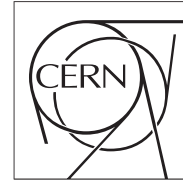




The Compact Muon Solenoid Experiment
Conference Report

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Triggering on electrons, jets and tau leptons with the CMS upgraded calorimeter trigger for the LHC RUN II

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Abstract

The Compact Muon Solenoid (CMS) experiment has implemented a sophisticated two-level online selection system that achieves a rejection factor of nearly 10^5 . During Run II, the LHC will increase its centre-of-mass energy up to 13 TeV and progressively reach an instantaneous luminosity of $2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. In order to guarantee a successful and ambitious physics programme under this intense environment, the CMS Trigger and Data acquisition (DAQ) system has been upgraded. A novel concept for the L1 calorimeter trigger is introduced the Time Multiplexed Trigger (TMT). In this design, nine main receive each all of the calorimeter data from an entire event provided by 18 preprocessors. This design is not different from that of the CMS DAQ and HLT systems. The advantage of the TMT architecture is that a global view and full granularity of the calorimeters can be exploited by sophisticated algorithms. The goal is to maintain the current thresholds for calorimeter objects and improve the performance for their selection. The performance of these algorithms will be demonstrated, both in terms of efficiency and rate reduction. The challenging aspects of the pile-up mitigation and firmware design will be presented.

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Triggering on electrons, jets and tau leptons with the CMS upgraded calorimeter trigger for the LHC RUN II

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ABSTRACT: The Compact Muon Solenoid (CMS) experiment has implemented a sophisticated two-level online selection system that achieves a rejection factor of nearly 10^5 . During Run II, the LHC will increase its centre-of-mass energy up to 13 TeV and progressively reach an instantaneous luminosity of $2 \times 10^{34} \text{cm}^{-2} \text{s}^{-1}$. In order to guarantee a successful and ambitious physics programme under this intense environment, the CMS Trigger and Data acquisition (DAQ) system has been upgraded. A novel concept for the L1 calorimeter trigger is introduced: the Time Multiplexed Trigger (TMT). In this design, nine main receive each all of the calorimeter data from an entire event provided by 18 preprocessors. This design is not different from that of the CMS DAQ and HLT systems. The advantage of the TMT architecture is that a global view and full granularity of the calorimeters can be exploited by sophisticated algorithms. The goal is to maintain the current thresholds for calorimeter objects and improve the performance for their selection. The performance of these algorithms will be demonstrated, both in terms of efficiency and rate reduction. The challenging aspects of the pile-up mitigation and firmware design will be presented.

KEYWORDS: Trigger; Algorithm; Upgrade.

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1. Introduction

The search for new physics crucially relies on the performance of the trigger system used to select the most interesting collisions amongst the millions occurring per second [1]. The CMS trigger system is organised in two consecutive steps: the hardware-based Level-1 (L1) trigger utilises coarse energy deposits in the calorimeters and signals in the muon systems to reduce the rate from about 40 MHz to 100 kHz; this is followed by the software-based High Level Trigger (HLT), implementing selection algorithms based on finer granularity and higher resolution information from all sub-detectors in regions of interest identified at L1. The output rate of the HLT is about 1000 Hz. The CMS electromagnetic calorimeter (ECAL) provides a precise measurement of the energies and positions of incident electrons and photons for both triggering and offline analysis purposes. ECAL and HCAL energies are combined to reconstruct hadronically decaying τ leptons, particle jets and energy sums. Run I of the LHC (2010-2012) already reached an instantaneous luminosity of nearly $8 \times 10^{33} \text{cm}^{-2} \text{s}^{-1}$ with p-p collisions at $\sqrt{s}=8$ TeV, 50 ns bunch-spacing and up to about 40 pile-up events per bunch-crossing, almost double the design pile-up. Nevertheless, the trigger system performed extremely well. Run II started in Spring 2015 with p-p collisions at $\sqrt{s}=13$ TeV, with 25 ns bunch-spacing. In 2016, the instantaneous luminosity is expected to reach up to $2 \times 10^{34} \text{cm}^{-2} \text{s}^{-1}$ and the number of pile-up events may be up to 70 per bunch-crossing. To avoid a significant increase in triggering energy thresholds, which would be detrimental for physics, an upgrade of the L1 trigger system is required [2].

2. Upgrade of the Level-1 trigger system for Run II

As the LHC restarts and operates at higher luminosity, the current CMS trigger system will not be capable of maintaining the thresholds required for the CMS physics programme. For example, a double-electron trigger, with thresholds of 13 GeV and 7 GeV in E_T for the two electrons respectively, had a L1 rate of 5 kHz in 2012; this would increase to about 50 kHz for the expected Run II conditions. A single electron trigger of 18 GeV threshold would give about 40 kHz, compared to 6 kHz during Run I. In these intense conditions, the implementation of pile-up mitigation techniques is required already at L1 to reach acceptable performance.

Modern technologies offer an effective solution to achieve these goals. The trigger primitives generated by the detector will be transmitted by newly installed high-speed optical links (4.8 to 6.4 Gb/s) replacing the existing copper cables (1.2 Gb/s), to a new system based on the μ TCA electronics standard. The system is based on custom designed AMC (Advanced Mezzanine Card) with Xilinx Virtex-7 FPGAs. In the TMT approach [3], these FPGAs use 10 Gb/s transceivers to gather information from the entire calorimeter for each event in a single FPGA, where sophisticated algorithms may be implemented. The complete view of the calorimeter will allow the trigger to compute global quantities such as the average energy density that can be used to estimate the pile-up level. The full calorimeter granularity provided to the algorithms is used to enhance the energy and position resolution of Level-1 candidates. The upgraded calorimeter trigger architecture is described in [4].

3. Improved selection algorithms at Level-1

The algorithms described here have been designed to thoroughly exploit the global view of the calorimeters and the full trigger tower (TT) granularity provided by the upgraded trigger system. The main goal is to get as close to the offline selection performance as possible by introducing innovative reconstruction techniques at firmware level. The Level-1 trigger system is a synchronous electronics system and therefore all algorithms implemented in the electronics must be of fixed latency. This is in contrast with offline reconstruction algorithms which are typically iterative. Lepton signatures are reconstructed with a dynamic clustering approach instead of a sliding window while particle jets are reconstructed with optimum size. The improved response results in sharper efficiency curves and cutting on dedicated isolation variables helps controlling the rate. Along with better reconstruction, identification and isolation techniques, the pile-up energy must be evaluated as well. What is required for the calorimeter trigger is an estimate of the energy that should be subtracted from the measured energy for each calorimeter object and additionally a means to remove objects which originate from pile-up particles. For the Phase 1 upgrade no information from tracking detectors is available to assist in these tasks so they must be accomplished using only the calorimeter data. Several approaches are possible and have been explored. The selection techniques must be able to evolve with the CMS physics programme and be flexible enough to remain robust in any changing conditions. The e/γ , τ lepton and jet algorithms will be presented along with their performance.

3.1 Electron and photon trigger algorithm

The improved e/γ algorithm [5] is described in Fig. 1. Clusters are seeded by local maxima of energy above a fixed threshold. The maximum size of the clusters is limited (at most 8 trigger towers can be clustered) in order to minimize the impact of pile-up energy deposits while including most of the electron or photon energy. An extended region in the ϕ -direction is used to obtain a better containment of the shower since electron and photon showers spread mostly along the ϕ -direction due to the magnetic field. Benefitting from the enhanced granularity, the e/γ candidate position can be computed as an energy-weighted average centered on the seed tower. As seen on Fig. 1, this results in a factor 4 improvement with respect to the Run I algorithm (centre of a 4×4 TT region) Additional background rejection is achieved by introducing a shape veto on the large variety of clusters produced by the dynamic clustering. The sum of the ECAL E_T of the seed and clustered towers is taken as the raw E_T of the cluster. A calibration derived from $Z\rightarrow ee$, is applied to this raw energy with factors depending on the η -position of the seed tower, the shape and the cluster E_T .

The E_T deposited in a 5×9 TTs isolation region displayed on Fig. 1, is computed excluding the footprint of the e/γ candidate. The threshold is a function of η and depends on an estimator for the number of pile-up interactions. This estimator corresponds to the number of TTs above a certain E_T threshold produced in the 8 central η rings of the calorimeters. The isolation threshold is currently tuned to reach 90% of trigger efficiency, constant as a function of pile-up and η .

The performance of this algorithm is compared with the Run I trigger using a $Z\rightarrow ee$ 2012 data sample (events selected with a tag-and-probe method) for the efficiency and a zero bias sample from a special high pile-up fill for the rates. The efficiency curves as well as the expected rate are shown in Fig. 2. The turn-on curve obtained is sharper than that of the Run 1 algorithm due to the recovery of energy lost through bremsstrahlung using a dynamic clustering at the TT level. The energy deposited by electrons is better clustered and leads to a better energy resolution of about 30%. In the endcaps, the upgrade improvement comes from the ability of the clustering to adapt to the peculiar geometry along with a more precise energy calibration. The rate of the single electron trigger can be reduced further by using the isolation criteria with a 10% efficiency loss. The shape veto can discriminate between e/γ and jet clusters reducing the fake rate, whilst keeping the efficiency loss negligible.

3.2 Selecting tau leptons

The algorithm developed is aiming at reconstructing efficiently hadronically decaying τ leptons at hardware trigger level [6]. Depending on the decay mode, several decay products may be producing more than one cluster spatially separated along the ϕ direction due to the magnetic field. Although the footprint of the τ lepton energy deposit is larger than that of an electron, the dynamic clustering developed for e/γ is perfectly adapted to reconstruct these individual clusters which can subsequently be merged. The calibration scheme is using separately the ECAL and HCAL energies and combines them linearly to optimize the response. The parameters are computed for various bins in p_T and η and stored in a Look-Up-Table (LUT). The isolation energy is derived in a similar way than for e/γ candidates. The τ footprint is subtracted from the isolation energy which is then

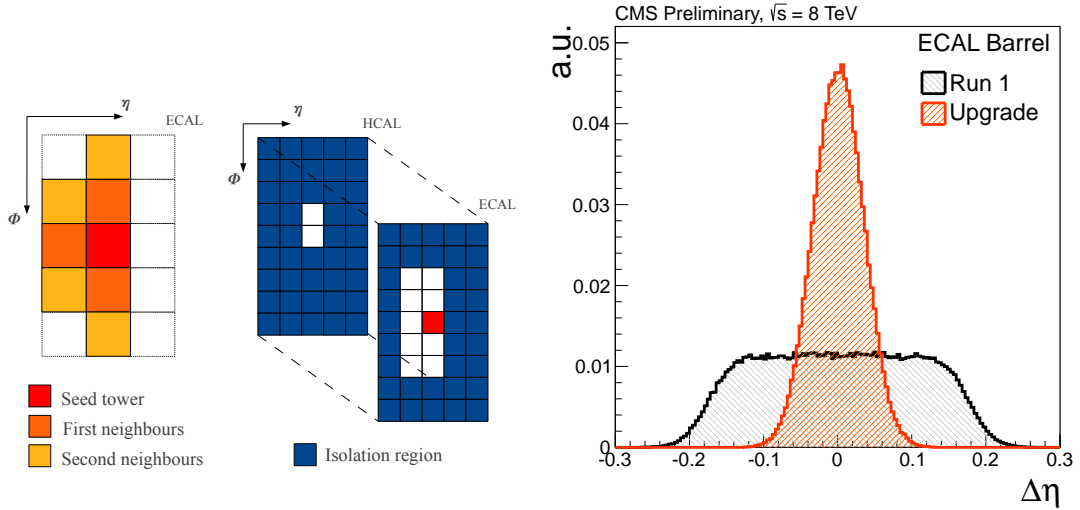


Figure 1. Left: The L1 e/γ clustering and isolation. A candidate is formed by clustering neighbour towers (orange and yellow) if they are linked to the seed tower (red). A candidate is considered as isolated if the E_T in the isolation region (blue) is smaller than a given value. The position resolution of the e/γ L1 trigger candidates in the ECAL barrel (right).

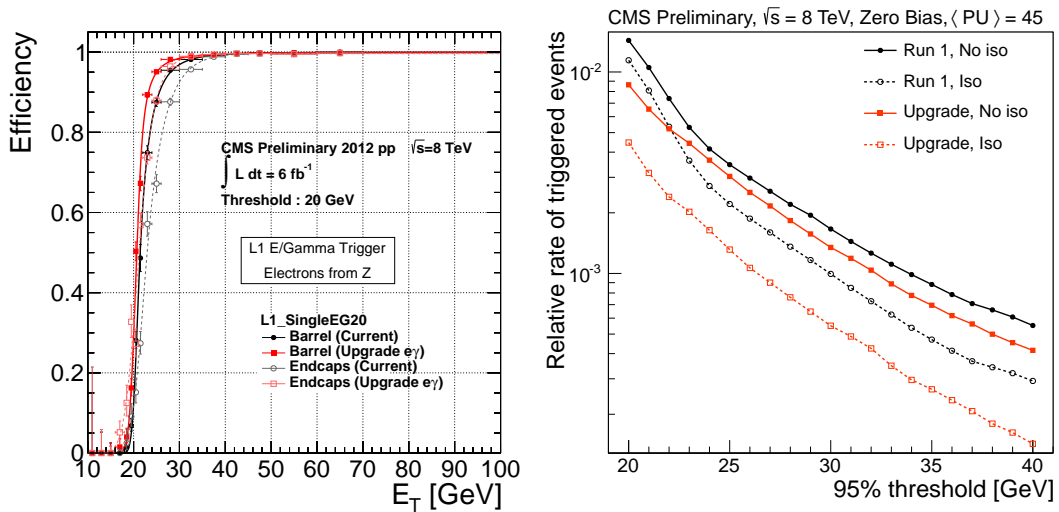


Figure 2. Electron trigger efficiency for a 20 GeV threshold at L1 as a function of the offline reconstructed E_T in the EB and EE (left). The efficiencies obtained with the current and the upgraded algorithms are shown with and without isolation criteria. Relative rate (right) of triggered events from 8 TeV zero bias data for an average pile-up of 45, obtained with the current (Run I) and the upgraded algorithms, both with and without their respective isolation requirements.

1 compared to a threshold depending on η . An additional shape veto LUT is also produced in order
 2 to discard background-like clusters from the list of possible τ candidates.

3 The performance of the τ lepton finder algorithm has been assessed on Monte Carlo simulation
 4 samples produced with $\sqrt{s} = 13$ TeV, a bunch spacing of 25 ns and 40 average pile-up interactions.

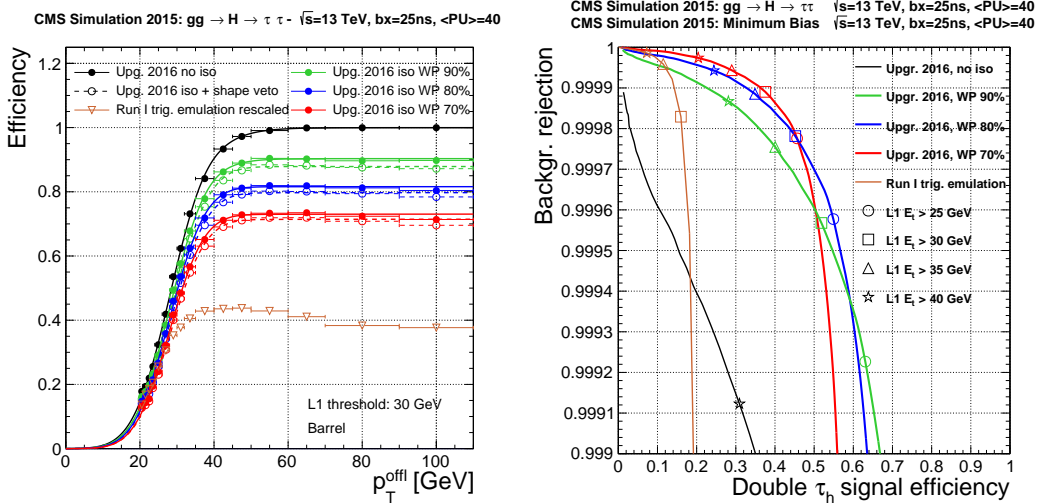


Figure 3. The Level-1 2016 upgraded trigger efficiency (left). The black line is computed for the upgrade trigger in absence of isolation, and the green, blue and red lines correspond to the 90%, 80% and 70% isolation efficiency working points (WP). The dashed lines correspond to additional requirements on the shape veto. The orange line denotes the Run I algorithm. Rate against efficiency for the double τ hadronic trigger at Level-1 (right). Various WP are shown as well as the Run 1 curve in orange.

1 The performance of the 2016 upgraded L1 algorithm is evaluated with respect to hadronic τ decays
 2 reconstructed offline using a particle flow based technique. The trigger efficiency as a function of
 3 the offline τ reconstructed p_T is displayed in Fig. 3 The 2016 upgrade algorithm shows superior
 4 triggering efficiency performance than Run 1 which does not reach a 100% plateau due to the im-
 5 plementation of a strict veto requirement and isolation on its candidates. Fig. 3 shows the trigger
 6 efficiency on signal against the background rejection. For the same background rejection a signif-
 7 icantly higher efficiency is achieved. Considering a target rate of 3 kHz for the double hadronic τ
 8 trigger, the L1 thresholds for the upgrade algorithm range between 30 (29) GeV and 42 (40) GeV
 9 depending on the working point of the isolation without (with) shape veto.

10 3.3 Jet and energy sums algorithms and their performance

11 The jet reconstruction algorithm described here is based on a similar square-jet approach as used
 12 in Run I but considering a 9×9 TT sliding window centered on a local maxima [7]. The window
 13 size chosen here is coherent with a cone radius of 0.4 for the offline jet reconstruction algorithm.
 14 In order to avoid double counting of jets the central TT energy is required to satisfy the inequalities
 15 illustrated in Fig. 4. The jet candidate energy is the sum of all 9×9 TTs energies. As seen on Fig. 4,
 16 the algorithm shows excellent agreement when compared with the anti- k_t algorithm used for offline
 17 reconstruction. Global quantities are also computed on full calorimeter granularity such as HT, the
 18 jet-based equivalent of the total E_T .

19 The pile-up subtraction is designed to operate on an event-by-event basis, allowing the per-
 20 formance to change dynamically based on differing pile-up conditions. A local pile-up correction
 21 technique called “chunky donut” was selected and proven to be efficient. Fig. 4 illustrates the
 22 chunky donut area that is used to estimate the local pile-up energy density to be subtracted from

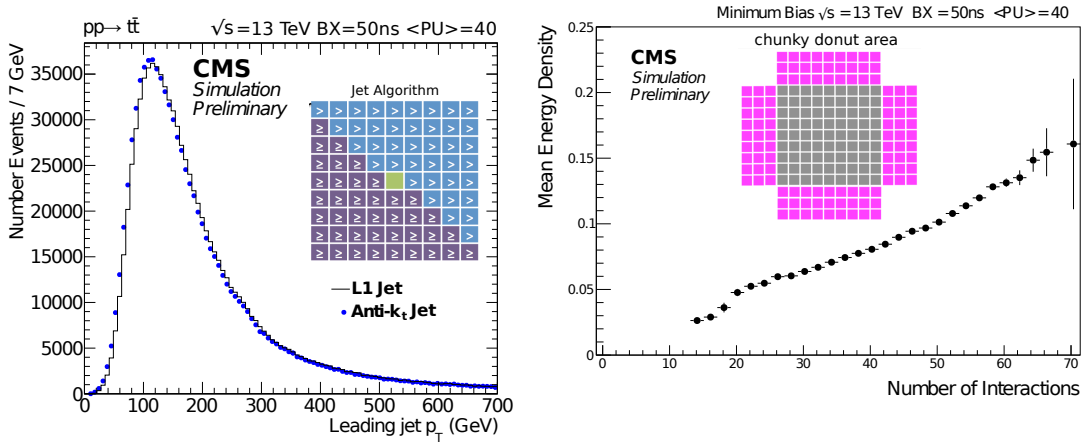


Figure 4. The 9×9 sliding window (left) used in the Level-1 jet algorithms. Any TT that satisfies these conditions is considered as a potential jet center. The histogram is a comparison between anti- k_t and the Level-1 jet algorithms. Anti- k_t is using a distance parameter of 0.4 and for the sake of performing a fair comparison; it is run on the same Level-1 TT inputs. The calorimeter area around the Level-1 jet considered for the chunky donut algorithm that is used to estimate the local energy density from pile-up (left). Correlations of the donut energy with the number of interactions is displayed.

1 the jet energy. Fluctuations that originate from the presence of other jets in the vicinity or the ab-
 2 sence of pile-up is mitigated by the extension of the donut area to four 9×3 strips around the jet
 3 and dropping the highest and lowest energy sides. Fig. 4 clearly illustrates the nice correlation of
 4 the donut energy with the number of interactions and consequently to the level of pile-up.

5 Jet energy calibration is expected to depend on the jet p_T and η . A dedicated LUT is derived
 6 from QCD di-jet MC by matching Level-1 jets to offline ones. The performance studies of the jet
 7 and energy sum algorithms were based on $t\bar{t}$ MC samples with an average of 40 pile-up interactions.
 8 A comparison is made to the same quantities produced by the Run 1 system as it was in 2012.
 9 Fig. 5 shows the Level-1 jet trigger efficiency as a function of the fourth-leading offline jet p_T
 10 for a threshold of 50 GeV. This trigger efficiency is shown against the rate (for an instantaneous
 11 luminosity of $7 \times 10^{33} \text{cm}^{-2} \text{s}^{-1}$ with 50 ns bunch crossing interval) on the same figure. The rate
 12 vs efficiency is shown for the global quantities HT on Fig. 5. The Level-1 upgraded jets algorithm
 13 with pile-up subtraction gives substantially better performance than the Run 1 trigger.

14 4. Firmware implementation and first commissioning results

15 The firmware implementation is particularly challenging as the electron finder along with the τ
 16 lepton and jet finders must fit within a single XC7V690T Xilinx FPGA. The firmware also includes
 17 core firmware, which comprises all the necessary logic to control the 72 input/output optical serial
 18 links. It also includes the configuration registers, the input pattern buffers and output spy buffers
 19 that should be accessible by software. The core firmware represents at this stage a total of 22%
 20 of the chip. The software interface is based on the IPBUS standard using libraries such as μHAL
 21 developed at CERN.

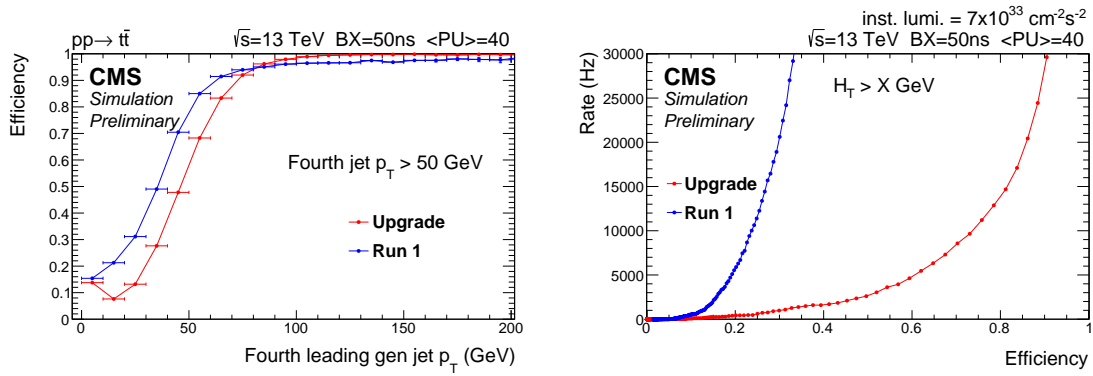


Figure 5. Level-1 Jet trigger efficiency as a function of the offline generated jet p_T for the fourth-leading jets in MC $t\bar{t}$ simulated events (left). Level-1 HT trigger rate against efficiency for an instantaneous luminosity $7 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$.

1 The calorimeters have a granularity of 72 TTs in the ϕ direction and 41 TTs in each of the
 2 positive and negative η direction (including here ECAL, HCAL and the forward calorimeter).
 3 The TMT architecture allows therefore the data to be rearranged in geometrical order (spatially-
 4 pipelined). The algorithms are thus fully pipelined and process the data at the incoming rate starting
 5 on the reception of the first data word. For the 32 bits received on each link, the internal computing
 6 frequency achieved is 240 MHz.

7 Upon reception, the TTs are combined to form basic blocks of 3×1 TTs which are then used
 8 to form larger objects such as 3×3 , 3×9 and 9×9 . These objects are the base components of the jet
 9 finder algorithm, the donut pile-up estimator and the lepton isolation region. A TT energy thresh-
 10 old is implemented to identify potential cluster seeds while a specific veto procedure is used to
 11 perform that particular step for jets as seen on Fig. 4. Global quantities are calculated as TTs are
 12 being received. At this early stage the 8 central η rings are summed to estimate the pile-up level
 13 for lepton isolation. In order to reduce the amount of resources required by the implementation of
 14 several adders in the lepton cluster logic, only “quality flags” are set for each of the incoming TT.
 15 These flags are determined by predefined criteria such as seeding, sharing or trimming. The cluster
 16 energy is thus computed as the sum of TTs with non-zero quality flag. The identification, H/E,
 17 shape identification criteria and energy calibration are implemented as LUTs. τ lepton candidates
 18 are built from e/γ -type clusters. Another LUT is consulted to perform cluster merging if required.
 19 A precise floor planning scheme has been developed to efficiently perform a place and route pro-
 20 cess and to guarantee that timing constraints will be satisfied after modification of VHDL sources
 21 during the development process. The fully pipelined firmware approach provides an efficient way
 22 to localize the processing, reduce the size and number of fan-outs, minimize routing delays and
 23 eliminates register duplication. The algorithm firmware is now loaded into the Layer-2 processors
 24 and comparisons of the outputs produced with the emulator are provided on Fig. 6. These show an
 25 excellent agreement of p_T and η distributions with the expected outputs from simulated MC data
 26 patterns.

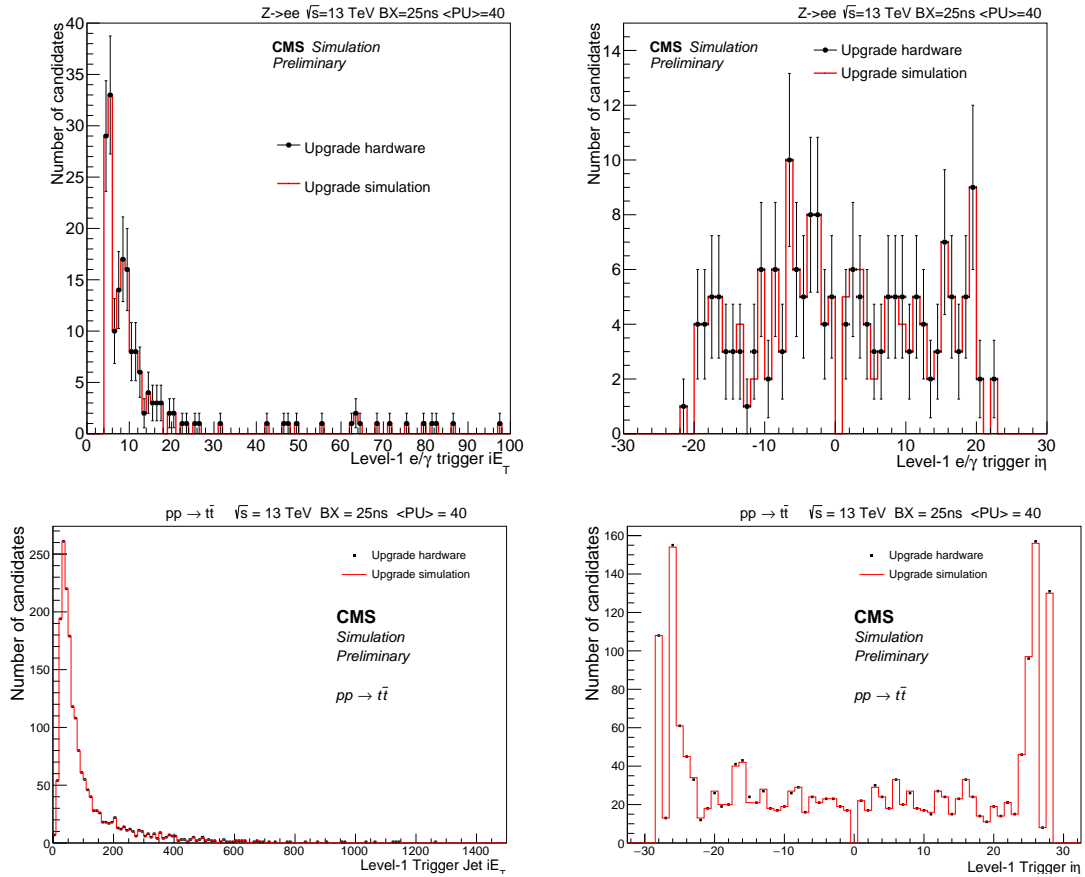


Figure 6. Level-1 e/γ upgraded algorithm firmware output distributions for p_T (top left) and η (top right) produced on $Z \rightarrow ee$ MC events by the Layer-2 firmware. Level-1 jet upgraded algorithm firmware output distributions for p_T (bottom left) and η (bottom right) produced on $t\bar{t}$ MC events. All results are compared to the expected outputs from the simulation.

1 5. Conclusion

2 The TMT approach provides an efficient way to build sophisticated algorithms as described in this
 3 document. The Level-1 e/γ , tau leptons, jets and energy sums algorithms have been presented
 4 and the performance reached are substantially higher than that of the Run 1 system. The upgraded
 5 trigger system is now fully deployed in the CMS service cavern and all the links amongst sub-
 6 components have been installed and validated. The algorithm firmware is now deployed on the
 7 production system. The outputs obtained from injected MC patterns were compared with the ex-
 8 pected simulation and shown to be in excellent agreement. The system is currently running in
 9 parallel alongside the existing calorimeter trigger for full commissioning with real collision data.

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