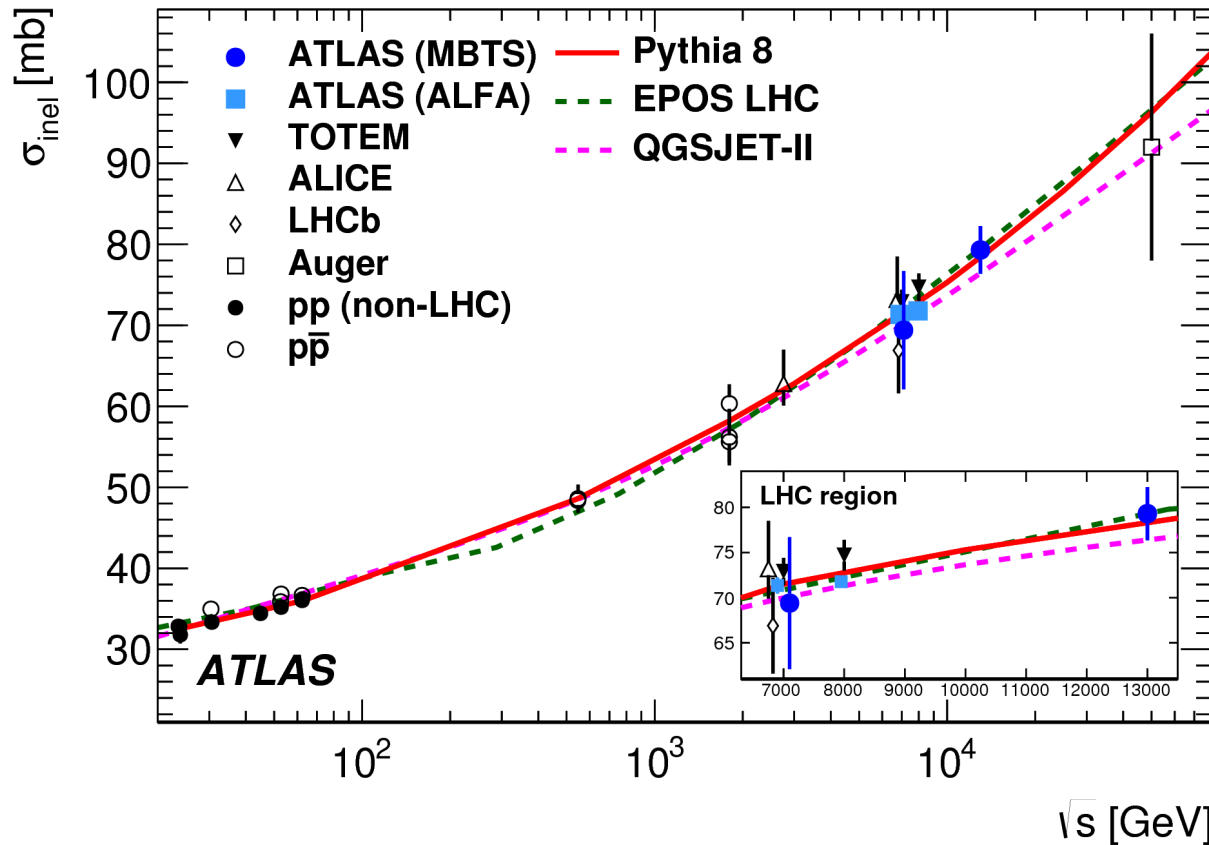
previous fiducial MBTS result

$$\frac{\sigma^{\text{MC}}(\xi < 10^{-6})}{\sigma^{7 \text{ TeV, MC}}(\xi < 5 \times 10^{-6})} = 1.015 \pm 0.081$$

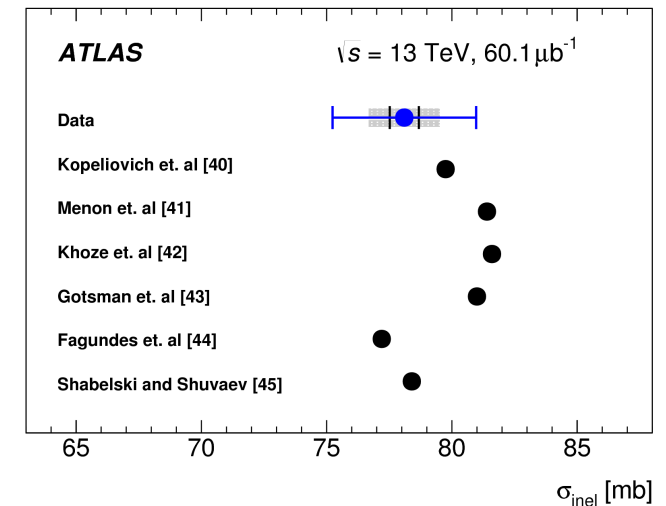
... extrapolation uncertainty is 2.5 mb

Extrapolation to Full Inelastic Cross Sec

$$\sigma_{inel} = 79.3 \pm 0.6(\text{exp}) \pm 1.3(\text{lum}) \pm 2.5(\text{extr}) \text{ mb}$$



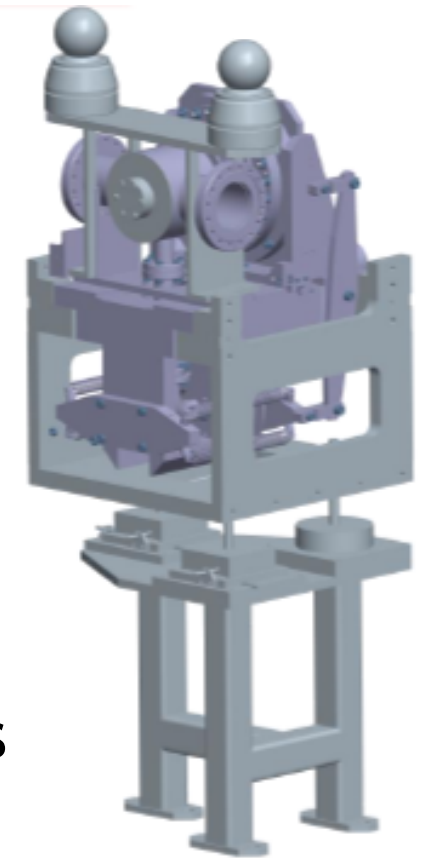
Consistent with continued growth with \sqrt{s}



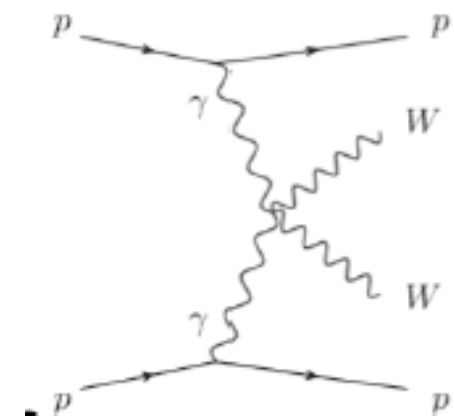
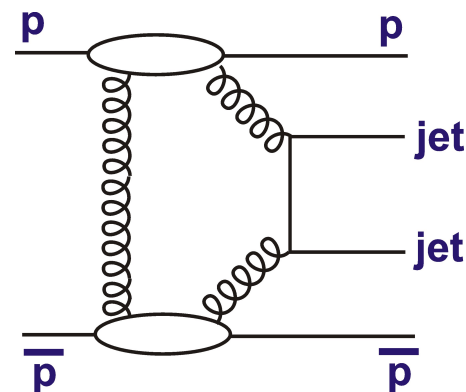
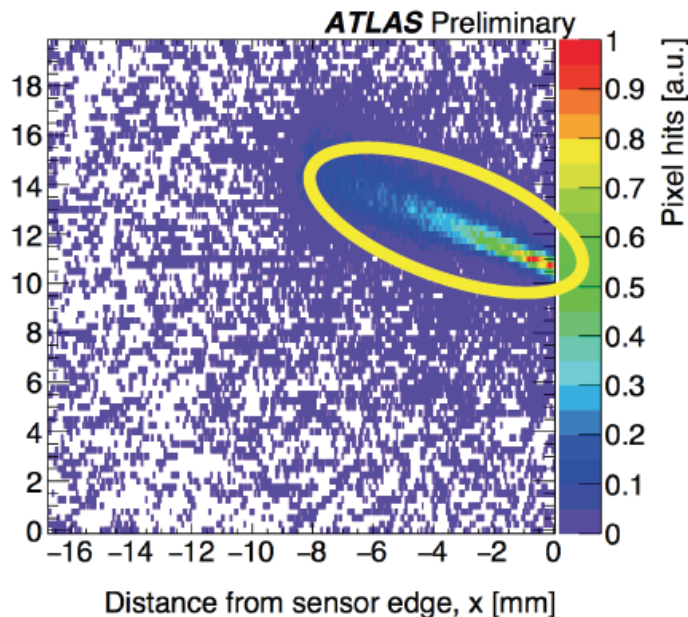
Within current uncertainties, result is consistent with indicative selection of models based on Regge phenomenology, eikonal approaches and other models of non-perturbative strong interactions

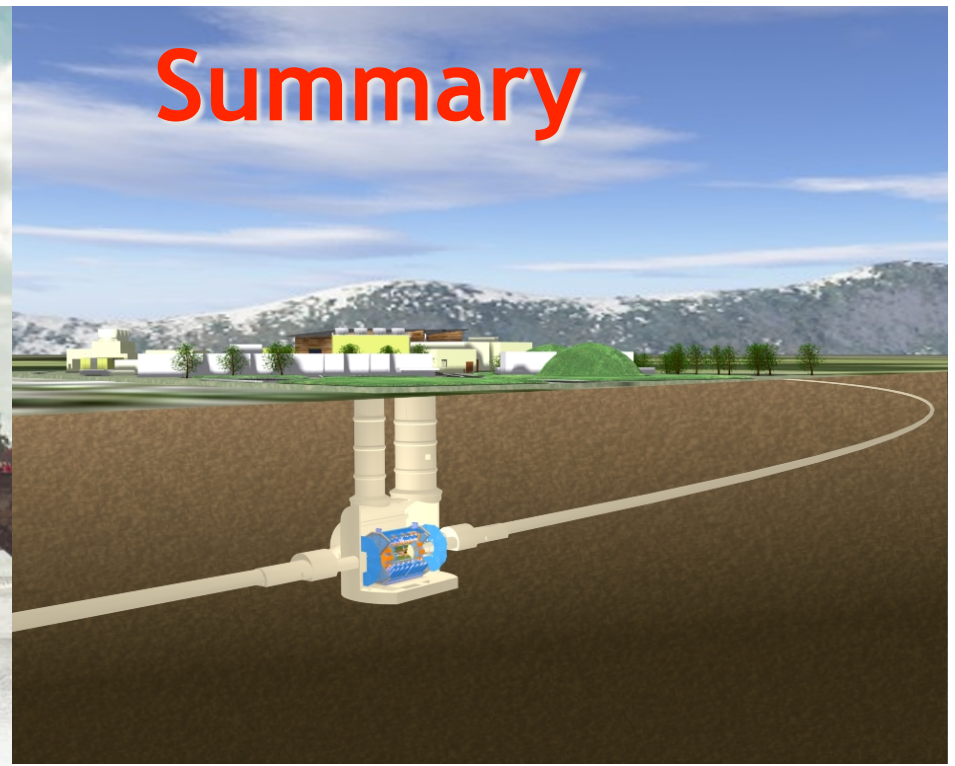
Future at ATLAS

- 13 TeV σ_{tot} ALFA measurement still ongoing
- Diffractive studies still ongoing
- ALFA being preserved for 14 TeV running
- Meanwhile, focus of Roman pots in ATLAS switches towards AFP & high lumi \rightarrow rare exclusive



(or exotic) processes ...





Summary

Proton-tagged elastic and total cross sections at $\sqrt{s} = 8$ TeV

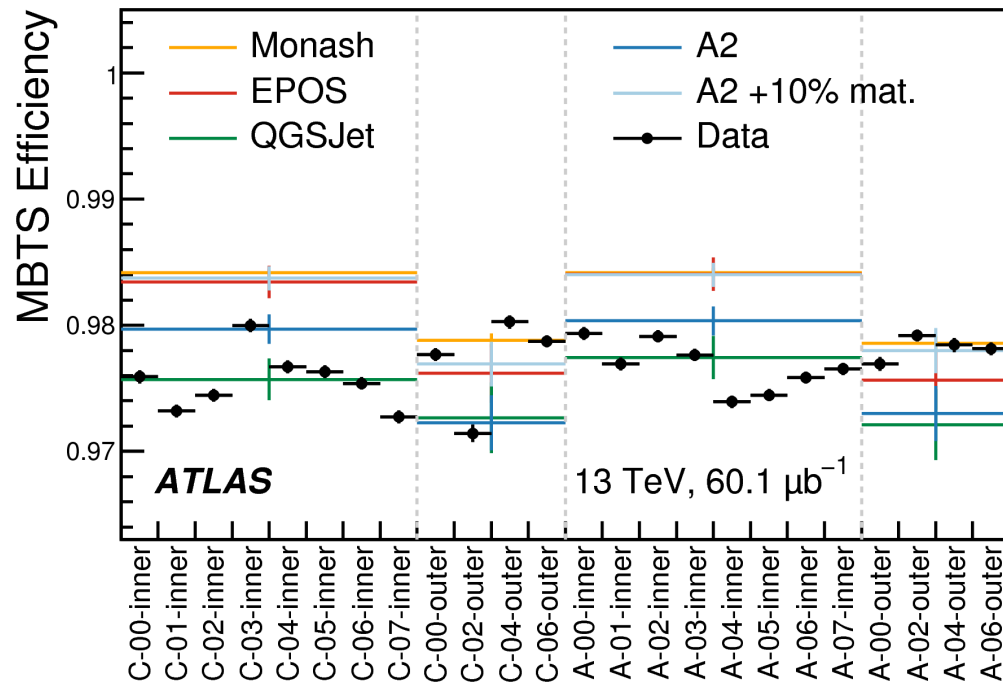
- Highly precise measurements of σ_{tot} ($\pm 1\%$), σ_{el} , σ_{inel} , B_{el}
- Continued cross sec growth and elastic peak shrinkage with \sqrt{s}
- Compatibility with TOTEM at 2σ level

Direct Inelastic Cross Section Measurement at $\sqrt{s} = 13$ TeV

- Improvement in precision ($\rightarrow 2\%$) over previous data.
- Some discrimination between models
- Consistent with ATLAS-ALFA extractions using optical theorem

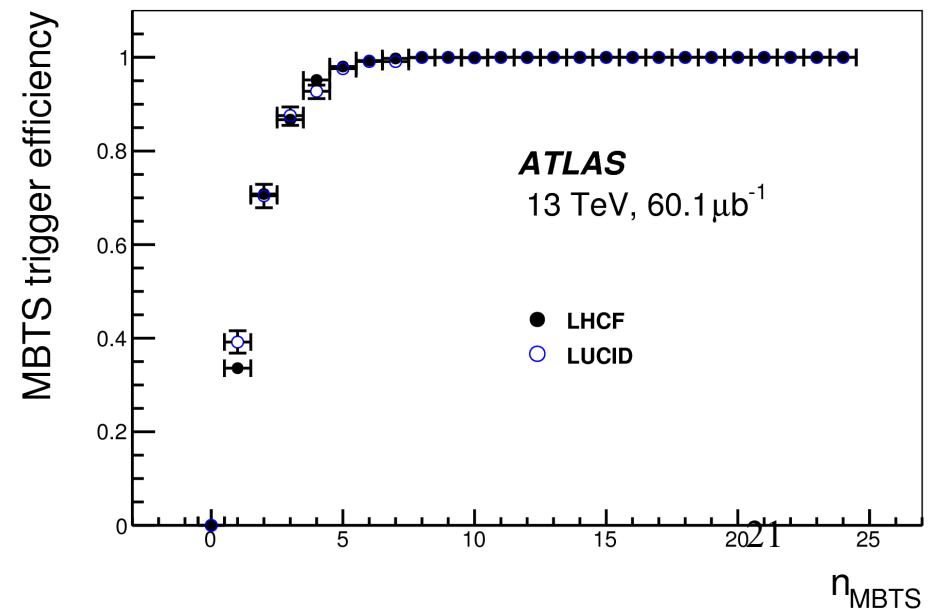
Back-ups / Working material

Controlling MBTS Efficiency with Data



- Efficiencies for each MBTS counter measured relative to tracks in inner detector where possible and calorimeter clusters where not.
 → MC efficiencies tuned accordingly

- Trigger efficiencies monitored relative to independent LUCID and LHCf triggers
- Efficiency / acceptance depends on MBTS segment multiplicity
- Analysis selection is $N_{\text{MBTS}} \geq 2$



Fiducial Cross Section Extraction

$$\sigma_{\text{inel}}^{\text{fid}}(\xi > 10^{-6}) = \frac{N - N_{\text{BG}}}{\epsilon_{\text{trig}} \times \mathcal{L}} \times \frac{1 - f_{\xi < 10^{-6}}}{\epsilon_{\text{sel}}} \quad \text{DROP???$$

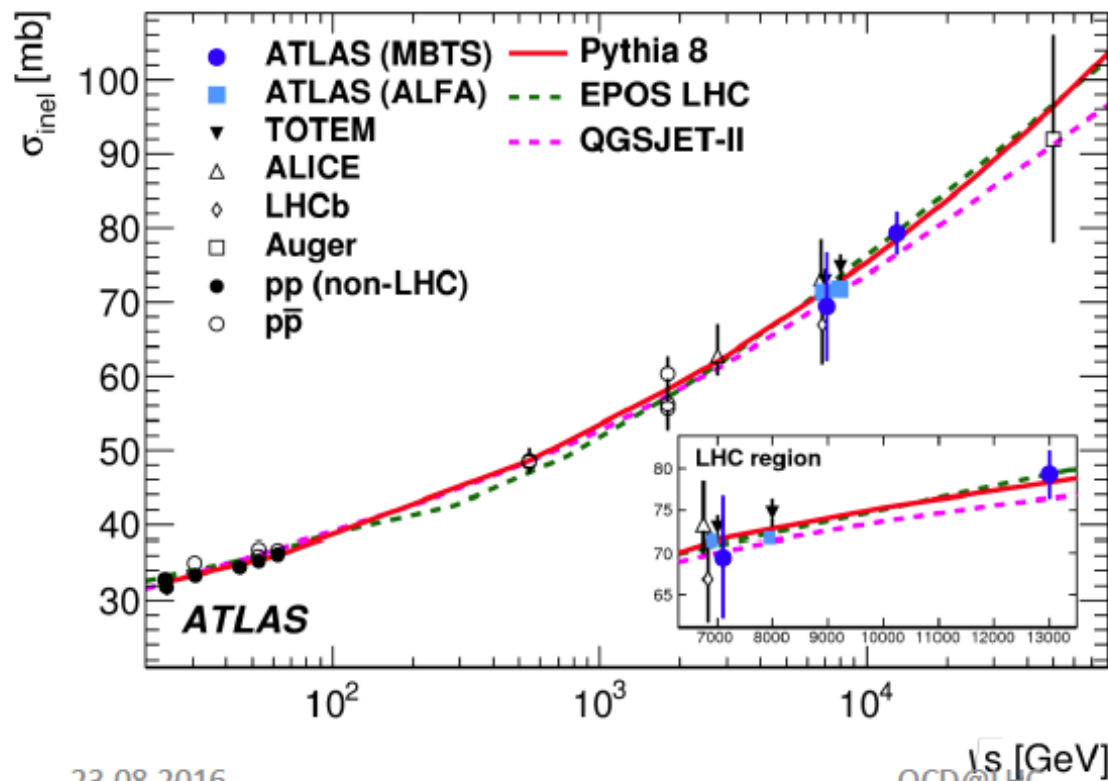
- N_{BG} : Small background from beam-gas, radiation & activation, determined using triggers in non-colliding bunches
- $(1 - f_{\xi < 10^{-6}}) / \epsilon_{\text{sel}} = C_{\text{MC}} = \text{MC acceptance and migration correction}$
- Luminosity from final calibration of Van der Meer scan
 $\rightarrow 1.9\%$ error

Factor	Value	Rel. uncertainty
Number of events passing the inclusive selection (N)	4159074	—
Number of background events (N_{BG})	51187	$\pm 50\%$
Integrated luminosity [μb^{-1}] (\mathcal{L})	60.1	$\pm 1.9\%$
Trigger efficiency (ϵ_{trig})	99.7%	$\pm 0.3\%$
MC correction factor (C_{MC})	99.3%	$\pm 0.5\%$

$$\sigma_{\text{inel}}^{\text{fid}}(13 \text{ TeV}) = 68.1 \pm 0.6(\text{exp.}) \pm 1.3(\text{lumi}) \text{ mb}^{22}$$

Summary of Total Inelastic Cross Sections

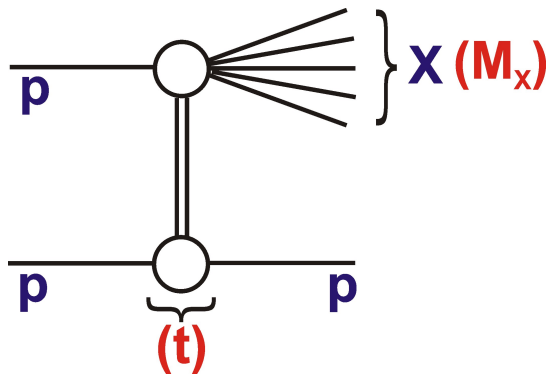
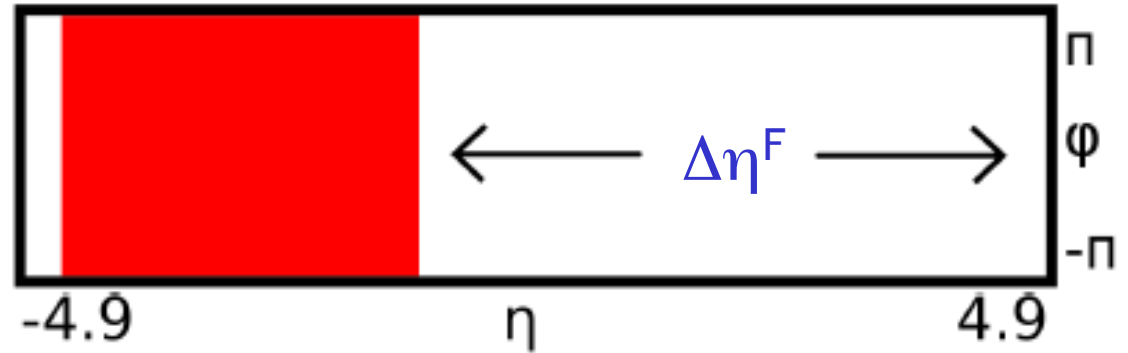
	7 TeV	8 TeV	13 TeV	Comments
MBTS	69.1 ± 7.3 mb		79.3 ± 2.9 mb	Main error contribution extrapolation
ALFA	71.3 ± 0.9 mb	71.7 ± 0.7 mb		Small errors due to precise lumin.
TOTEM	72.9 ± 1.5 mb	74.7 ± 1.7 mb		Based on elastic & inelastic rates
CMS			71.3 ± 3.5 mb	Preliminary, based on HF calorimeters



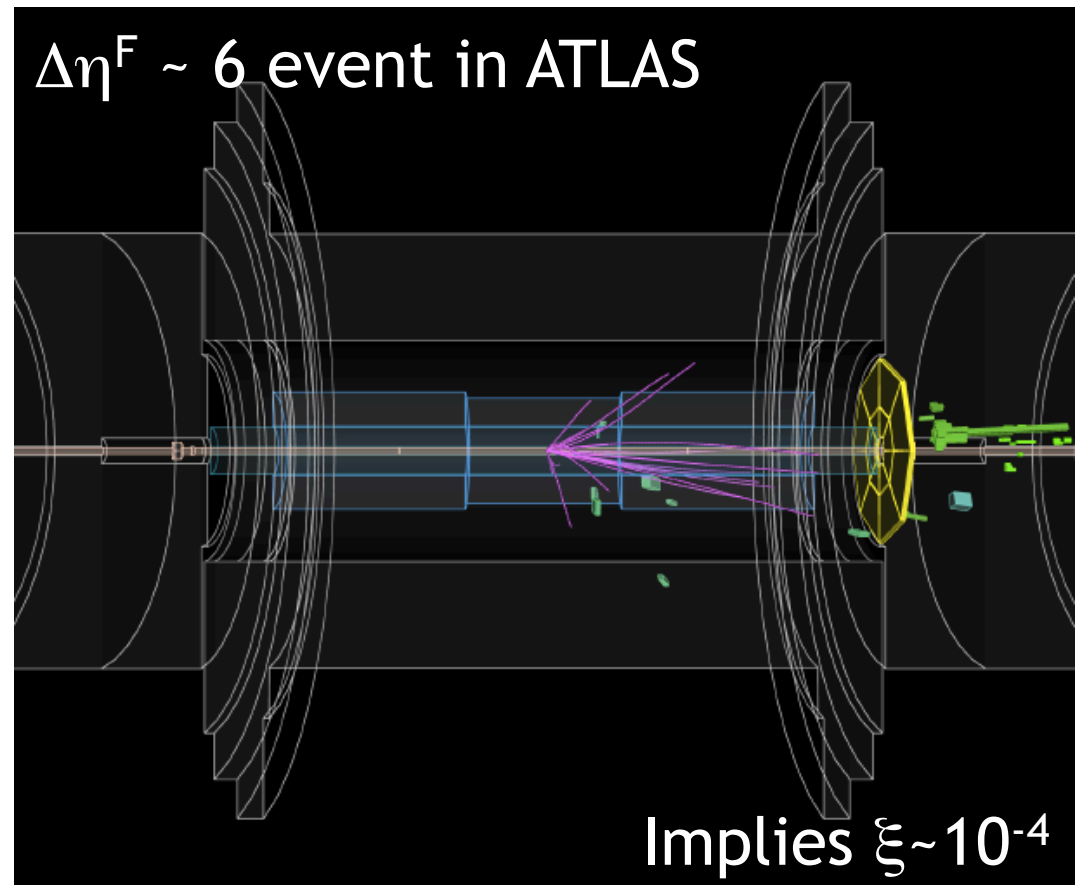
- In general increase with \sqrt{s} visible as expected.
- Values at 7 TeV and 8 TeV agree within errors.
- At 13 TeV some discrepancy, but large errors.

Rapidity gap cross-sections

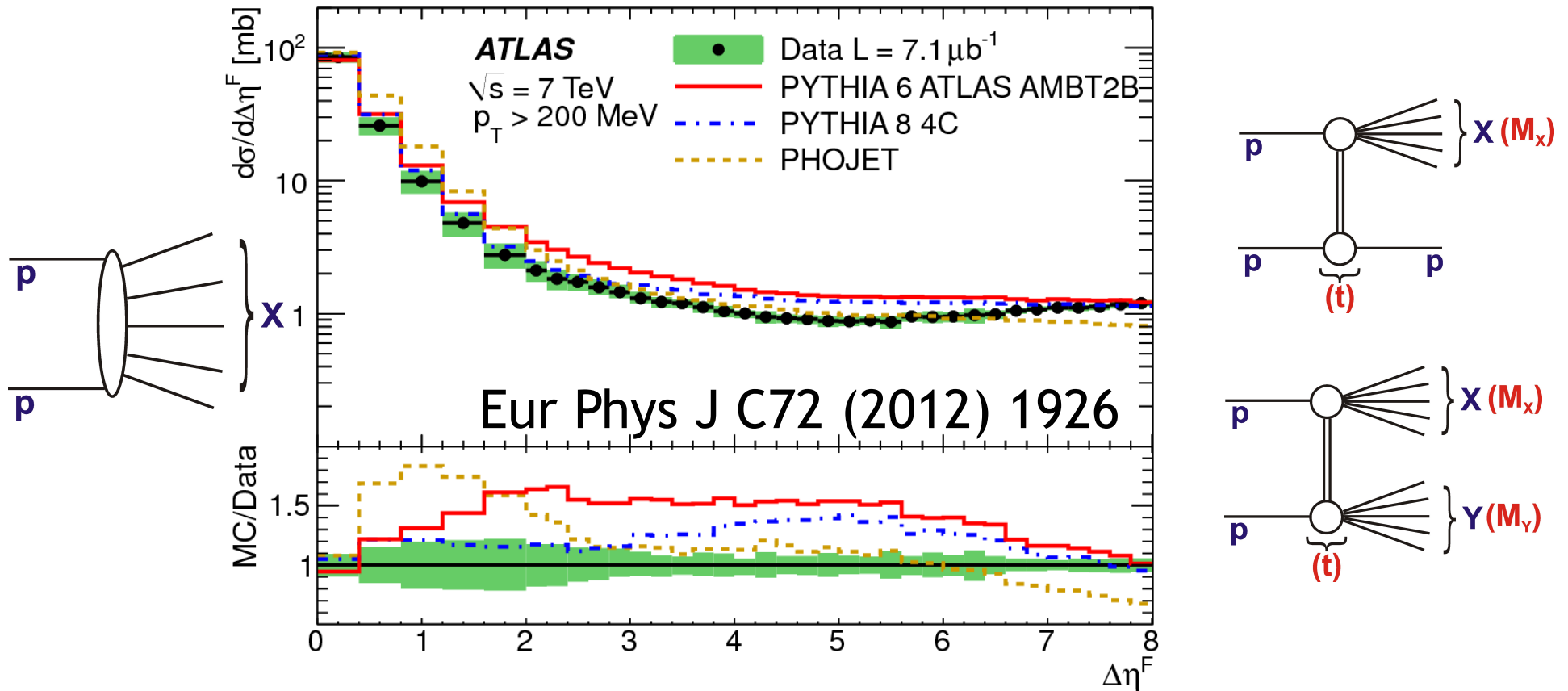
Method developed by ATLAS to measure hadron Level cross section as a function of $\Delta\eta^F$: forward or backward rapidity gap extending to limit of instrumented range: i.e. including $\eta = \pm 4.9$



... no statement on $\eta > \sim 4.9$
 ... large $\Delta\eta^F$ sensitive to SD + low M_Y DD



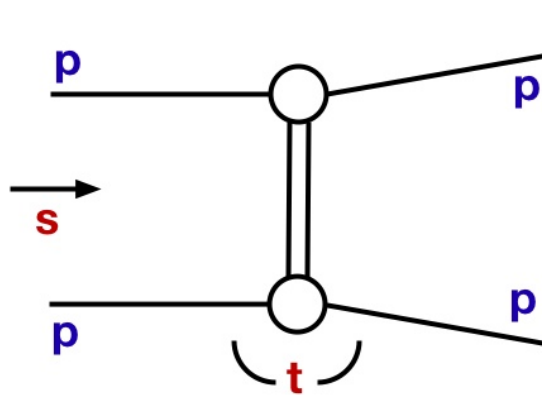
Inclusive Differential Gap Cross Section



- Large $\Delta\eta^F$: Diffractive plateau with $\sim 1\text{mb}$ per unit gap size, consistent with soft pomeron ($\alpha_{\text{IP}}(0) = 1.058 \pm 0.036$)
- Small $\Delta\eta^F$: sensitive to hadronisation fluctuations / MPI in ND

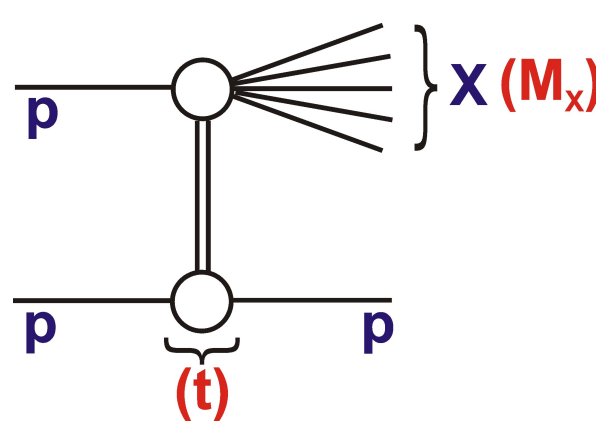
Can the same method be applied to hard diffractive processes?...

Decomposing the pp Cross Section



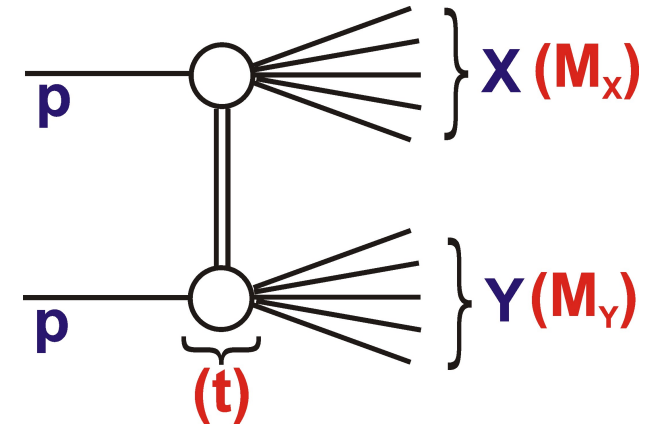
Elastic

1 degree of freedom
 → scattering angle / t



Single diffractive dissociation

Also M_X

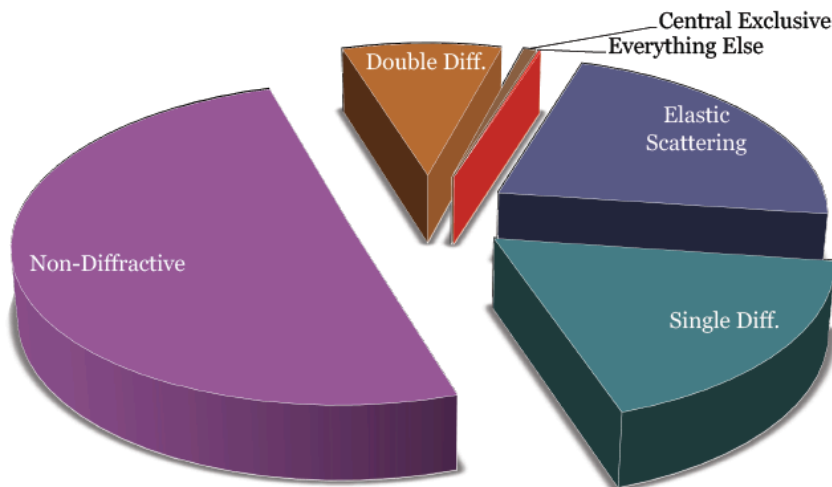


Double diffractive dissociation

Also M_Y

$$\xi = \frac{M_X^2}{s} = 1 - \frac{E'_p}{E_p}$$

$$\xi_Y = \frac{M_Y^2}{s}$$



At LHC, M_X, M_Y can be as large as 1 TeV → plenty of phase space to produce jets and other hard probes

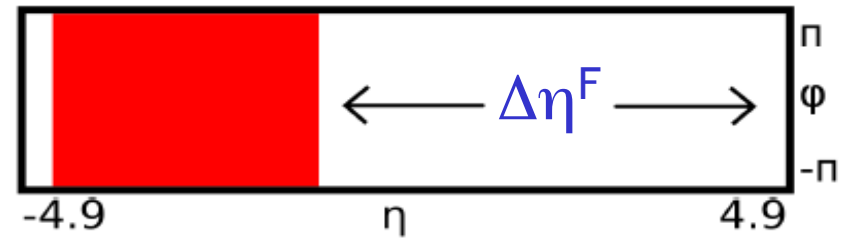
Probably for the back-up

$$\frac{d\sigma_{el}}{dt_i} = \frac{1}{\Delta t_i} \times \frac{\mathcal{M}^{-1} [N_i - B_i]}{A_i \times \epsilon^{\text{reco}} \times \epsilon^{\text{trig}} \times \epsilon^{\text{DAQ}} \times L_{\text{int}}}$$

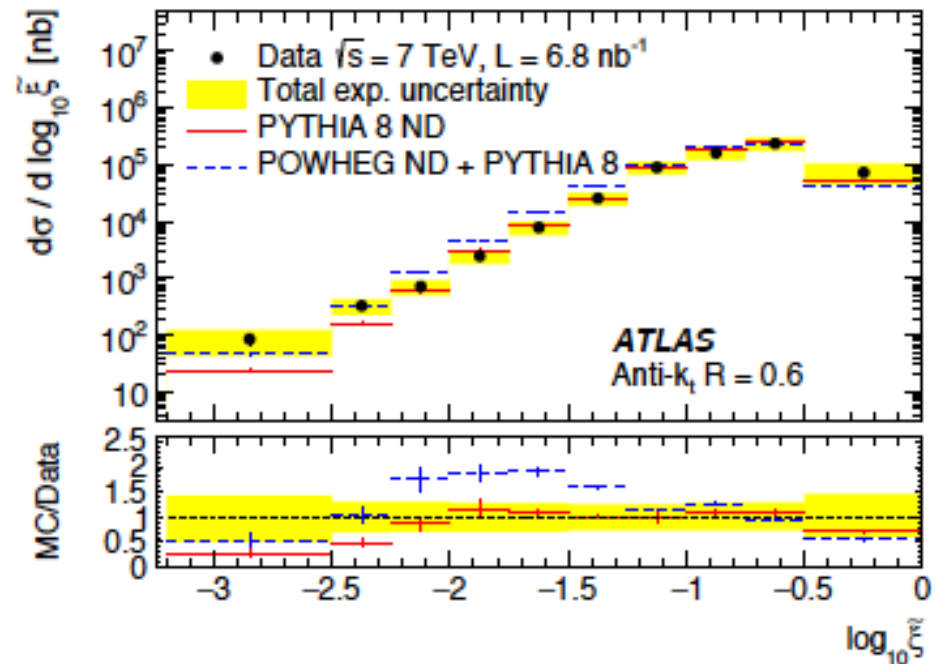
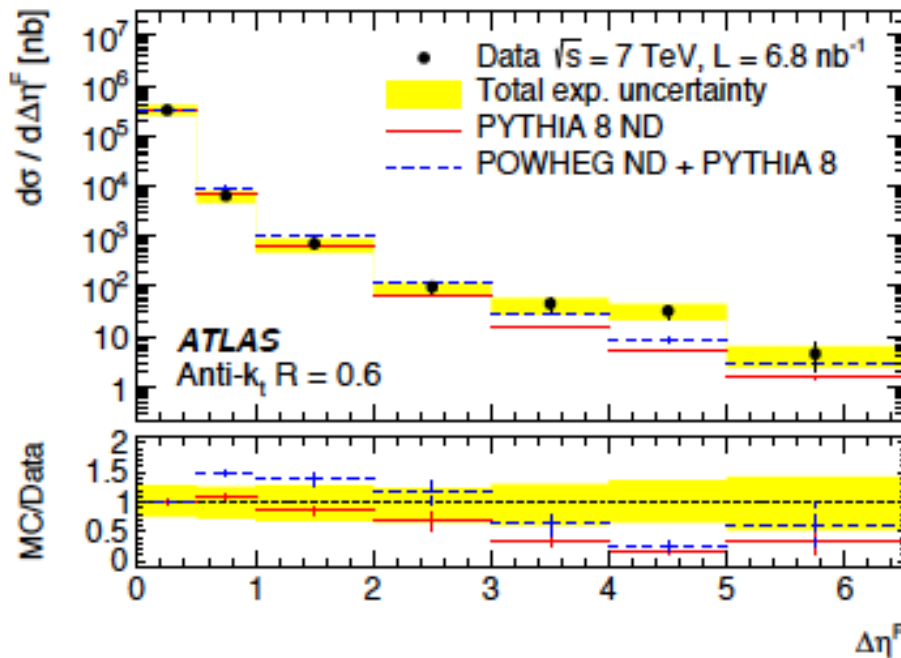
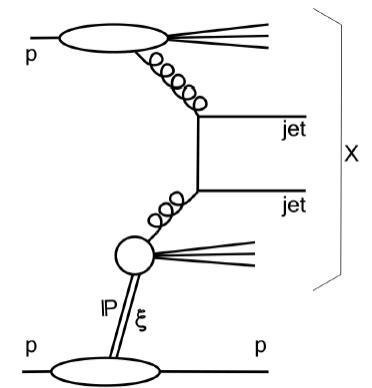
The diagram illustrates the components of the differential cross-section equation and their associated experimental parameters:

- Unfolding procedure** (black box) points to the \mathcal{M}^{-1} term.
- Events per bin** (green box) points to the N_i term.
- Background (0.12%)** (red box) points to the B_i term.
- Acceptance (fit range > 10%)** (blue box) points to the A_i term.
- Reconstruction efficiency ($\approx 90\%$)** (purple box) points to the ϵ^{reco} term.
- Trigger efficiency (> 99.9%)** (yellow box) points to the ϵ^{trig} term.
- Dead-time fraction (0.44%)** (red box) points to the ϵ^{DAQ} term.
- Luminosity 1.5% error** (dark blue box) points to the L_{int} term.

Corrected Diffractive Jet Data v Non-Diffractive Models



- Kinematic suppression of large gaps \rightarrow no clear diffractive plateau (unlike minimum bias case)
- ND models matched to data at small gap sizes give contributions compatible with data up to largest $\Delta\eta^F$ and smallest ξ ... **no clear diffractive signal ...**



Evidence for Diffractive Contribution

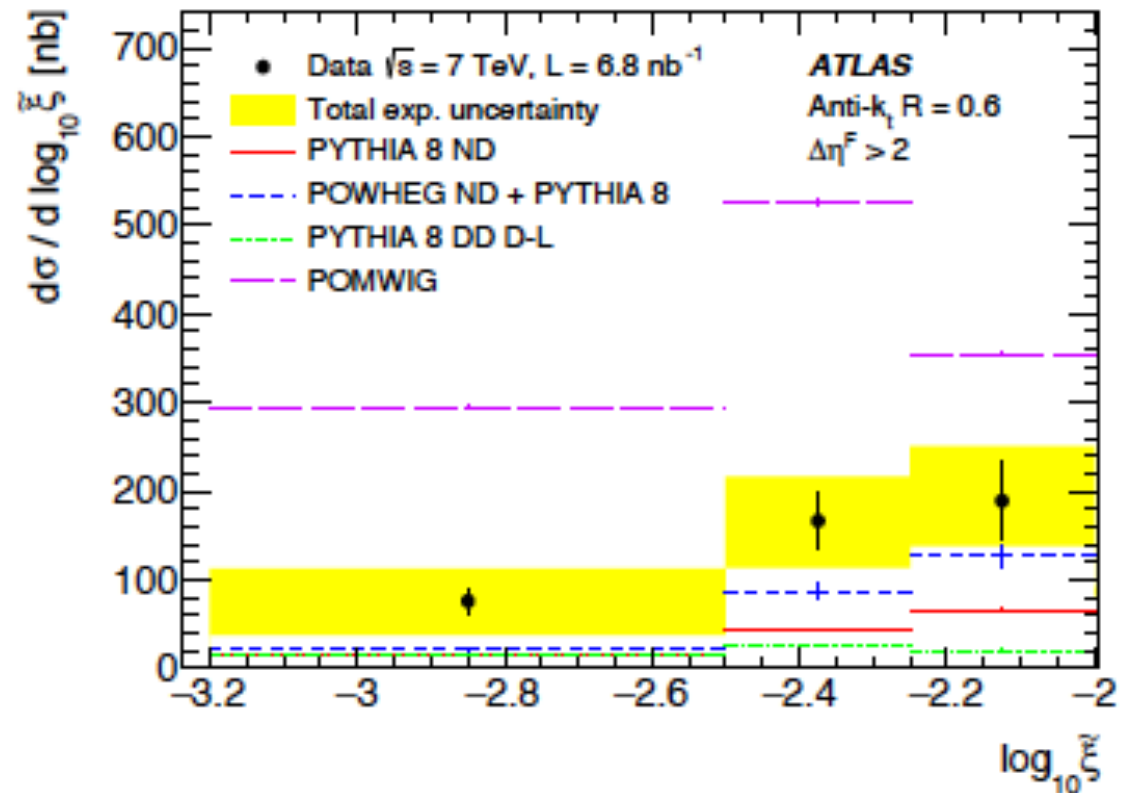
Focusing on small ξ , whilst simultaneously requiring large gap size ($\Delta\eta_F > 2$) gives best sensitivity to diffractive component

→ Models with no SD jets are below data by factor $> \sim 3$

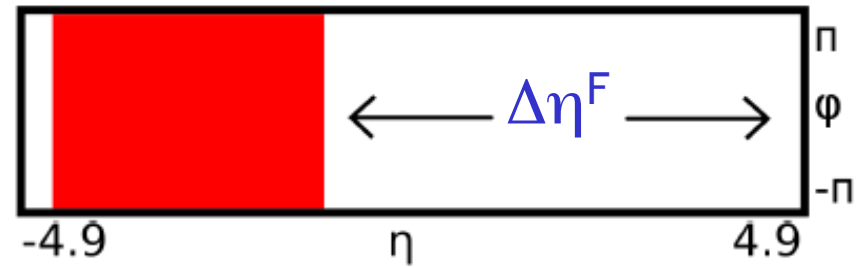
→ Comparison of smallest ξ with DPDF-based model (POMWIG) leads to rapidity gap survival probability estimate ...

- Model dependence not investigated in detail
- In context of POMWIG, using anti- k_T with $R=0.6$:

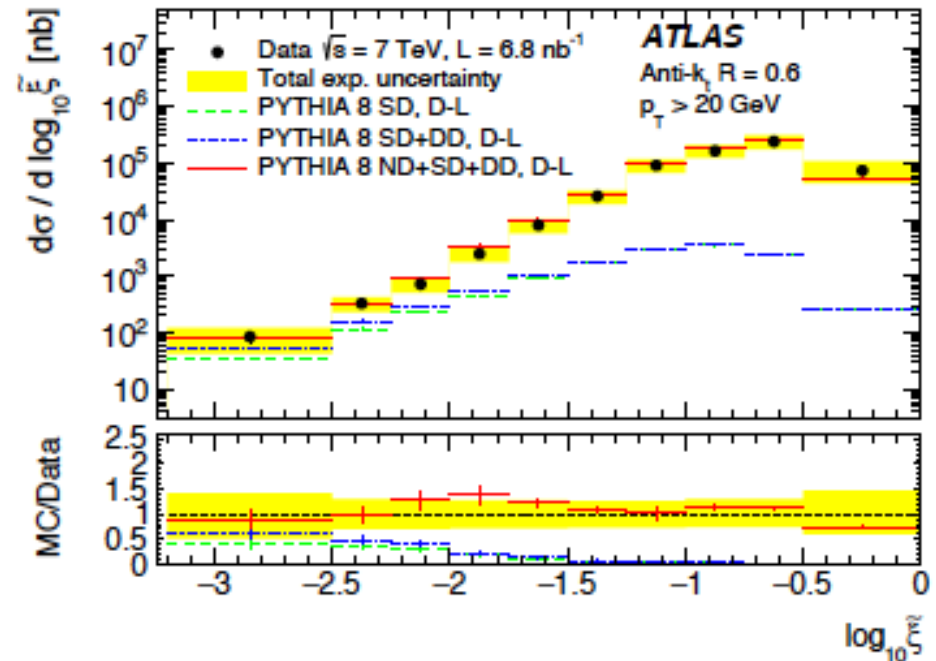
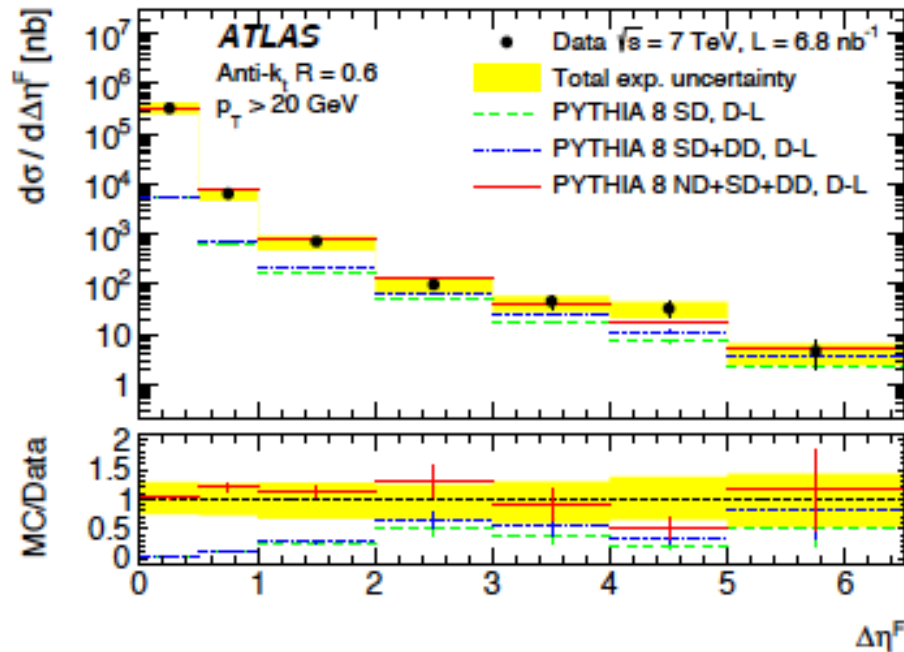
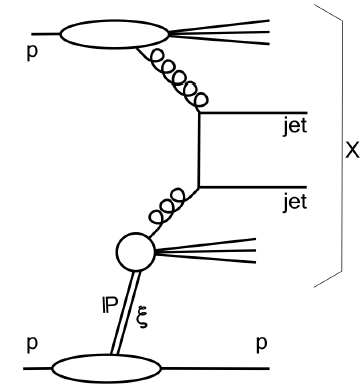
$$S^2 = 0.16 \pm 0.04 \text{ (stat.)} \pm 0.08 \text{ (exp. syst.)} ,$$



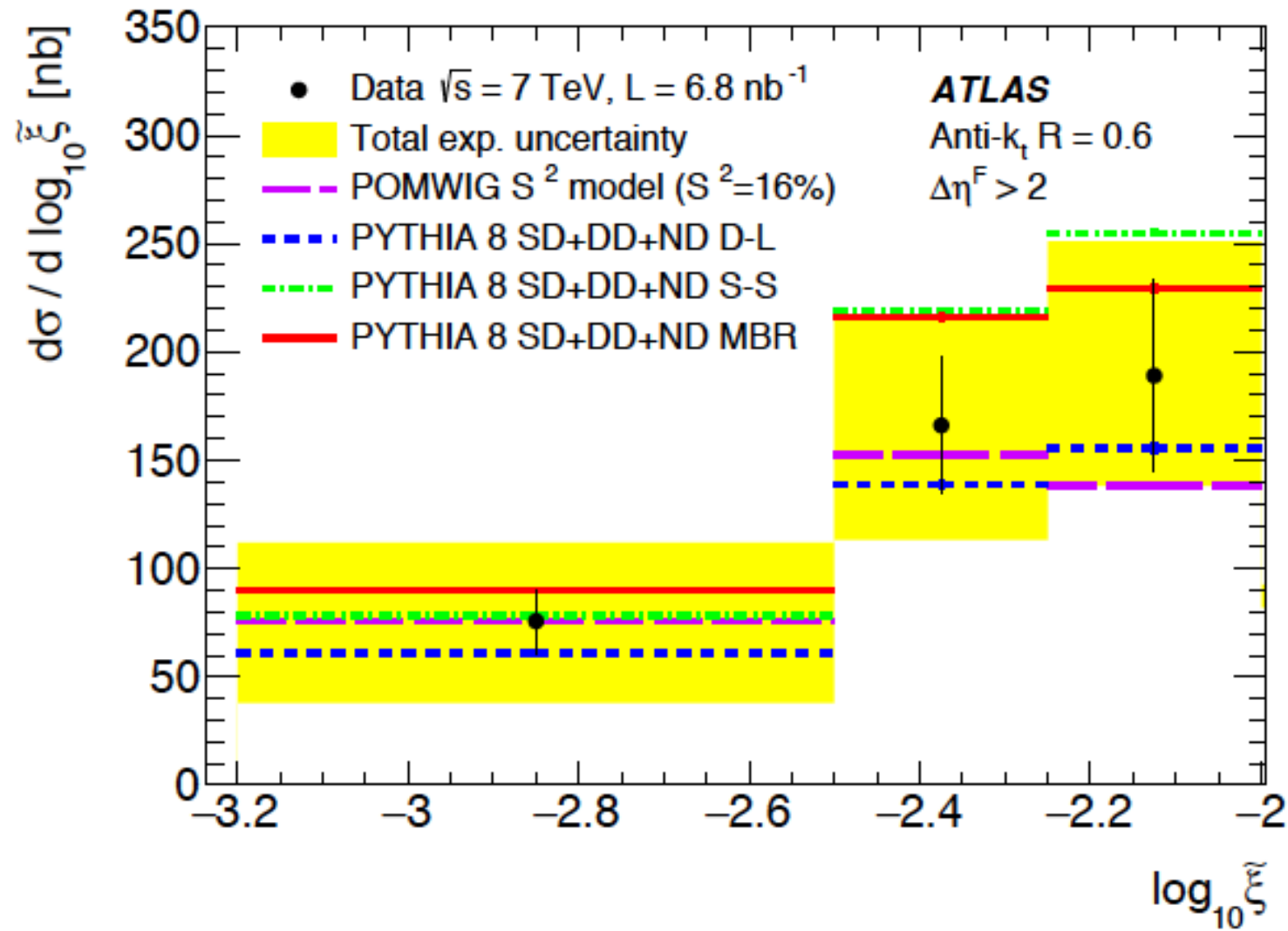
Comparison with Full PYTHIA8



‘Off-the shelf’ PYTHIA8 ND*0.71 + SD + DD does a good job at all $\Delta\eta_F$ and ξ , with no need for a gap survival factor (though ND dominates, so compatible with a wide range of S^2 values).

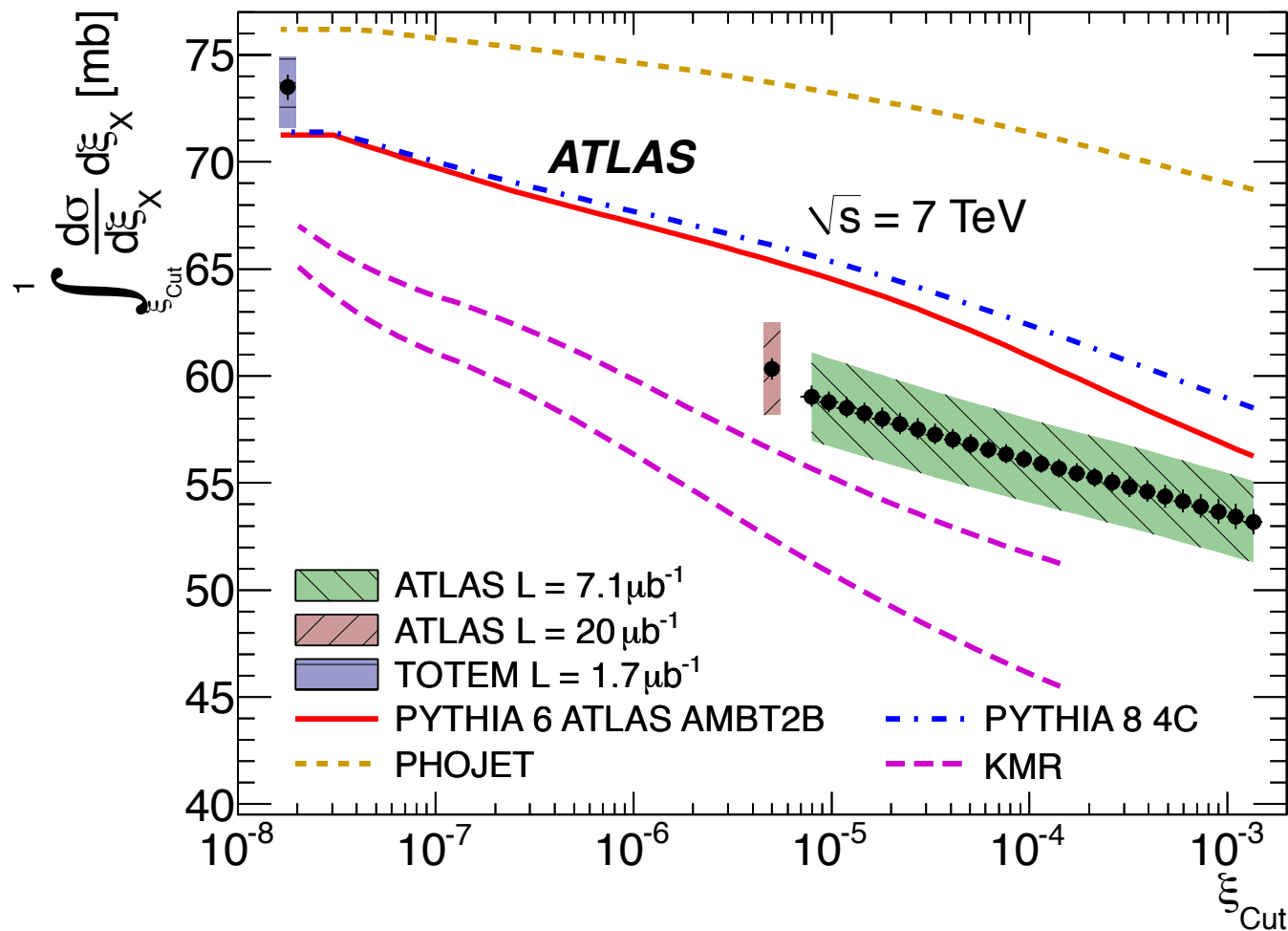


Diffractive Models Focusing on $\Delta\eta^F > 2$



Phys Lett B754 (2016), 214

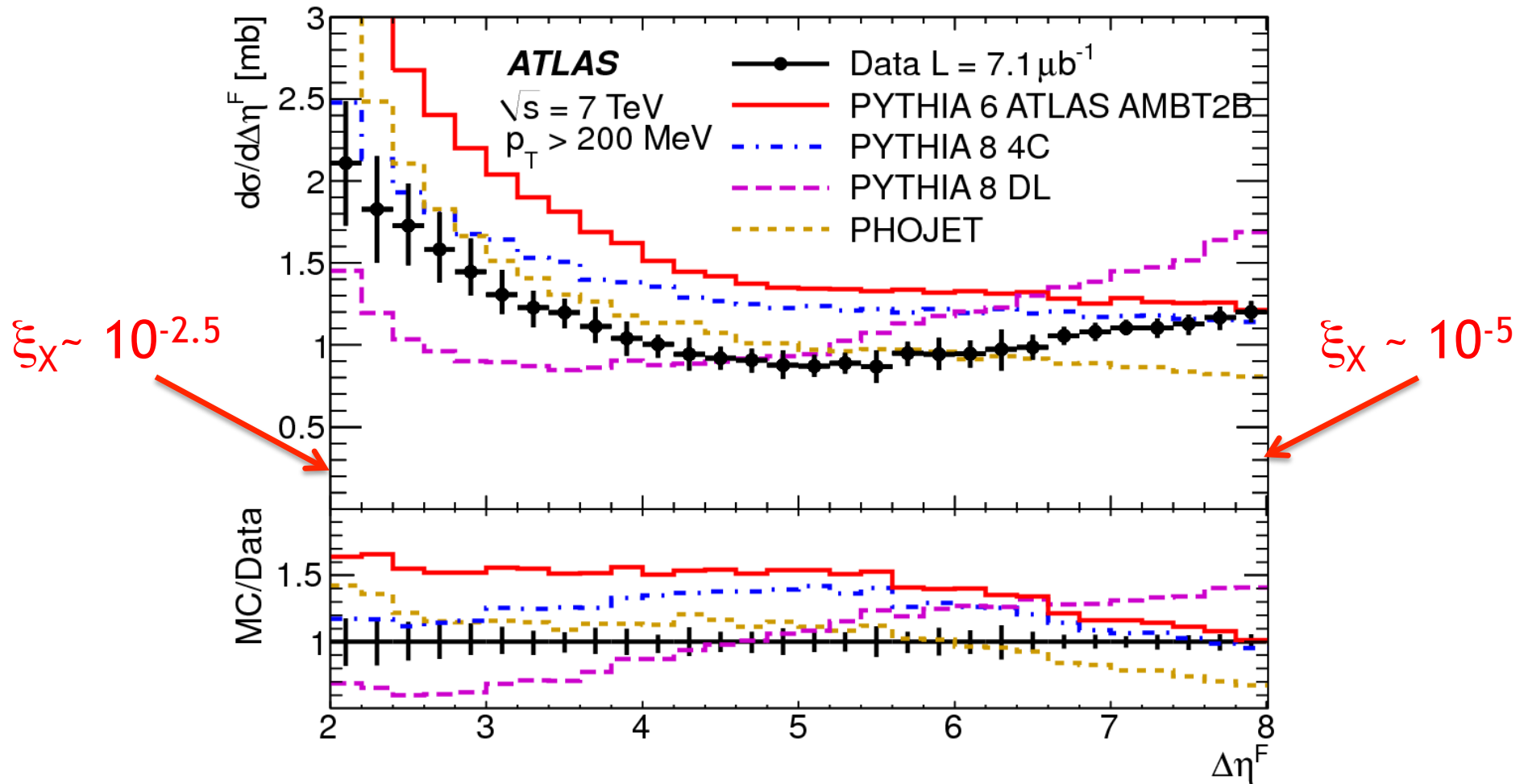
Investigating Low Mass Extrapolations



[Inelastic cross section excluding diffractive channels with $\xi < \xi_{\text{cut}}$]

- Integrating ATLAS gap cross section up to some max $\Delta\eta^F$ (equivalently min ξ_x) and comparing with TOTEM indicates that small ξ_x region underestimated in PHOJET and PYTHIA:
- 14 mb with $\xi < 10^{-5}$, compared to 6 (3) mb in PYTHIA (PHOJET)

Large Gaps and Diffractive Dynamics



- Default PHOJET, PYTHIA have $\alpha_{\text{IP}}(0) = 1$; DL has $\alpha_{\text{IP}}(0) = 1.085$
- Fit to large $\Delta\eta^F$ data: $\alpha_{\text{IP}}(0) = 1.058 \pm 0.003$ (stat) ± 0.036 (syst)
- CMS also favour intermediate value of $\alpha_{\text{IP}}(0)$

Large Gaps and Diffractive Dynamics

- Diffractive plateau with ~ 1 mb per unit of gap size for $\Delta\eta^F > 3$
- Broadly described by models
- $\alpha_{\text{IP}}(0) = 1.058 \pm 0.036$ (ATLAS)
- Further significant progress will require proton tagging to unfold SD from DD and ND

