IPP meeting

• Inclusive D⁰ and D⁺ to K⁰ and K⁺ branching ratios: HVFL versus PDG - Paul Colas.

In connection with the K^0 -ln(x) spectrum Gerald suggested to look at the $D^+ \to K$ decay, as incorrect branching ratios there may be a source of difference between the data and the HVFL03 MC. Paul found that the decay $D^+ \to K^-$ may have a little more kaonless decays (20.8 PDG versus 28.5 sum of exclusive decays HVFL04). A big difference is seen on $D^0 \to K^0$: 42 ± 5 % (PDG) versus 27.2 (HVFL04) and on $D^0 \to K^-$: 46 ± 4 % (PDG) versus 64.1 % (HVFL04). Obtaining matching between the sum of the exclusive decay channels and the inclusive one is a non-trivial excercise. Work on this point is in progres.

• Photons from conversion pairs - Frank Stephan.

In preparation for the DPG details were shown on this analysis. Pairfinding based on QPAIRF (for details on selection see transparencies) is shown to give a detailed picture of the material distribution inside the detector. Also the crude modelling of the VDET material in the 1991 Monte Carlo shows up clearly when a MC-data comparison is made. A photon energy resolution of The effiency correction is evaluated for low and high radius separately. On average the data has $6.92 \pm 2.5\%$ (stat.) $\pm 5.0\%$ (sys.) photons per event. This is to be compared to 7.100 in HVFL02, 7.108 in Jetset and 7.085 in Herwig; these numbers have a statistical error of 0.03 % on them. The comparison of the corrected data spectrum (-ln(x)) to the Monte Carlo models shows a discrepancy at the low momentum side. It is not clear whether this represent a problem.

∑⁰ and ρ⁰ studies - Armin Böhrer.
For the Σ⁰ the decay Σ⁰ → Λ⁰ + γ is studied. Armin started by expressing some worries having used the 1991 Monte Carlo too, which has Vdet material modelling problem. In the 1991 + 1992 data he finds 90 ± 17 (stat.) ± 12 (sys.) Σ⁰'s which corresponds to 0.072 ± 0.014 (stat.) ± 0.010 (sys.) per event. This is to be compared to 0.075 in HVFL03, 0.060 in Jetset, 0.062 in Herwig. Comparing the Chi (dss) rate as measured

by Robert Johnson (0.0273 \pm 0.0013 \pm 0.0016) with the Σ^0 (uds) rate, he extracts a strangeness suppression of 0.38 \pm 0.08 (stat.).

On the side of the ρ^0 studies Armin showed his improved understanding of the implementation of the Bose-Einstein correlations. Including the $\operatorname{lin}(Q)$ terms in the form of an interference term is very important. This interference term is a function of x, it is 1 at low x and drops to 0 at high x. The stability of the ρ^0 fit improves greatly with respect to the case where $\operatorname{lin}(Q)$ was not taken into account. The number of ρ^0 's is stable as a function of x when comparing a Breit-wigner shape with the Lund shape. Also the ratio of ρ^0 -rates in data and MC is now yields compatible results when the omega is coupled to the ρ or left free in the fit. And last but not least, Armin showed the understanding now in terms of the $\operatorname{mass}(\pi^+\pi^-)$ rather than the mass difference spectrum and gave both the ρ^0 and f⁰ spectra as a function of x. These correspond to an integrated ρ^0 -rate of 1.57 \pm 0.06 \pm 0.24 (m $_{\rho}$ < 2 GeV) and a f⁰-rate of 0.15 \pm 0.02.