



The Compact Muon Solenoid Experiment
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Probing Majorana Neutrinos at the CMS

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Abstract

Observing a Lepton Number Violation process ($\Delta L = 2$) is a clear way to establish the Majorana nature of the neutrino mass. We search for a signature of same-sign dileptons and two jets from heavy neutrino decay. The latest search results for heavy Majorana neutrinos at the LHC using the CMS detector will be presented.

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Probing Majorana Neutrinos at the CMS

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We report a search for heavy Majorana neutrinos (N) decaying into a W boson and a lepton using the CMS detector at the Large Hadron Collider. We search for a signature of two same-sign electrons and two jets or two same-sign muons and two jets or a same-sign electron-muon pair and two jets using 19.7 fb^{-1} of data collected during 2012 in proton-proton collisions at a center-of-mass energy of 8 TeV. We find the data to be consistent with the expected standard model background and upper limits are set on $|V_{iN}|^2$, as a function of heavy Majorana neutrino mass from 40-500 GeV/c^2 , where V_{iN} is the mixing element of the heavy Majorana neutrino with the standard model neutrino. The upper limits obtained for $(|V_{eN}|^2, |V_{\mu N}|^2, |V_{eN}V_{\mu N}^*|)$ are (0.022, 0.00470, 0.009) for $m_N = 50 \text{ GeV}/c^2$, (0.016, 0.0123, 0.009) for $m_N = 200 \text{ GeV}/c^2$, and (0.72, 0.583, 0.47) for $m_N = 500 \text{ GeV}/c^2$. These results extend considerably the regions excluded by previous direct searches.

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1. Introduction

The discovery of neutrino oscillation established that neutrinos have non-zero masses. Results from various neutrino oscillation experiments and cosmological constraints showed the upper limit of neutrino mass sum is very small, less than 0.13 eV.

The so-called "SeeSaw" mechanism is able to explain this smallness of neutrino masses in natural way by introducing a heavy neutrino state N . Here, the mass of standard neutrino is given by $m_\nu \sim y_\nu^2 v^2 / m_N$, where y_ν is a Yukawa coupling, v is the Higgs vacuum expectation value in the standard model (SM) and m_N is the mass of the heavy Majorana neutrino state. Direct searches of heavy Majorana neutrinos are very important to understand the nature of neutrinos and the origin of their masses. We considered a heavy Majorana neutrino mixes with the SM neutrinos, with m_N and V_{lN} as free parameters of the model. Here V_{lN} is a mixing element describing the mixing between the heavy Majorana neutrino and the SM neutrino of flavour l .

In the process we performed the search, the heavy Majorana neutrino is produced by s-channel production of a W boson, which decays via $W^\pm \rightarrow Nl^\pm$. The N is assumed to be a Majorana particle, so it can decay via $N \rightarrow W^\mp l^\pm$ with $W^\mp \rightarrow q\bar{q}'$, resulting in a $l^\pm l^\pm + q\bar{q}'$ final state.

2. CMS detector and Event Selection

The central feature of the CMS detector is a superconducting solenoid with an internal radius of 3 m. The solenoid provides a magnetic field of 3.8 T along the direction of the counter-clockwise rotating beam. Within the superconducting solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Muons are measured in gas ionization detectors embedded in the steel return yoke. Extensive forward calorimetry complements the coverage provided by the barrel and endcap detectors. The first level of the CMS trigger system, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select up to 100 kHz of the most interesting events. The High Level Trigger (HLT) processor farm uses information from all CMS subdetectors to further decrease the event rate to roughly 300 Hz before data storage. A more detailed description of the CMS detector can be found in [3].

In the analysis, we use several samples of simulated SM process which include a full treatment of the pp collisions. These samples include the full simulation of the CMS detector based on GEANT4 and are reconstructed using the same CMS software as used for data. In order to simulate heavy Majorana neutrino events, we use a leading-order (LO) event generator describes in [4], which was implemented in ALPGEN v2.14 [5] with the CTEQ6M parton distribution functions (PDFs) [7]. We assumed $|V_{lN}|^2 = 1$, and considered mass of heavy Majorana neutrino from 40 to 500 GeV/c² [11]. We scale the LO cross section by a factor of 1.34 to account for higher-order corrections, based on the next-to-next-to-leading-order calculation in FEWZ [9, 10]. Parton showering

and hadronization are simulated using PYTHIA v6.4.22 [8].

Additional selections are applied to the triggered events to ensure the presence of well-identified dileptons and jets. In this analysis, we concentrate on the same-sign (SS) dilepton and additional two jets without missing any final state energy from the heavy Majorana neutrino decay. Lepton candidates are required to pass CMS official tight ID. Additionally, we used more tight relative isolation cut to reduce fake lepton background from QCD processes. The invariant mass of the dilepton pair is required to be above $10 \text{ GeV}/c^2$ and outside the Z mass peak for the dielectron channel, to remove a significant amount of background from charge mismeasurement. We also veto third lepton to suppress background from diboson production events. Jets are selected with CMS official selection. To reduce background from top quark decays, b-quark jet veto was applied where the medium working point of the combined secondary vertex tagger [12] has been used.

Depending on the heavy Majorana neutrino mass, signal events from heavy Majorana neutrino decays have different kinematic properties. In the low neutrino mass search region ($m_N < m_W$), the W-boson propagator is on-shell and the final state system of $M(l^\pm l^\pm jj)$ should be close to the W mass. In the high neutrino mass search region ($m_N > m_W$), the W-boson propagator is off-shell but the W-boson from the heavy Majorana neutrino decay is on-shell, so the $M(jj)$ will be close to the W mass. We require missing transverse energy (MET) to be less than 30 (35) GeV in low-mass (high-mass) search region because we have no neutrino at final state of signal process. We applied additional cuts to reduce background without losing many portion of signal events. Detailed information about event selection can be found in [1, 2].

3. Background Estimation and Systematic Uncertainties

There are three major sources of same-sign dilepton backgrounds. The first and the most important background events are originated from misidentification (fake) of objects as prompt leptons. These "prompt leptons" originate from B hadron decays, light-quark or gluon jets as well as from photon conversions, and are typically not well isolated. The simulation is not reliable because of lack of statistics comes from small misidentification probability. Data driven method from QCD enriched data sample was used to estimate fake lepton background [1, 2].

The second contribution is from SM production of two genuine isolated leptons of the same charge. This irreducible background consist of diboson events as well as triboson and $t\bar{t}$ plus boson, other rare processes include Higgs event where Higgs decays into neutral bosons, or associated Higgs production VH and $t\bar{t}H$ in which case also $H \rightarrow WW$ contribute. We estimate the background using PYTHIA genertator and MADGRAPH v5.1.3.30 [6].

The third background source is from opposite-sign event in which the charge of one of the electrons is mismeasured. We used both simulated sample of Drell-Yan process and data to evaluate the charge mismeasurement probability as a function of Pt and η of the electron. The charge mismeasurement background of muon is ignored because of much low probability [1, 2].

The main source of systematic uncertainties are associated with the background estimates. The overall systematic uncertainty on the misidentified lepton background is 40%, 28% and 35% for dielectron, dimuon and electron-muon channels respectively. An uncertainty of 12% assigned to the background from charge mis-measurement for dielectron channel. The remaining candidates are cross section uncertainty of MC samples and detector effect. Detailed information about systematic uncertainty can be found in [1, 2].

4. Results and Summary

We see no evidence for a significant excess in the data beyond the backgrounds predicted from SM and we set 95% confidence level (CL) exclusion limits on the square of the Majorana neutrino mixing element and cross section times branching fraction for $pp \rightarrow Nl^\pm \rightarrow l^\pm l^\pm q\bar{q}'$ as a function of m_N . We use Poisson distributions for the signal and log-normal distributions for the nuisance parameters. For simple interpretation, each channel of dielectron, dimuon, electron-muon channels is considered as independent analysis to each other. We set limits on $|V_{eN}|^2$, $|V_{\mu N}|^2$, and $|V_{eN}V_{\mu N}^*|$ as a function of m_N .

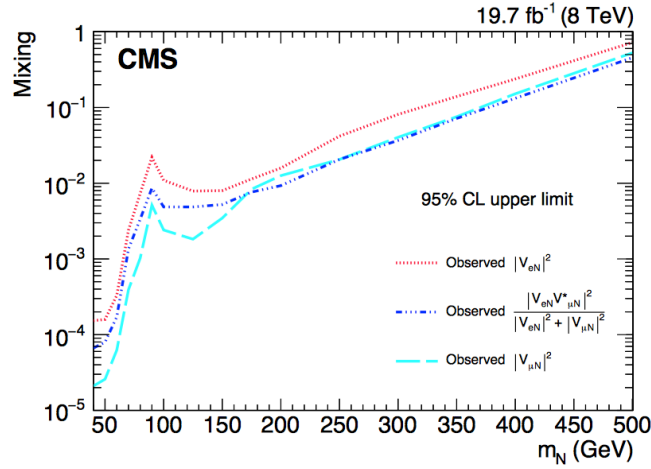


Figure 1: 95% CL limits on mixing parameter between the heavy Majorana neutrino and SM neutrinos

We have performed a search for heavy Majorana neutrinos decaying into same-sign leptons and two jets using 19.7 fb^{-1} of data collected during 2012 in pp collisions at a center-of-mass energy of 8 TeV. We do not observe any excess of events beyond the expected standard model background and set 95% CL upper limits on the mixing elements, $|V_{eN}|^2$, $|V_{\mu N}|^2$, and $|V_{eN}V_{\mu N}^*|$, of the Majorana neutrino as a function of Majorana neutrino mass from $40 \text{ GeV}/c^2$ to $500 \text{ GeV}/c^2$.

The results significantly improve the previous LHC limits on $|V_{lN}|^2$ as well as the studied mass hypothesis range.

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