[1] A. Falkowski *et al*, *Hidden Higgs Decaying to Lepton Jets*, *JHEP* **05** (2010) 077 [arXiv:1002.2952] **[2]** ATLAS EXOT-2013-22 **[3]** ATLAS EXOT-2013-22 Aux **[4]** LHC Higgs Cross Section Working Group 2014

RESULTS

A search for dLJ pairs was performed using the full 2012 ATLAS dataset at $\sqrt{s} = 8 \text{ TeV}$ (20.3 fb⁻¹). The search was kept largely model-independent until the limit-setting stage. To facilitate re-casting of the results, t rigger and reconstruction efficiency tables were produced as a function of dark photon $c\tau$ and p_T using a dedicated "LJ gun" Monte Carlo tool. A similar, complementary search for "prompt" LJ pairs (where y_d has small or zero cτ, and therefore decay vertex compatible with the event's primary vertex) was also performed.

For Run 2 (13 TeV), the analysis will be extended to:

Additional topologies, e.g. new narrow-scan trigger will enable searches for only one LJ in the final state

- Tri-muon: 3 MSonly tracks, $p_T > 6$ GeV (for pair of Type 0 dLJs)
- Narrow-scan: 2 MSonly tracks in Δ R = 0.5 cone, leading $p_T > 20$ GeV, sub-leading > 6 GeV (for Type 0, 1 dLJs). New in Run 2
- CaloRatio: jet $p_T > 30$ GeV with low EM fraction (for Type 1, 2 dLJs)

• Higher Higgs masses (for BSM extended Higgs sectors) and associated Higgs production

.

TRIGGERS

Cuts defined using data-driven method for
 QCD multi-jet contamination:

Simplified matrix (ABCD) method assumes multi-

jet background factorizable in 2D plane
 E 10
 E 10
 $\frac{20}{20}$ 10
 E 10
 E 10
 E 10
 QCD multi-jet contamination:

The ATLAS trigger system has multiple levels. L1 is hardware-based, using the calorimeters and MS to define regions-of-interest. The next two levels, combined into a High-Level Trigger (HLT) in Run 2, are software-based and are fed from all subdetectors.

The following HLT triggers are used for dLJs:

 Lightest unstable hidden states in MeV to GeV range \rightarrow typically produced with large boost \rightarrow highly-collimated decay products Decay back to SM with high branching fraction \rightarrow e⁺e⁻, μ⁺μ⁻, or π⁺π⁻ in final state

SELECTION REQUIREMENTS

Non-negligible lifetime \rightarrow decay vertex displaced relative to primary vertex of event

> **Cuts defined to optimize signal significance, with variables ordered by separation power:**

Smaller mixing parameter ϵ yields longer y_d lifetime. BRs of the y_d depend on its mass. Much of the parameter space in such models features long-lived, boosted low-mass $\gamma_d \rightarrow l^+l^-$ yielding final-state dLJs.

- Jet Vertex Tagger: Rejects QCD jets (Type 2 dLJs)
- Jet Width: Rejects QCD (Type 2)
- Jet EM Fraction: Rejects QCD (Type 2)
- Muon timing (∆t between RPC layers): Rejects muon bundles (Types 0, 1)
- Jet timing: Rejects mis-reconstructed cosmics (Types 1, 2)
- Beam-Induced Background tagging: Rejects fake BIB jets accompanied by φ–matched muon segments parallel to beampipe (Type 2)

We target γ^{decays} beyond the Inner Detector (ID) **up to the Muon Spectrometer (MS).**

Simplified matrix (ABCD) method assumes multijet background factorizable in 2D plane

- ∑p_T: Scalar sum of transverse momentum of ID tracks (pT>0.5 GeV belonging to primary vertex of event) in Δ R = 0.5 cone around LJ centre
- max(|φ|): between leading LJ in event and LJ that is farthest from it in φ
- dLJs back-to-back from FRVZ processes have high $|\varphi|$ and low Σp_T

LJ-finding uses a clustering algorithm with Δ R = 0.5 cone, and special reconstruction considerations:

SEARCH FOR LONG-LIVED NEUTRAL PARTICLES DECAYING INTO LEPTON-JETS WITH THE ATLAS DETECTOR IN PROTON-PROTON COLLISION DATA

MOTIVATION

A wide range of Beyond the Standard Model (BSM) theories predict a hidden sector, weakly coupled to the visible sector. One particularly promising, yet challenging, collider search strategy targets processes with:

dLJs IN DARK PHOTON MODELS

dLJs are a particularly distinctive "smoking gun" signature of the dark photon $(A'$ or γ_d), the heavy gauge boson of an additional U(1)'. In "vector portal" models, the y_d can kinetically mix with the SM photon:

 $\mathcal{L} \supset \frac{\epsilon}{2} F^{\mu\nu} A_{\mu\nu} + m_{\gamma_d}^2 A^{\gamma 2}$

BENCHMARK MODELS

As benchmarks, we take two hidden sector models (Falkowsky-Ruderman-Volansky-Zupan, FRVZ):

H (SM-like Higgs or a BSM heavy neutral variety) decays to heavy hidden fermions fd2 . Each fd2 may decay directly to a γ^d and Hidden Lightest Stable Particle (left), or through a hidden scalar sd1 (right). The γ^d 's decay to LJs, which usually come off back-to-back. [1]

dLJ SEARCH STRATEGY

- Muon pairs appear in spectrometer as "MSonly" tracks (no associated ID tracks)
- Electron / pion pairs appear in calorimeters as "LJCalTracks" (narrow isolated jets, with much less energy deposition in EMCal than in HCal)

- Collimated final-state particles difficult to reconstruct due to detector granularity
- Reconstruction of tracks with displaced decay

vertices require removal of primary vertex constraints used in standard algorithms, and is especially difficult outside ID

LJ TYPE2

 \Box ID

 \Box HCAL

 \Box MS

ATLAS Simulation

TYPE2 LJ

 $m_v = 0.05 \text{ GeV}$

 $m_v = 0.15 \text{ GeV}$

 $m_v = 0.40 \text{ GeV}$

• $m_v = 0.90 \text{ GeV}$

• $m_{\gamma_d} = 1.50 \text{ GeV}$

 $L_{xy}^{t_d}$ [m]

EMCAL

Miriam Diamond, University of Toronto, on behalf of the Exotics Group of the ATLAS Collaboration

Displaced Lepton-Jets (dLJs): Collimated jet-like structures, produced far from the primary vertex, containing pair(s) of muons, electrons, and/or light hadrons

SYSTEMATIC UNCERTAINTIES

(In order of Run 1 importance)

Cosmic ray background estimation

QCD multijet background estimation

Trigger efficiency

 \bullet γ_d detection efficiency and p_T resolution

Higgs production $\sigma_{gg\;{\rm fusion}}$

Muon reconstruction efficiency

Effect of pile-up on Σp_T

Normalization of integrated luminosity

Jet energy scale

BACKGROUND SOURCES

- QCD multi-jet**:** γγ, γ+jets, tt, single top, Drell-Yan e **+**e **-** / μ **+**μ **-** , Z/W+jets, diboson
- Cosmic-ray muon energy deposits in calorimeters (for Types 1, 2 dLJs)**:** misreconstructed as jets
- Cosmic muon bundles (for Types 0, 1)**:** mainly concentrated in barrel
- Beam-induced background (for Type 2)**:** highenergy muon longitudinally crossing detector, with bremsstrahlung in HCal barrel

Exclusion contours were established in the plane of kinetic

With the FRVZ models (assuming SM Higgs produced via gluon-gluon fusion) as benchmarks, limits were established on σ **x BR(H** \rightarrow **xy_d + X), thereby setting BR-dependent exclusion bounds on cτ(γ^d) in these models. Bounds for FRVZ process (A), H → γd γd +X , shown here. [2]**

mixing parameter vs γ_d mass, in the context of vector portal models. The prompt and displaced LJ contours cover an area of parameter space untouched by other experiments, although they do depend upon an additional parameter, BR(H \rightarrow γ_d **+ X). [2]**

Simultaneous counting experiment in control and signal regions, with $N_A = N_D \times N_B/N_C$, provides estimate of **background contamination in signal region A. [2]**

 H igher $γ_d$ masses

Higher Higgs production cross-sections at 13 TeV, which will greatly benefit dLJ search sensitivity. [4]

