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## Transverse momentum dependence of inclusive primary charged-particle production in p–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV

The ALICE Collaboration\*

### Abstract

The transverse momentum ( $p_{\text{T}}$ ) distribution of primary charged particles is measured at midrapidity in minimum-bias p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV with the ALICE detector at the LHC in the range  $0.15 < p_{\text{T}} < 50$  GeV/ $c$ . The spectra are compared to the expectation based on binary collision scaling of particle production in pp collisions, leading to a nuclear modification factor consistent with unity for  $p_{\text{T}}$  larger than 2 GeV/ $c$ . The measurement is compared to theoretical calculations and to data in Pb–Pb collisions at  $\sqrt{s_{\text{NN}}} = 2.76$  TeV.

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\*See Appendix A for the list of collaboration members



Measurements of particle production in proton-nucleus collisions at high energies enable the study of fundamental properties of Quantum Chromodynamics (QCD) over a broad range of parton fractional momentum  $x$  and parton densities (see [1] for a review). They also provide reference measurements for the studies of deconfined matter created in nucleus-nucleus collisions [2].

The first measurements of charged-particle production in minimum-bias p–Pb collisions at the LHC at a centre-of-mass energy per nucleon-nucleon pair of  $\sqrt{s_{\text{NN}}} = 5.02$  TeV [3, 4] showed that: i) the charged particle multiplicity density at midrapidity scales approximately with the number of participating nucleons ( $\langle N_{\text{part}} \rangle = 7.9 \pm 0.6$  for minimum-bias collisions) calculated in a Glauber model [5] and ii) the transverse momentum ( $p_{\text{T}}$ ) spectrum, measured in the range 0.5–20 GeV/ $c$  [4], exhibits binary collision scaling above a few GeV/ $c$ , as expected in the absence of any significant nuclear modification effect. The latter is quantified by the nuclear modification factor,  $R_{\text{pPb}}$ , the ratio of the  $p_{\text{T}}$  spectrum in p–Pb collisions and a reference obtained by scaling the measurement in pp collisions with the number of binary nucleon-nucleon collisions in p–Pb. The preliminary result by the CMS collaboration [6] hints at an enhancement of particle production in p–Pb collisions above binary collision scaling, leading to  $R_{\text{pPb}} > 1$ , for  $p_{\text{T}}$  exceeding about 30 GeV/ $c$ . The preliminary result by the ATLAS collaboration [7] exhibits, for collisions corresponding to 0–90% centrality,  $R_{\text{pPb}}$  values slightly above unity and independent of  $p_{\text{T}}$  in the range 4–22 GeV/ $c$ .

In this letter we present an update of our previously published  $p_{\text{T}}$  spectra of primary charged particles [4] based on the 60 times larger sample size collected with the ALICE detector [8] in 2013 in minimum-bias collisions. These data allow a significant extension of the transverse momentum range. The present analysis is essentially identical to the previous and therefore we update only the information related to the enlarged data set; the reader is referred to the earlier publications [4, 9–11] for a more detailed and complete description.

**Table 1:** Systematic uncertainties on the  $p_{\text{T}}$ -differential yields in p–Pb and pp collisions. The quoted ranges span the  $p_{\text{T}}$  dependence of the uncertainties in the measured range, 0.15–50 GeV/ $c$ . Normalization uncertainties are also quoted.

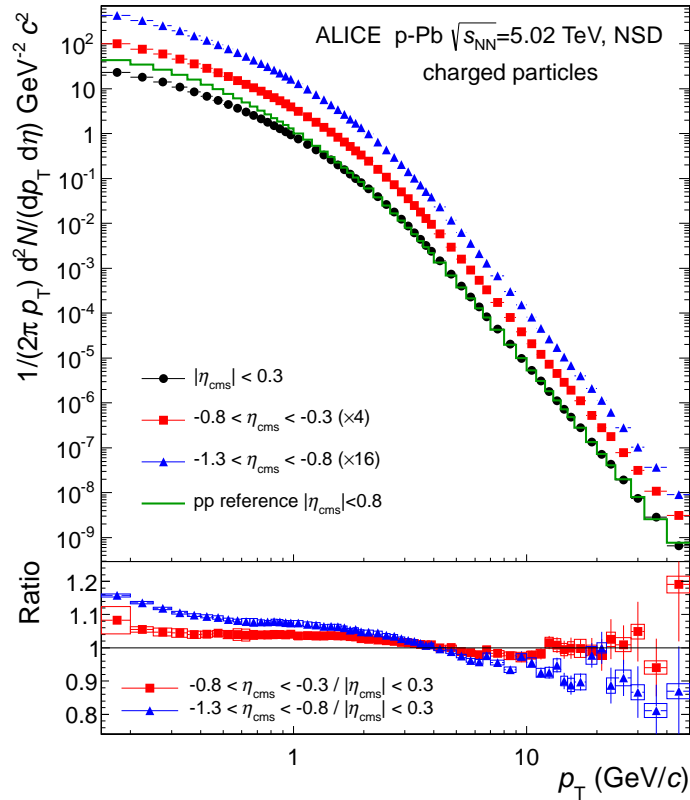
Uncertainty	Value
Event selection	0.6%
Track selection	1.0–5.5%
Tracking efficiency	3.0%
$p_{\text{T}}$ resolution	0–1.3%
$p_{\text{T}}$ scale	0–1.5%
Particle composition	0.1–0.4%
MC generator used for correction	1.0%
Secondary particle rejection	0.5–4.0%
Material budget	0.2–1.5%
Total for p–Pb, $p_{\text{T}}$ -dependent	3.4–6.7%
Normalization p–Pb	3.1%
Total for pp, $p_{\text{T}}$ -dependent	6.8–8.2%
Normalization pp	3.6%
Nuclear overlap $\langle T_{\text{pPb}} \rangle$	3.6%

The ALICE minimum-bias trigger is defined by a coincidence of signals in detectors covering in pseudorapidity<sup>1</sup>  $2.8 < \eta < 5.1$  (VZERO-A) and  $-3.7 < \eta < -1.7$  (VZERO-C). In the 2013 data sample, 106 million events (corresponding to an integrated luminosity of  $50.7 \pm 1.6 \mu\text{b}^{-1}$ ) satisfy the trigger and offline event-selection criteria, which select essentially non-single-diffractive (NSD) minimum-bias col-

<sup>1</sup>In the laboratory frame  $\eta = -\ln[\tan(\vartheta/2)]$ , with  $\vartheta$  the polar angle between the charged particle and the beam axis; the proton beam has negative  $\eta$ .

lisions. The centre-of-mass pseudorapidity is defined as  $\eta_{\text{cms}} = -\eta - |y_{\text{NN}}|$ , with the proton beam at positive rapidity;  $|y_{\text{NN}}| = 0.465$  is the rapidity of the centre-of-mass for nucleon-nucleon collisions. This equation is exact only for massless or very high  $p_{\text{T}}$  particles. The spectra are corrected on a statistical basis using the measurements by ALICE in p–Pb collisions of the  $\eta$  distribution of inclusive charged particle production [3] and of the pion, kaon, and proton yields [12]; this correction depends on the  $\eta_{\text{cms}}$  range and on  $p_{\text{T}}$ , reaching about 20% for the lowest  $p_{\text{T}}$  bin. The systematic uncertainty of the particle composition [12] leads to a systematic uncertainty in our spectra of up to 0.4%.

The systematic uncertainties on the spectra are evaluated as in previous analyses of pp [10], Pb–Pb [9], and p–Pb [4] data. The  $p_{\text{T}}$  scale uncertainty is negligible below 20 GeV/ $c$  and reaches 1.5% at 50 GeV/ $c$ . The main contributions and the total uncertainties are listed in Table 1.



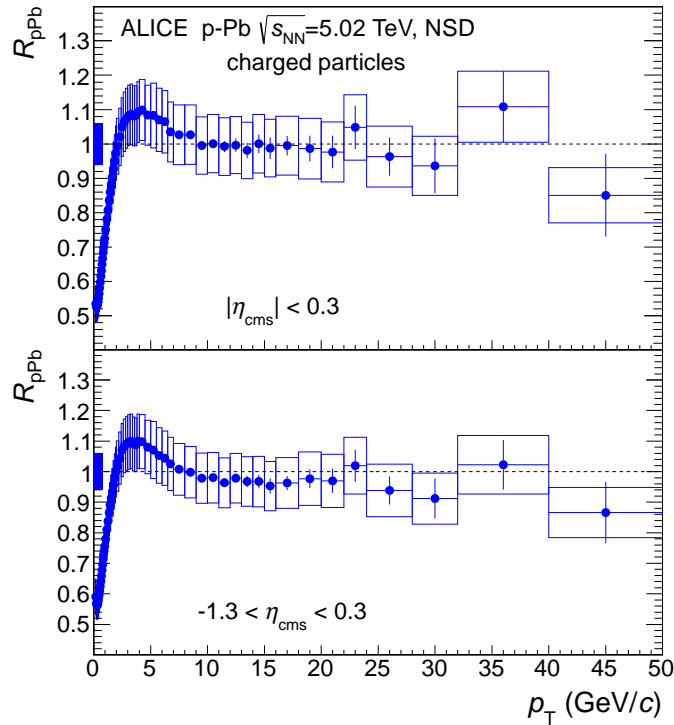
**Fig. 1:** Transverse momentum distributions of charged particles in minimum-bias (NSD) p–Pb collisions for different pseudorapidity ranges (upper panel). The spectra are scaled by the factors indicated. The histogram represents the reference spectrum (cross section scaled by the nuclear overlap function,  $T_{\text{pPb}}$ ) in inelastic pp collisions, determined in  $|\eta| < 0.8$ . The lower panel shows the ratio of spectra in p–Pb at backward pseudorapidities to that at  $|\eta_{\text{cms}}| < 0.3$ . The vertical bars (boxes) represent the statistical (systematic) uncertainties.

The  $p_{\text{T}}$  spectra of charged particles measured in minimum-bias (NSD) p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV are shown in Fig. 1 for the ranges  $|\eta_{\text{cms}}| < 0.3$ ,  $-0.8 < \eta_{\text{cms}} < -0.3$ , and  $-1.3 < \eta_{\text{cms}} < -0.8$ . The pp reference spectrum,  $\langle T_{\text{pPb}} \rangle d^2\sigma^{\text{pp}}/d\eta dp_{\text{T}}$ , is also included.  $\langle T_{\text{pPb}} \rangle$  is the average nuclear overlap function, calculated using the Glauber model [13], which gives  $\langle T_{\text{pPb}} \rangle = \langle N_{\text{coll}} \rangle / \sigma_{\text{NN}} = 0.0983 \pm 0.0035$  mb $^{-1}$ , with  $\langle N_{\text{coll}} \rangle = 6.9 \pm 0.6$  and  $\sigma_{\text{NN}} = 70 \pm 5$  mb. Since the data in pp collisions [10] do not indicate any  $\eta$  dependence of the  $p_{\text{T}}$  spectrum in the range measured by ALICE ( $|\eta| < 0.8$ ) our current reference spectrum is, differently than in [4, 10], for  $|\eta| < 0.8$ . It was obtained by data interpolation at low  $p_{\text{T}}$  and by scaling the measurement at  $\sqrt{s} = 7$  TeV with the ratio of spectra calculated with NLO pQCD at

$\sqrt{s} = 5.02$  and 7 TeV [10].

In the lower panel of Fig. 1 the ratios of the spectra for backward ( $-0.8 < \eta_{\text{cms}} < -0.3$  and  $-1.3 < \eta_{\text{cms}} < -0.8$ ) pseudorapidity ranges to that at  $|\eta_{\text{cms}}| < 0.3$  are shown. The indication of a slight softening of the  $p_T$  spectrum when going from central to backward (Pb-side) pseudorapidity, observed already in the pilot-run data of 2012 [4] (note opposite  $\eta_{\text{cms}}$  sign convention in [4]) is confirmed with better significance and extended in  $p_T$  down to 0.15 GeV/ $c$ .

A good description of our earlier measurement of spectra in p–Pb collisions [4] was achieved in the EPOS3 model [14] including a hydrodynamical description of the collision, while the PHSD model [15] significantly underestimated the spectra for  $p_T$  values of several GeV/ $c$ .



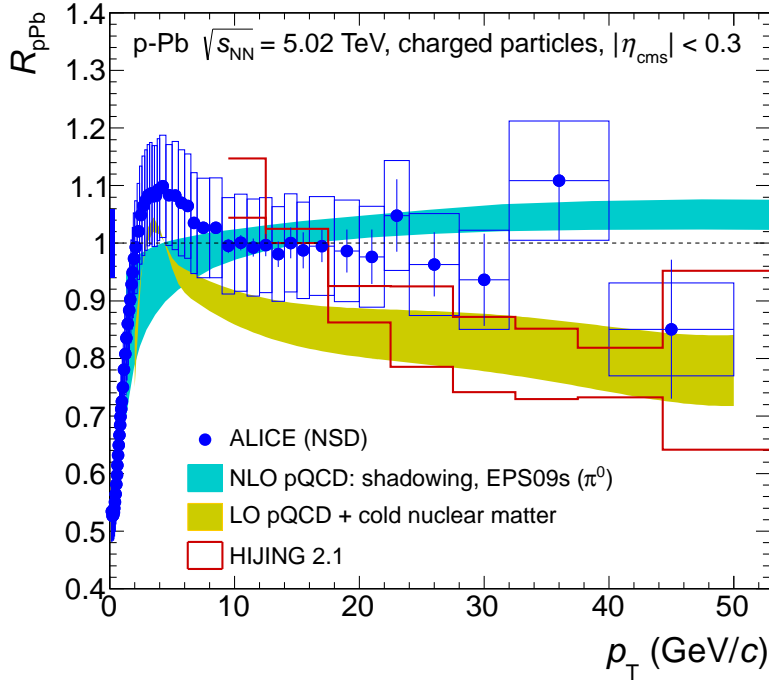
**Fig. 2:** The nuclear modification factor of charged particles as a function of transverse momentum, measured in minimum-bias (NSD) p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV in two pseudorapidity ranges,  $|\eta_{\text{cms}}| < 0.3$  and  $-1.3 < \eta_{\text{cms}} < 0.3$ . The statistical errors are represented by vertical bars, the systematic errors by boxes around data points. The relative systematic uncertainties on the normalization are shown as boxes around unity near  $p_T = 0$ .

In order to quantify nuclear effects in p–Pb collisions, the  $p_T$ -differential yield relative to the pp reference, the nuclear modification factor, is calculated as:

$$R_{\text{pPb}}(p_T) = \frac{d^2 N^{\text{pPb}}/d\eta dp_T}{\langle T_{\text{pPb}} \rangle d^2 \sigma^{\text{pp}}/d\eta dp_T}, \quad (1)$$

where  $N^{\text{pPb}}$  is the charged particle yield in p–Pb collisions.

The measurement of the nuclear modification factor  $R_{\text{pPb}}$  for charged particle production in  $|\eta_{\text{cms}}| < 0.3$  and  $-0.3 < \eta_{\text{cms}} < 1.3$  is shown in Fig. 2. The uncertainties of the p–Pb and pp spectra are added in quadrature, separately for the statistical and systematic uncertainties. The systematic uncertainties are largely correlated between adjacent  $p_T$  bins. The total systematic uncertainty on the normalization, the quadratic sum of the uncertainty on  $\langle T_{\text{pPb}} \rangle$ , the normalization of the pp reference spectrum and



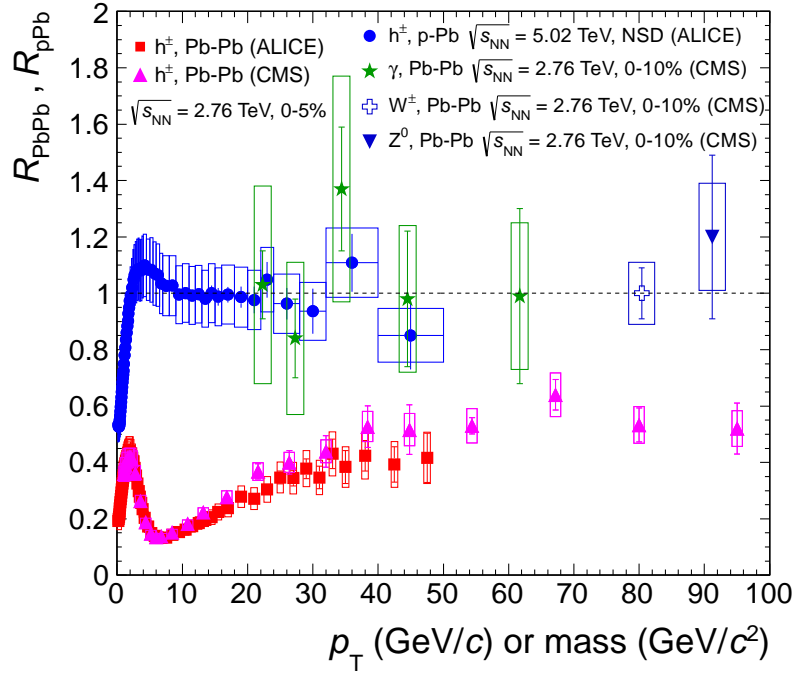
**Fig. 3:** Transverse momentum dependence of the nuclear modification factor  $R_{pPb}$  of charged particles measured in minimum-bias (NSD) p–Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. The ALICE data in  $|\eta_{cms}| < 0.3$  (symbols) are compared to model calculations [18–20] (bands, see text for details). The vertical bars (boxes) show the statistical (systematic) uncertainties. The relative systematic uncertainty on the normalization is shown as a box around unity near  $p_T = 0$ .

the normalization of the p–Pb data, amounts to 6.0%. The  $R_{pPb}$  factor is consistent with unity up to  $p_T = 50$  GeV/c. The average values of  $R_{pPb}$  in  $|\eta_{cms}| < 0.3$  are  $0.995 \pm 0.007$  (stat.)  $\pm 0.084$  (syst.) for the  $p_T$  range 10–20 GeV/c,  $0.990 \pm 0.031$  (stat.)  $\pm 0.090$  (syst.) in the range 20–28 GeV/c and  $0.969 \pm 0.056$  (stat.)  $\pm 0.090$  (syst.) in the range 28–50 GeV/c. The systematic uncertainties are weighted averages of the values in  $p_T$  bins, with statistical uncertainties as inverse square weights; all values carry in addition the common overall normalization uncertainty of 6%.

The data indicate a small enhancement,  $R_{pPb}$  above unity, barely significant within systematic errors, around 4 GeV/c, i.e. in the  $p_T$  region where the much stronger Cronin enhancement is seen at lower energies [16, 17].

The p–Pb data provide important constraints to models of nuclear modification effects. As an illustration, in Fig. 3 the measurement of  $R_{pPb}$  at  $|\eta_{cms}| < 0.3$  is compared to theoretical model predictions. The predictions for shadowing [18], calculated at next-to-leading order (NLO) with the EPS09s nuclear modification of parton distribution functions describe the data well for  $p_T \gtrsim 6$  GeV/c (the calculations are for  $\pi^0$ ). The LO pQCD model including cold nuclear matter effects [19] exhibits a distinct trend of decreasing  $R_{pPb}$ , which is not supported by the data. The prediction with the HIJING 2.1 model, shown for two fragmentation schemes [20], exhibits a more pronounced trend of decreasing  $R_{pPb}$  at high  $p_T$ . It is interesting to note that calculations with the EPOS LHC model [21], not included here, show a similar trend. Several predictions based on the saturation (Color Glass Condensate) model are available [22–24]; they were shown previously [4] to describe the  $R_{pPb}$  data for  $p_T$  of a few GeV/c.

In Fig. 4 we compare the measurement of the nuclear modification factor for inclusive primary charged-particle ( $h^\pm$ ) production in p–Pb collisions to that in central (0–5% centrality) Pb–Pb collisions [9, 25]. The p–Pb data demonstrate that the suppression of hadron production at high  $p_T$  in Pb–Pb collisions,



**Fig. 4:** Transverse momentum dependence of the nuclear modification factor  $R_{pPb}$  of charged particles ( $h^\pm$ ) measured in minimum-bias (NSD) p–Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV in comparison to data on the nuclear modification factor  $R_{PbPb}$  in central Pb–Pb collisions. The Pb–Pb data are for charged particle [9, 25], direct photon [26],  $Z^0$  [27] and  $W^\pm$  [28] production. All data are for midrapidity.

understood in theoretical models as a consequence of parton energy loss in (deconfined) QCD matter (see [9] and references therein), has no contribution from initial state effects. The ALICE p–Pb data show no sign of nuclear matter modification of hadron production at high  $p_T$  and are therefore fully consistent with the observation of binary collision scaling in Pb–Pb of observables which are not affected by hot QCD matter (direct photons [26] and vector bosons [27, 28])

In summary, we have extended our measurements of the charged-particle  $p_T$  spectra and nuclear modification factor in minimum-bias (NSD) p–Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. The results, covering a substantially-extended  $p_T$  range,  $0.15 < p_T < 50$  GeV/c, exhibit, within uncertainties, no deviation from binary collision scaling at high  $p_T$ ; the nuclear modification factor remains consistent with unity for  $p_T \gtrsim 2$  GeV/c. The data are described by a prediction based on NLO pQCD calculations with PDF shadowing and further underline our earlier observation [4] that initial state effects do not contribute to the strong suppression of hadron production at high  $p_T$  observed at the LHC in Pb–Pb collisions.

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## A The ALICE Collaboration

B. Abelev<sup>69</sup>, J. Adam<sup>37</sup>, D. Adamová<sup>77</sup>, M.M. Aggarwal<sup>81</sup>, M. Agnello<sup>105,88</sup>, A. Agostinelli<sup>26</sup>, N. Agrawal<sup>44</sup>, Z. Ahammed<sup>124</sup>, N. Ahmad<sup>18</sup>, I. Ahmed<sup>15</sup>, S.U. Ahn<sup>62</sup>, S.A. Ahn<sup>62</sup>, I. Aimo<sup>105,88</sup>, S. Aiola<sup>129</sup>, M. Ajaz<sup>15</sup>, A. Akhmedov<sup>53</sup>, S.N. Alam<sup>124</sup>, D. Aleksandrov<sup>94</sup>, B. Alessandro<sup>105</sup>, D. Alexandre<sup>96</sup>, A. Alici<sup>12,99</sup>, A. Alkin<sup>3</sup>, J. Alme<sup>35</sup>, T. Alt<sup>39</sup>, S. Altinpinar<sup>17</sup>, I. Altsybeev<sup>123</sup>, C. Alves Garcia Prado<sup>113</sup>, C. Andrei<sup>72</sup>, A. Andronic<sup>91</sup>, V. Anguelov<sup>87</sup>, J. Anielski<sup>49</sup>, T. Antičić<sup>92</sup>, F. Antinori<sup>102</sup>, P. Antonioli<sup>99</sup>, L. Aphecetche<sup>107</sup>, H. Appelshäuser<sup>48</sup>, S. Arcelli<sup>26</sup>, N. Armesto<sup>16</sup>, R. 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A. Nyanin<sup>94</sup>, J. Nystrand<sup>17</sup>, H. Oeschler<sup>87</sup>, S. Oh<sup>129</sup>, S.K. Oh<sup>vi,40</sup>, A. Okatan<sup>63</sup>, L. Olah<sup>128</sup>, J. Oleniacz<sup>126</sup>, A.C. Oliveira Da Silva<sup>113</sup>, J. Onderwaater<sup>91</sup>, C. Oppedisano<sup>105</sup>, A. Ortiz Velasquez<sup>32</sup>, A. Oskarsson<sup>32</sup>, J. Otwinowski<sup>91</sup>, K. Oyama<sup>87</sup>, P. Sahoo<sup>45</sup>, Y. Pachmayer<sup>87</sup>, M. Pachr<sup>37</sup>, P. Pagano<sup>29</sup>, G. Paic<sup>58</sup>, F. Painke<sup>39</sup>, C. Pajares<sup>16</sup>, S.K. Pal<sup>124</sup>, A. Palmeri<sup>101</sup>, D. Pant<sup>44</sup>, V. Papikyan<sup>1</sup>, G.S. Pappalardo<sup>101</sup>, P. Pareek<sup>45</sup>, W.J. Park<sup>91</sup>, S. Parmar<sup>81</sup>, A. Passfeld<sup>49</sup>, D.I. Patalakha<sup>106</sup>, V. Paticchio<sup>98</sup>, B. Paul<sup>95</sup>, T. Pawlak<sup>126</sup>, T. Peitzmann<sup>52</sup>, H. Pereira Da Costa<sup>14</sup>, E. Pereira De Oliveira Filho<sup>113</sup>, D. Peresunko<sup>94</sup>, C.E. Pérez Lara<sup>75</sup>, A. Pesci<sup>99</sup>, V. Peskov<sup>48</sup>, Y. Pestov<sup>5</sup>, V. Petráček<sup>37</sup>, M. Petran<sup>37</sup>, M. Petris<sup>72</sup>, M. Petrovici<sup>72</sup>, C. Petta<sup>27</sup>, S. Piano<sup>104</sup>, M. Pikna<sup>36</sup>, P. Pillot<sup>107</sup>, O. Pinazza<sup>99,34</sup>, L. Pinsky<sup>115</sup>, D.B. Piyarathna<sup>115</sup>, M. Płoskoń<sup>68</sup>, M. Planinic<sup>121,92</sup>, J. Pluta<sup>126</sup>, S. Pochybova<sup>128</sup>, P.L.M. Podesta-Lerma<sup>112</sup>, M.G. Poghosyan<sup>34</sup>, E.H.O. Pohjoisaho<sup>42</sup>, B. Polichtchouk<sup>106</sup>, N. Poljak<sup>92</sup>, A. Pop<sup>72</sup>, S. Porteboeuf-Houssais<sup>64</sup>, J. Porter<sup>68</sup>, B. Potukuchi<sup>84</sup>, S.K. Prasad<sup>127</sup>, R. Preghenella<sup>99,12</sup>, F. Prino<sup>105</sup>, C.A. Pruneau<sup>127</sup>, I. Pshenichnov<sup>51</sup>, G. Puddu<sup>23</sup>, P. Pujahari<sup>127</sup>, V. Punin<sup>93</sup>, J. Putschke<sup>127</sup>, H. Qvigstad<sup>21</sup>, A. Rachevski<sup>104</sup>, S. Raha<sup>4</sup>, J. Rak<sup>116</sup>, A. Rakotozafindrabe<sup>14</sup>, L. Ramello<sup>30</sup>, R. Raniwala<sup>85</sup>, S. Raniwala<sup>85</sup>, S.S. Räsänen<sup>42</sup>, B.T. Rascanu<sup>48</sup>, D. Rathee<sup>81</sup>, A.W. Rauf<sup>15</sup>, V. Razazi<sup>23</sup>, K.F. Read<sup>118</sup>, J.S. Real<sup>65</sup>, K. Redlich<sup>vii,71</sup>, R.J. Reed<sup>129</sup>, A. Rehman<sup>17</sup>, P. Reichelt<sup>48</sup>, M. Reicher<sup>52</sup>, F. Reidt<sup>34</sup>, R. Renfordt<sup>48</sup>, A.R. Reolon<sup>66</sup>, A. Reshetin<sup>51</sup>, F. Rettig<sup>39</sup>, J.-P. Revol<sup>34</sup>, K. Reygers<sup>87</sup>, V. Riabov<sup>79</sup>, R.A. Ricci<sup>67</sup>, T. Richert<sup>32</sup>, M. Richter<sup>21</sup>, P. Riedler<sup>34</sup>, W. Riegler<sup>34</sup>, F. Riggi<sup>27</sup>, A. Rivetti<sup>105</sup>, E. Rocco<sup>52</sup>, M. Rodríguez Cahuantzi<sup>2</sup>, A. Rodríguez Manso<sup>75</sup>, K. Røed<sup>21</sup>, E. Rogochaya<sup>61</sup>, S. Rohni<sup>84</sup>, D. Rohr<sup>39</sup>, D. Röhrich<sup>17</sup>, R. Romita<sup>76</sup>, F. Ronchetti<sup>66</sup>, L. Ronflette<sup>107</sup>, P. Rosnet<sup>64</sup>, A. Rossi<sup>34</sup>, F. Roukoutakis<sup>82</sup>, A. Roy<sup>45</sup>, C. Roy<sup>50</sup>, P. Roy<sup>95</sup>, A.J. Rubio Montero<sup>10</sup>, R. Rui<sup>24</sup>, R. Russo<sup>25</sup>, E. Ryabinkin<sup>94</sup>, Y. Ryabov<sup>79</sup>, A. Rybicki<sup>110</sup>, S. Sadovsky<sup>106</sup>, K. Šafařík<sup>34</sup>, B. Sahlmuller<sup>48</sup>, R. Sahoo<sup>45</sup>, P.K. Sahu<sup>56</sup>, J. Saini<sup>124</sup>, S. Sakai<sup>68</sup>, C.A. Salgado<sup>16</sup>, J. Salzwedel<sup>19</sup>, S. Sambyal<sup>84</sup>, V. Samsonov<sup>79</sup>, X. Sanchez Castro<sup>50</sup>, F.J. Sánchez Rodríguez<sup>112</sup>, L. Šándor<sup>54</sup>, A. Sandoval<sup>59</sup>, M. Sano<sup>120</sup>, G. Santagati<sup>27</sup>, D. Sarkar<sup>124</sup>, E. Scapparone<sup>99</sup>, F. Scarlassara<sup>28</sup>, R.P. Scharenberg<sup>89</sup>, C. Schiaua<sup>72</sup>, R. Schicker<sup>87</sup>, C. Schmidt<sup>91</sup>, H.R. Schmidt<sup>33</sup>, S. Schuchmann<sup>48</sup>, J. Schukraft<sup>34</sup>, M. Schulc<sup>37</sup>, T. Schuster<sup>129</sup>, Y. Schutz<sup>107,34</sup>, K. Schwarz<sup>91</sup>, K. Schweda<sup>91</sup>, G. Scioli<sup>26</sup>, E. Scomparin<sup>105</sup>, R. Scott<sup>118</sup>, G. Segato<sup>28</sup>, J.E. Seger<sup>80</sup>, Y. Sekiguchi<sup>119</sup>, I. Selyuzhenkov<sup>91</sup>, J. Seo<sup>90</sup>, E. Serradilla<sup>10,59</sup>, A. Sevcenco<sup>57</sup>, A. Shabetai<sup>107</sup>, G. Shabratova<sup>61</sup>, R. Shahoyan<sup>34</sup>, A. Shangaraev<sup>106</sup>, N. Sharma<sup>118</sup>, S. Sharma<sup>84</sup>, K. Shigaki<sup>43</sup>, K. Shtejer<sup>25</sup>, Y. Sibiraki<sup>94</sup>, S. Siddhanta<sup>100</sup>, T. Siemiarczuk<sup>71</sup>, D. Silvermyr<sup>78</sup>, C. Silvestre<sup>65</sup>, G. Simatovic<sup>121</sup>, R. Singaraju<sup>124</sup>, R. Singh<sup>84</sup>, S. Singha<sup>124,73</sup>, V. Singhal<sup>124</sup>, B.C. Sinha<sup>124</sup>, T. Sinha<sup>95</sup>, B. Sitar<sup>36</sup>, M. Sitta<sup>30</sup>, T.B. Skaali<sup>21</sup>, K. Skjerdal<sup>17</sup>, M. Slupecki<sup>116</sup>, N. Smirnov<sup>129</sup>, R.J.M. Snellings<sup>52</sup>, C. Søgaard<sup>32</sup>, R. Soltz<sup>69</sup>, J. Song<sup>90</sup>, M. Song<sup>130</sup>, F. Soramel<sup>28</sup>, S. Sorensen<sup>118</sup>, M. Spacek<sup>37</sup>, E. Spiriti<sup>66</sup>, I. Sputowska<sup>110</sup>, M. Spyropoulou-Stassinaki<sup>82</sup>, B.K. Srivastava<sup>89</sup>, J. Stachel<sup>87</sup>, I. Stan<sup>57</sup>, G. Stefanek<sup>71</sup>, M. Steinpreis<sup>19</sup>, E. Stenlund<sup>32</sup>, G. Steyn<sup>60</sup>, J.H. Stiller<sup>87</sup>, D. Stocco<sup>107</sup>, M. Stolpovskiy<sup>106</sup>, P. Strmen<sup>36</sup>, A.A.P. Suaide<sup>113</sup>, T. Sugitate<sup>43</sup>, C. Suire<sup>46</sup>, M. Suleymanov<sup>15</sup>, R. Sultanov<sup>53</sup>, M. Šumbera<sup>77</sup>, T. Susa<sup>92</sup>, T.J.M. Symons<sup>68</sup>, A. Szabo<sup>36</sup>, A. Szanto de Toledo<sup>113</sup>, I. Szarka<sup>36</sup>, A. Szczepankiewicz<sup>34</sup>, M. Szymanski<sup>126</sup>, J. Takahashi<sup>114</sup>, M.A. Tangaro<sup>31</sup>, J.D. Tapia Takaki<sup>viii,46</sup>, A. Tarantola Peloni<sup>48</sup>, A. Tarazona Martinez<sup>34</sup>, M.G. Tarczila<sup>72</sup>, A. Tauro<sup>34</sup>, G. Tejeda Muñoz<sup>2</sup>, A. Telesca<sup>34</sup>, C. Terrevoli<sup>23</sup>, J. Thäder<sup>91</sup>, D. Thomas<sup>52</sup>, R. Tieulent<sup>122</sup>, A.R. Timmins<sup>115</sup>, A. Toia<sup>102</sup>, V. Trubnikov<sup>3</sup>, W.H. Trzaska<sup>116</sup>, T. Tsuji<sup>119</sup>, A. Tumkin<sup>93</sup>, R. Turrisi<sup>102</sup>, T.S. Tveter<sup>21</sup>, K. Ullaland<sup>17</sup>, A. Uras<sup>122</sup>, G.L. Usai<sup>23</sup>, M. Vajzer<sup>77</sup>, M. Vala<sup>54,61</sup>, L. Valencia Palomo<sup>64</sup>, S. Vallero<sup>87</sup>, P. Vande Vyvre<sup>34</sup>, J. Van Der Maarel<sup>52</sup>, J.W. Van Hoorne<sup>34</sup>, M. van Leeuwen<sup>52</sup>, A. Vargas<sup>2</sup>, M. Vargyas<sup>116</sup>, R. Varma<sup>44</sup>, M. Vasileiou<sup>82</sup>, A. Vasiliev<sup>94</sup>, V. Vechernin<sup>123</sup>, M. Veldhoen<sup>52</sup>, A. Velure<sup>17</sup>, M. Venaruzzo<sup>24,67</sup>, E. Vercellin<sup>25</sup>, S. Vergara Limón<sup>2</sup>, R. Vernet<sup>8</sup>, M. Verweij<sup>127</sup>, L. Vickovic<sup>109</sup>, G. Viesti<sup>28</sup>, J. Viinikainen<sup>116</sup>, Z. Vilakazi<sup>60</sup>, O. Villalobos Baillie<sup>96</sup>, A. Vinogradov<sup>94</sup>, L. Vinogradov<sup>123</sup>, Y. Vinogradov<sup>93</sup>, T. Virgili<sup>29</sup>, Y.P. Viyogi<sup>124</sup>, A. Vodopyanov<sup>61</sup>, M.A. Völkl<sup>87</sup>, K. Voloshin<sup>53</sup>, S.A. Voloshin<sup>127</sup>, G. Volpe<sup>34</sup>, B. von Haller<sup>34</sup>, I. Vorobyev<sup>123</sup>, D. Vranic<sup>91,34</sup>, J. Vrláková<sup>38</sup>, B. Vulpescu<sup>64</sup>, A. Vyushin<sup>93</sup>, B. Wagner<sup>17</sup>, J. Wagner<sup>91</sup>, V. Wagner<sup>37</sup>, M. Wang<sup>7,107</sup>, Y. Wang<sup>87</sup>, D. Watanabe<sup>120</sup>, M. Weber<sup>115</sup>, J.P. Wessels<sup>49</sup>, U. Westerhoff<sup>49</sup>, J. Wiechula<sup>33</sup>, J. Wikne<sup>21</sup>, M. Wilde<sup>49</sup>, G. Wilk<sup>71</sup>, J. Wilkinson<sup>87</sup>, M.C.S. Williams<sup>99</sup>, B. Windelband<sup>87</sup>, M. Winn<sup>87</sup>, C.G. Yaldo<sup>127</sup>, Y. Yamaguchi<sup>119</sup>, H. Yang<sup>52</sup>, P. Yang<sup>7</sup>, S. Yang<sup>17</sup>, S. Yano<sup>43</sup>, S. Yasnopolskiy<sup>94</sup>, J. Yi<sup>90</sup>, Z. Yin<sup>7</sup>, I.-K. Yoo<sup>90</sup>, I. Yushmanov<sup>94</sup>, V. Zaccolo<sup>74</sup>, C. Zach<sup>37</sup>, A. Zaman<sup>15</sup>, C. Zampolli<sup>99</sup>, S. Zaporozhets<sup>61</sup>, A. Zarochentsev<sup>123</sup>, P. Závada<sup>55</sup>, N. Zaviyalov<sup>93</sup>, H. Zbroszczyk<sup>126</sup>, I.S. Zgura<sup>57</sup>, M. Zhalov<sup>79</sup>, H. Zhang<sup>7</sup>, X. Zhang<sup>7,68</sup>, Y. Zhang<sup>7</sup>, C. Zhao<sup>21</sup>, N. Zhigareva<sup>53</sup>, D. Zhou<sup>7</sup>, F. Zhou<sup>7</sup>, Y. Zhou<sup>52</sup>, Zhou, Zhuo<sup>17</sup>, H. Zhu<sup>7</sup>, J. Zhu<sup>7</sup>, X. Zhu<sup>7</sup>, A. Zichichi<sup>12,26</sup>, A. Zimmermann<sup>87</sup>, M.B. Zimmermann<sup>49,34</sup>, G. Zinovjev<sup>3</sup>, Y. Zoccarato<sup>122</sup>, M. Zyzak<sup>48</sup>

## Affiliation notes

<sup>i</sup> Deceased



- ii Also at: St. Petersburg State Polytechnical University
- iii Also at: Department of Applied Physics, Aligarh Muslim University, Aligarh, India
- iv Also at: M.V. Lomonosov Moscow State University, D.V. Skobeltsyn Institute of Nuclear Physics, Moscow, Russia
- v Also at: University of Belgrade, Faculty of Physics and "Vinča" Institute of Nuclear Sciences, Belgrade, Serbia
- vi Permanent Address: Permanent Address: Konkuk University, Seoul, Korea
- vii Also at: Institute of Theoretical Physics, University of Wrocław, Wrocław, Poland
- viii Also at: University of Kansas, Lawrence, KS, United States

## Collaboration Institutes

- 1 A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia
- 2 Benemérita Universidad Autónoma de Puebla, Puebla, Mexico
- 3 Bogolyubov Institute for Theoretical Physics, Kiev, Ukraine
- 4 Bose Institute, Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Kolkata, India
- 5 Budker Institute for Nuclear Physics, Novosibirsk, Russia
- 6 California Polytechnic State University, San Luis Obispo, CA, United States
- 7 Central China Normal University, Wuhan, China
- 8 Centre de Calcul de l'IN2P3, Villeurbanne, France
- 9 Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba
- 10 Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
- 11 Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
- 12 Centro Fermi - Museo Storico della Fisica e Centro Studi e Ricerche "Enrico Fermi", Rome, Italy
- 13 Chicago State University, Chicago, USA
- 14 Commissariat à l'Energie Atomique, IRFU, Saclay, France
- 15 COMSATS Institute of Information Technology (CIIT), Islamabad, Pakistan
- 16 Departamento de Física de Partículas and IGFAE, Universidad de Santiago de Compostela, Santiago de Compostela, Spain
- 17 Department of Physics and Technology, University of Bergen, Bergen, Norway
- 18 Department of Physics, Aligarh Muslim University, Aligarh, India
- 19 Department of Physics, Ohio State University, Columbus, OH, United States
- 20 Department of Physics, Sejong University, Seoul, South Korea
- 21 Department of Physics, University of Oslo, Oslo, Norway
- 22 Dipartimento di Fisica dell'Università 'La Sapienza' and Sezione INFN Rome, Italy
- 23 Dipartimento di Fisica dell'Università and Sezione INFN, Cagliari, Italy
- 24 Dipartimento di Fisica dell'Università and Sezione INFN, Trieste, Italy
- 25 Dipartimento di Fisica dell'Università and Sezione INFN, Turin, Italy
- 26 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Bologna, Italy
- 27 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Catania, Italy
- 28 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Padova, Italy
- 29 Dipartimento di Fisica 'E.R. Caianiello' dell'Università and Gruppo Collegato INFN, Salerno, Italy
- 30 Dipartimento di Scienze e Innovazione Tecnologica dell'Università del Piemonte Orientale and Gruppo Collegato INFN, Alessandria, Italy
- 31 Dipartimento Interateneo di Fisica 'M. Merlin' and Sezione INFN, Bari, Italy
- 32 Division of Experimental High Energy Physics, University of Lund, Lund, Sweden
- 33 Eberhard Karls Universität Tübingen, Tübingen, Germany
- 34 European Organization for Nuclear Research (CERN), Geneva, Switzerland
- 35 Faculty of Engineering, Bergen University College, Bergen, Norway
- 36 Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia
- 37 Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
- 38 Faculty of Science, P.J. Šafárik University, Košice, Slovakia
- 39 Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 40 Gangneung-Wonju National University, Gangneung, South Korea

- 41 Gauhati University, Department of Physics, Guwahati, India
- 42 Helsinki Institute of Physics (HIP), Helsinki, Finland
- 43 Hiroshima University, Hiroshima, Japan
- 44 Indian Institute of Technology Bombay (IIT), Mumbai, India
- 45 Indian Institute of Technology Indore, Indore (IITI), India
- 46 Institut de Physique Nucléaire d'Orsay (IPNO), Université Paris-Sud, CNRS-IN2P3, Orsay, France
- 47 Institut für Informatik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 48 Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 49 Institut für Kernphysik, Westfälische Wilhelms-Universität Münster, Münster, Germany
- 50 Institut Pluridisciplinaire Hubert Curien (IPHC), Université de Strasbourg, CNRS-IN2P3, Strasbourg, France
- 51 Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
- 52 Institute for Subatomic Physics of Utrecht University, Utrecht, Netherlands
- 53 Institute for Theoretical and Experimental Physics, Moscow, Russia
- 54 Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia
- 55 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- 56 Institute of Physics, Bhubaneswar, India
- 57 Institute of Space Science (ISS), Bucharest, Romania
- 58 Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 59 Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 60 iThemba LABS, National Research Foundation, Somerset West, South Africa
- 61 Joint Institute for Nuclear Research (JINR), Dubna, Russia
- 62 Korea Institute of Science and Technology Information, Daejeon, South Korea
- 63 KTO Karatay University, Konya, Turkey
- 64 Laboratoire de Physique Corpusculaire (LPC), Clermont Université, Université Blaise Pascal, CNRS-IN2P3, Clermont-Ferrand, France
- 65 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS-IN2P3, Grenoble, France
- 66 Laboratori Nazionali di Frascati, INFN, Frascati, Italy
- 67 Laboratori Nazionali di Legnaro, INFN, Legnaro, Italy
- 68 Lawrence Berkeley National Laboratory, Berkeley, CA, United States
- 69 Lawrence Livermore National Laboratory, Livermore, CA, United States
- 70 Moscow Engineering Physics Institute, Moscow, Russia
- 71 National Centre for Nuclear Studies, Warsaw, Poland
- 72 National Institute for Physics and Nuclear Engineering, Bucharest, Romania
- 73 National Institute of Science Education and Research, Bhubaneswar, India
- 74 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- 75 Nikhef, National Institute for Subatomic Physics, Amsterdam, Netherlands
- 76 Nuclear Physics Group, STFC Daresbury Laboratory, Daresbury, United Kingdom
- 77 Nuclear Physics Institute, Academy of Sciences of the Czech Republic, Řež u Prahy, Czech Republic
- 78 Oak Ridge National Laboratory, Oak Ridge, TN, United States
- 79 Petersburg Nuclear Physics Institute, Gatchina, Russia
- 80 Physics Department, Creighton University, Omaha, NE, United States
- 81 Physics Department, Panjab University, Chandigarh, India
- 82 Physics Department, University of Athens, Athens, Greece
- 83 Physics Department, University of Cape Town, Cape Town, South Africa
- 84 Physics Department, University of Jammu, Jammu, India
- 85 Physics Department, University of Rajasthan, Jaipur, India
- 86 Physik Department, Technische Universität München, Munich, Germany
- 87 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- 88 Politecnico di Torino, Turin, Italy
- 89 Purdue University, West Lafayette, IN, United States
- 90 Pusan National University, Pusan, South Korea
- 91 Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt, Germany
- 92 Rudjer Bošković Institute, Zagreb, Croatia

- 93 Russian Federal Nuclear Center (VNIIEF), Sarov, Russia
- 94 Russian Research Centre Kurchatov Institute, Moscow, Russia
- 95 Saha Institute of Nuclear Physics, Kolkata, India
- 96 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- 97 Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
- 98 Sezione INFN, Bari, Italy
- 99 Sezione INFN, Bologna, Italy
- 100 Sezione INFN, Cagliari, Italy
- 101 Sezione INFN, Catania, Italy
- 102 Sezione INFN, Padova, Italy
- 103 Sezione INFN, Rome, Italy
- 104 Sezione INFN, Trieste, Italy
- 105 Sezione INFN, Turin, Italy
- 106 SSC IHEP of NRC Kurchatov institute, Protvino, Russia
- 107 SUBATECH, Ecole des Mines de Nantes, Université de Nantes, CNRS-IN2P3, Nantes, France
- 108 Suranaree University of Technology, Nakhon Ratchasima, Thailand
- 109 Technical University of Split FESB, Split, Croatia
- 110 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Cracow, Poland
- 111 The University of Texas at Austin, Physics Department, Austin, TX, USA
- 112 Universidad Autónoma de Sinaloa, Culiacán, Mexico
- 113 Universidade de São Paulo (USP), São Paulo, Brazil
- 114 Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
- 115 University of Houston, Houston, TX, United States
- 116 University of Jyväskylä, Jyväskylä, Finland
- 117 University of Liverpool, Liverpool, United Kingdom
- 118 University of Tennessee, Knoxville, TN, United States
- 119 University of Tokyo, Tokyo, Japan
- 120 University of Tsukuba, Tsukuba, Japan
- 121 University of Zagreb, Zagreb, Croatia
- 122 Université de Lyon, Université Lyon 1, CNRS/IN2P3, IPN-Lyon, Villeurbanne, France
- 123 V. Fock Institute for Physics, St. Petersburg State University, St. Petersburg, Russia
- 124 Variable Energy Cyclotron Centre, Kolkata, India
- 125 Vestfold University College, Tonsberg, Norway
- 126 Warsaw University of Technology, Warsaw, Poland
- 127 Wayne State University, Detroit, MI, United States
- 128 Wigner Research Centre for Physics, Hungarian Academy of Sciences, Budapest, Hungary
- 129 Yale University, New Haven, CT, United States
- 130 Yonsei University, Seoul, South Korea
- 131 Zentrum für Technologietransfer und Telekommunikation (ZTT), Fachhochschule Worms, Worms, Germany