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For the LHCb collaboration

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- Outlook for Run-II (period after LS1 – 2015–2018)

- LHC running parameters
- LHCb improvements

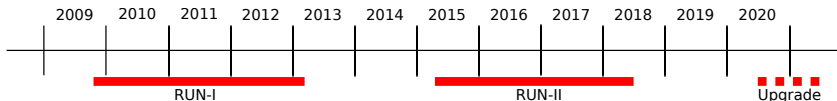
- Upgrade (beyond LS2 – 2020–)

- New detectors
- Trigger and DAQ upgrade
- Higher luminosity

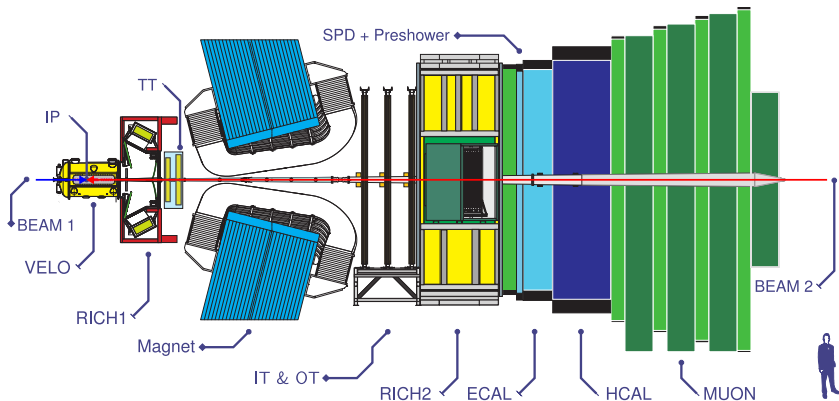
- Benchmark DCPV measurements at LHCb

- γ measurements (tree processes)
- EW penguins (Rare decays): $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $B^+ \rightarrow K^+ \mu^+ \mu^-$
- EW penguins (Radiative decays) $B^0 \rightarrow K^{*0} \gamma$
- Gluonic penguins: ($B \rightarrow VV$ decays) $B_s^0 \rightarrow \phi \phi$

First discuss Run-II and Upgrade, and look at performance for “benchmark channels”



- LHCb: single arm spectrometer with acceptance $2 < \eta < 5$
 - optimized for b physics (rare decays, CPV, spectroscopy)
 - also well suited for c physics



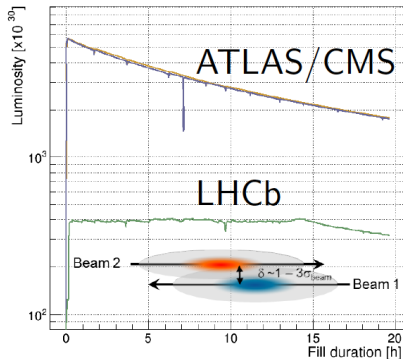
LHCb Run-II prospects 2015–2018

Changes with respect to Run-I

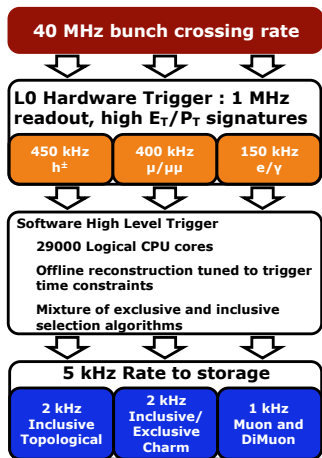
- luminosity at LHCb $L = 4 \times 10^{32} \text{cm}^{-2} \text{s}^{-1}$
 - kept constant within $\pm 5\%$
 - compare to $L = 2 - 4 \times 10^{32} \text{cm}^{-2} \text{s}^{-1}$ for Run-I (2011–2012)
- higher $b\bar{b}$ cross-section at 13 TeV
 - assume factor 1.6
- higher multiplicity and inelastic cross-section
- 25 ns bunch spacing

Run-II: expect $\approx 5 \text{fb}^{-1}$

- Optimize trigger scheme
 - Push decision to later stages in the trigger (improve purity)
 - Effectively more CPU power for the trigger ($\approx 2\times$)



Architecture of LHCb trigger for Run-I



Three stages:

Hardware (“L0”) max output 1 MHz

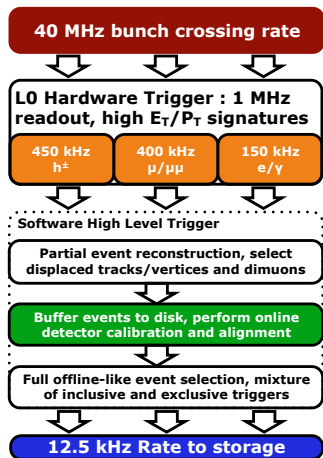
Software (“HLT1”) fast reconstruction and initial selection

Software (“HLT2”) almost off-line reconstruction and final selection

Output rate up to 5 kHz
(design 2 kHz)

Hlt1 and Hlt2 are run in sequence on the same event in one task

Architecture of LHCb trigger for Run-II



Three stages:

Hardware (“L0”) max output 1 MHz

Software (“HLT1”) fast reconstruction and initial selection

Software (“HLT2”) “off-line” reconstruction and final selection

Output rate of HLT2 12.5 kHz

HLT1 writes events to local storage

Do calibration/alignment before HLT2

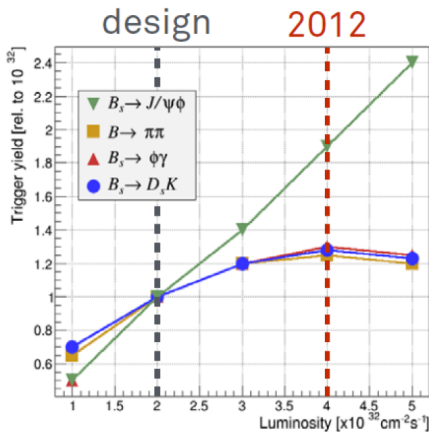
HLT2 can reach off-line performance

enriched selection!

UPGRADE:

Higher luminosity
 Main trigger limitation: 1MHz hardware trigger limit
 Rate, occupancy and radiation damage ...

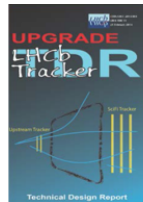
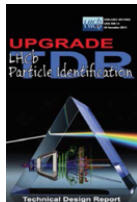
- Remove limitation of 1 MHz hardware trigger
 - All reconstructed in software trigger – 30 MHz
 - Trigger redesign
 - New front end electronics and DAQ system
- Luminosity $L = 2 \times 10^{33} \text{cm}^{-2} \text{s}^{-1}$
- High particle densities
 - Replace sub-detectors to reduce occupancy
- Radiation damage issues



Charm trigger will also profit!

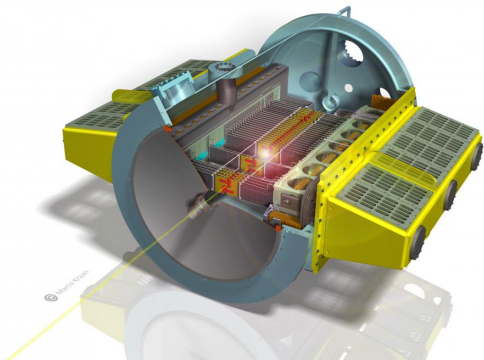
LHCb Upgrade documentation

- LHCb upgrade (LOI): CERN/LHCC 2011-001
- Framework TDR for the LHCb Upgrade (TDR): CERN/LHCC 2012-007
- LHCb VELO Upgrade (TDR): CERN/LHCC 2013-021
- LHCb PID Upgrade (TDR): CERN/LHCC 2013-022
- LHCb Tracker Upgrade (TDR): CERN/LHCC 2014-001
- LHCb Trigger and Online Upgrade (TDR): CERN/LHCC 2014-016



Upgraded Vertex Locator

- 40 MHz readout
- $55 \times 55 \mu\text{m}^2$ pixel size (VeloPix: development from TimePix)
- increased radiation hardness
- microchannel CO_2 cooling

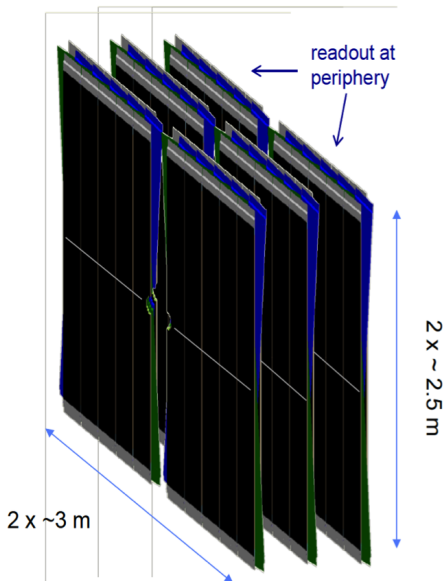


Vertex Detector: replace $r - \phi$ strip geometry with pixel geometry



Scintillating Fibres (SciFi) with Silicon Photo-Multiplier (SiPM) readout

- 250 μm diameter fibres
- fast track reconstruction in trigger
- radiation damage: fibres OK up to 50 fb^{-1}
- dark current: SiPM cooled to -40°



Selected DCPV measurements

A visualization of particle tracks, likely from a detector, showing a central point with many tracks radiating outwards. The tracks are colored in shades of blue and green, and some have small markers along their paths.

Recent LHCb results:

- γ measurements (tree processes) Dalitz plot analysis
- EW penguins (Rare decays): $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $B^+ \rightarrow K^+ \mu^+ \mu^-$
- EW penguins (Radiative decays): $B^0 \rightarrow K^{*0} \gamma$
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γ measurement: Dalitz analysis

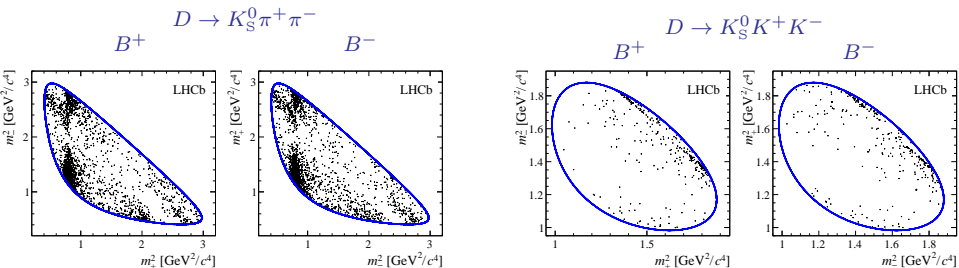
Use predefined binning of Dalitz plot in $B^\pm \rightarrow DK^\pm$ decays with $D \rightarrow K_S^0 h^+ h^-$
Model independent analysis using strong phase difference measured by CLEO-c

[Phys. Rev. D82 (2010) 112006, arXiv:1010.2817]

Measured inputs from $B^0 \rightarrow D^* \mu X$ with $D^* \rightarrow D\pi$ to remove detection asymmetries

Remove model-dependent inputs

Recent LHCb result: [subm. JHEP, \[arXiv:1408.2748\]](#)



m_+^2 and m_-^2 : m^2 of $K_S^0 h^-$ and $K_S^0 h^+$ combinations

Results of γ measurement

Simultaneous fit to all bins to x_{\pm} and y_{\pm} :

$$x_{\pm} \equiv r_B \cos(\delta_B \pm \gamma) \quad \text{and} \quad y_{\pm} \equiv r_B \sin(\delta_B \pm \gamma)$$

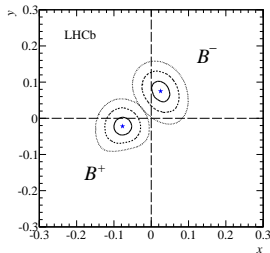
$$x_+ = (-7.7 \pm 2.4 \pm 1.0 \pm 0.4) \times 10^{-2}$$

$$x_- = (2.5 \pm 2.5 \pm 1.0 \pm 0.5) \times 10^{-2}$$

$$y_+ = (-2.2 \pm 2.5 \pm 0.4 \pm 1.0) \times 10^{-2}$$

$$y_- = (7.5 \pm 2.9 \pm 0.5 \pm 1.4) \times 10^{-2}$$

First two uncertainties are statistical and systematic
third is due to strong phases from CLEO-c

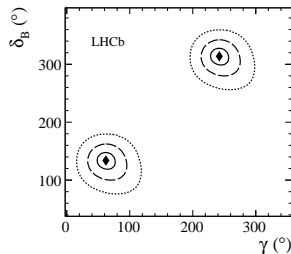
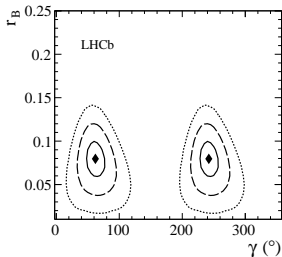


$$r_B = 0.080^{+0.019}_{-0.021}$$

$$\gamma = (62^{+15}_{-14})^\circ$$

$$\delta_B = (134^{+14}_{-15})^\circ$$

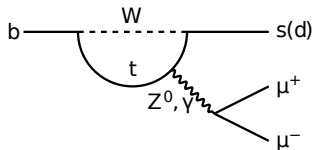
Most precise γ from single
DCPV measurement



Electroweak Penguins

$$B^0 \rightarrow K^{*0} \mu^+ \mu^- \quad \text{and} \quad B^+ \rightarrow K^+ \mu^+ \mu^-$$
$$b \rightarrow s(d) l^+ l^- \text{ decays}$$

flavour changing neutral currents (FCNC)
forbidden at tree level in the SM – only loops and boxes
NP effect can be relatively large



Observables:

Studies of isospin asymmetry LHCb paper [JHEP, arXiv:1403.8044], [JHEP 06 (2014) 133]

Zero crossing point, LHCb: $\mu^+ \mu^-$ forward-backward asymmetry [JHEP 05 (2014) 082] [JHEP 08 (2013) 131]

new measurement of the direct CP asymmetry \mathcal{A}^{CP}

$$\mathcal{A}^{CP} \equiv \frac{\Gamma(\bar{B} \rightarrow \bar{K}^{(*)} \mu^+ \mu^-) - \Gamma(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\Gamma(\bar{B} \rightarrow \bar{K}^{(*)} \mu^+ \mu^-) + \Gamma(B \rightarrow K^{(*)} \mu^+ \mu^-)}$$

LHCb: [subm. JHEP, \[arXiv:1408.0978\]](#)

Electroweak Penguins

$$\mathcal{A}^{CP} \text{ in } B^0 \rightarrow K^{*0} \mu^+ \mu^- \text{ and } B^\pm \rightarrow K^\pm \mu^+ \mu^-$$

Counting experiment: Measure: CP asymmetries

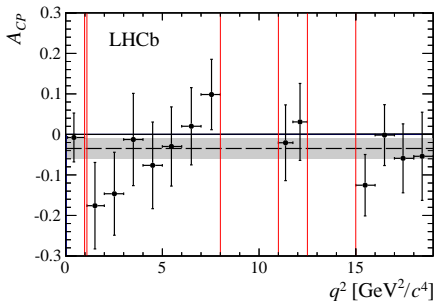
$$\mathcal{A}^{CP}(B^0 \rightarrow K^{*0} \mu^+ \mu^-) = -0.035 \pm 0.024 \pm 0.003,$$

$$\mathcal{A}^{CP}(B^\pm \rightarrow K^\pm \mu^+ \mu^-) = 0.012 \pm 0.017 \pm 0.001,$$

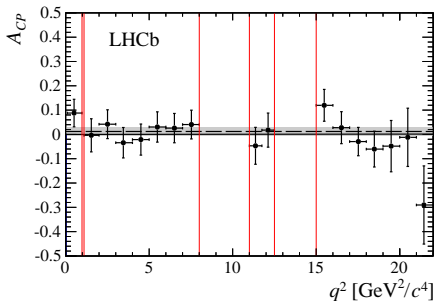
Systematic uncertainties much smaller than statistical errors

SM expects a very small asymmetry: sensitivity to NP

$$B^0 \rightarrow K^{*0} \mu^+ \mu^-$$

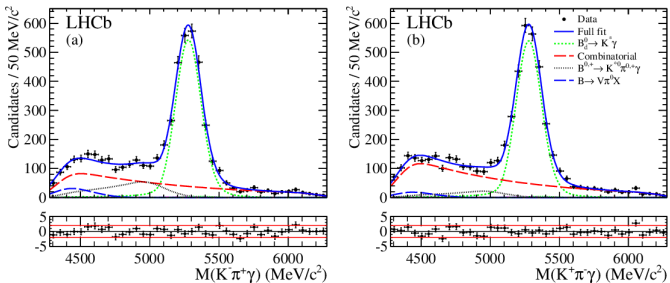


$$B^\pm \rightarrow K^\pm \mu^+ \mu^-$$



Electroweak Penguins Radiative decays

Direct CP asymmetry in $B^0 \rightarrow K^{*0} \gamma$ decay
Electroweak penguin
Sensitive to photon polarization



$$A_{CP}(B^0 \rightarrow K^{*0} \gamma) = (0.8 \pm 1.7 (\text{stat.}) \pm 0.9 (\text{syst.}))\%$$

Experimental syst. error will also improve with larger data samples

LHCb: [Nucl. Phys. B867 (2013) 1-18] (1 fb^{-1} data set)

$$\text{SM prediction: } A_{SM}(B^0 \rightarrow K^{*0} \gamma) = (-0.61 \pm 0.43)\%$$

M. Matsumori, A. I. Sanda, and Y. Y. Keum, Phys. Rev. D72 (2005) 014013, arXiv:hep-ph/0406055

Gluonic Penguins

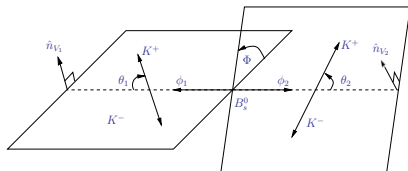
Direct CPV in $B_s^0 \rightarrow \phi \phi$

Triple-product asymmetries in $B_s^0 \rightarrow \phi \phi$ decay:
T-odd combinations

$$U = \sin(\Phi) \cos(\Phi)$$

$$V = \sin(\pm\Phi)$$

$$V \text{ positive if } \cos(\theta_1) \cos(\theta_2) \geq 0$$



CP violation in a decay time integrated method
asymmetries from mixing strongly suppressed: (M. Gronau and J. L. Rosner, Phys. Rev. D84 (2011) 096013, arXiv:1107.1232)
SM close to zero – sensitivity to NP: (A. Datta, M. Duraisamy, and D. London, Phys. Rev. D86 (2012) 076011, arXiv:1207.4495)

new LHCb results: [subm. JHEP, \[arXiv:1407.2222\]](#)

The triple-product asymmetries (decay time integrated fit)

$$A_U = -0.003 \pm 0.017(\text{stat}) \pm 0.006(\text{sys})$$

$$A_V = -0.017 \pm 0.017(\text{stat}) \pm 0.006(\text{sys})$$

Measurement still dominated by statistical uncertainties

Physics prospects for the Upgrade

Type	Observable	LHC Run 1	LHCb 2018	Upgrade	Theory
B_s^0 mixing	$\phi_s(B_s^0 \rightarrow J/\psi \phi)$ (rad)	0.050	0.025	0.009	~ 0.003
	$\phi_s(B_s^0 \rightarrow J/\psi f_0(980))$ (rad)	0.068	0.035	0.012	~ 0.01
	$A_{sl}(B_s^0)$ (10^{-3})	2.8	1.4	0.5	0.03
Gluonic penguin	$\phi_s^{\text{eff}}(B_s^0 \rightarrow \phi \phi)$ (rad)	0.15	0.10	0.023	0.02
	$\phi_s^{\text{eff}}(B_s^0 \rightarrow K^{*0} \bar{K}^{*0})$ (rad)	0.19	0.13	0.029	< 0.02
	$2\beta^{\text{eff}}(B^0 \rightarrow \phi K_S^0)$ (rad)	0.30	0.20	0.04	0.02
Right-handed currents	$\phi_s^{\text{eff}}(B_s^0 \rightarrow \phi \gamma)$	0.20	0.13	0.030	< 0.01
	$\tau^{\text{eff}}(B_s^0 \rightarrow \phi \gamma) / \tau_{B_s^0}$	5%	3.2%	0.8%	0.2 %
Electroweak penguin	$S_3(B^0 \rightarrow K^{*0} \mu^+ \mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.04	0.020	0.007	0.02
	$q_0^2 A_{FB}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)$	10%	5%	1.9%	$\sim 7\%$
	$A_I(K \mu^+ \mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.09	0.05	0.017	~ 0.02
	$\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-) / \mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)$	14%	7%	2.4%	$\sim 10\%$
Higgs penguin	$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$ (10^{-9})	1.0	0.5	0.19	0.3
	$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) / \mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$	220%	110%	40%	$\sim 5\%$
Unitarity triangle	$\gamma(B \rightarrow D^{(*)} K^{(*)})$	7°	4°	1.1°	negligible
angles	$\gamma(B_s^0 \rightarrow D_s^\mp K^\pm)$	17°	11°	2.4°	negligible
	$\beta(B_s^0 \rightarrow J/\psi K_S^0)$	1.7°	0.8°	0.31°	negligible
Charm CP violation	$A_\Gamma(D^0 \rightarrow K^+ K^-)$ (10^{-4})	3.4	2.2	0.5	-
	ΔA_{CP} (10^{-3})	0.8	0.5	0.12	-

LHCb Run 1: 3 fb^{-1}

LHCb 2018: assuming 5 fb^{-1}

LHCb Upgrade: assuming 50 fb^{-1}

Estimate of the theoretical uncertainty



Conclusions

Conclusions

- Many direct CPV measurements are statistically limited after LHCb Run-I
- Run-II will provide three times more $b\bar{b}$ – reduce errors by factor two
- Running after the Upgrade will provide a further order of magnitude more data
- Many systematic uncertainties will also be reduced with more data
- Many of these measurements will not be systematics-limited
- Theoretical errors also small