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# CMS Conference Report

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## Small- $x$ QCD studies with CMS at the LHC

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### Abstract

The capabilities of the CMS experiment to study the low- $x$  parton structure and QCD evolution in the proton and the nucleus at LHC energies are presented through four different measurements, to be carried out in Pb-Pb at  $\sqrt{s_{NN}} = 5.5$  TeV: (i) the charged hadron rapidity density  $dN_{ch}/d\eta$  and (ii) the ultraperipheral (photo)production of  $\Upsilon$ ; and in p-p at  $\sqrt{s} = 14$  TeV: (iii) inclusive forward jets and (iv) Mueller-Navelet dijets (separated by  $\Delta\eta \gtrsim 8$ ).

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# Introduction

At high energies, the cross-sections of all *hadronic objects* (protons, nuclei, or even photons “fluctuating” into  $q\bar{q}$  vector states) are dominated by scatterings involving gluons. Gluons clearly outnumber quarks in the small momentum fraction (low- $x$ ) range of the parton distribution functions (PDFs) as a consequence of the QCD parton splitting probabilities described by the DGLAP [1] and BFKL [2] evolution equations. The fast growth of the gluon densities  $xG(x, Q^2)$  for decreasing  $x$  conspicuously observed in DIS  $ep$  at HERA [3], cannot however continue indefinitely since this would violate unitarity even for scatterings with  $Q^2 \gg \Lambda_{\text{QCD}}^2$ . For small enough  $x$  values, gluons must start to recombine in a process known as gluon saturation [4]. This phenomenon occurs when the size occupied by the partons becomes similar to the size of the hadron  $\pi R^2$ , or in terms of the *saturation momentum*  $Q_s$  when:  $Q^2 \lesssim Q_s^2(x) \simeq \alpha_s xG(x, Q^2)/\pi R^2$ .  $Q_s$  grows with the number  $A$  of nucleons in the “target”, the collision energy  $\sqrt{s}$ , and the rapidity of the gluon  $y = \ln(1/x)$ , according to:  $Q_s^2 \sim A^{1/3} x^{-0.3} \sim A^{1/3} (\sqrt{s})^{0.3} \sim A^{1/3} e^{0.3y}$ . The  $A$  dependence implies that, at equal energies, saturation effects will be enhanced by factors as large as  $A^{1/3} \approx 6$  in a heavy nucleus ( $A = 208$  for Pb) compared to protons. Theoretically, the regime of low- $x$  QCD can be effectively described in the “Color Glass Condensate” (CGC) framework, where all gluon fusions and multiple scatterings are “resummed” into classical high-density gluon wavefunctions [5]. The corresponding evolution is given in this case by the BK/JIMWLK [6] *non-linear* equations.

Experimentally, the most direct way to access the low- $x$  PDFs in hadronic collisions is by measuring perturbative probes (heavy- $Q$ , jets, high- $p_T$  hadrons, prompt  $\gamma$ , ...) at large  $\sqrt{s}$  and *forward* rapidities [7]. For a  $2 \rightarrow 2$  parton scattering, the *minimum*  $x$  probed in a process with a particle of momentum  $p_T$  produced at pseudo-rapidity  $\eta$ , is  $x_2^{\text{min}} = x_T e^{-\eta}/(2 - x_T e^{\eta})$  where  $x_T = 2p_T/\sqrt{s}$ . Thus,  $x_2^{\text{min}}$  decreases by a factor of  $\sim 10$  every 2 units of rapidity. The experimental capabilities of the CMS experiment are extremely well adapted for the study of low- $x$  phenomena with proton and ion beams. The acceptance of the CMS/TOTEM system is the largest ever available in a collider, and the detector is designed to measure different particles with excellent momentum resolution [8]: jets ( $|\eta| < 6.6$ ),  $\gamma$  and  $e^\pm$  ( $|\eta| < 3$ ), muons ( $|\eta| < 2.5$ ), hadrons ( $|\eta| < 6.6$ ), plus neutrals in the Zero-Degree Calorimeters (ZDCs,  $|\eta| > 8.3$ ). We present four observables measurable in CMS which are sensitive to parton saturation effects in the proton and nucleus wave-functions at LHC energies.

## 1 Measurements in PbPb collisions at $\sqrt{s_{NN}} = 5.5$ TeV

### (1) Charged hadron PbPb rapidity density: $dN_{ch}/d\eta$

In high-energy heavy-ion collisions, the *hadron* rapidity density  $dN/d\eta$  is directly related to the number of initially released *partons* at a given  $\eta$ . CGC approaches which effectively take into account a reduced initial parton flux in the nuclear PDFs, reproduce successfully the absolute hadron yields (as well as their centrality and  $\sqrt{s_{NN}}$  dependences) at SPS – RHIC energies [9, 10]. At LHC, the expected PbPb multiplicities at  $\eta = 0$  are  $dN/d\eta|_{\eta=0} \approx 2000$  (Fig. 1, left). CMS simulation studies from hit counting in the innermost Si pixel layer ( $|\eta| < 2.5$ ) indicate that the occupancy remains less than 2% and that, on an event-by-event basis, the reconstructed  $dN_{ch}/d\eta$  is within  $\sim 2\%$  of the true primary multiplicity (Fig. 1, right) [11].

### (2) $\Upsilon$ photoproduction in ultra-peripheral PbPb ( $\rightarrow \gamma \text{Pb}$ ) $\rightarrow \Upsilon + \text{Pb}^* \text{Pb}^{(*)}$ collisions

Ultra-peripheral collisions (UPCs) of heavy ions generate strong electromagnetic fields (equivalent to a flux of quasi-real photons) which can be used to constrain the low- $x$  behaviour of  $xG(x, Q^2)$  via  $Q\bar{Q}$  photoproduction [13]. Lead beams at 2.75 TeV have Lorentz factors  $\gamma = 2930$  leading to maximum photon energies  $\omega_{\text{max}} \approx \gamma/R \sim 100$  GeV, and c.m. energies  $W_{\gamma\gamma}^{\text{max}} \approx 160$  GeV and  $W_{\gamma A}^{\text{max}} \approx 1$  TeV. The  $x$  values probed in  $\gamma\text{Pb} \rightarrow \Upsilon\text{Pb}$  processes at  $y = 2.5$  can be as low as  $x \sim 10^{-4}$ . Full simulation+reconstruction [11] of input distributions from the STARLIGHT MC [14] show that CMS can measure  $\Upsilon \rightarrow e^+e^-, \mu^+\mu^-$  within  $|\eta| < 2.5$ , in UPCs tagged with neutrons detected in the ZDCs. Figure 2 shows the reconstructed  $dN/dm_{l^+l^-}$  around the  $\Upsilon$  mass for an integrated PbPb luminosity of  $0.5 \text{ nb}^{-1}$ . With a total yield of  $\sim 400$   $\Upsilon$ , detailed  $p_T$  and  $\eta$  studies can be carried out, to constrain the low- $x$  gluon density in the Pb nucleus.

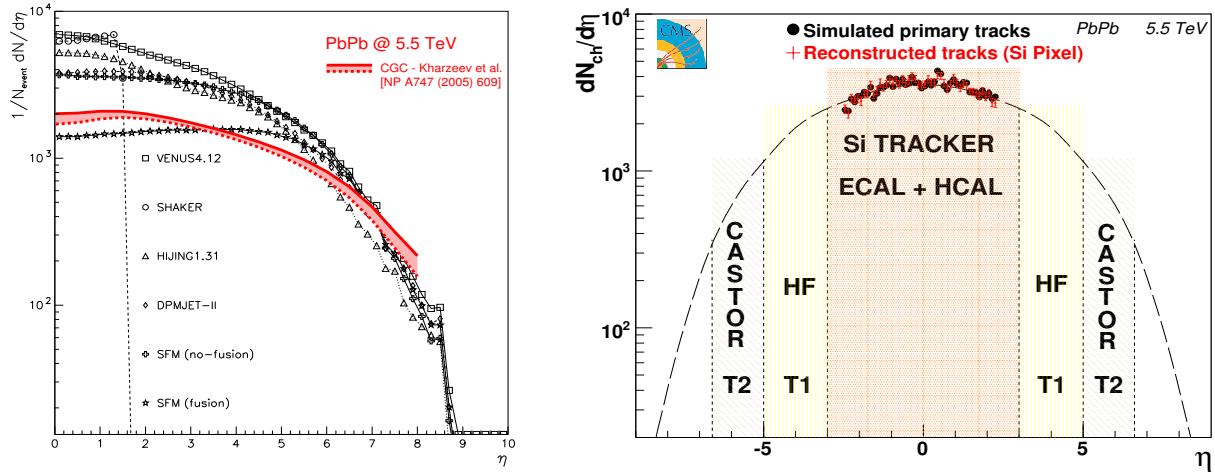


Figure 1: Left: Model predictions for  $dN/d\eta$  in central PbPb at the LHC [9, 12]. Right: Range of charged particle rapidities covered by CMS (Si tracker, HF, CASTOR) and TOTEM (T1, T2 trackers). The HIJING PbPb distribution of primary simulated tracks within  $|\eta| < 2.5$  (black dots) is compared to the reconstructed hits in layer 1 of the Si tracker (red crosses) [11].

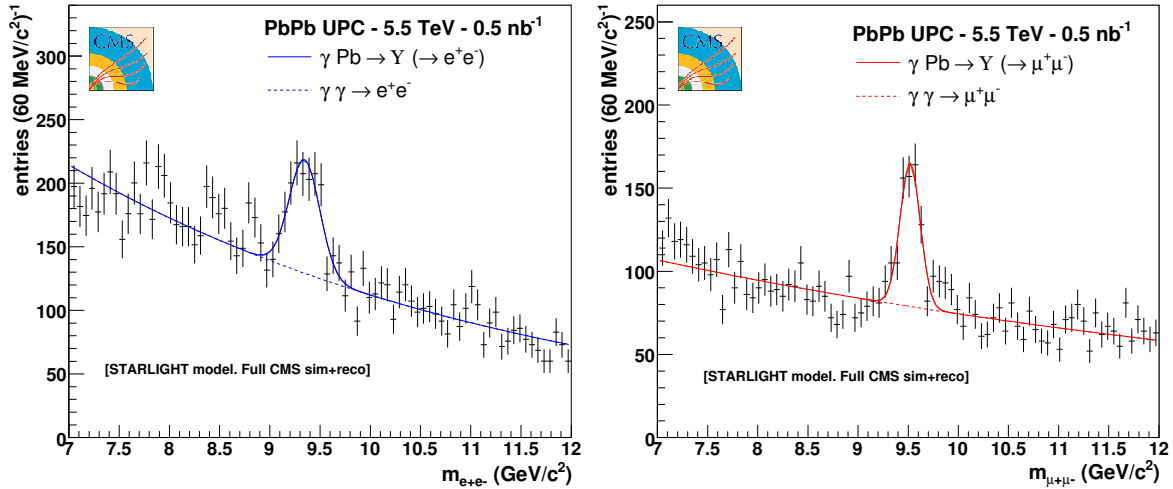


Figure 2: Expected  $e^+e^-$  (left),  $\mu^+\mu^-$  (right) invariant mass distributions from  $\gamma Pb \rightarrow Y Pb^*$  ( $Y \rightarrow l^+ l^-$ , signal) and  $\gamma\gamma \rightarrow l^+ l^-$  (background) in UPC PbPb at  $\sqrt{s_{NN}} = 5.5$  TeV in CMS.

## 2 Measurements in pp collisions at $\sqrt{s} = 14$ TeV

### (3) Inclusive forward jet production: $pp \rightarrow \text{jet}+X$ , with $3 < |\eta_{\text{jet}}| < 5$

Jet measurements at Tevatron have provided valuable information on the proton PDFs. At 14 TeV, the production of jets with  $E_T \approx 20\text{--}100$  GeV in the CMS forward calorimeters (HF and CASTOR) probes the PDFs down to  $x_2 \approx 10^{-6}$  [7]. Figure 3-left shows the single inclusive jet spectrum in both HF's ( $3 < |\eta| < 5$ ) expected for a short first run with just  $1 \text{ pb}^{-1}$  integrated luminosity. The spectrum has been obtained from a preliminary study using PYTHIA 6.403 with jet reconstruction at the *particle-level* (i.e. no detector effects are included apart from the HF tower  $\eta$ - $\phi$  granularity) [15]. Although at such low  $E_T$ 's systematic uncertainties can be as large as  $\sim 30\%$ , the available statistics for this study is very high.

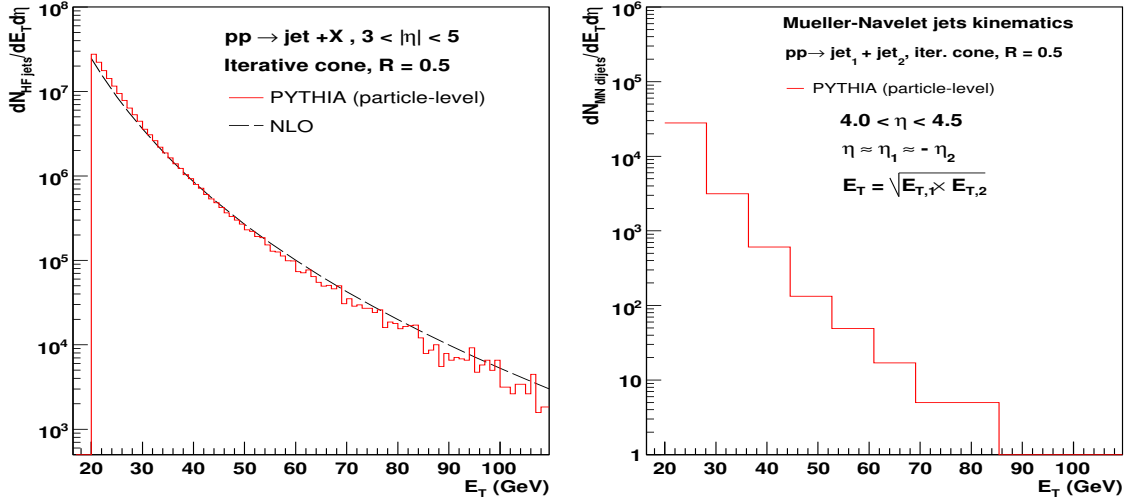


Figure 3: Expected jet yields in pp at  $\sqrt{s} = 14$  TeV ( $1 \text{ pb}^{-1}$ ) obtained from PYTHIA 6.403 at the particle-level (*no* full detector response, underlying event or hadronization corrections). Left: Single inclusive jets in HF ( $3 < |\eta| < 5$ ) (compared to a NLO calculation with scale  $\mu = E_T$  [16]). Right: Dijets separated by  $\Delta\eta = 8-9$  with the Müller-Navelet kinematics cuts described in the text, as a function of  $E_T \equiv \sqrt{E_{T,1} \times E_{T,2}}$ .

#### (4) Mueller-Navelet dijets: $pp \rightarrow \text{jet}_1 + \text{jet}_2$ , with large $\Delta\eta = \eta_2 - \eta_1$

Inclusive dijet production at large pseudorapidity intervals – Müller-Navelet (MN) jets – has been considered an excellent testing ground for BFKL [17] and non-linear QCD [18] evolutions. The large rapidity separation between partons enhances the available longitudinal momentum phase space for BFKL radiation. Gluon saturation effects are expected to reduce the (pure BFKL) MN cross section by a factor of  $\sim 2$  for jets separated by  $\Delta\eta \approx 9$  [18]. In order to estimate the expected statistics for  $1 \text{ pb}^{-1}$  (first run without pile-up), we have selected the PYTHIA events which pass the MN kinematics cuts:  $|E_{T,1} - E_{T,2}| < 2.5$  GeV,  $|\eta_1| - |\eta_2| < 0.25$ , and  $\Delta\eta = 6 - 10$  [15]. Figure 3-right shows the results for  $\Delta\eta = 8-9$ . The expected dijet yields for this  $\eta$  separation indicate that these studies are statistically feasible at the LHC.

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