MOTIVATION A wide range of Beyond the Standard Model (BSM) theories predict a hidden sector, weakly coupled to the visible sector. One particularly promising, yet challenging, collider search strategy targets processes with:

Lightest unstable hidden states in MeV to GeV range \rightarrow typically produced with large boost \rightarrow highly-collimated decay products Decay back to SM with high branching fraction \rightarrow e⁺e⁻, µ⁺µ⁻, or π⁺π⁻ in final state

Displaced Lepton-Jets (dLJs): Collimated jet-like structures, produced far from the primary vertex, containing pair(s) of muons, electrons, and/or light hadrons

dLJs IN DARK PHOTON MODELS

dLJs are a particularly distinctive "smoking gun" signature of the dark photon (A' or γ_d), the heavy gauge boson of an additional U(1)'. In "vector portal" models, the γ_d can kinetically mix with the SM photon:

lifetime. BRs of the γ_d depend on its mass. Much

of the parameter space in such models features

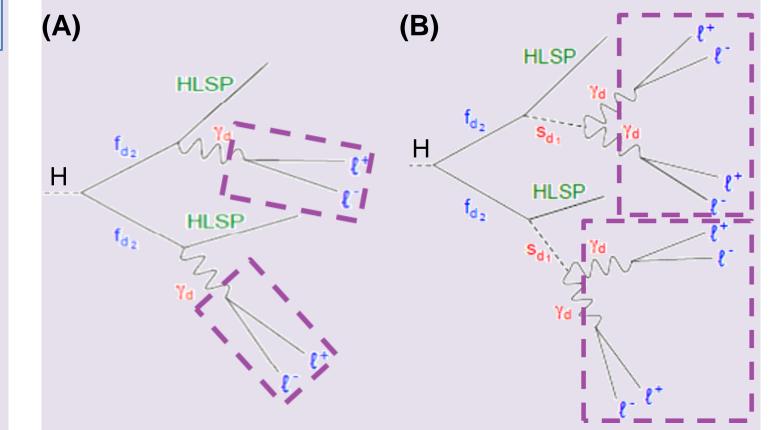
 $\mathcal{L} \supset rac{\epsilon}{2} F^{\mu
u} A_{\mu
u} + m_{\gamma_d}^2 A^{\prime 2}$

Smaller mixing parameter ϵ yields longer γ_{d}

SEARCH FOR LONG-LIVED NEUTRAL PARTICLES **DECAYING INTO LEPTON-JETS WITH THE ATLAS DETECTOR IN PROTON-PROTON COLLISION DATA**

BENCHMARK MODELS

As benchmarks, we take two hidden sector models (Falkowsky-Ruderman-Volansky-Zupan, FRVZ):



dLJ SEARCH STRATEGY

We target γ_d decays beyond the Inner Detector (ID) up to the Muon Spectrometer (MS).

- Muon pairs appear in spectrometer as "MSonly" tracks (no associated ID tracks)
- Electron / pion pairs appear in calorimeters as "LJCalTracks" (narrow isolated jets, with much less energy deposition in EMCal than in HCal)

LJ-finding uses a clustering algorithm with $\Delta R = 0.5$ cone, and special reconstruction considerations:

- Collimated final-state particles difficult to reconstruct due to detector granularity
- Reconstruction of tracks with displaced decay

Non-negligible lifetime \rightarrow decay vertex displaced relative to primary vertex of event

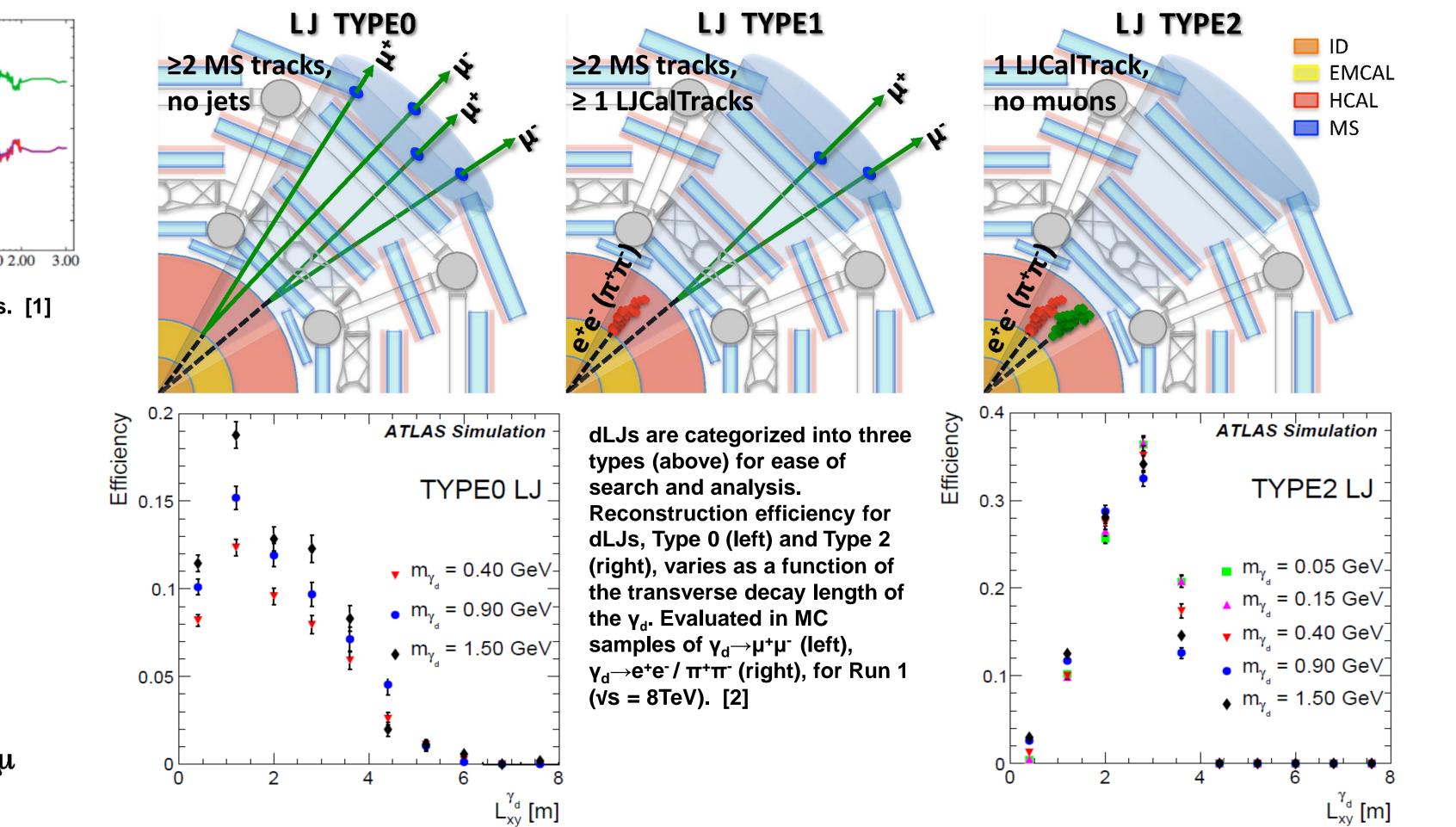
γα→μμ

0.05

long-lived, boosted low-mass $\gamma_d \rightarrow l^+l^-$ yielding final-state dLJs. 0.50 e^+e^- Event display of a simulated dLJ. [3] 2 0.30 B Hadrons 0.10 0.15 0.20 0.30 0.50 0.70 1.00 1.50 2.00 3.00 γ_d Mass [GeV] Dependence of γ_d BRs on mass. [1] ATLAS Simulation ____m, = 0.40 GeVm_{γ_} = 0.90 GeV ••••m_γ = 1.50 GeV⁻ 0.5 $\sqrt{(\eta_{1} - \eta_{2})^{2} + (\phi_{1} - \phi_{2})^{2}}$ **Opening angle of constituents of** dLJs from $\gamma_d \rightarrow \mu^+ \mu^-$. Evaluated in MC γα→μμ samples for Run 1 ($\sqrt{s} = 8$ TeV). [2]

H (SM-like Higgs or a BSM heavy neutral variety) decays to heavy hidden fermions f_{d2} . Each f_{d2} may decay directly to a y_d and Hidden Lightest Stable Particle (left), or through a hidden scalar s_{d1} (right). The γ_d 's decay to LJs, which usually come off back-to-back. [1]

vertices require removal of primary vertex constraints used in standard algorithms, and is especially difficult outside ID



BACKGROUND SOURCES

SELECTION REQUIREMENTS

RESULTS

- QCD multi-jet: $\gamma\gamma$, γ +jets, tt, single top, Drell-Yan $e^+e^-/\mu^+\mu^-$, Z/W+jets, diboson
- Cosmic-ray muon energy deposits in calorimeters (for Types 1, 2 dLJs): misreconstructed as jets
- Cosmic muon bundles (for Types 0, 1): mainly concentrated in barrel
- Beam-induced background (for Type 2): highenergy muon longitudinally crossing detector, with bremsstrahlung in HCal barrel

TRIGGERS

The ATLAS trigger system has multiple levels. L1 is hardware-based, using the calorimeters and MS to define regions-of-interest. The next two levels, combined into a High-Level Trigger (HLT) in Run 2, are software-based and are fed from all subdetectors.

The following HLT triggers are used for dLJs:

- Tri-muon: 3 MSonly tracks, $p_T > 6$ GeV (for pair of Type 0 dLJs)
- Narrow-scan: 2 MSonly tracks in $\Delta R = 0.5$ cone, leading $p_T > 20$ GeV, sub-leading > 6 GeV (for Type 0, 1 dLJs). New in Run 2
- CaloRatio: jet $p_T > 30$ GeV with low EM fraction (for Type 1, 2 dLJs)

$H \to 2\gamma_d + X$	m _H =125GeV	m _H =800GeV

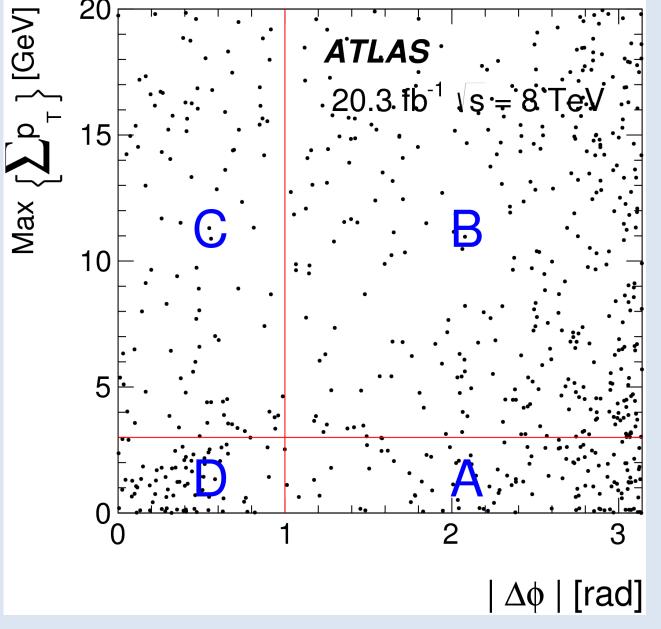
Cuts defined to optimize signal significance, with variables ordered by separation power:

- Jet Vertex Tagger: Rejects QCD jets (Type 2 dLJs)
- Jet Width: Rejects QCD (Type 2)
- Jet EM Fraction: Rejects QCD (Type 2)
- Muon timing (Δt between RPC layers): Rejects muon bundles (Types 0, 1)
- Jet timing: Rejects mis-reconstructed cosmics (Types 1, 2)
- Beam-Induced Background tagging: Rejects fake BIB jets accompanied by φ -matched muon segments parallel to beampipe (Type 2)

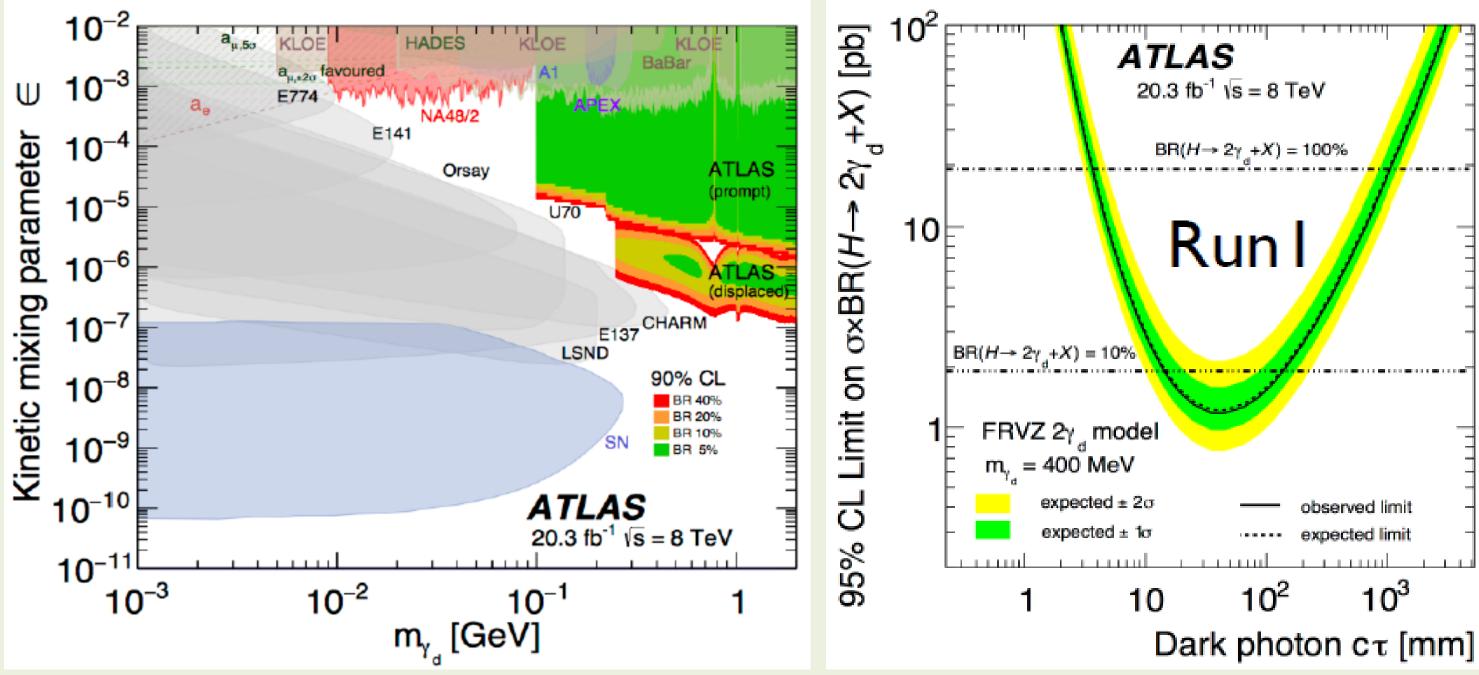
Cuts defined using data-driven method for QCD multi-jet contamination:

Simplified matrix (ABCD) method assumes multijet background factorizable in 2D plane

- $\sum p_{T}$: Scalar sum of transverse momentum of ID tracks (pT>0.5 GeV belonging to primary vertex of event) in $\Delta R = 0.5$ cone around LJ centre
- $max(|\phi|)$: between leading LJ in event and LJ that is farthest from it in ϕ
- dLJs back-to-back from FRVZ processes have high $|\phi|$ and low $\sum p_T$



A search for dLJ pairs was performed using the full 2012 ATLAS dataset at $v_s = 8$ TeV (20.3 fb⁻¹). The search was kept largely model-independent until the limit-setting stage. To facilitate re-casting of the results, trigger and reconstruction efficiency tables were produced as a function of dark photon $c\tau$ and p_{τ} using a dedicated "LJ gun" Monte Carlo tool. A similar, complementary search for "prompt" LJ pairs (where γ_d has small or zero $c\tau$, and therefore decay vertex compatible with the event's primary vertex) was also performed.



With the FRVZ models (assuming SM Higgs produced via gluon-gluon fusion) as benchmarks, limits were established on $\sigma x BR(H \rightarrow x\gamma_d + X)$, thereby setting BR-dependent exclusion bounds on $c\tau(\gamma_d)$ in these models. Bounds for FRVZ process (A), $H \rightarrow \gamma_d \gamma_d + X$, shown here. [2]

Exclusion contours were established in the plane of kinetic mixing parameter vs γ_d mass, in the context of vector portal models. The prompt and displaced LJ contours cover an area of parameter space untouched by other experiments, although they do depend upon an additional parameter, BR(H $\rightarrow \gamma_d$ + X). [2]

For Run 2 (13 TeV), the analysis will be extended to:

Additional topologies, e.g. new narrow-scan trigger will enable searches for only one LJ in the final state

	Run 2	Run 1	Run 2
Tri-muon	2.0%	2.9%	2.4%
Narrow-scan	10.6%		23.0%
CaloRatio	0.3%	2.3%	9.7%
OR of all	11.9%	4.6%	32.0%
$\textbf{H} \rightarrow \textbf{4}\textbf{\gamma}_{d} \textbf{+} \textbf{X}$	m _H =125GeV		m _H =800GeV
	Run 2	Run 1	Run 2
Tri-muon	Run 2 4.9%	Run 1 5.8%	Run 2 7.8%
Tri-muon Narrow-scan			
	4.9%	5.8%	7.8%
Narrow-scan	4.9% 8.3%	5.8%	7.8% 38.4%

Simultaneous counting experiment in control and signal regions, with $N_A = N_D \times N_B/N_C$, provides estimate of background contamination in signal region A. [2]

Higher γ_d masses

Higher Higgs masses (for BSM extended Higgs sectors) and associated Higgs production

SYSTEMATIC UNCERTAINTIES

(In order of Run 1 importance)

Cosmic ray background estimation

QCD multijet background estimation

Trigger efficiency

 γ_d detection efficiency and p_T resolution

Higgs production $\sigma_{gg fusion}$

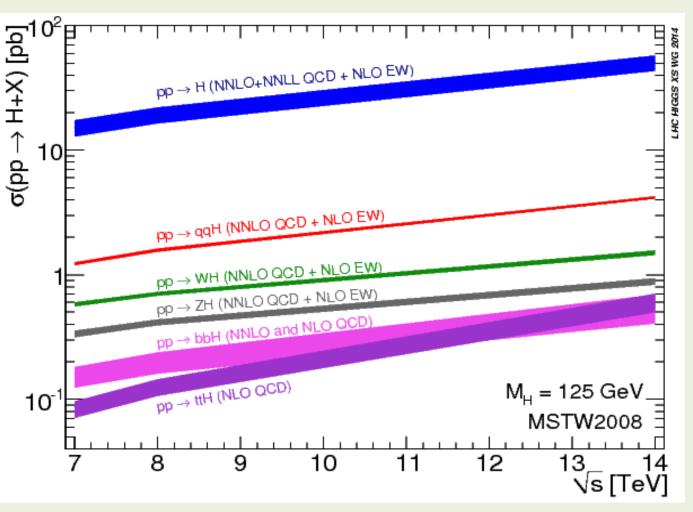
Muon reconstruction efficiency

Effect of pile-up on $\sum p_{T}$

Normalization of integrated luminosity

Jet energy scale

[1] A. Falkowski et al, Hidden Higgs Decaying to Lepton Jets, JHEP 05 (2010) 077 [arXiv:1002.2952] [4] LHC Higgs Cross Section Working Group 2014 [2] ATLAS EXOT-2013-22 [3] ATLAS EXOT-2013-22 Aux



Higher Higgs production cross-sections at 13 TeV, which will greatly benefit dLJ search sensitivity. [4]



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