

Supersonic jets of hydrogen and helium for laser wakefield acceleration

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The properties of laser wakefield accelerated electrons in supersonic gas flows of hydrogen and helium are investigated. At identical backing pressure, we find that electron beams emerging from helium show large variations in their spectral and spatial distributions, whereas electron beams accelerated in hydrogen plasmas show a higher degree of reproducibility. In an experimental investigation of the relation between neutral gas density and backing pressure, it is found that the resulting number density for helium is $\sim 30\%$ higher than for hydrogen at the same backing pressure. The observed differences in electron beam properties between the two gases can thus be explained by differences in plasma electron density. This interpretation is verified by repeating the laser wakefield acceleration experiment using similar plasma electron densities for the two gases, which then yielded electron beams with similar properties.

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The development of bright and ultrashort sources of particles and x rays is an important area of research. Such sources are of interest in many domains, including materials science, chemistry, biology, and medicine. Currently, emerging sources based on laser-plasma acceleration [1] are attracting significant attention. The accelerator can be very compact, and the particle beams have several unique characteristics. Recent achievements include the generation of electron beams with high energies (few GeV) [2], short pulse duration (few femtoseconds) [3], high peak current (few kA) [4], low energy spread ($< 1.5\%$) [5], and low emittance (few $\text{mm} \times \text{mrad}$) [6]. For most demanding applications, however, the stability of the source is also very important. A critical issue for laser wakefield accelerator (LWFA) research is to find ways to decrease shot-to-shot fluctuations.

In a typical LWFA, an intense laser pulse is focused in a neutral gas medium and atoms, or molecules, are rapidly ionized by the leading edge of the laser pulse. The main part of the pulse interacts with a plasma, and free electrons are displaced by the laser ponderomotive force which leads to a significant charge separation and a copropagating plasma wave. Strong accelerating electric fields ($\sim 100 \text{ GV/m}$) are present in the plasma wave, and copropagating electrons can be accelerated to high energies if they have sufficient initial kinetic energy and are located in an appropriate phase of the plasma wave. In the so-called bubble regime [7], the injection of electrons can be

achieved by driving the plasma wave to such a high amplitude that the wave breaks. This occurs as the velocity of the electrons exceeds the phase velocity of the plasma wave and results in self-injection of electrons from the background plasma into the accelerating phase of the plasma wave.

The threshold for wave breaking can be described as a laser power threshold [8] as well as a laser energy threshold [9] for a given plasma electron density n_e . Thus, for a given set of laser parameters, the self-injection threshold can be found by adjusting n_e . Assuming ideal gas behavior, the neutral gas number density n in a supersonic jet is proportional to the pressure p_0 supplied to the nozzle and for a fully ionized gas $n_e = N_e n$, where N_e is the number of electrons per atom, or molecule, depending on the gas species. Thus, for fully ionized gases, $n_e \propto p_0$. In this article, we present, to our knowledge, the first comparative study of electron beams emerging from supersonic jets of H_2 and He. These gases were chosen since they will be fully ionized for the present experimental conditions.

The experimental investigations were conducted using the multiterawatt laser at the Lund Laser Centre. This Ti:sapphire-based system produced 37 fs duration laser pulses with 650 mJ of energy on target during the present study. An $f/15$ parabolic mirror focused the laser pulse to a $16 \mu\text{m}$ (FWHM) spot measured in vacuum, which yielded a peak intensity of $5.7 \times 10^{18} \text{ W/cm}^2$. The beam waist was positioned, within one Rayleigh length, at the front edge of a supersonic gas flow released from a 2 mm diameter nozzle. Behind the interaction medium, along the laser propagation axis, a permanent dipole magnet dispersed the accelerated electrons according to energy. The dispersed electron beams impacted on a scintillating screen, imaged using a 16-bit digital camera. The integrated charge above

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the spectrometer threshold energy (40 MeV) was also estimated using the measured response of the scintillator screen [10,11].

In Fig. 1, two five-image sequences of electron beams accelerated in 2 mm gas jets of H₂ and He operated at $p_0 = 9.5$ bar are presented. It is apparent that electron beams originating from H₂ [Fig. 1(a)] were, compared to those accelerated in He plasmas [Fig. 1(b)], more stable in terms of maximum electron energy, position, and spatial divergence, as well as integrated beam charge. Most electron energy spectra contained a single peak with a relatively large energy spread, corresponding to the dispersed electron beams shown in Fig. 1(a). Also, the individual images shown in Fig. 1(a) are similar to the average of the full sequence, consisting of ten images, which is shown in Fig. 1(c). However, the electron beams emerging from He [Fig. 1(b)] fluctuated significantly and suffered from filamentation, which was not the case for

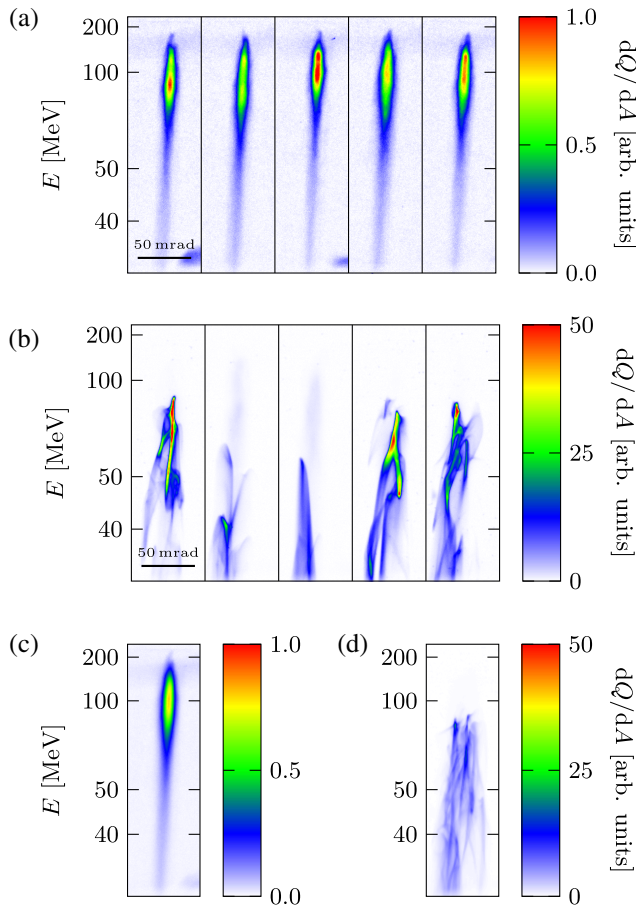


FIG. 1. False-color images of five electron beams emerging from (a) H₂ and (b) He dispersed by a permanent dipole magnet. In both cases, a 2 mm nozzle was used at a fixed backing pressure of 9.5 bar. The reproducibility of the data is shown by the average of ten individual images of electron spectra for beams emerging from (c) H₂ and (d) He. All color scales are normalized to the maximum signal in (a).

beams from H₂. Most of the energy spectra of the beams originating in He had multiple peaks, each often having very small energy spreads. It is also apparent that the sequence average [see Fig. 1(d)] is not similar to any of the individual images shown in Fig. 1(b). When comparing the two series, it can also be deduced that the integrated charge of beams accelerated in He is significantly larger than those accelerated in H₂.

The integrated beam charge was measured in a sequence of pulses while varying the pressure in the range 3–15 bar, and the results are shown in Fig. 2. As can be seen, the threshold for self-injection, which is the point where beam charge increases rapidly, is at 9 bar for He but occurs at 11 bar for H₂, indicating differences between the two media.

We have evaluated several phenomena in order to explain our observations, such as differences in the neutral gas ionization and the corresponding ionization-induced defocusing [12]. However, the intensity needed [13] for He \rightarrow He⁺ is 1.4×10^{15} W/cm², and for He⁺ \rightarrow He²⁺ is 8.8×10^{15} W/cm², which are at least 2 orders of magnitude below the peak laser intensity used in this experiment. Thus, this effect should have been noticeable only at the front of the laser pulse and in the wings. Simulations of the laser-pulse evolution performed using the code WAKE [14], which included ionization of neutral gases, did not show any significant differences in pulse characteristics when propagating through H₂ compared to He at identical n_e .

Another possible cause for the behavior in Fig. 2 could be weaker accelerating fields for H₂ than for He. Since H₂ is a molecular gas, the background of positively charged ions in the bubble behind the laser pulse might not be homogenous, as is expected for monatomic gases such as He. Assuming that the protons of the fully ionized H₂ ions

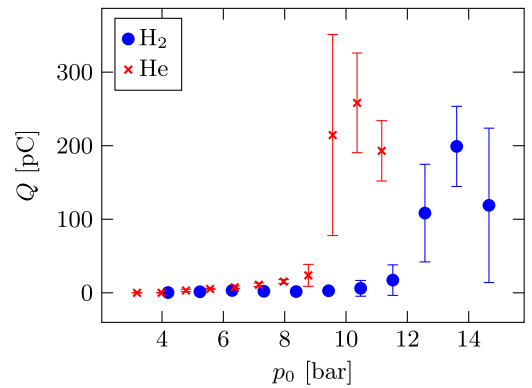


FIG. 2. Measured charge Q in the electron beams accelerated in a 2 mm gas jet over the scanned pressure range 3–15 bar in H₂ (blue circles) and He (red crosses) plotted as functions of the backing pressure. Each point represents the average of ten individual measurements with error bars indicating one standard deviation in each direction. Note that only electrons with an energy exceeding the cutoff (40 MeV) contributes to Q in this figure.

are separated by their molecular bond distance (0.074 nm), Coulomb repulsion will cause an explosion. However, in a H₂ plasma, the initial ion speed (~ 4.5 nm/fs) is too small to have a noticeable effect on the ion density in the bubble.

Finally, the differences between the gases seen in Figs. 1 and 2 can be due to fluid mechanical differences between the gases. To determine the magnitude of such an influence on the resulting n_e , a simple model of a converging-diverging nozzle was investigated. The relation between the nozzle throat sizes and flow Mach number M is [15]

$$\left(\frac{r_0}{r^*}\right)^2 = \frac{1}{M} \left[\frac{2 + (\kappa - 1)M^2}{(\kappa + 1)} \right]^{\frac{\kappa + 1}{2(\kappa - 1)}}, \quad (1)$$

where r_0 is the nozzle exit radius, r^* is the critical radius where the flow reaches sonic speeds inside the nozzle, and κ is the ratio of specific heats of the gas with numerical values 1.41 for H₂ and 1.66 for He [16]. For the specified $r^* = 0.39$ mm of the 2.0 mm diameter nozzle used in the experiments, Eq. (1) yields $M = 3.5$ and $M = 4.2$, for H₂ and He, respectively. Assuming that the gas can be described as an ideal gas, it is also possible to express the density at the nozzle exit, n_{exit} , as [17]

$$n_{\text{exit}} = \frac{p_0}{k_B T_0} \left[1 + \frac{\kappa - 1}{2} M^2 \right]^{-\frac{1}{\kappa - 1}}, \quad (2)$$

where k_B is Boltzmann's constant and $T_0 = 293$ K the temperature. As the flow exits the nozzle, it will diverge with half-angle φ given by $\varphi = \alpha + \theta$, where $\alpha = \arcsin M^{-1}$ is the Mach cone half-angle and θ the nozzle expansion angle. This means that, using cylindrical symmetry, the radius of the gas flow can be written as $r = r_0 + h \tan \varphi$, where h is the vertical distance from the nozzle exit. Assuming that φ remains constant, the gas density at a specific h close to the nozzle exit can be estimated as $n = n_{\text{exit}}(r_0/r)^2$.

As is seen from Eqs. (1) and (2), there is a nontrivial relation between n_{exit} and the gas-species-dependent κ . Therefore, characterizing the relation between p_0 and n for both gases released from the nozzle was necessary and performed experimentally. The phase shift introduced by He at $n = 5 \times 10^{18} \text{ cm}^{-3}$ over 2 mm for 633 nm light is 0.14 rad (corresponding to a 14 nm optical path length difference), which is difficult to measure with an ordinary interferometer. Therefore, n was measured as a function of p_0 with a setup consisting of an expanded HeNe-laser beam and a wave-front sensor [18], which is sensitive enough to determine the phase shift introduced by He. By assuming full ionization, n_e is then plotted as a function of p_0 for the two gases in Fig. 3, which clearly shows that they resulted in different n_e at all p_0 . Using r^* as a fitting parameter in Eq. (1) to simultaneously fit the theoretical model to experimental results obtained for both H₂ and He yielded $r^* \approx 0.35$ mm, which is close to the specified critical radius

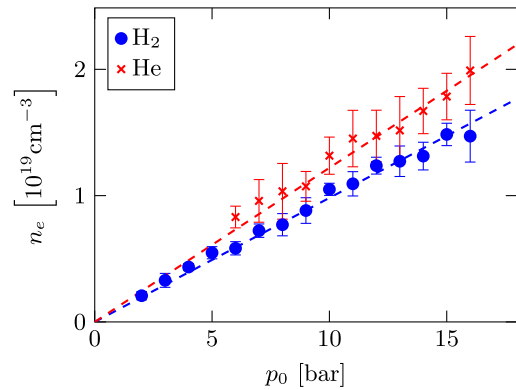


FIG. 3. The plasma electron number density n_e , 1 mm from the nozzle orifice (2 mm diameter) as a function of the applied backing pressure (p_0) for H₂ (blue circles) and He (red crosses). Assuming full ionization, the plateau electron number density n_e along the center axis in the laser propagation direction is determined from measurements of the neutral gas number density (n) using a setup consisting of an expanded HeNe beam and a wave-front sensor. Each point represents the average of 10–20 individual measurements, and the error bars indicate one standard deviation in each direction. The dashed lines are the theoretical results fitted with regards to r^* .

of the nozzle. The fitted results, shown as dashed lines in Fig. 3, are in excellent agreement with the experimental data.

From the theoretical model, it was found that $n_{\text{He}} \approx 1.3n_{\text{H}_2}$. Thus, n_e in He is $\sim 30\%$ higher than for H₂ at any specific p_0 . Compensating for this difference and plotting the data in Fig. 2 as a function of n_e instead of p_0 results in Fig. 4(a). Now it can be seen that the rapid increase in Q occurs at the same n_e for both gas species. The effect observed in Fig. 1 is therefore not significantly due to any of the previously discussed differences between the two gas species but can be explained by the relation between n_e and p_0 in Fig. 3. In Fig. 1, the electrons were accelerated in gas jets with $p_0 = 9.5$ bar which corresponds to $n_e = 9.3 \times 10^{18} \text{ cm}^{-3}$ for H₂ and $n_e = 1.2 \times 10^{19} \text{ cm}^{-3}$ for He. Using a similar n_e for He as for H₂ in Fig. 1(a) results in Fig. 4(b). Now, the accelerated electron beams emerging from He are very similar to the ones from H₂, which is also seen when comparing the averages of ten individual images in Figs. 1(c) (H₂) and 4(c) (He). Laser self-focusing inside the plasma becomes stronger with increasing n_e , resulting in a smaller spot size w_0 and a higher normalized vector potential a_0 for He than for H₂. For small w_0 and high a_0 , it is expected that self-injection LWFA results in unstable, high charge electron beams, since transversal injection dominates over longitudinal injection [19]. When longitudinal injection is the dominant injection mechanism (large w_0 and small a_0), the accelerated electron beams becomes very stable, but with low charge. Thus, the differences seen in Fig. 1 can be explained by the differences in n_e between H₂ and He at

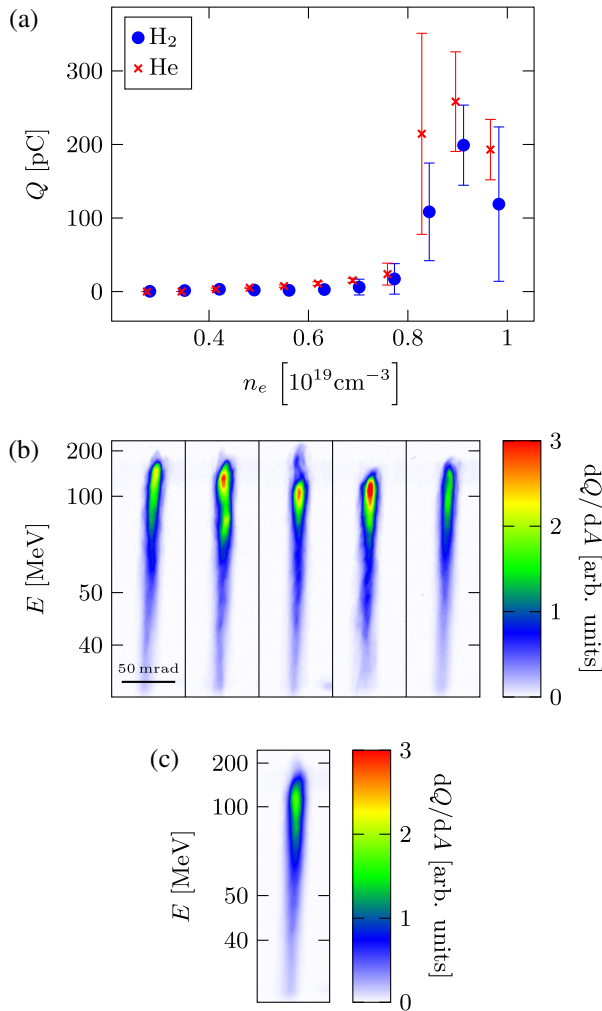


FIG. 4. (a) Measured charge Q from Fig. 2 plotted against plasma electron density n_e (assuming full ionization) instead of backing pressure p_0 , by using the results of Fig. 3. Again, only electrons with energies exceeding the cutoff energy of the spectrometer setup contribute to Q . Each point represents the average of ten individual measurements with error bars indicating one standard deviation. (b) False-color images of five electron beams emerging from He at $n_e = 9.8 \times 10^{18} \text{ cm}^{-3}$ dispersed by a permanent dipole magnet, showing a high resemblance to electron beams accelerated in H_2 at similar n_e . (c) Average of the full ten-image sequence which is partly shown in (b). The color scales are normalized to the maximum signal in Fig. 1(a).

identical p_0 , since the two series have different injection mechanisms.

In this study, we have shown that electron beams emerging from H_2 and He at identical nozzle backing pressures have different properties. This is found to be primarily a result of the supersonic gas jet number density dependence on a specific heat ratio which, generally, differs between gas species. Repeating the experiment using similar n_e for both gases confirms these findings, since the resulting beams of accelerated electrons then showed similar properties regardless of gas species. Thus, both

gases resulted in stable, low charge electron beams for $n_e < 8 \times 10^{18} \text{ cm}^{-3}$, which can be deduced from Fig. 4. It is also believed that this effect can have implications when using gas mixtures as an acceleration medium and should be studied further.

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