Neutrino Factory Downstream Systems

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Abstract

We describe the Neutrino Factory accelerator systems downstream from the target and capture area. These include the bunching and phase rotation, cooling, acceleration, and decay ring systems. We also briefly discuss the R&D program under way to develop these systems, and indicate areas where help from CERN would be invaluable.

1 Introduction

A muon-based Neutrino Factory (NF) will be a powerful tool for the experimentalists. Design and performance estimates for such a facility have been ongoing for about 10 years. This effort is fully international, and includes The Neutrino Factory and Muon Collider Collaboration in the U.S., the UK Neutrino Factory group and Work Package 3 of the EUROnu Design Study in Europe, and the Japan Neutrino Factory Working Group in Asia.

Here, we will consider the Neutrino Factory systems downstream of the target and capture region, namely the bunching, phase rotation, cooling, acceleration, and decay ring systems. The upstream systems are covered in the paper by Kirk [1].

2 Muon Accelerator Description

2.1 Advantages

Muon-based facilities can address several of the key outstanding particle physics questions that can be addressed with accelerator-based experiments. In the neutrino sector, the neutrino beam is derived from the decays of either μ^+ or μ^- circulating in a decay ring. The kinematics of such decays is well known, and there are minimal hadronic uncertainties in the spectrum and flux. The neutrino beam contains high-energy electron neutrinos, above the τ production threshold, so this channel can be observed in the experiments. Oscillations from $\nu_e \rightarrow \nu_\mu$ give rise to easily detectable "wrong-sign" muons (i.e., muons whose sign is opposite to that of the circulating beam in the decay ring).

At the energy frontier, the fact that the muon is a point particle means that the full beam energy is available for particle production. Moreover, because the muon mass is much greater than that of the electron, a Muon Collider has almost no synchrotron radiation. This results in a narrower energy spread at the interaction point than occurs at an electron–positron collider, and it permits a circular Muon Collider that uses its expensive rf equipment more efficiently and has a small footprint that can fit on an existing laboratory site.

2.2 Challenges

There are two main challenges associated with producing a muon beam:

1) muons are created as a tertiary beam $(p \rightarrow \pi \rightarrow \mu)$

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2) muons have a very short lifetime, 2.2 µs at rest

The first challenge results in a low production rate, and thus requires a multi-MW proton source, and a target capable of handling this power. Moreover, the production process gives rise to a beam with a very large energy spread and transverse emittance. This, in turn, means that solenoidal focusing is preferred for the initial portions of the facility (since solenoids focus in both planes simultaneously), that some form of emittance cooling is needed, and that a high-acceptance acceleration system and decay ring are required.

The second challenge means that rapid beam manipulations are mandatory. This implies that the presently untested ionization cooling technique is needed to reduce the beam emittance, which creates a need for high-gradient rf cavities that can operate in a strong solenoidal field. It also implies the need for a rapid acceleration system after the cooling channel. There is an additional aspect of the rapid decay of muons that affects the facility design. Decay electrons emitted by the beam in the muon decay ring create a large heat load in the mid-plane of the superconducting magnets.

2.3 Ingredients of a Neutrino Factory

A NF comprises a number of sections, including the proton driver, the target, capture, and decay section, the bunching and phase rotation section, the cooling section, the acceleration section, and finally the decay ring. Figure 1 shows the layout representing the baseline NF design of the International Design Study of a Neutrino Factory (IDS-NF) [2]. In what follows we describe the downstream systems, starting from the bunching and phase rotation.

2.3.1 Bunching and Phase Rotation

Because the beam from the upstream target and capture section is unsuitable for the downstream accelerators, it must be "conditioned" before it can be transported further. The conditioning involves a rotation in longitudinal phase space to reduce the energy spread and creation of a 201-MHz bunch train suitable for capture in the cooling and acceleration systems to follow. The task is accomplished with a series of rf cavities whose frequencies decrease from about 325 to 201 MHz along the channel in a prescribed way. Optimization of the length and performance of this section is in progress.

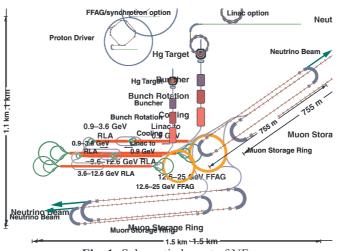


Fig. 1: Schematic layout of NF.

2.3.2 Ionization Cooling

As part of the International Scoping Study of a Neutrino Factory [3], the performance of the various cooling channel designs was compared. As a result of that comparison, the so-called "Feasibility Study 2" (FS2) channel [4] was found to perform the best. When coupled with a proton driver providing 4 MW of 5–15 GeV protons (in 2 ns bunches), this channel, illustrated schematically in Fig. 2, meets the requirement of 10²¹ useful muon decays per 10⁷ s year. The channel transmits both muon signs simultaneously, interleaved at different phases of the 201-MHz rf system. It is worth noting that the actual implementation of such a channel, such as that being built for the Muon Ionization Cooling Experiment (MICE) [5], is considerably more complicated than the simple "physicist's view" illustrated in Fig. 2.

2.3.3 Acceleration

The baseline scheme adopted for the IDS-NF (see Fig. 3) comprises a pre-acceleration linac that increases the beam energy to 0.9 GeV. This is followed by a pair of "dog-bone" Recirculating Linear Accelerators (RLAs) that raise the energy from 0.9–3.6 GeV and from 3.6–12.6 GeV, respectively. The last stage of acceleration, from 12.6–25 GeV, is provided by a non-scaling Fixed-Field, Alternating Gradient (FFAG) ring. Optics for the linac, the RLAs, the required injection chicanes, and the transfer lines have been completed [6].

The last system, the non-scaling FFAG ring, presents several technical challenges. First, there is a strong coupling between the longitudinal and transverse dynamics that is important for the emittance values of interest to a muon FFAG. Different amplitude particles have different flight paths and thus different flight times. The result is that large-amplitude particles slip out of phase with the rf and are not fully accelerated. Using sextupoles to effect a partial chromatic correction has been shown [7] to ameliorate the problem and appears workable.

The second challenge concerns injection and extraction into the densely packed FFAG ring. Progress on this topic has recently been reported by Pasternack [8], though the kicker and septum magnet specifications remain daunting.

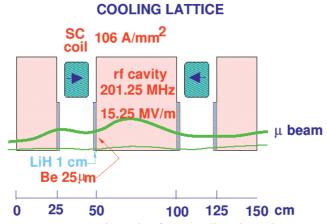


Fig. 2: Schematic of FS2 cooling channel.

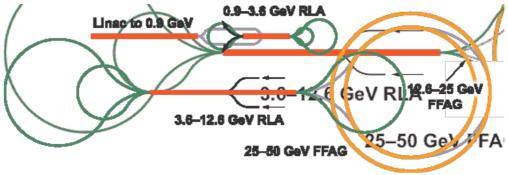


Fig. 3: Schematic of IDS-NF baseline acceleration system.

2.3.3.1 Racetrack Ring

As suggested by the name, a racetrack ring has two long straight section that can be aimed at a single detector site. If there are two baselines, there are two independent rings. The facility could be operated with one muon species in each ring, with the two species interchanged periodically, or it could be operated with both species counter-rotating in each ring. Although it is likely to be more expensive, this configuration is very flexible compared with the triangle case (see Section 2.3.3.2), in the sense that it can be used for two detector sites with no spatial constraints, and can be used with the full beam even if only one detector is operating. For these reasons, it has been adopted as the IDS-NF baseline [2,3].

2.3.3.2 Triangle Ring

In this configuration [3], two rings would be stacked side-by-side in a single tunnel, with one ring storing each muon sign. A typical layout is shown in Fig. 4. Two detectors can be illuminated with interleaved trains of positive and negative muon decay products. In terms of the percentage of circumference available for decays, the triangle ring is somewhat more efficient than the racetrack design. However, its geometry constrains the locations of the two baselines. If the two sites are in the same direction from the ring, or if only a single site is used, the racetrack is preferred.

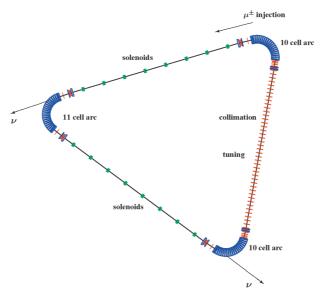


Fig. 4: Schematic of NF triangle ring.

3 R&D Program

A substantial R&D program is now under way to validate the various design choices for a NF. Broadly, the program can be separated into three categories:

- Simulations
- Component Development
- System Tests

This program is being carried out worldwide, managed via loose, but effective, international coordination.

Simulations include studies and optimization of the accelerator design. At present, this activity is being carried out under the auspices of the IDS-NF. Component development includes the development of rf cavities, magnets, and liquid-hydrogen absorbers suitable for use in a NF. The key issue at present is the degradation in maximum gradient observed for room-temperature cavities immersed in a strong axial magnetic field. Finally, system tests involve proof-of-concept tests to validate the overall performance and cost of critical subsystems. Because such tests require substantial resources, they are typically carried out by means of international collaborations.

3.1 Overview of Technical Issues

Each R&D category has its own issues to examine. In the simulations area, the main tasks include optimization of the subsystem designs and eventually end-to-end tracking of the entire facility. For component development, the critical R&D topics include development of normal conducting rf cavities capable of operating in a strong axial magnetic field, development of low-frequency superconducting rf (SRF) cavities, development of fast, wide aperture kicker magnets for the FFAG ring, and development of decay ring magnets that can tolerate the high mid-plane heat load resulting from muon decay electrons. System tests include the recently completed MERcury Intense Target (MERIT) experiment at CERN [9], the Muon Ionization Cooling Experiment (MICE) [5], presently under way at RAL, and the non-scaling FFAG experiment EMMA [10], under way at Daresbury Laboratory. At some future time, a 6D cooling experiment to serve as proof-of-concept for a Muon Collider cooling channel will undoubtedly be carried out.

3.2 IDS-NF

The goal of the IDS-NF program is to deliver, in a 2013 time frame, a Reference Design Report (RDR) for a NF. In the RDR, the physics performance of the facility will be detailed, and the specification of each of the required accelerator, diagnostics, and detector systems will be defined. A cost estimate for the facility described in the RDR will be included in the document. The present baseline NF design shown in Fig. 1 is a result of this design effort. The EU contributes strongly to this effort via its EUROnu design study program.

3.3 Normal Conducting RF

As mentioned earlier, the main challenge for operation of a NF cooling channel (and also the bunching and phase rotation sections) is the operation of rf cavities in a strong solenoidal focusing field. As shown in Ref. [11], for vacuum rf cavities the maximum stable gradient decreases as the magnetic field is increased. Interestingly, an rf cavity filled with high-pressure hydrogen gas shows no such degradation [12]. Present plans call for investigating different materials, such as beryllium, for the vacuum cavities, and measuring the response of a cavity filled with high-pressure hydrogen to an intense beam of ionizing radiation.

3.4 MICE

The MICE experiment [5], now under way at RAL, aims to design, engineer, fabricate, and test with muons, a section of a realistic NF cooling channel. Detailed comparisons with simulation codes will be made to validate the codes as a design tool for an eventual facility. The cooling channel components are being fabricated now, and are already providing information on both the cost and complexity of a muon cooling channel.

3.5 EMMA

EMMA, being fabricated at Daresbury Laboratory in the UK as a primarily UK effort, will test an electron model of a non-scaling FFAG ring. This represents the first test of such a device and will serve to demonstrate the feasibility of the concept. While it is not designed to be a scale model of a muon FFAG, the EMMA ring will serve to investigate longitudinal dynamics, transmission, emittance growth, and the influence of resonances on the beam. As shown in Fig. 5, the components are mostly fabricated and are now being installed. Commissioning is anticipated to commence early in 2010.

4 Possibilities for CERN Participation

There are many areas where expertise from CERN could—and hopefully will—make substantial and necessary contributions to the R&D program outlined in this paper. We list these below:

- 1. Target facility design. CERN staff have experience in making estimates of the shielding needed for a 4-MW target facility. They also have the necessary skills to explore the safety and environmental aspects of the proposed mercury-jet target and beam dump system. Developing a robust target facility design from a safety perspective would greatly benefit from CERN involvement, both in defining the requirements and in preparing a design to satisfy them.
- 2. *Proton driver design*. The design of this portion of the NF facility is expected to be site dependent. CERN could develop a site-specific design based on the SPL operating at 4 MW with ~2 ns bunches.
- 3. Engineering and costing of key components. CERN engineering staff, especially those completing the LHC project, are world experts in normal conducting and superconducting rf systems, cryogenic systems, superconducting magnets, and fast kicker magnets. All these components are needed for the NF design.



Fig. 5: EMMA accelerator components installed on a girder.

It is worth noting that CERN staff made key contributions to both the NF design and the design of the MICE experiment in the "early days." The intellectual effort provided then is greatly missed and is badly needed.

In the longer term, CERN participation in an international 6D cooling experiment for a MC would be of great value to the scientific community. A Muon Collider is a larger and more complicated facility than a Neutrino Factory, and its construction would surely be carried out as an international endeavor .

5 Summary

R&D toward both a NF and a MC is making steady progress, with strong EU contributions in many areas. MERIT has established the ability of a mercury jet to tolerate 4 MW of protons. MICE is progressing toward a demonstration of muon ionization cooling in a few years, and EMMA will begin commissioning shortly. In order to develop realistic technical designs and their corresponding cost estimates, CERN help in component engineering and cost estimating will be critical. Development of muon-based accelerator facilities offers great scientific promise and remains a worthy—and challenging—goal to pursue. It seems only natural that CERN, representing the premier particle physics laboratory in the world, would play a significant role in such a program.

Acknowledgments

I am grateful for the continued R&D efforts of my colleagues in the Neutrino Factory and Muon Collider Collaboration and the IDS-NF whose work has resulted in the design and R&D progress summarized here.

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Synergies between the needs of LHC, neutrinos and Radio-active Ion Beams

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Abstract

An extensive upgrade of the LHC injectors is planned for reaching the ultimate potential of the collider during the second phase of its life, after 2020. As a result, the beam delivered by the SPS for fixed target physics will also be improved, in particular to the benefit of conventional neutrino beams. Moreover, the Superconducting Proton Linac (SPL) which is under design can be at the core of a multi-MW proton driver at 5 and/or 2.5 GeV, serving a neutrino facility, and/or a Radioactive Ion Beam facility of the next generation. The future accelerator complex is described and its potential for other applications than LHC is detailed.

1 Introduction

The LHC is designed for colliding proton beams at a centre of mass energy of 14 TeV with a nominal luminosity of 10^{34} cm⁻²s⁻¹ in two interaction regions [1]. After modification, the existing complex of accelerators has been made capable to serve as injector and provide the beam required for that purpose [2]. Starting for these favorable conditions, the LHC should reach nominal performance within a few years. Without upgrade, the peak luminosity will then saturate and the halving time of the statistical experimental errors will become very large. An analysis has therefore been made of the possible scenarios for achieving a much higher rate of increase of the integrated luminosity (e.g., a factor of 10 would enlarge the discovery range for new particles by about 25 % in mass [3]). Beyond the need for improvements in the LHC itself, the characteristics of the injected beam have to be upgraded and the reliability of the whole complex has to be significantly increased [4, 5].

2 Plans for future LHC injectors

2.1 Motivation

The present LHC injectors have been built more than 30 years ago (the PS is the oldest, celebrating in November 2009 the 50th anniversary of its first beam). Since their design, the state of knowledge in accelerator physics and technology has drastically progressed. Even though they have been regularly adapted and their operating mode has been sophisticated to cope with new needs and achieve unexpected levels of performance, they work nowadays at their absolute limit for the LHC. As the operation of the collider will become more regular, the (lack of) reliability of the injectors is likely to negatively impact on the integrated luminosity. Hence the plan to build new injectors operating simply and satisfying the future needs of the LHC with adequate performance margin.

Luminosity is directly related to beam brightness, the ratio N_b / ε_N of the number of protons in a bunch N_b to its normalized transverse emittance ε_N (in the approximation of same emittances in both transverse phase planes) which is then the main beam characteristic to improve. In the case of protons where synchrotron radiation is too weak to provide sufficient cooling, beam brightness is, at best,

preserved during acceleration. The maximum achievable brightness is therefore limited by the tolerable space charge induced tune spread ΔQ_{sc} at low energy. In a given synchrotron, it is governed by the relation in Eq. (1):

$$\Delta Q_{sc} \propto -\frac{N_b}{\varepsilon_N} \frac{R}{\beta \gamma^2} \tag{1}$$

where R is the average radius and β, γ are the usual relativistic beam parameters.

2.2 Description

For the design of the future injector chain, the target brightness corresponds to bunches of 3.4 10¹¹ protons within nominal emittances of 3.75 mm.mrad (25 ns time interval between bunches) circulating at 7 TeV in the LHC. The proposed complex [5] is sketched in Fig. 1, together with the present machines.

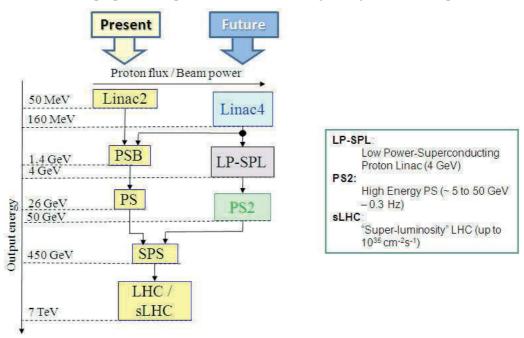


Fig. 1: Present and future accelerators

A new 50 GeV synchrotron (PS2) [6] will replace the existing 26 GeV PS and provide beam with the specified brightness (including $\sim 20\%$ margin) at injection in the SPS. Its main characteristics are given in Table 1. The increased injection energy in the SPS will reduce space charge and improve beam stability. Combined with other upgrades concerning mainly RF and the coating of the vacuum chamber to reduce secondary electron yield, this will allow the SPS to preserve beam characteristics and deliver the specified beam to the LHC.

Because of its superiority in terms of beam performance and potential for future applications, a superconducting linac (the Low Power SPL or LP-SPL) has been selected as the injector for the 50 GeV proton synchrotron PS2 [7]. Its low energy front-end (up to 160 MeV), called Linac4 [8], will first be used to replace Linac2 as injector of the PS Booster, providing beam with twice the $\beta \gamma^2$ of the present 50 MeV Linac2 and hence doubling the brightness of the PSB beam. The layout of these new accelerators on the CERN site [9] is shown in Fig. 2.

Table 1: PS2 characteristics

| Physical quantity | Value | Reason |
|--|--|--|
| Injection energy | 4 GeV | Space charge in PS2 (from Eq. (1)) |
| Ejection energy | 50 GeV | SPS improvement |
| Nb of protons/bunch with 25 ns bunch spacing for LHC (total in 168 bunches) | 4 10 ¹¹ p/b (6.7 10 ¹³ p) | Brightness goal for LHC upgrade (with 15 % margin) |
| Nb of protons/bunch with 25 ns bunch spacing for SPS fixed target (total in 168 bunches) | 6 10 ¹¹ p/b (1 10 ¹⁴ p) | SPS flux for fixed target physics |
| Cycling period to 50 GeV | 2.4 s | LHC filling time & SPS flux |
| Transverse emittances at ejection of LHC-type beam (normalized – 1 sigma) | 3π mm.mrad | SPS requirement for LHC |
| Longitudinal emittance/bunch (25 ns bunch spacing) | 0.35 eVs | SPS requirement for LHC |
| Circumference (ratio PS2/SPS) | 1346.4 m (15/77) | LHC need for 25, 50 and 75 ns bunch spacing |
| RF harmonics for 25 ns bunch spacing (resp. 50 or 75 ns) | 180 (90 or 60) | LHC need for 25, 50 and 75 ns bunch spacing |

For the needs of the LHC, the LP-SPL will only deliver a moderate beam power (~140 kW) at 4 GeV. However, its superconducting accelerating structures are capable to operate at much higher duty cycle than required by the LHC. At the cost of some upgrade of RF and infrastructure, this capability could be later exploited to increase beam power up to 4 MW at 5 and/or 2.5 GeV for fulfilling the needs of a neutrino facility and/or a Radioactive Ion Beam facility of the next generation [10]. The initial and potential evolution of the linac specifications are summarized in Table 2.

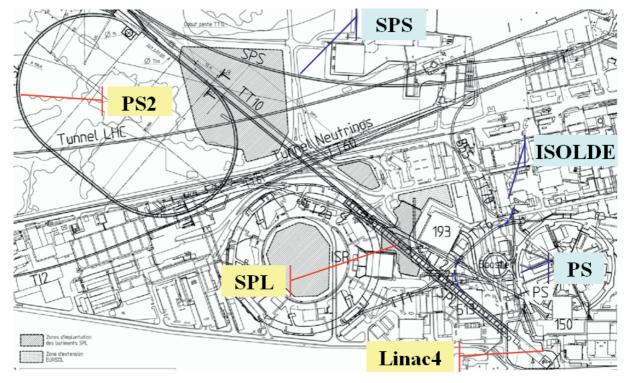


Fig. 2: Layout on the CERN site

Table 2: Low Power and High Power SPL beam characteristics

| | LP-SPL | HP-SPL Option 1 | HP-SPL Option 2 |
|--|--------|-----------------------|----------------------|
| Maximum kinetic energy [GeV] | 4 | 4 or 5 ^a | 4 or 5 ^a |
| Average beam current during pulse [mA] | 20 | 20 | 40 ^b |
| Pulsing rate [Hz] | 2 | 50 | 50 |
| Pulse duration [ms] | 0.9 | 0.9 | 1.2 ^b |
| Beam power [MW] | 0.14 | 2.25 @ 2.5 GeV | 5 @ 2.5 GeV |
| | | or 4.5 MW at 5 GeV | and 4 MW at 5 GeV |

^a Required for a neutrino factory.

2.3 Implementation plan

2.3.1 Linac4

The construction of the LP-SPL front-end, Linac4, has been approved at the end of 2007, as an efficient means to increase the performance of the PS complex and to prepare for the LP-SPL itself. As a result of this first step, the brightness per pulse of the PS Booster will be doubled, Linac4 providing beam at 160 MeV, with twice the $\beta \gamma^2$ of the present 50 MeV Linac2. The PS will be filled by a single pulse from the PSB, minimizing the time spent with a large space charge on the PS injection flat porch and hence suppressing the associated emittance blow-up and beam loss.

Linac4 is equipped with normal conducting accelerating structures. The frequency of 352.2 MHz has been chosen for its convenience for accelerating protons in this energy range, and because of the availability of a large inventory of LEP RF equipment. The block diagram of the accelerator is shown in Fig. 3. Thirteen former-LEP klystrons and six new pulsed devices are used to excite the four different types of accelerating structures. A four vane RFQ bunches and accelerates the beam up to 3 MeV where a wideband / high speed chopper (rise and fall time <2 ns) tailors the bunch train to the needs of the following synchrotron. An Alvarez-DTL equipped with permanent quadrupole magnets brings the energy up to 50 MeV. Cavity-Coupled DTL (CCDTL) structures are used in the energy range from 50 to 102 MeV followed by Pi Mode Structures (PIMS) for acceleration up to 160 MeV.



Fig. 3: Block diagram of Linac4

Civil Engineering work will finish at the end of 2010. Installation of the infrastructure will then take place in 2011, followed by the installation of the accelerator itself. Linac beam commissioning will start in the middle of 2012 and last until the third quarter of 2013, when all accelerators will be stopped and the PSB will be modified. After 3 months of beam commissioning, operation will resume in April 2014 and Linac4 will become the source of all protons for high energy physics at CERN.

2.3.2 New LHC injectors

The study of the LP-SPL, PS2 and SPS upgrade has also been approved at the end of 2007, with the goal of preparing a project proposal for mid-2012.

^b Required for 2 simultaneous users of high beam power or for 5 MW at 2.5 GeV

The LP-SPL will deliver beam at 2Hz with the characteristics shown in Table 2. It is made up of 2 sections of superconducting cavities accelerating the H $^-$ beam from 160 MeV to 4 GeV [9]. Both sections operate at 704 MHz and use 5-cell elliptical cavities differing by the use of different geometric β (respectively 0.65 and 1.0). Challenging peak surface fields of 50 MV/m and 100 mT are assumed, corresponding to accelerating gradients of 19.3 and 25 MV/m. Medium β cavities are grouped in 10 cryomodules of 11.5 m length containing 6 cavities and 2 quadrupole doublets. High β cavities are grouped in 18 cryomodules of 14.3 m length containing 8 cavities and 1 quadrupole doublet. As shown in the block diagram of Fig. 4, beam extraction is foreseen at ~1.4 GeV for supplying particles to the ISOLDE experimental area (Fig. 2).

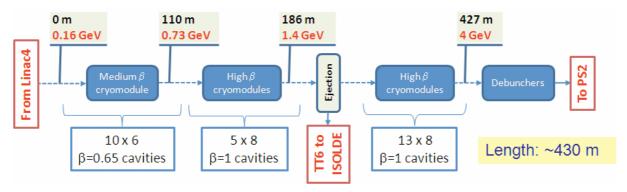


Fig. 4: Block diagram of the LP-SPL

The main characteristics of PS2 and its beam are summarized in Table 1. The basic design choices for the accelerator have recently been made [6]. To satisfy the integration requirements, a race track shaped is preferred. Charge exchange injection is implemented to accumulate the H beam from the LP-SPL. Fast injection is foreseen for heavy ions from LEIR. Beam can be extracted at the highest energy (50 GeV) either in a single turn (for LHC), in five turns (for SPL fixed target physics) or by slow resonant ejection (for a potential experimental area). The lattice is of the Negative Momentum Compaction (NMC) type, to avoid the crossing of transition during acceleration (γ_t =37i). The RF system operates over the frequency range from 18 to 40 MHz. The spacing between bunches and the overall time structure of the circulating beam is obtained directly at injection, through an adequate chopping in the LP-SPL front end. Compared to the complexity of the scheme presently used in the PSB and PS, this simple mode of operation is expected to be very robust and reliable.

If the decision to build the new injectors is taken at the end of 2012, their construction will start at the beginning of 2013. Once the tunnels are available, the LP-SPL will be progressively installed and beam commissioned at increasingly higher energies. PS2 will first be tested with a fast injected proton beam from the PS and afterwards with the H beam from the LP-SPL using charge exchange injection. Both accelerators will be fully beam commissioned without interfering with the LHC physics program. Once the foreseen beam characteristics are obtained at the exit of PS2, the connection will be made with the SPS in 2019-2020.

2.3.3 High Power SPL options

The beam power of the LP-SPL can later be upgraded to multi-MW by increasing the cycling rate to 50 Hz (Table 2), which implies the replacement/upgrade of power supplies and a major upgrade of the infrastructure (electricity, water cooling and cryogenics).

If a beam power larger than 2 MW is needed at 2.5~GeV, or if two high power users require $\sim 4~\text{MW}$ simultaneously, the beam current during the pulse will have to be doubled to 40~mA, doubling the number of high power RF amplifiers. Enough space will be reserved in the linac tunnel to eventually add accelerating structures and bring the beam energy to 5~GeV, as required for a neutrino factory.

3 Potential for neutrino experiments

3.1 PS2

As a result of the requirements of the LHC upgrade, PS2 will be capable to deliver up to 10^{14} protons/pulse at 50 GeV every 2.4 s to the SPS for fixed target physics. This corresponds to an average beam power of 330 kW or to $\sim 4 \times 10^{20}$ pot/year (~ 6 times the capability of today's PS).

To meet the needs of neutrino physicists who are expecting 3×10^{21} pot/year [11], a new design is therefore needed. Beam loss analysis and management should be studied in great detail to make sure that hands-on maintenance will remain possible. The only means to increase beam power being linked to the intensity per pulse and the repetition rate, it can be predicted that a MW-class PS2 will require larger magnet apertures, more RF voltage and RF power, more powerful magnets power supplies and much more involved collimation and beam abort systems.

3.2 SPS with the new injectors

With respect to the PS, PS2 will deliver to the SPS for fixed target physics a beam of much better characteristics (Table 3).

| | PS | PS2 | Advantage of PS2 | Benefit for SPS |
|--------------------------------------|-----------------------------|---------------|---|--|
| Injection energy in the SPS [GeV] | 14 | 26-50 | Injection above transition Smaller beam size | No transition crossing Less loss at injection |
| Number of pulses for filling the SPS | 2 | 1 | No SPS injection flat porch | Less loss on flat porch Gain of 1.2s in SPS cycle |
| Longitudinal beam characteristics | Partly bunched | Fully bunched | Matched bunch to bucket transfer | Less capture loss |
| Number of protons/pulse | $2 \times 3 \times 10^{13}$ | 10^{14} | Higher intensity | Potential for higher intesnity |

Table 3: Comparison between PS and PS2 for fixed target physics with the SPS

Assuming that the SPS is upgraded and made capable to accelerate and transfer the beams specified for the LHC upgrade, the beam for fixed target physics will be limited in the SPS mostly by RF (voltage and power) and activation due to beam loss. If beam is sent onto the CNGS target [12], the design parameters of the CNGS facility will be the actual limits (maximum intensity per pulse for target and horn, maximum flux to target/horn and hadron stop). Moreover, all radiation protection calculation shall be redone with the new parameters, more shielding is likely to be needed and a new INB approval by the French IRSN has to be obtained. This has been studied in 2006, and the results published in Ref. [12] show that the potential proton flux on target at 400 GeV can be brought up to 2 or 3 times the nominal value for CNGS (Table 4).

Table 4: Protons on target per year $[\times 10^{19}]$ for 200 days of SPS operation with 80 % machine availability

| SPS cycle length | 6 | s | 4.8 | s |
|--|----------|------|----------|------|
| Injection momentum | 14 GeV/c | | 26 GeV/c | |
| Beam sharing | 0.45 | 0.85 | 0.45 | 0.85 |
| Present injectors (4.8×10 ¹³ p/p) | 5 | 9.4 | | |
| Future injectors + SPS RF upgrade for LHC (7×10 ¹³ p/p) | | | 9 | 17.1 |
| Future injectors + new SPS RF system (10 ¹⁴ p/p) | | | 12.9 | 24.5 |

3.3 High power SPL

A superconducting proton linac is a safe and reliable solution for the supply of a high power proton beam at a few GeV. As such, the LP-SPL has the potential to be upgraded to multi-MW of beam power at up to 5 GeV. For neutrino applications, however, the proton beam on target must also have a time structure which a linac cannot deliver: for a superbeam, a pulse of a few µs is required, while for a neutrino factory a few bunches of ~2ns rms bunch length are mandatory. The linac beam has therefore to be complemented with one or two rings.

3.3.1 Application to a Neutrino Factory

A neutrino factory has very demanding needs (Table 5), as defined by the ISS working group [13].

| Parameter Value | | |
|--|------------|--|
| rarameter | vaiue | |
| Average beam power [MW] | 4 | |
| Pulse repetition frequency [Hz] | 50 | |
| Beam kinetic energy [GeV] | 10 ± 5 | |
| Bunch length (rms) [ns] | 2 ± 1 | |
| Number of proton bunches | 3 or 5 | |
| Sequential extraction delay [µs] | ≥ 17 | |
| Pulse duration (liquid Hg target) [µs] | ≤ 40 | |
| Pulse duration (solid target) [µs] | ≥ 20 | |

Table 5: Proton driver requirements of a neutrino factory [13]

A scheme has been designed which can potentially meet these specifications [14], using the 5 GeV high power SPL and two fixed energy rings of approximately 300 m. In the first ring (the accumulator), the chopped linac beam is accumulated in a few long bunches, using charge-exchange injection. The accumulator is isochronous to preserve the time structure of the linac beam, and it has no RF system to minimize the impedance. Once accumulation is finished, bunches are transferred one by one to the second ring (the compressor) where they are rotated in the longitudinal phase plane and ejected to the target when their length is minimum. This principle is sketched in Fig. 5 in the case of 6 bunches. Bunch rotation takes place with the energy stored in the cavities of the low frequency RF system.

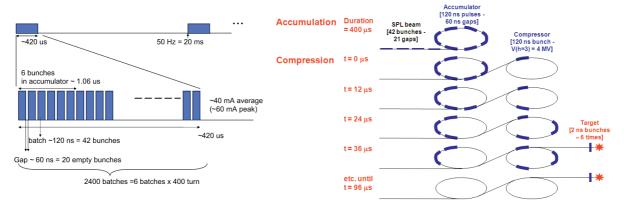


Fig. 5: Principle of bunch generation for a neutrino factory

The accumulator and compressor rings have been designed and particle tacking simulations have shown that bunches of 2 ns rms length can indeed be generated (Fig.6). The study of collective effects in

the accumulator has revealed that the impedances required for stability are within reach. More difficult scenarios have been developed for generating less than 6 bunches. The main parameters of the rings in these different cases are shown in Table 6.

| Ring | Parameter | 6 bunches case | 3 & 1 bunch cases |
|-------------|-------------------------------|----------------|-------------------|
| Accumulator | Circumference [m] | 318.5 | 185.8 |
| | Nb. of accumulation turns | 400 | 640 / 1920 |
| | Type of magnets | NC | SC |
| Compressor | Circumference [m] | 314.2 | 200 |
| | Nb. of compression turns | 36 | 86 |
| | RF voltage on h=3 (MV] | 4 | 1.7 |
| | Transition gamma | 2.3 | 2.83 |
| | Type of magnets | SC | NC |
| | Interval between bunches [µs] | 12 | 30 |

Table 6: Main parameters of the accumulator and compressor for a neutrino factory

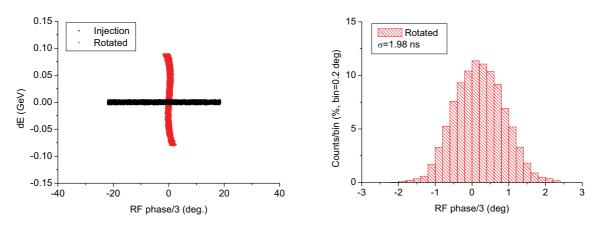


Fig. 6: Bunch rotation in the longitudinal phase plane (left) and longitudinal density at ejection (right)

3.3.2 Application to a low energy neutrino "super-beam"

To generate a low energy neutrino superbeam from π decay, the time structure of the beam provided by the accumulator is perfectly adequate and the compressor is not necessary.

3.3.3 Application to a "beta-beam" facility

In the context of the EURISOL Design Study supported within the 6^{th} Framework Program of the European Union [15], the possibility has been studied to generate electron neutrinos and antineutrinos from the beta decay of 6 He or 18 Ne radioactive ions. It has been shown that the necessary flux of 3×10^{13} 6 He per second can be obtained irradiating a Beryllium Oxide target with the spallation neutrons resulting from a proton beam of ~ 200 kW at 1-2 GeV on a converter target. The SPL would be perfectly able to satisfy this need. However, no satisfying solution has yet been found for getting the necessary flux of 2×10^{13} 18 Ne ions/s, although a promising technique based on 70 MeV proton irradiating an Aluminium Oxide target has recently been proposed.

4 Synergy with a Radioactive Ion Beam Facility

With its 50 Hz rate, the high power SPL is a competitive proton driver for a radioactive ion beam facility of the next generation (EURISOL-like) [15]. The beam from the SPL can directly be used, and even preferably the H⁻ ions, without the need for an intermediate ring. Figure 7 shows a possible layout on the CERN site.

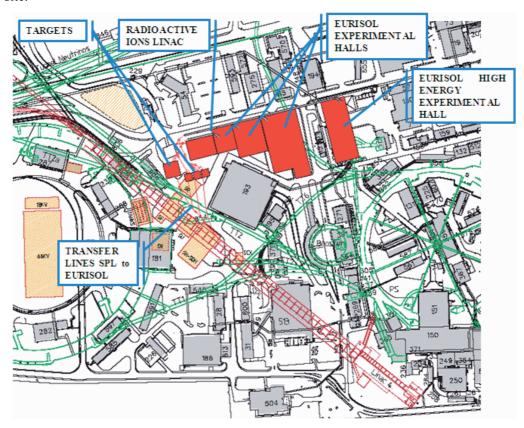


Fig. 7: ISOLDE and EURISOL layouts on the CERN site

In a first stage, the present ISOLDE area will receive ~ 1.4 GeV protons from the LP-SPL in proton pulses of 20 mA / 0.9 ms length at an average rate of 1.25 Hz (3 pulses out of 4, every 0.6 s), which corresponds to a beam power of 31 kW. If ISOLDE users are interested in a lower beam energy of 1 GeV and a reduced pulse length of 0.38 ms, the rate could be increased to ~ 3 Hz for the same beam power, by reducing the field in the cavities and accelerating a slightly higher current in the LP-SPL (28 mA instead of 20 mA).

In a second stage, the LP-SPL could be upgraded to high beam power by increasing its repetition rate to 50 Hz (HP-SPL Option 1 in Table 2). The EURISOL experimental area could then receive up to 2.25 MW of beam power at 2.5 GeV, if there is no other user of the SPL high power beam. To get 5 MW, or to be able to operate simultaneously with another user of high beam power, the beam current should be increased by a factor of two to 40 mA, by doubling the number of klystrons in the SPL (HP-SPL Option 2 in Table 2).

5 Summary

The chain of accelerators made up of the LP-SPL, PS2 and the upgraded SPS will provide a lot of flexibility for tailoring the beam injected inside the LHC to whatever solution will finally be implemented for upgrading the integrated luminosity. Users of ISOLDE and SPS fixed target beams will

immediately benefit from these new accelerators. The LP-SPL and the SPS have the potential to be upgraded for supplying a higher beam power. If this capability is needed for PS2, its design should be revisited accordingly. A superconducting proton linac is especially worthwhile in the CERN context because of its capability to be upgraded into the proton driver for a neutrino facility and/or a EURISOL-like radioactive ion beam facility.

Acknowledgement

This document summarizes the work of multiple teams (HIP and PAF working groups; PS2, SPL and SPS upgrade study teams; neutrino study team) and of numerous workshops (especially the HHH and NuFact series). The references are an attempt at pointing to the main sources of information.

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