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Review of Particle Properties Particle Data Group

ALAN RITTENBERG, ANGELA BARBARO-GALTIERI, THOMAS LASINSKI, ARTHUR H. ROSENFELD,
and THOMAS G. TRIPPE

Lawrence Radiation Laboratory, University of California, Berkeley, California 94720*

MATTS ROOS† and CLAUDE BRICMAN

CERN, Geneva, Switzerland

PAUL SÖDING

DESY, Hamburg, Germany

NAOMI BARASH-SCHMIDT

Brandeis University, Waltham, Massachusetts 02154

CHARLES G. WOHL

Oxford University, Oxford, England

This review of the properties of leptons, mesons, and baryons is an updating of Review of Particle Properties, Particle Data Group [Phys. Letters **33B**, 1 (1970)]. Data are evaluated, listed, averaged, and summarized in Tables. A data booklet is also available.

CONTENTS

I. Introduction, Credits, Consultants.....	S1	<i>SU(3)</i> Conventions.....	S26
II. Selection and Organization of Data.....	S2	<i>SU(3)</i> Isoscalar Factors.....	S27
III. Criteria for Resonances.....	S2	c.m. Energy and Momentum vs Beam Momentum.....	S28
IV. Parameters and Conventions.....	S3	Special Relativity, Phase Space, and Cross Sections.....	S29
A. Quantum Numbers.....	S3	Confidence Level vs χ^2 for n_D Degrees of Freedom.....	S30
B. Particle Names.....	S4	Gaussianlike Distributions.....	S30
C. Masses and Widths.....	S4	Atomic and Nuclear Properties of Materials.....	S31
D. Muon-Decay Parameters.....	S6	Multiple Coulomb Scattering.....	S31
E. <i>K</i> -Decay Parameters.....	S7	Radioactivity and Radiation Protection.....	S31
F. Baryon-Decay Parameters.....	S8	Range and Energy Loss in Copper.....	S32
V. Statistical Procedures.....	S9	Range and Energy Loss in Liquid Hydrogen.....	S33
A. Confidence Levels and Errors.....	S9	Data Card Listings.....	S34
B. Unconstrained Averaging, Scale Factors.....	S10	Illustrated Key.....	S34
C. Constrained Fits.....	S10	Stable Particles.....	S36
VI. Particle Data Group Publications.....	S10	Mesons.....	S61
Acknowledgments.....	S11	Baryons.....	S97
References (for above sections).....	S11	Appendix I. Test of $\Delta I = \frac{1}{2}$ Rule for <i>K</i> Decays.....	S147
Particle Properties Tables.....	S12	Appendix II. <i>SU(3)</i> Classification of Resonances.....	S148
Stable Particles.....	S12		
Addendum.....	S15		
Mesons.....	S16		
Baryons.....	S21		
Miscellaneous Tables, Figures, and Formulae.....	S25		
Physical and Numerical Constants.....	S25		
Clebsch-Gordan Coefficients and Spherical Harmonics.....	S26		

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† On leave of absence from the Institute of Nuclear Physics, University of Helsinki, Helsinki, Finland.

I. INTRODUCTION, CREDITS, CONSULTANTS

This review is an updating through January 1971 of our previous review (Particle Data Group, 1970).

In this version of the text we concentrate on topics that are either new or essential. For complementary information on our standard procedures the reader is referred to our 1969 article (Particle Data Group, 1969).

We remind the reader that it is inappropriate to make reference to this compilation instead of to an

original work. In our Data Card Listings we provide references to all original data quoted by us.

The responsibilities of the authors of this compilation (who prefer to be quoted as the "Particle Data Group") can roughly be broken down as follows:

Stable Particles: A. Barbaro-Galtieri, N. Barash-Schmidt, T. G. Trippe. Our eta meson consultant is L. R. Price.

Meson Resonances: M. Roos, A. H. Rosenfeld, P. Söding.

Baryon Resonances: A. Barbaro-Galtieri, C. Bricman, T. Lasinski, C. G. Wohl.

General: All Berkeley authors.

Consultants: The three teams just mentioned must come to a consensus on how to treat the data, and must write a number of mini-reviews. It is impractical to spread this responsibility over more than a few people in each team, and still expect to meet publication deadlines. Hence we limit our number of authors (to ten in this edition), but thereby leave gaps in our coverage, both intellectual and geographical. To help us overcome this deficiency, we have solicited the help of consultants:

- Chih-Yung Chien (Johns Hopkins University) and
- LeRoy R. Price (Univ. of California at Irvine).

Mark Sakitt (BNL) has recently agreed to aid us in the future, and we hope to add a few more consultants.

The usefulness of this compilation depends in large part on the interaction among the users and authors and consultants. We appreciate comments, criticisms, and suggestions for improvements of all stages of data retrieval, processing, and presentation.

II. SELECTION AND ORGANIZATION OF DATA

All particles are considered to fall into one of the three groups:

- (1) Stable Particles, immune to decay via the strong interaction.
- (2) Meson Resonances.
- (3) Baryon Resonances.

These groups are maintained within the two main parts of the compilation:

- (1) Tables of particle properties
- (2) Data Card Listings.

The Data Card Listings contain the original information (data, references, etc.), weighted averages, comments and "mini-reviews." Immediately preceding the Data Card Listings is an illustrated key thereto.

Data that have not been used in the averages are set off in parentheses; our reasoning is then often given in a footnote below the data. If the reason is not given,

it is one of the following:

- The quantity was presented with no error stated.
- The result comes from a preprint or conference report and has not been verified by its authors.
- It involves some assumptions that we do not wish to incorporate.
- It is of poor quality, e.g., bad signal-to-noise ratio.
- Two experiments give contradictory results, so that it is senseless to average the data.

When the data for a particle have received special treatment or when they present special problems, this is noted in a mini-review in the Data Card Listings.

The Tables represent the output of weighted averages and some critical judgment. The extent to which "blind" averaging has been tempered with judgment is explained in footnotes to the Tables. In general, however, the footnotes are less complete than are the collection of notes and mini-reviews in the Data Card Listings. The reader is therefore encouraged to familiarize himself with the Data Card Listings, and ultimately with the original experiments.

Further details about our selection of data can be found in the text of the 1969 edition, (Particle Data Group, 1969, Sec. III).

III. CRITERIA FOR RESONANCES

An experimentalist who finds some evidence for a peak in a mass spectrum will of course want to know what has been seen in that region in the past; hence, we strive to have the Data Card Listings serve as an archive for any substantial claim or evidence for a new state.

For the Tables, on the other hand, we wish to be more conservative, and to include only those peaks or resonances that have a $\gtrsim 90\%$ chance of survival. One's betting odds for survival are of course completely subjective; they are influenced mainly by the amount of information available (such as partial-wave analyses), and somewhat by the degree of controversy over interpretation and how long it will be before more information is available.

More details on our acceptance criteria are as follows:

- (1) Consider, first, peaks where there is *no* partial-wave information:

(1a) Most mesons and Ξ^* peaks, plus N^* and Y^* peaks above 2200 MeV. Unless the peak is experimentally shaky, we put it in the table.

(1b) Peaks not far above some threshold, e.g., $A1 \rightarrow \rho\pi$ or $A3 \rightarrow f\pi$, and their K^* counterparts Q and L . These are at least partly explained by the "Deck effect," or its more modern version "double-Regge-pole (DRP) exchange." But the notion of duality tells us that "to explain is not to explain away," i.e., the

“explanation of a peak” with the DRP model need not contradict its interpretation as an s -channel resonance. So we put these four peaks on the Meson Table, but add comments like “ $A1$ interpretation not clear; 160 MeV above $\rho\pi$ threshold.” For more discussion of this problem, see the $A1$ mini-review in the listings.

(2) Consider, next, peaks for which there *are* partial-wave analyses. We can then check that the Argand plot follows roughly a left-hand circle, and that the “speed” of the amplitude also shows a maximum near the resonance energy; further, there should be data well above the resonance, showing that the speed again decreases. Indeed proper behavior of the partial-wave amplitude will accredit a resonance even if its elasticity is too small to make a noticeable peak in the cross section.

Of course, even if Argand plots are available, it may still be a matter of opinion as to what behavior constitutes a resonance. Such an example is the K^+p peak (near $K\Delta$ threshold) called Z_1 (1900), which is discussed in a mini-review in the listings. The P_{13} Argand plots are displayed there, and most suggest a resonance; however, there is disagreement between the various analyses as to the speed of the amplitude, i.e., as to whether it has a Breit–Wigner type of behavior. In addition, the errors on the amplitudes are still large, and we prefer to wait a bit longer before we put Z_1 on the table.

Approximate Breit–Wigner behavior of a partial-wave amplitude appears to us to be the most satisfactory test for a resonance, since after all a Breit–Wigner amplitude is the Fourier transform of an exponential decay of a state with a finite lifetime. We are aware that this approximate Breit–Wigner behavior could be accidental, but can only hope that such an accident is improbable.

We now ask “How likely is it that peaks of class (1) above (no way to check them with partial-wave analysis) will eventually be confirmed as resonances?” We know of no experimentally convincing peak that has been shown to having *nothing* to do with a resonance. But be warned that broad peaks may be misleading: they may contain several resonances, or they may include a resonance narrower than the peak, plus some other complications.

●Before 1966 we might have tabulated the πp bumps at 1512 and 1688 MeV as single resonances, whereas partial-wave analysis shows that each contains several resonances.

●Abrams *et al.* have reported broad peaks in the $\bar{N}N$ total cross section at 2190, 2345, and 2380 MeV. One of these indeed overlaps our tabulated $U(2375)$ meson, whose width is given as 30 MeV. But $\bar{p}p$ s -channel experiments, which can study all the final states formed at c.m. energies near 2345 MeV, show that the rest of the $\bar{p}p(2345)$ peak, which is 140-MeV wide, must be attributed to a jumble of many different, unresolved

effects. So we do not accept as a single resonance the whole $\bar{N}N(2345)$ peak, whose width is measured as 140 MeV.

Despite these cases where a resonance has turned out to be narrower than a bump, and only partly to explain the bump, *most* baryon enhancements have been confirmed by partial-wave analysis. A relevant example is the following:

●Before the $N^*(1470, P_{11})$ was confirmed in partial-wave analyses, it was seen as a missing mass or $p\pi\pi$ peak produced peripherally in high-energy pp collisions, and (like $A1$, Q , $A3$, and L) was partly explained by the Deck effect and later by double-Regge-pole exchange. Thus nowadays we say that the s -channel resonance *confirms* the production peak, and that this supports one’s confidence in duality.

In summary, we enter onto the tables experimentally convincing peaks unless there is contradictory information; and we expect that most of these peaks will eventually be confirmed as one or more resonances.

IV. PARAMETERS AND CONVENTIONS

A. Quantum Numbers

The symbols $I^G(J^P)C$ represent:

- I = isospin
- G = G parity
- J = spin
- P = space parity
- C = charge conjugation parity.

Mesons

The charge conjugation operator C turns particle into antiparticle and has eigenvalues ± 1 only for neutral states, so it is useful to define an extension G which has eigenvalues for charged states too. It is usually* defined by

$$G = C \exp(i\pi I_y). \quad (1)$$

A neutral nonstrange state is an eigenstate of $\exp(i\pi I_y)$ with eigenvalue $(-1)^I$. Then we can write the eigenvalue equation for the whole multiplet as

$$G = C_n (-1)^I, \quad (2)$$

where C_n (n for neutral) is the eigenvalue C would have if applied to the neutral member of the multiplet. Thus, for a π^0 , C has the eigenvalue $+1$, and since $I=1$, $G=-1$. For the charged pion, there are no eigenvalues corresponding to C and to the isospin rotation, but Eqs. (1) and (2) still give $G=-1$.

Consider a meson as a bound state of fermion–antifermion, e.g., $\bar{q}q$, with orbital angular momentum l , and with the two quark spins coupling to give a spin

*Most texts define it as in Eq. (1); see, e.g., Gasiorowicz (1966); however, sometimes the rotation is taken about I_x . The difference between the two conventions is mentioned in a footnote in Källén (1964).

S. Then one can show that the charge-conjugation eigenvalue [defined in Eq. (2)] is

$$C_n = (-1)^{t+s}. \quad (3)$$

Equations (2) and (3) combine to give

$$G = (-1)^{t+s+I}. \quad (4)$$

The parity is

$$P = -(-1)^t. \quad (5)$$

Equations (3) and (5) combine to give

$$C_n P = -(-1)^s$$

so all singlets ($^1S_0, ^1P_1, \dots$) have $C_n P = -1$, and all triplets ($^3S_1, \dots$) have $C_n P = +1$.

If, instead of $\bar{q}q$, we consider the meson as a state of *boson-antiboson* (e.g., $A_2 \rightarrow \bar{K}K$), it turns out that some signs cancel, and Eqs. (3) and (4) [not (5)!] apply *unchanged*. Of course the mesons are usually spinless and S is zero, but the equations are more general. Equations (3) and (4) can be considered as selection rules forbidding many decays.

For proofs see our 1969 text (Particle Data Group, 1969), and Appendix by C. Zemach. We repeat here as the summary Table I that we used in 1969 as Table II.

Baryons and Mesons

Well-established quantum numbers are underlined (except for stable particles, where most of the quantum numbers are established). We have used flimsy evidence to guess many of the remaining ones, and we have indicated with “?” the ones for which there is almost no evidence.

As is customary, we define antiparticles as the result of operating with *CPT* on particles, so both share the same spins, masses, and mean lives. Whenever there is a particularly interesting test of *CPT* invariance we include it in the Stable Particles Table.

B. Particle Names

If a *meson* has a well-accepted colloquial name, we use it. If not, we name it by a single symbol which specifies its atomic mass number A ($=0$ for mesons), its hypercharge Y , its isospin I , and, for a nonstrange meson, its G parity. We choose

$$I=0; \quad \eta \quad \text{if } G \text{ is even, } \phi \text{ if it is odd;}$$

$$I=1; \quad \rho \quad \text{if } G \text{ is even, } \pi \text{ if it is odd;}$$

$$I=\frac{1}{2}; \quad K$$

$$I=\frac{3}{2}; \quad L \quad (\text{if ever established}).$$

To crowd even more information onto the symbol, we add a subscript giving J^P . Thus η_{0+} (1070). If J^P is not

known, but must be “normal” ($0^+, 1^-, 2^+, \dots$), e.g., because $K\pi$ decays are seen, we use the subscript N . Thus $K_N(1420)$. If such modes are *not* seen (and are not otherwise forbidden), we *guess* that it is because J is “abnormal,” and we write, for example, $K_A(1320)$.

For *baryons*, no attempt has been made to attach a subscript about J and P . The symbols are

$$Z_I \quad \text{for } Y=2, \quad I=0, 1;$$

$$N \quad \text{for } Y=1, \quad I=\frac{1}{2};$$

$$\Delta \quad \text{for } Y=1, \quad I=\frac{3}{2};$$

$$\Lambda \quad \text{for } Y=0, \quad I=0;$$

$$\Sigma \quad \text{for } Y=0, \quad I=1;$$

$$\Xi \quad \text{for } Y=-1, \quad I=\frac{1}{2};$$

$$\Omega \quad \text{for } Y=-2, \quad I=0.$$

For stable baryons of each Y and I we use the symbol standing alone; for resonances, the mass is in parentheses [i.e., $N(1688)$, $\Lambda(1405)$, $\Sigma(1765)$, etc.]. The J^P assignment is reported in the Table as $\frac{1}{2}^+$, $\frac{3}{2}^-$, $\frac{5}{2}^+$, etc., and also by the symbols P_{11} , D_{13} , F_{15} , which refer to the πp , $K p$, or $K \bar{p}$ partial-wave amplitude where the resonant state occurs (the first subscript refers to the isospin state).

When two *baryons* have identical quantum numbers we warn the reader by adding a *prime* to the symbol for the heavier one, e.g., p , $N'(1470, \frac{1}{2}^+)$. In the case of baryon resonances described by Argand diagrams which exhibit more than one resonance, we use one prime for the first, two for the second, \dots ; thus the series of which the proton is the stable member becomes: p , $N'(1470, \frac{1}{2}^+)$, $N''(1780, \frac{1}{2}^+)$.

If there is only one resonance on an Argand plot, and thus no need for distinctions, we use no primes.

For *some* pairs of *mesons* with identical quantum numbers, we also use primes; e.g., η , η' ; f , f' .

C. Masses and Widths

In the Tables, columns are headed Mass M , and Width Γ . We speak loosely of M as the position of a resonant peak, and of Γ as its full width at half-maximum. We now want to make these statements more precise. Values of M and Γ are extracted from data via *models*, and we cannot average these values if the models are dissimilar. In the next few paragraphs we discuss this point in slightly more detail, using the example of an s -channel resonance.

An *elastic* nonrelativistic Breit-Wigner T -matrix element is usually written

$$T_{11} = \frac{\Gamma/2}{M - m - i\Gamma/2}. \quad (6)$$

Here $\Gamma(m)$ is the width for decay into the channel 1, with angular momentum l . It contains barrier-penetra-

TABLE I. $I^G(J^P)$ of mesons from $\bar{q}q$ model. For the distinction between abnormal J^P and abnormal C, see text. $I=\frac{1}{2}$ states share the same values of J^P as the $I=0$ and 1 states shown, but are not eigenstates of G . The middle column, which gathers together $(J^P)_{N \text{ or } A}$ CP, is a redundant intermediate step intended to make the Table easier to read.

Parity	$\bar{q}q$ State		$(J^P)_{CP}$ Normal or abnormal	$I^G(J^P)C_n$	Examples and comments
	CP	CP			
	-	+			
Parity -	$1S_0$		$(0^-)_{A^-}$	$\begin{cases} 0^+(0^-)+ \\ 1^-(0^-)+ \end{cases}$	η, η' π
			$3S_1$	$(1^-)_{N^+}$	$\begin{cases} 0^-(1^-)- \\ 1^+(1^-)- \end{cases}$
Parity +	$1P_1$		$(1^+)_{A^-}$	$\begin{cases} 0^-(1^+)- \\ 1^+(1^+)- \end{cases}$	B
		$3P_0$	$(0^+)_{N^+}$	$\begin{cases} 0^+(0^+)+ \\ 1^-(0^+)+ \end{cases}$	$\eta_0(1060)$ $\pi_N(1016)$
		$3P_1$	$(1^+)_{A^+}$	$\begin{cases} 0^+(1^+)+ \\ 1^-(1^+)+ \end{cases}$	A1
		$3P_2$	$(2^+)_{N^+}$	$\begin{cases} 0^+(2^+)+ \\ 1^-(2^+)+ \end{cases}$	f, f' A2
Parity -	$1D_2$		$(2^-)_{A^-}$	$\begin{cases} 0^+(2^-)+ \\ 1^-(2^-)+ \end{cases}$	Regge recurrence of $1S_0, 0^-$
		$3D_1$	$(1^-)_{N^+}$	same as $3S_1$	
		$3D_2$	$(2^-)_{A^+}$	$\begin{cases} 0^-(2^-)- \\ 1^+(2^-)- \end{cases}$	Regge recurrence of top abnormal-C state below: $(J^P)C_n = (0^-)-$
		$3D_3$	$(3^-)_{N^+}$	$\{J > 2\}$	
Parity +	$1F_3$		$(3^+)_{A^-}$	$\{J > 2\}$	
		$3F_2$	$(2^+)_{N^+}$	same as $3P_2$	Another A2?
		$3F_3$	$(3^+)_{A^+}$	$\{J > 2\}$	
		$3F_4$	$(4^+)_{N^+}$	etc.	

ABNORMAL C STATES THAT CANNOT COME FROM $\bar{q}q$ MODEL

Abnormal C states Have no $\bar{q}q$ model	$(0^-)_{A^+}$	$\begin{cases} 0^-(0^-)- \\ 1^+(0^-)- \end{cases}$	All except $J^P = 0^-$ are $J^P = \text{normal},$ $CP = -1$
	$(1^-)_{N^-}$	$\begin{cases} 0^+(1^-)+ \\ 1^-(1^-)+ \end{cases}$	
	$(0^+)_{N^-}$	$\begin{cases} 0^-(0^+)- \\ 1^+(0^+)- \end{cases}$	
	$(2^+)_{N^-}$	$\begin{cases} 0^-(2^+)- \\ 1^+(2^+)- \end{cases}$	
	$(3^-)_{N^-}$	$\begin{cases} 0^+(3^-)+ \\ 1^-(3^-)+ \end{cases}$	

tion factors which can vary rapidly with energy; near threshold $\Gamma(m)$ should start up as q^{2l+1} , and then level off. Various m dependences are used, mostly variants of the general form

$$\Gamma(m) \propto \{q^2/[1+(qR)^2]\}^l (q/m). \quad (7)$$

For a choice of forms, see Jackson (1967), Pisut and Roos (1968), and Barbaro-Galtieri (1968). Of course the detailed shape of the amplitude will depend on the form chosen. So, although Γ is *related* to the full width, it can be measured in terms of the behavior of T at resonance. It is easy to show (Herndon *et al.*, 1970) that the relation is

$$\text{"Speed"}(\text{res}) = \left. \frac{dT}{dm} \right|_{m=M} = \frac{x_e}{\Gamma(M)/2}, \quad (8)$$

where the elasticity, $x_e = \Gamma_e/\Gamma$, is introduced next. Further properties of "Speed" are discussed in the baryon mini-review at the front of the baryon Data Card Listings.

For an *inelastic* resonance feeding into channel β ,

$$T_{1\beta} = \frac{\frac{1}{2}(\Gamma_1\Gamma_\beta)^{1/2}}{M-m-i\Gamma/2} = (x_1x_\beta)^{1/2} \frac{\Gamma/2}{M-m-i\Gamma/2}, \quad (9)$$

where

$$\Gamma = \sum_1^N \Gamma_\beta, \quad x_\beta = \Gamma_\beta/\Gamma, \quad (10)$$

and x_1 (called the elasticity) is often written x_e .

The channel cross section $\sigma_{1\beta}$ for the reaction $1 \rightarrow \beta$ is

$$\sigma_{1\beta} = 4\pi\lambda^2 (J+\frac{1}{2}) |T_{1\beta}|^2, \quad (11)$$

where $J = l \pm \frac{1}{2}$.

Resonances seen in production are even more complicated. $(\Gamma_1)^{1/2}$ disappears from T , and must be replaced with some model-dependent parametrization of the production process.

In conclusion, we have seen that because of the energy dependence of Γ even the amplitude T for a resonance does not have a full width at half-maximum equal to Γ (but it does peak at or near M). Then kinematic factors enter into the cross section for formation [Eq. (11)] or production, and displace the observed peak away from M . For quantitative examples, see Barbaro-Galtieri (1968).

Most of the useful information on the N , Δ , Λ , and Σ baryon resonances with $M < 2000$ MeV has come from partial-wave analysis. Masses and widths of most of these states are dependent on the model, as well as on the data used by the different groups that performed these analyses; therefore, the masses in the main baryon Tables are not averages, but plausible guesses, and the errors are "external errors" based on the consistency among different analyses. For the procedures adopted, different from resonance to resonance,

see the appropriate mini-review in the Data Card Listings.

Resonances with mass $M > 2000$ MeV have been detected primarily in total-cross-section experiments.

We can use Eq. (11) to relate the height of the peak at resonance σ_{res} to the elasticity x_e . At resonance, the *channel* cross section is

$$\sigma_{\text{res}}(1 \rightarrow \beta) = 4\pi\lambda^2 (J+\frac{1}{2}) x_e x_\beta, \quad (12)$$

and the *total* cross section is

$$\sigma_{\text{res}}^{(\text{total})} = 4\pi\lambda^2 (J+\frac{1}{2}) x_e. \quad (13)$$

If J is known, we can solve for x_e . If J is not known, the product $(J+\frac{1}{2})x_e$ is given in the baryon table.

D. Muon-Decay Parameters

The μ -decay parameters describe the momentum spectrum (ρ and η), the asymmetry (ξ and δ), and the helicity (h) of the electron in the process $\mu^\pm \rightarrow e^\pm + \nu + \bar{\nu}$. Assuming a local and lepton-conserving interaction, the matrix element may be written as

$$\sum_i \langle \bar{e} | \Gamma_i | \mu \rangle \langle \bar{\nu} | \Gamma_i (C_i + C_i' \gamma_5) | \nu \rangle, \quad (14)$$

where the summation is taken over $i = S, V, T, A, P$. Using the definitions and sign conventions of Kinoshita and Sirlin (1957), we have, for the momentum parameters:

$$\rho = [3g_A^2 + 3g_V^2 + 6g_T^2]/D, \quad (15)$$

$$\eta = [g_S^2 - g_P^2 + 2g_A^2 - 2g_V^2]/D; \quad (16)$$

for the asymmetry parameters:

$$\xi = [+6g_S g_P \cos \phi_{SP} - 8g_A g_V \cos \phi_{AV} + 14g_T^2 \cos \phi_{TT}]/D, \quad (17)$$

$$\delta = [-6g_A g_V \cos \phi_{AV} + 6g_T^2 \cos \phi_{TT}]/D\xi; \quad (18)$$

and for the parameter describing the helicity of the electron:

$$h = \pm [2g_S g_P \cos \phi_{SP} - 8g_A g_V \cos \phi_{AV} - 6g_T^2 \cos \phi_{TT}]/D. \quad (19)$$

Here

$$D = g_S^2 + g_P^2 + 4g_A^2 + 4g_V^2 + 6g_T^2, \quad (20)$$

$$g_i^2 = |C_i|^2 + |C_i'|^2, \quad (21)$$

and

$$\cos \phi_{ij} = \text{Re}(C_i^* C_j' + C_i' C_j^*). \quad (22)$$

The quantities g_i are defined to be real nonnegative numbers, and the ϕ_{ij} are phase angles between the i -type and j -type interactions. Under the assumption of two-component neutrinos, $C_i' = -C_i$, and $C_j' = -C_j$, the S , P , and T terms vanish, and ϕ_{AV} is the phase angle between C_A and C_V in the complex plane.

By using the above equations and the experimental determinations of ρ , η , ξ , δ , and h , limits can be placed on g_S/g_V , g_A/g_V , g_T/g_V , g_P/g_V , and ϕ_{AV} . The results, listed in the data cards, assume neither two-component neutrinos nor time-reversal invariance. If, however, two-component neutrinos are assumed, then $\sin \phi_{AV}$ is the amplitude of time-reversal violation. Note that most experiments study only the upper end of the spectrum where ρ and η are highly correlated, so they can only report ρ for $\eta \equiv 0$ and η for $\rho \equiv \frac{3}{4}$. The values for ρ and η we use here were obtained by combining measurements of both upper and lower ends of the spectrum and turn out to be nearly uncorrelated.

Note also that the radiative corrections are unambiguous only when $g_S = g_T = g_P = 0$. The same limits on g_A/g_V and ϕ_{AV} are obtained, however, as when g_S , g_T , and g_P are left free.

E. K -Decay Parameters

1. Dalitz Plot for $K \rightarrow 3\pi$ Decays

The small deviation from uniformity of the Dalitz plot for the 3π decay of the K meson is usually described by a "slope parameter" (Dalitz, 1956). For the τ and τ' decays of the charged K 's, and the τ^0 decay mode of the K_L^0 , we parameterize the Dalitz plot distribution by the expression

$$|M|^2 \propto 1 + g(s_3 - s_0)/m_{\pi^+}{}^2 + h[(s_3 - s_0)/m_{\pi^+}{}^2]^2 + j(s_2 - s_1)/m_{\pi^+}{}^2 + \dots, \quad (23)$$

where $m_{\pi^+}{}^2$ has been introduced so as to make the coefficients g , h , and j dimensionless, and

$$s_i = (\mathbf{p}_K - \mathbf{p}_i)^2 = (m_K - m_i)^2 - 2m_K T_i \quad i = 1, 2, 3, \quad (24)$$

$$s_0 = \frac{1}{3} \sum_i s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2). \quad (25)$$

Here the \mathbf{p}_i are 4-vectors, m_i and T_i are mass and kinetic energy of the i th pion, and the index 3 is used for the odd pion.

The coefficient g is a measure of the slope in the variable s_3 (or T_3) of the Dalitz plot, while h measures the quadratic dependence on s_3 . The coefficient j is related to the asymmetry of the plot and must be zero if CP invariance holds. Note also that if CP is good, g must be the same for τ^+ and τ^- , and similarly for h .

At present there is no compelling experimental evidence for either the h or the j term [for upper limits on the j term, see Section E.2(b) below]. Thus we stop the above expansion at the first term and list only g . Since different experiments use different forms for $|M|^2$, in order to compare the experiments we have converted to g whatever coefficients have been measured. See the *listings* for details on this point.

Relations among τ^\pm , τ'^\pm , and τ^0 are predicted by the $\Delta I = \frac{1}{2}$ rule. See Appendix I for these relations and a discussion of this rule.

2. CP Violation in K^0 Decays

We list parameters for four different reactions in which CP can be tested [Okun and Rubbia (1967), Steinberger (1969)].

(a) $K_S \rightarrow \pi^+ \pi^- \pi^0$. The quantity measured here is the ratio of amplitudes

$$A_S(K_S \rightarrow \pi^+ \pi^- \pi^0) / A_L(K_L \rightarrow \pi^+ \pi^- \pi^0) \equiv x + iy. \quad (26)$$

If CPT invariance holds and there is no $I=3$ state present, then x is zero, and one measures y , which is the CP -violating part. In the K_S^0 *listings*, we give the results for (26) under Branching Ratio $R4$. Our procedure is to assume that $x=0$, and to list $(A_S/A_L)^2$ in the form of a branching ratio.

(b) *Charge asymmetry in $K_L \rightarrow 3\pi$ decays*. As mentioned above, the presence of a term in $(s_2 - s_1)$ in Expression (23) describing the Dalitz plot distribution for τ^\pm , τ^0 decays of K mesons would be an indication of CP violation. Rather than listing values of the $(s_2 - s_1)$ coefficient j in Eq. (23), we choose to list σ_\pm from the equivalent expression

$$|M|^2 \propto 1 + \sigma_\pm \frac{2}{3} (T_+ - T_-) / T_{\pm \max} + (CP \text{ nonviolating terms}), \quad (27)$$

where T_\pm are the kinetic energies of the charged pions. We have momentarily abandoned the form involving the Mandelstam variables s_i in favor of (27) because the latter has been consistently used by experimenters searching for CP violation. We list σ_\pm among the CP -violating parameters at the back of the K_L^0 data cards. Note that only upper limits have been reported for this quantity.

(c) *Asymmetry in the $K_L \rightarrow \pi^\mp l^\pm \nu$ decays*. The quantity measured and compiled here is

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^- l^+ \nu) - \Gamma(K_L \rightarrow \pi^+ l^- \nu)}{\Gamma(K_L \rightarrow \pi^- l^+ \nu) + \Gamma(K_L \rightarrow \pi^+ l^- \nu)}. \quad (28)$$

This asymmetry violates CP invariance. If CPT is good, for a pure K_L^0 beam, δ can be written as

$$\delta = 2 \frac{1 - |x|^2}{|1 - x|^2} \text{Re } \epsilon, \quad (29)$$

where x is the $\Delta S = \Delta Q$ -violating parameter defined in Sec. E.3, and ϵ is the parameter of the expansion

$$|K_L\rangle = \{[(1 + \epsilon) |K\rangle - (1 - \epsilon) |\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}\}, \quad (30a)$$

$$|K_S\rangle = \{[(1 + \epsilon) |K\rangle + (1 - \epsilon) |\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}\}. \quad (30b)$$

(d) $K_L \rightarrow 2\pi$ decay. The relevant parameters are:

$$\eta_{+-} = A(K_L \rightarrow \pi^+\pi^-) / A(K_S \rightarrow \pi^+\pi^-) = |\eta_{+-}| \exp(i\phi_{+-}), \quad (31)$$

$$\eta_{00} = A(K_L \rightarrow \pi^0\pi^0) / A(K_S \rightarrow \pi^0\pi^0) = |\eta_{00}| \exp(i\phi_{00}), \quad (32)$$

ϵ , defined in Eqs. (30) above, and

$$\epsilon' = (i/\sqrt{2}) \exp i(\delta_2 - \delta_0) \text{Im}(A_2/A_0). \quad (33)$$

Here A_i and δ_i are the amplitude and phase of $\pi\pi$ scattering at the K mass, defined by

$$\langle I=0 | T | K \rangle = \exp(i\delta_0 A_0), \quad (34a)$$

$$\langle I=2 | T | K \rangle = \exp(i\delta_2 A_2). \quad (34b)$$

Wu and Yang (1964) have derived the relationships

$$\eta_{+-} = \epsilon + \epsilon', \quad (35a)$$

$$\eta_{00} = \epsilon - 2\epsilon'. \quad (35b)$$

At present many models have been proposed to explain the experimental results on CP violation, but more data are needed before the cause of CP violation can be ascertained.

We list η_{+-} , η_{00} , ϕ_{+-} , and ϕ_{00} . The phases are measured directly, whereas the magnitudes η_{+-} and η_{00} are derived parameters. We use, as far as we can, the directly measured quantities as input, and calculate η_{+-} and η_{00} from the values given by our constrained fits. Therefore, if one looks at the Data Card Listings, most of the $|\eta|$ measurements appear in the form of branching ratios, with appropriate comments. We then give the values of η_{+-} and $|\eta_{00}|^2$ (Particle Data Group, 1969) in a separate list at the end of the CP -violating parameters section of the K_L^0 Data Card Listings.

3. $\Delta S = \Delta Q$ Rule in K^0 Decays.

The relative amount of $\Delta S \neq \Delta Q$ component present is measured by the parameter x , defined as

$$x = A(\bar{K}^0 \rightarrow \pi^- \ell^+ \nu) / A(K^0 \rightarrow \pi^- \ell^+ \nu). \quad (36)$$

We list $\text{Re}\{x\}$ and $\text{Im}\{x\}$.

4. Form Factors in K_{l3} Leptonic Decays.

Assuming that only the vector current contributes to these decays, we write the matrix element as

$$M \propto f_+(q^2) [(P_K + P_\pi)_\mu \bar{u} \ell \gamma_\mu (1 + \gamma_5) u_\nu] + f_-(q^2) [m_\ell \bar{u} \ell (1 + \gamma_5) u_\nu], \quad (37)$$

where P_K and P_π are the four momenta of K and π mesons, respectively; m_ℓ is the lepton mass; f_+ and f_- are dimensionless form factors which can depend only on $q^2 = (P_K - P_\pi)^2$, the square of the momentum transfer to the leptons. The parameters we list are:

λ_\pm , the energy dependence of the $f_\pm(q^2)$ form factor

$$f_\pm(q^2) = f_\pm(0) [1 + \lambda_\pm (q/m_\pi)^2]; \quad (38)$$

and ξ , the ratio of the two form factors

$$\xi = f_- / f_+. \quad (39)$$

The quantity ξ can be determined in different ways:

(1) by measuring the $K_{\mu 3} / K_{e 3}$ branching ratio and comparing it with the theoretical ratio as given in terms of $\xi(0) = f_-(0) / f_+(0)$.

$$\begin{aligned} \Gamma(K_{\mu 3}^\pm) / \Gamma(K_{e 3}^\pm) &= 0.6457 + 0.1264 \text{Re } \xi + 0.0192 |\xi|^2 \\ &\quad + 1.4115 \lambda_+ + 0.4754 \lambda_- \text{Re } \xi + 0.0080 \lambda_+ \text{Re } \xi, \\ \Gamma(K_{\mu 3}^0) / \Gamma(K_{e 3}^0) &= 0.6452 + 0.1246 \text{Re } \xi + 0.0186 |\xi|^2 \\ &\quad + 1.3162 \lambda_+ + 0.4370 \lambda_- \text{Re } \xi + 0.0064 \lambda_+ \text{Re } \xi. \end{aligned} \quad (40)$$

See CABIBBO 66 and FEARING 70 (for the charge-dependent formulae) in the card listings. Note that the first constant has been changed to 0.6457; the earlier value was a misprint,* which we copied from CABIBBO 66.

(2) by studying the Dalitz plot of the $K_{\mu 3}$ decay. The $K_{e 3}$ Dalitz plot distribution is only dependent upon the λ_+ parameter, whereas the $K_{\mu 3}$ distribution is dependent upon λ_- , λ_+ , ξ . Often experimenters have measured only the momentum spectrum of either the π or the lepton and compared it with the predicted spectrum. See note in the K^+ Data Card Listings for a discussion of this method. For a formula relating the Dalitz plot variables to ξ , see for example, BRENE 61 in the K^\pm card listings.

(3) by measuring the muon polarization in $K_{\mu 3}$ decay. In the rest frame of the K , the μ is expected to be polarized in the direction \mathbf{A} with $\mathbf{P} = \mathbf{A} / |\mathbf{A}|$, where \mathbf{A} is given (CABIBBO 64 in K^+ card listings) by

$$\begin{aligned} A &= \alpha_1(\xi) \mathbf{p}_\mu - \alpha_2(\xi) \{ (\mathbf{p}_\mu / m_\mu) [(m_K - E_\pi) \\ &\quad + (\mathbf{p}_\pi \cdot \mathbf{p}_\mu) (E_\mu - m_\mu) / |\mathbf{p}_\mu|^2] + \mathbf{p}_\pi \} \\ &\quad + m_K \text{Im } \xi(q^2) (\mathbf{p}_\pi \times \mathbf{p}_\mu). \end{aligned} \quad (41)$$

If time-reversal invariance holds, ξ is real, and thus there is no polarization perpendicular to the K -decay plane.

See the note in the Listings, after K^+ decays, for discussions of experimental results.

F. Baryon-Decay Parameters

A/V ratio for baryon leptonic decays. The baryon part of the matrix element for these decays may be written as

$$\langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) | B_i \rangle, \quad (42)$$

where B_i and B_f represent initial and final baryons, and

* We thank Drs. H. W. Fearing and J. Smith for calling this mistake to our attention.

g_A and g_V the axial and vector coupling constants. Here the Pauli representation is used for the γ matrices. The definition of g_A/g_V is

$$g_A/g_V = |g_A/g_V| e^{i\delta}, \quad (43)$$

where δ is $0+n\pi$ if time-reversal invariance holds (see JACKSON 57 in neutron card listings).

In neutron beta decay the measurements are consistent with time reversal, so g_A/g_V therefore is nearly real and has been considered to be such in all the baryon leptonic decays. Notice that by using the above definition of the matrix element with the Pauli representations, the value of g_A/g_V in neutron beta decay is negative.

We compile the ratio g_A/g_V with its sign, for those decays for which it has been measured. For the neutron beta decay we compile also the phase δ .

Asymmetry parameters in nonleptonic hyperon decays. The transition matrix for the hyperon decay may be written as

$$M = s + p(\boldsymbol{\sigma} \cdot \mathbf{q}), \quad (44)$$

where s and p are the parity-changing and the parity-conserving amplitudes, respectively, $\boldsymbol{\sigma}$ is the Pauli spin operator, and \mathbf{q} is a unit vector along the direction of the decay baryon in the hyperon rest frame.

The asymmetry parameters are defined by the relations

$$\alpha = 2 \operatorname{Re} (s^*p) / (|s|^2 + |p|^2), \quad (45a)$$

$$\beta = 2 \operatorname{Im} (s^*p) / (|s|^2 + |p|^2), \quad (45b)$$

$$\gamma = (|s|^2 - |p|^2) / (|s|^2 + |p|^2). \quad (45c)$$

With the transition matrix (44), the angular distribution of the decay baryon, in the hyperon rest system, is of the form

$$I = 1 + \alpha \mathbf{P}_Y \cdot \mathbf{q}, \quad (46)$$

where $\mathbf{P}_Y = \langle Y | \boldsymbol{\sigma} | Y \rangle$ is the hyperon polarization.

In the notation of Lee and Yang (1957), the polarization \mathbf{P}_B of the decay baryon is*

$$\mathbf{P}_B = \frac{[(\alpha + \mathbf{P}_Y \cdot \mathbf{q})\mathbf{q} + \beta(\mathbf{P}_Y \times \mathbf{q}) + \gamma\mathbf{q} \times (\mathbf{P}_Y \times \mathbf{q})]}{(1 + \alpha \mathbf{P}_Y \cdot \mathbf{q})}, \quad (47)$$

where \mathbf{P}_B is defined in that rest system of the baryon obtained by a Lorentz transformation along \mathbf{q} from the hyperon rest system in which \mathbf{q} and \mathbf{P}_Y are defined. Note that α is the helicity of the decay baryon for unpolarized hyperons.

The three parameters α , β , and γ satisfy the relation

$$\alpha^2 + \beta^2 + \gamma^2 = 1. \quad (48)$$

* Note that Ref. 13 contains a misprint. The minus sign in the definition of β should be replaced by a 2. In addition, our unit vector \mathbf{q} is the direction of the baryon, whereas their unit vector \mathbf{p} is the direction of the pion.

It is then convenient to describe hyperon nonleptonic decays in terms of the two independent parameters α and the angle ϕ defined by

$$\beta = (1 - \alpha^2)^{1/2} \sin \phi, \quad (49a)$$

$$\gamma = (1 - \alpha^2)^{1/2} \cos \phi, \quad (49b)$$

which has a more nearly Gaussian distribution than β or γ . Evidently

$$-\frac{1}{2}\pi \leq \phi \leq \frac{1}{2}\pi \quad \text{for } \gamma > 0, \quad (50a)$$

$$+\frac{1}{2}\pi \leq \phi \leq \frac{3}{2}\pi \quad \text{for } \gamma < 0. \quad (50b)$$

In discussing time-reversal invariance, the quantity of interest is Δ , defined by

$$\alpha = 2 |s| |p| \cos \Delta / (|s|^2 + |p|^2), \quad (51a)$$

$$\beta = -2 |s| |p| \sin \Delta / (|s|^2 + |p|^2); \quad (51b)$$

that is, Δ is the phase angle of s relative to p . Evidently

$$-\frac{1}{2}\pi \leq \Delta \leq \frac{1}{2}\pi \quad \text{for } \alpha > 0, \quad (52a)$$

$$+\frac{1}{2}\pi \leq \Delta \leq \frac{3}{2}\pi \quad \text{for } \alpha < 0. \quad (52b)$$

Under the assumption of time-reversal invariance, the angle Δ must satisfy the relation

$$\Delta = \delta_s - \delta_p, \quad (53)$$

modulo π , where δ_s and δ_p are the pion-baryon scattering phase shifts at the appropriate energy and for the appropriate isospin state. For Λ decay, assuming the validity of the $|\Delta I| = \frac{1}{2}$ rule,

$$\Delta = \delta_s - \delta_p = (6.8 \pm 2.0) \text{ deg.}^*$$

On the data cards we list α and ϕ for each decay since they are the most closely related to the experiments and are essentially uncorrelated. Whenever necessary we have changed the signs of the reported values, so as to agree with our conventions. In the Stable Particles Table we give α , ϕ , and Δ with errors; and for convenience we also give the central value of γ , without an error.

V. STATISTICAL PROCEDURES

This section is a much abbreviated version of Sec. IX in the text of our January 1970 edition (Particle Data Group, 1970), to which the reader is referred details. See also the mini-review on K^* masses and mass differences in the Data Card Listings.

A. Confidence Levels and Errors

Quoted errors represent one standard deviation (σ). Upper and lower limits represent 68.3% confidence bounds (1σ), unless otherwise stated.

* This value for $\delta_s - \delta_p$ is derived from the phase-shift analyses by Roper *et al.* (1965). The error is our estimation of the uncertainty.

The errors in the Tables and the errors of the averages in the *listings* often include a scale factor S ; see Sec. V.B below.

Quantities that have changed more than 1σ since our January 1970 edition (Particle Data Group, 1970) are italicized in the Tables. For a discussion see Sec. V.B in the text of that edition.

B. Unconstrained Averaging, Scale Factors

In the absence of constraints, we calculate a weighted average

$$\bar{x} \pm \delta\bar{x} = (\sum \omega_i x_i / \sum \omega_i) \pm [1 / (\sum \omega_i)^{1/2}];$$

$$\omega_i = [1 / (\delta x_i)^2], \quad (54)$$

where the sums run over N experiments. We also calculate χ^2 and compare it with its expectation value of $N-1$.

If $\chi^2 > N-1$, we increase the error $\delta\bar{x}$ in Eq. (54) by a factor

$$S = [\chi^2 / (N-1)]^{1/2}. \quad (55)$$

It is easy to design statistical tests for determining whether one experiment (or a group of experiments) is consistent with the other experiments. However, statistics do not tell us who is wrong in case of contradictions. When $S \gg 1$, one can conclude either that:

- (1) some (or all) experiments are wrong or
- (2) some (or all) experiments have underestimated their errors, or
- (3) the experiments do not measure the same quantity (systematic errors).

We do our best to resolve these cases. If we cannot, we *assume* that *all* experimentalists underestimated their errors by the same scale factor. If we scale up all input errors by this factor, χ^2 returns to $N-1$, and of course the output error scales up by the same factor.

If all the experiments have errors of about the same size, the above procedure is straightforward. If, however, there are both precise and imprecise (large errors) measurements of a particular quantity, one must be very careful not to permit the imprecise ones to "dilute" the scale factor. See our January 1970 edition (Particle Data Group, 1970) for the prescription we use to handle this effect.

We often plot an ideogram to guide the reader in deciding which data he might reject before making his own selected average.

For further discussion of ideograms and scale factors, we refer the reader to Sec. IX of our January 1970 edition (Particle Data Group, 1970).

C. Constrained Fits

The information on partial-decay fractions P_i and partial widths $\Gamma_i = P_i \Gamma_{\text{total}}$ is frequently given by

branching ratios R_j , say, $R_1 = P_1 / (P_1 + P_2)$, $R_2 = P_2 / P_3$, $R_3 = P_1 / P_2$, $R_4 = P_3 / (P_1 + P_2 + P_3)$, etc.*

The number of experimental inputs R_j is often greater than the number of decay modes. In these cases we fit all available information on the P_i , Γ_i , and R_j subject to the constraint $\sum P_i = 1$. When, in addition, the input R_j are contradictory so that scale factors may have to be introduced, one has to resort to iterative procedures.

The Data Card Listings give the values of the fitted R_j , P_i , and Γ_i , together with the error matrices of the P_i and of the Γ_i . For details about this procedure the reader is referred to the text of the January 1970 edition (Particle Data Group, 1970), Sec. IV.B.

VI. PARTICLE DATA GROUP PUBLICATIONS

To obtain a reprint of this report, or any of the items listed below, write either Scientific Information Service, CERN, or Technical Information Division, Lawrence Radiation Laboratory, Berkeley, whichever is closer.

A. Pocket-Sized Particle Data Booklet

In addition to the present version of the Review of Particle Properties available from CERN and LRL, a pocket booklet is available. This contains the tables only, plus some additional useful information, selected for high-energy physicists. The complete pocket version comprises the data booklet, a 16-month diary, a mini-atlas, and a plastic cover. Any of these items that you have requested in the past will automatically be sent to you, but please note that our mailing lists are self-cancelling; unless you return the request card that is sent with each mailing, your name will be removed from our mailing list. If you wish to order any items in bulk we must charge 25 cents (US) for each of the pocket-sized items.

B. Other Compilations

The following reports of the Particle Data Group and of the CERN HERA Group are available from both CERN and LRL:

1. Cross Sections

A Compilation of *K⁺N Reactions*, L. R. Price, N. Barash-Schmidt, O. Benary, R. W. Bland, A. H. Rosenfeld, C. G. Wohl, UCRL-20000 *K⁺N* (September 1969).

A Compilation of *YN Reactions*, O. Benary, N. Barash-Schmidt, L. R. Price, A. H. Rosenfeld, G. Alexander, UCRL-20000 *YN* (January 1970).

Compilation of *Elastic Scattering Data*, G. C. Fox and C. Quigg, UCRL-20001 (January 1970).

NN and ND Interactions (above 0.5 GeV/c)—A Compilation, O. Benary, L. R. Price, G. Alexander, UCRL-20000 *NN* (August 1970).

* We are also able to fit *products* of rates from formation experiments as given in Eq. (12).

A Compilation of *Pion-Nucleon Scattering Data*, G. Giacomelli, P. Pini, and S. Stagni, CERN/HERA 69-1.

Data Compilation of *Antiproton-proton Reactions into Antihyperon-hyperon*, B. Sadoulet, CERN/HERA 69-2.

A Compilation of *Total and Total Elastic Cross Sections*, G. Giacomelli, CERN/HERA 69-3.

A Collection of *Pion Photoproduction Data*. I—From the Threshold to 1.5 GeV, P. Spillantini and V. Valente, CERN/HERA 70-1.

Compilation of Cross Sections. I—*Proton Induced Reactions*, J. D. Hansen, D. R. O. Morrison, N. Tovey, E. Flaminio, CERN/HERA 70-2.

Compilation of Cross Sections. II—*Antiproton Induced Reactions*, E. Flaminio, J. D. Hansen, D. R. O. Morrison, N. Tovey, CERN/HERA 70-3.

Compilation of Cross Sections. III—*K⁺ Induced Reactions*, E. Flaminio, J. D. Hansen, D. R. O. Morrison, N. Tovey, CERN/HERA 70-4.

Compilation of Cross Sections. IV— *π^+ Induced Reactions*, E. Flaminio, J. D. Hansen, D. R. O. Morrison, N. Tovey, CERN/HERA 70-5.

2. Partial-wave amplitudes

πN Partial-Wave Amplitudes—A Compilation, D. J. Herndon, A. Barbaro-Galtieri, A. H. Rosenfeld, UCRL-20000 πN (February 1970).

Each of these reports in turn lists relevant references. A further list of references can be found in: "Compilation of Coupling Constants and Low-Energy Parameters," by G. Ebel *et al.*, [Nucl. Phys. **B17**, 1 (1970), and Springer Tracts in Modern Physics, Vol. 55, Sept. 1970 edition, edited by G. Höhler.]

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PARTICLE PROPERTIES TABLES

April 1971

"Now go, write it before them in a table, and note it in a book,
that it may be for the time to come for ever and ever."

Isaiah 30:7-8

". . . or at least until the next edition."

Particle Data Group

N. Barash-Schmidt, A. Barbaro-Galtieri, C. Bricman, T. Lasinski,
A. Rittenberg, M. Roos, A. H. Rosenfeld, P. Söding, T. G. Trippe and C. G. Wohl

(Closing date for data: Feb. 1, 1971)

Stable Particles

For additional parameters, see Addendum to this table.

Quantities in italics have changed by more than one (old) standard deviation since January 1970.

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean life (sec) $c\tau$ (cm)	Partial decay mode		p or Pmax ^b (MeV/c)	
				Mode	Fraction ^a		
γ	$0, 1(1^-)^-$	$0(<2.)10^{-21}$	stable	stable			
ν	ν_e $J = \frac{1}{2}$ ν_μ	$0(<60 \text{ eV})$ $0(<1.6)$	stable	stable			
e	$J = \frac{1}{2}$	0.5110041 $\pm.0000016$	stable ($>2 \times 10^{21} \text{ y}$)	stable			
μ	$J = \frac{1}{2}$ $m^2 = 0.0112$ $m_\mu - m_{\pi^\pm} = -33.916$ $\pm.011$	105.6599 $\pm.0014$	2.1983×10^{-6}	$e\nu\bar{\nu}$	100	53	
			$\pm.0008$	$e\gamma\gamma$	(<1.6)	10^{-5}	53
			$c\tau = 6.590 \times 10^4$	$3e$	(<1.3)	10^{-7}	53
				$e\gamma$	(<2.2)	10^{-8}	53
π^\pm	$1^-(0^-)$ $m^2 = 0.0195$	139.576 $\pm.011$	2.6024×10^{-8}	$\mu\nu$	100 %	30	
			$\pm.0024$	$e\nu$	(1.24 ± 0.03)	10^{-4}	70
			$c\tau = 780.2$	$\mu\nu\gamma$	^c (1.24 ± 0.25)	10^{-4}	30
			$(\tau^+ - \tau^-)/\bar{\tau} =$ $(0.05 \pm 0.07)\%$	$\pi^0 e\nu$	(1.02 ± 0.07)	10^{-8}	5
			(test of CPT)	$e\nu\gamma$ e^+e^-	^c (3.0 ± 0.5)	10^{-8}	70
			(<3.4)	10^{-8}	70		
π^0	$1^-(0^-)^+$ $m^2 = 0.0182$ $m_{\pi^\pm} - m_{\pi^0} = 4.6041$ $\pm.0037$	134.972 $\pm.012$	0.84×10^{-16}	$\gamma\gamma$	(98.84 ± 0.04)%	67	
			$\pm.10 \text{ S} = 2.1^*$	γe^+e^-	(1.16 ± 0.04)%	67	
			$c\tau = 2.5 \times 10^{-6}$	$\gamma\gamma\gamma$	(<5)	10^{-6}	67
				$e^+e^-e^+e^-$	^d (3.47)	10^{-5}	67

Stable Particles

Particle	$I(GJP)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean life (sec) $c\tau$ (cm)	Partial decay mode				
				Mode	Fraction ^a	p or Pmax ^b (MeV/c)		
K^\pm	$\frac{1}{2}(0^-)$	493.84 ± 0.11 $m^2 = 0.244$ $m_{K^\pm} - m_{K^0} = -3.95$ ± 0.13 $S = 1.1^*$	1.2371×10^{-8} ± 0.0026 $S = 1.9^*$ $c\tau = 370.8$ $(\tau^+ - \tau^-) / \bar{\tau} =$ $(.11 \pm .09)\%$ (test of CPT) $S = 1.2^*$	$\mu\nu$	(63.77 \pm 0.28)%	$S = 1.1^*$	236	
				$\pi\pi^0$	(20.92 \pm 0.29)%	$S = 1.2^*$	205	
				$\pi\pi^-\pi^+$	(5.58 \pm 0.03)%	$S = 1.1^*$	126	
				$\pi\pi^0\pi^0$	(1.68 \pm 0.04)%		133	
				$\mu\pi^0\nu$	(3.20 \pm 0.11)%	$S = 1.8^*$	215	
				$e\pi^0\nu$	(4.86 \pm 0.07)%	$S = 1.1^*$	228	
				$\pi\pi^\mp e^\pm\nu$	(3.3 \pm 0.3)	10^{-5}	204	
				$\pi\pi^\pm e^\mp\nu$	(<7)	10^{-7}	204	
				$\pi\pi^\mp\mu^\pm\nu$	(0.9 \pm 0.4)	10^{-5}	151	
				$\pi\pi^\pm\mu^\mp\nu$	(<3)	10^{-6}	151	
				$e\nu$	(1.30 \pm 0.18)	10^{-5}	247	
				$\pi\pi^0\gamma$	c (<1.9)	10^{-4}	205	
				$\pi\pi^+\pi^-\gamma$	c (10 \pm 4)	10^{-5}	126	
				$\pi e\nu\gamma$	c (6 \pm 4)	10^{-4}	227	
				πe^+e^-	(<0.4)	10^{-6}	227	
				$\pi\mu^+\mu^-$	(<2.4)	10^{-6}	172	
				$\pi\gamma\gamma$	c (<0.4)	10^{-4}	227	
				$\pi\nu\bar{\nu}$	(<1.2)	10^{-6}	227	
$\pi\gamma$	(<4)	10^{-6}	227					
K^0	$\frac{1}{2}(0^-)$	497.79 ± 0.15 $S = 1.1^*$	50% K_{Short}^0 , 50% K_{Long}^0					
K_S^0	$\frac{1}{2}(0^-)$	$S = 1.1^*$ $m^2 = 0.248$	0.862 $\times 10^{-10}$ ± 0.006 $S = 1.2^*$ $c\tau = 2.58$	$\pi^+\pi^-$	(68.7 \pm 0.5)%	$S = 1.3^*$	206	
				$\pi^0\pi^0$	(31.3 \pm 0.5)%		209	
				$\mu^+\mu^-$	(<.7)	10^{-5}	225	
				e^+e^-	(<35)	10^{-5}	249	
K_L^0	$\frac{1}{2}(0^-)$	$m_{K_L^0} - m_{K_S^0} = 0.5398 \times 10^{10} \hbar \text{ sec}^{-1}$ ± 0.0033	5.172×10^{-8} ± 0.043 $c\tau = 1550$	$\pi^0\pi^0\pi^0$	(21.4 \pm 0.7)%	$S = 1.1^*$	139	
				$\pi^+\pi^-\pi^0$	(12.6 \pm 0.3)%		133	
				$\pi\mu\nu$	(26.8 \pm 0.6)%		216	
				$\pi e\nu$	(38.9 \pm 0.6)%	$S = 1.1^*$	229	
				$\pi^+\pi^-$	(0.157 \pm 0.005)%		206	
				$\pi^0\pi^0$	(0.094 \pm 0.019)%	$S = 1.5^*$	209	
				$\pi^+\pi^-\gamma$	c (<0.4)	10^{-3}	206	
				$\gamma\gamma$	(5.6 \pm 0.5)	10^{-4}	249	
				$e\mu$	(<1.6)	10^{-9}	238	
				$\mu^+\mu^-$	(<1.9)	10^{-9}	225	
e^+e^-	(<1.6)	10^{-9}	249					
η	$0^+(0^-)^+$	548.8 ± 0.6 $S = 1.4^*$ $m^2 = 0.301$	$\Gamma = (2.63 \pm 0.59) \text{ keV}$ Neutral decays 72.2%	$\gamma\gamma$	(38.6 \pm 1.1)%	$S = 1.2^*$	274	
				$\pi^0\gamma\gamma$	(3.3 \pm 1.1)%		258	
				$3\pi^0$	(30.3 \pm 1.1)%		180	
				$\pi^+\pi^-\pi^0$	(23.1 \pm 1.0)%		175	
				$\pi^+\pi^-\gamma$	(4.7 \pm 0.2)%		236	
				$\pi^0e^+e^-$	(<0.03)%		258	
				$\pi^+\pi^-e^+e^-$	(0.1 \pm 0.1)%		236	
				$\pi^+\pi^-\pi^0\gamma$	(<0.2)%		175	
				$\pi^+\pi^-\gamma\gamma$	(<0.2)%		236	
				$\mu^+\mu^-$	(2 \pm 1)		10^{-5}	253
				$\mu^+\mu^-\pi^0$	(<5)		10^{-4}	211
				p	$\frac{1}{2}(\frac{1}{2}^+)$		938.2592 ± 0.0052 $m^2 = 0.8803$	stable ($> 2 \times 10^{28} \text{ y}$)
n	$\frac{1}{2}(\frac{1}{2}^+)$	939.5527 ± 0.0052 $m^2 = 0.8828$ $m_p - m_n = -1.29344$ ± 0.00007	$e(0.932 \pm 0.014) 10^3$ $c\tau = 2.79 \times 10^{13}$			$\mu^- \nu$		100

Stable Particles

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean life (sec) $c\tau$ (cm)	Partial decay mode		p or P _{max} ^b (MeV/c)
				Mode	Fraction ^a	
Λ	$0(\frac{1}{2}^+)$ $S=1.3^*$ $m^2=1.245$	1115.59 ± 0.06	2.517×10^{-10} $\pm .024$ $S=1.2^*$ $c\tau = 7.55$	$p\pi^-$	(64.0 ± 0.7)%	$S=1.1^*$ 100
				$n\pi^0$	(36.0 ± 0.7)%	104
				$p e \nu$	(0.80 ± 0.06) 10^{-3}	163
				$p \mu \nu$	(1.35 ± 0.60) 10^{-4}	131
Σ^+	$1(\frac{1}{2}^+)$ $S=1.7^*$ $m^2=1.415$ $m_{\Sigma^+} - m_{\Sigma^-} = -7.95$ $\pm .12$ $S=1.4^*$	1189.42 ± 0.11	0.800×10^{-10} $\pm .006$ $c\tau = 2.40$	$p\pi^0$	(51.7 ± 0.8)%	189
				$n\pi^+$	(48.3 ± 0.8)%	185
				$\gamma\gamma$	(1.24 ± 0.18) 10^{-3}	$S=1.4^*$ 225
				$n\pi^+\gamma$	(1.30 ± 0.24) 10^{-4}	185
				$\Lambda e^+\nu$	(2.02 ± 0.47) 10^{-5}	72
				$\Lambda e^+\nu$	(< 2.4) 10^{-5}	202
$\Lambda e^+\nu$	(< 1.0) 10^{-5}	224				
Σ^0	$1(\frac{1}{2}^+)$	1192.51 ± 0.10 $m^2=1.422$	< 1.0×10^{-14} $c\tau < 3 \times 10^{-4}$	$\Lambda \gamma$	100 %	74
				Λe^+e^-	d(5.45) 10^{-3}	74
Σ^-	$1(\frac{1}{2}^+)$ $S=1.1^*$ $m^2=1.434$ $m_{\Sigma^0} - m_{\Sigma^-} = -4.86$ $\pm .06$	1197.37 ± 0.07	1.489×10^{-10} $\pm .022$ $S=1.8^*$ $c\tau = 4.46$	$n\pi^-$	100 %	193
				$ne^-\nu$	(1.09 ± 0.05) 10^{-3}	230
				$n\mu^-\nu$	(0.45 ± 0.04) 10^{-3}	210
				$\Lambda e^-\nu$	(0.60 ± 0.06) 10^{-4}	79
				$n\pi^-\gamma$	c(1.0 ± 0.2) 10^{-4}	193
Ξ^0	$\frac{1}{2}(\frac{1}{2}^+)^f$ $m^2=1.729$ $m_{\Xi^0} - m_{\Xi^-} = -6.6$ $\pm .7$	1314.7 ± 0.7	3.03×10^{-10} $\pm .18$ $c\tau = 9.08$	$\Lambda\pi^0$	100 %	135
				$p\pi^-$	(< 0.9) 10^{-3}	299
				$pe^-\nu$	(< 1.3) 10^{-3}	323
				$\Sigma^+e^-\nu$	(< 1.5) 10^{-3}	119
				$\Sigma^-e^+\nu$	(< 1.5) 10^{-3}	112
				$\Sigma^+\mu^-\nu$	(< 1.5) 10^{-3}	64
				$\Sigma^-\mu^+\nu$	(< 1.5) 10^{-3}	49
				$p\mu^-\nu$	(< 1.3) 10^{-3}	309
Ξ^-	$\frac{1}{2}(\frac{1}{2}^+)^f$ $m^2=1.746$	1321.31 ± 0.17	1.660×10^{-10} $\pm .037$ $S=1.1^*$ $c\tau = 4.98$	$\Lambda\pi^-$	100 %	139
				$\Lambda e^-\nu$	g(0.67 ± 0.23) 10^{-3}	190
				$\Sigma^0e^-\nu$	(< 0.5) 10^{-3}	123
				$\Lambda\mu^-\nu$	(< 1.3) 10^{-3}	163
				$\Sigma^0\mu^-\nu$	(< 0.5) %	70
				$n\pi^-$	(< 1.1) 10^{-3}	303
				$ne^-\nu$	(< 1.0) %	327
Ω^-	$0(\frac{3}{2}^+)^f$ $m^2=2.797$	1672.5 ± 5	$1.3^{+0.4}_{-0.3} \times 10^{-10}$ $c\tau = 3.9$	$\Xi^0\pi^-$	Total of 28 events seen	294
				$\Xi^-\pi^0$		290
				ΛK^-		211

* S = Scale factor = $\sqrt{\chi^2/(N-1)}$, where $N \approx$ number of experiments. S should be ≈ 1 . If $S > 1$, we have enlarged the error of the mean, δx , i. e., $\delta x \rightarrow S\delta x$. This convention is still inadequate, since if $S \gg 1$, the experiments are probably inconsistent, and therefore the real uncertainty is probably even greater than $S\delta x$. See text and ideogram in data card listings.

a. Quoted upper limits correspond to a 90% confidence level.

b. In decays with more than two bodies, P_{max} is the maximum momentum that any particle can have.

c. See data card listings for energy limits used in measuring this branching ratio.

d. Theoretical value; see also data card listings.

e. See note in data card listings.

f. P for Ξ and J^P for Ω^- not yet measured. Values reported are SU(3) predictions.

g. Assumes rate for $\Xi^- \rightarrow \Sigma^0 e^-\nu$ small compared with $\Xi^- \rightarrow \Lambda e^-\nu$.

ADDENDUM TO
Stable Particles

Magnetic moment								
e	$1.001\ 159\ 644$ $\pm 0.000\ 000\ 007$	$\frac{e\hbar}{2m_e c}$	μ Decay parameters ^a					
μ	$1.001\ 166\ 16$ $\pm 0.000\ 000\ 31$	$\frac{e\hbar}{2m_\mu c}$	$\rho = 0.752 \pm 0.003$	$\eta = -0.12 \pm 0.21$	$\xi = 0.972 \pm 0.013$	$\delta = 0.755 \pm 0.009$	$h = 1.00 \pm 0.13$	
			$ g_A/g_V = 0.86^{+0.33}_{-0.11}$		$\phi = 180^\circ \pm 15^\circ$			
K^\pm	Mode	Partial rate (sec⁻¹)	$\Delta I = \frac{1}{2}$ rule for $K^\pm \rightarrow 3\pi$		Form factors for leptonic decays			
	$\mu\nu$	$(51.55 \pm 0.25)10^6$	$S=1.2^*$	$\pi^+\pi^+\pi^-$	$c_g = -.206 \pm .007$	See listings for λ, ξ		
	$\pi\pi^0$	$(16.91 \pm 0.24)10^6$	$S=1.2^*$	$\pi^-\pi^+\pi^-$	$c_g = -.194 \pm .007$			
	$\pi\pi^+\pi^-$	$(4.51 \pm 0.02)10^6$	$S=1.1^*$	$\pi^+\pi^0\pi^0$	$c_g = .527 \pm .017$			
	$\pi\pi^0\pi^0$	$(1.36 \pm 0.04)10^6$		See also listings and Appendix I				
$\mu\pi^0\nu$	$(2.59 \pm 0.09)10^6$	$S=1.8^*$						
$e\pi^0\nu$	$(3.92 \pm 0.06)10^6$	$S=1.2^*$						
K_S^0	$\pi^+\pi^-$	$(0.796 \pm .007)10^{10}$	$S=1.2^*$	CP violation parameters		$I = \frac{1}{2}$ rule for $K_L^0 \rightarrow 3\pi$		
	$\pi^0\pi^0$	$(0.364 \pm .006)10^{10}$	$S=1.3$	$ \eta_{+-} = (1.95 \pm 0.03)10^{-3}$	$\phi_{+-} = (44 \pm 3)^\circ$	$\pi^+\pi^-\pi^0$	$c_g = .60 \pm .03$	
K_L^0	$\pi^0\pi^0\pi^0$	$(4.14 \pm 0.13)10^6$	$S=1.1^*$	$ \eta_{00} = 2.2 \pm 0.2$	10^{-3}	$\phi_{00} = (51 \pm 30)^\circ$	See also listings & App. I	
	$\pi^+\pi^-\pi^0$	$(2.44 \pm 0.06)10^6$		Charge asymmetry:		$\Delta S = -\Delta Q$		
	$\pi\mu\nu$	$(5.19 \pm 0.12)10^6$		$S=1.5^*$	$\frac{\Gamma(K_L^0 \rightarrow \ell^+) - \Gamma(K_L^0 \rightarrow \ell^-)}{\Gamma(K_L^0 \rightarrow \ell^+) + \Gamma(K_L^0 \rightarrow \ell^-)}$	$Re\ x = -.003 \pm .031$	$S=1.7^*$	
	$\pi e\nu$	$(7.52 \pm 0.14)10^6$	$S=1.1^*$			$Im\ x = -.037 \pm .043$	$S=1.1^*$	
	$\pi^+\pi^-$	$(3.03 \pm 0.10)10^4$				Form factors for leptonic decays		
$\pi^0\pi^0$	$(1.82 \pm 0.37)10^4$	$S=1.5^*$			See listings for λ, ξ			
η	Mode	Asymmetry parameter						
	$\pi^+\pi^-\pi^0$	$(1.2 \pm 0.5)\%$ $S=1.3^*$						
	$\pi^+\pi^-\gamma$	$(1.1 \pm 1.3)\%$						
	Magnetic moment	Decay parameters ^b						
	$(e\hbar/2m_p c)$	Measured		Derived		g_A/g_V	g_V/g_A	
p	2.792782 ± 0.000047							
n	-1.913148 ± 0.000066	$pe^- \nu$				-1.231 ± 0.010	$[\delta = (181.1 \pm 1.3)^\circ]$	
Λ	-0.70 ± 0.07	$p\pi^-$	0.645 ± 0.016	$(-6.3 \pm 3.5)^\circ$	0.76	$(7.4^{+4.0}_{-4.1})^\circ$		
		$n\pi^0$	0.649 ± 0.046					
		$pe\nu$			-0.83 ± 0.18	$S=1.3^*$		
Σ^+	2.59 ± 0.46	$p\pi^0$	-0.991 ± 0.019	$(22 \pm 90)^\circ$	0.12	$(183^{+11}_{-12})^\circ$		
		$n\pi^+$	$+0.066 \pm 0.016$	$(167 \pm 20)^\circ$	-0.97	$(-73^{+136}_{-10})^\circ$		
		$p\gamma$	$-1.03^{+.52}_{-.42}$	$S=1.1^*$				
Σ^-		$n\pi^-$	-0.069 ± 0.008	$(10 \pm 15)^\circ$	0.98	$(249^{+12}_{-115})^\circ$	See listings	
		$\Delta e^- \nu$			0.35 ± 0.18			
Ξ^0		$\Lambda\pi^0$	-0.35 ± 0.08	$(25 \pm 21)^\circ$	0.85	$(228^{+16}_{-37})^\circ$		
Ξ^-		$\Lambda\pi^-$	-0.40 ± 0.03	$(-4 \pm 8)^\circ$	0.91	$(170^{+18}_{-17})^\circ$		

*S = scale factor. Quoted error includes scale factor; see footnote to main Stable Particles Table for definition.

a. $|g_A/g_V|$ defined by $g_A^2 = |C_A|^2 + |C'_A|^2$, $g_V^2 = |C_V|^2 + |C'_V|^2$, and $\Sigma \langle \bar{e} | \Gamma_i | \mu \rangle \langle \bar{\nu} | \Gamma_i (C_1 + C'_1 \gamma_5) | \nu \rangle$;

ϕ defined by $\cos \phi = -Re(C_A^* C'_V + C'_A C_V^*) / g_A g_V$ [for more details, see text].

b. The definition of these quantities is as follows [for more details on sign convention, see text]:

$$\alpha = \frac{2|s||p|\cos \Delta}{|s|^2 + |p|^2}; \quad \beta = \sqrt{1-\alpha^2} \sin \phi; \quad g_A/g_V \text{ defined by } \langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) | B_i \rangle;$$

$$\beta = \frac{-2|s||p|\sin \Delta}{|s|^2 + |p|^2}; \quad \gamma = \sqrt{1-\alpha^2} \cos \phi; \quad \delta \text{ defined by } g_A/g_V = |g_A/g_V| e^{i\delta}.$$

c. The definition of the slope parameter of the Dalitz plot is as follows: $|M|^2 = 1 + g \left(\frac{s_3 - s_0}{m_\pi^2} \right)$.

Mesons

April 1971

Quantities in italics have changed by more than one (old) standard deviation since January 1970

Name	$G \begin{matrix} & 0 & 1 \\ \hline - & \phi & \pi \\ + & \eta & \rho \end{matrix}$	$I^G(J^P)C_n$ estab.	Mass M (MeV)	Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		p or Pmax (MeV/c) ^(b)	
						Mode	Fraction %		
$\pi^{\pm}(140)$ $\pi^0(135)$		$1^-(0^-)^+$	139.58 134.97	0.0 7.2 eV ± 1.2 eV	0.019483 0.018217	See Stable Particles Table			
$\eta(549)$		$0^+(0^-)^+$	548.8 ± 0.6	2.70 keV ± 0.67 keV	0.301 ± 0.000	All neutral $\pi^+\pi^-\pi^0$ + $\pi^+\pi^-\gamma$	72 28	See Stable Particles Table	
η_0 or ϵ	$\begin{matrix} (700-) \\ (1000) \end{matrix}$	$0^+(0^+)^+$	≥ 750 δ_0^0 seems to stay near 90° from 750 to 1150 MeV; see note in listings	$\gg 100$	≥ 0.6	$\pi\pi$	100		
$\rho(765)$		$1^+(1^-)^-$	765 ± 10 (c)	125 ± 20 (c)	0.585 ± 0.095	$\pi\pi$ e^+e^- $\mu^+\mu^-$ For upper limits, see footnote (e)	≈ 100 .0060 \pm .0008 (d) .0067 \pm .0012 (d)	356 382 368	
$\omega(784)$		$0^-(1^-)^-$	783.9 ± 0.3 S=1.4*	11.4 ± 0.9	0.614 ± 0.009	$\pi^+\pi^-\pi^0$ $\pi^+\pi^-$ $\pi^0\gamma$ e^+e^- For upper limits, see footnote (g)	89.8 \pm 4.0 0.93 \pm 0.25 (f) 9.3 \pm 1.2 0.0066 \pm .0017 S=1.4*	328 366 380 392	
η' or X^0		$0^+(\bar{1}^-)^+$ $J^P = 0^-$ or 2^-	957.5 ± 0.8	< 4	0.917 < 0.004	$\eta\pi\pi$ $\pi^+\pi^-\gamma$ (mainly $\rho^0\gamma$) $\gamma\gamma$ [note (h)] For upper limits, see footnote (i)	64.0 \pm 5.0 S=1.1* 29.4 \pm 2.7 S=1.2* 6.6 \pm 3.7	231 458 479	
$\delta(962)$		≥ 1 ()	962 ± 5	< 5	0.925 < 0.005	Observed in missing mass	} Widths incompatible		
$\pi_N(975)$		$1^-(0^+)^+$	975 $\pm 10^S$	58 ± 11	0.950 ± 0.056	$\eta\pi$ possibly seen		315	
Ξ These three could be related; see note in listings.									
$\pi_N(1016)$ $\rightarrow KK$		$1^-(0^+)^+$	1016 ± 10	≈ 25 } if res.	1.032 ± 0.025	$K^{\pm}K^0$ $\eta\pi$	Only mode seen < 80	111 342	
Resonance, virtual bound state, or antibound state, still not distinguished.									
$\phi(1019)$		$0^-(1^-)^-$	1019.5 ± 0.6 S = 1.5*	4.0 ± 0.3	1.039 ± 0.004	K^+K^- $K_L K_S$ $\pi^+\pi^-\pi^0$ (incl. $\rho\pi$) e^+e^- $\mu^+\mu^-$ For upper limits, see footnote (j)	46.4 \pm 2.8 S=1.2* 35.4 \pm 4.0 S=1.6* 18.2 \pm 5.4 S=1.9* .035 \pm .003 .023 \pm .005	125 108 461 509 498	
η_0 or S^*	(1070)	$0^+(0^+)^+$	1070 $\pm 30^S$ if resonance	150-300 (k)	1.14	$\pi\pi$ KK	< 65 > 35	517 201	
Resonance and scattering length both possible (k); see notes in listing on η_0 (700-1000, $\pi\pi$) and η_N (1080, $J > 0$)									
$A_1(1070)$		$1^-(1^+)^+$	1070 $\pm 20^S$ $J^P = 2^-$ not excluded	50-200 dep. on bgr.	1.14	3π [see note (l) on $\rho\pi$] KK [G = (-1) ^{l+I} forbids $K\bar{K}$ for J = 1]	≈ 100 < 0.25	488 201	
Resonance interpretation slightly in doubt; only 160 MeV above $\rho\pi$ threshold. See note (t).									

Mesons

Name	$I^G(J^P)C_{\eta}$ estab.	Mass M (MeV)	Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		p or P _{max} ^(b) (MeV/c)
					Mode	Fraction %	
B(1235)	$\underline{1}^+(1^+)-$	1233 $\pm 10^S$	100 $\pm 20^S$	1.53 $\pm .13$	$\omega\pi$ $\pi\pi\pi$ $K\bar{K}$ For other upper limits, see footnote (m)	≈ 100 < 30 < 2 } Absence suggests J ^P = Abnormal	350 602 371
f(1260)	$0^+(2^+)+$	1269 $\pm 10^S$	154 $\pm 25^S$	1.60 $\pm .20$	$\pi\pi$ $2\pi^+2\pi^-$ $K\bar{K}$	≈ 80 7±2 ≈ 5	618 554 391
D(1285)	$0^+(A)+$ S=1.4* J ^P = 0 ⁻ , 1 ⁺ , 2 ⁻ , with 1 ⁺ favoured	1286 ± 4	33 ± 4	1.65 $\pm .04$	$K\bar{K}\pi$ [mainly $\pi_N(1016)$] $\pi\pi\eta$ $\pi_N(975)\pi$ $\pi\pi\rho$	Seen Possibly large Possibly seen Not seen	305 484 244 352
A ₂ (1310)	$1^-(2^+)+$ Either: A ₂ Wide ($\approx 1300, \Gamma \approx 80$) + A ₂ Narrow ($\approx 1300, \Gamma \approx 10$) or: A ₂ Low ($\approx 1290, \Gamma \approx 25$) + A ₂ High ($\approx 1310, \Gamma \approx 25$)	(1310)	(85)	(1.72)	< > = unsplit and unresolved avge. $\rho\pi$ $\eta\pi$ $K\bar{K}$ $\eta'(958)\pi$	$\langle \approx 76 \rangle$ $\langle \approx 18 \rangle$ $\langle \approx 5 \rangle$ $\langle \sim 1 \rangle$	409 523 420 274
Existing data can be explained by a pair of resonances which are either narrow-on-wide or possibly side-by-side; three recent high-resolution expts. show no split. See note (n).							
E(1422)	$0^+(0^-)+$ J ^P = 1 ⁺ not excluded See note in listings	1422 ± 4	69 ± 8	2.02 $\pm .10$	$K^*\bar{K} + \bar{K}^*K$ $\pi_N(1016)\pi$ $\pi\pi\eta$ $\pi\pi\rho$	50 ± 10 50 ± 10 < 60 Not seen	153 326 568 457
f'(1514)	$0^+(2^+)+$	1514 ± 5	73 ± 23 S=1.8*	2.29 $\pm .11$	$K\bar{K}$ $K^*\bar{K} + \bar{K}^*K$ $\pi\pi$ $\eta\pi\pi$ $\eta\eta$	72 ± 12 10 ± 10 < 14 18 ± 10 < 40 } (o)	570 294 744 624 521
$\pi/\rho(1540)$ or F ₁	$\underline{1}^+(A)$ Evidence based on only one experiment	1540 ± 5	40 ± 15	2.37 $\pm .06$	$K^*\bar{K} + \bar{K}^*K$	Only mode seen	321
$\pi_A(1640)$ or A ₃	$\underline{1}^-(A)+$ J ^P = 2 ⁻ preferred Resonance interpretation in doubt; only 240 MeV above f π threshold, see note (t)	1640 $\pm 10^S$	50-200 (t)	2.67 $\pm .15$	f π 3 π $\omega\pi\pi$ (p)	Dominant Possibly observed	310 792 597
$\rho_N(1650)$	$0^-(N)-$ S=1.2*	1664 ± 13	141 ± 17	2.83 $\pm .23$	$\rho\pi$ 3 π 5 π	Dominant Possibly observed 10±10	639 797
$\rho_N(1660)$ or g	$\underline{1}^+(N)-$ J ^P = 1 ⁻ , 3 ⁻ , ... with 3 ⁻ favoured	1660 $\pm 20^S$	≤ 200	2.76	2 π K \bar{K}	Dominant 7 ± 3	818 664
Nearby peaks in 4 π modes have been grouped under $\rho(1710)$ below. This does not imply that $\rho_N(1660)$ and $\rho(1710)$ are different resonances.							
$\rho(1710)$ + 4 π	$\underline{1}^+()-$	1712 $\pm 10^S$	125 ± 25	2.93 $\pm .21$	4 π + [$\pi^+A_2^0(\rightarrow\pi^+\pi^-\pi^0)/\text{all}(\pi^+\pi^+\pi^-\pi^0)$] + [$\pi^+\omega(\rightarrow\pi^+\pi^-\pi^0)/\text{all}(\pi^+\pi^+\pi^-\pi^0)$] + [$\rho^+\rho^0$] $\pi^+\phi/\pi^+\pi^+\pi^-\pi^0$ $\pi^+2\pi^+2\pi^-\pi^0/\pi^+\pi^+\pi^-\pi^0$ + [$\pi\pi\rho$]	40 ± 20 25 ± 10 Seen < 11 < 15 Seen]	797 339 667 382 539 704 649
See note under $\rho_N(1660)$ above. $\omega\pi$ mode has mass 1633 MeV, see listing.							

← See R-Region Fig. in listings before $\rho_N(1660)$

Mesons

Name	$I^G(J^P)C_n$ estab.	Mass M (MeV)	Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		
					Mode	Fraction %	p or Pmax ^(b) (MeV/c)
☞ See note (q) for states grouped as R(1750), η/ρ (1830), ϕ/π (1830)							
S(1930)	$1^+(\geq 2)^-$	1929 ± 14	20 $\pm 15^S$		$\pi\pi$ $p\bar{p}$	Probably seen Probably seen	955 226
Seen in $\pi^- p \rightarrow (MM)^- p$, appears confirmed in $\bar{p} p \rightarrow \bar{p} p$							
☞ See note (q) for states grouped as ρ (2100), T(2200), ρ (2275), and $N\bar{N}$ (2350).							
U(2375)	$\frac{1}{2}^-$	2374 ± 13	30 $\pm 20^S$	5.57 $\pm .07$	Seen in $\pi^- p \rightarrow (MM)^- p$ and $\bar{p} p \rightarrow \bar{K} K^* \pi \pi$.		
☞ See note (q) for heavier states.							
K^+ (494) K^0 (498)	$1/2(0^-)$	493.84 497.79		0.244 0.248	See Stable Particles Table		
K^* (892)	$1/2(1^-)$	892.6 ± 0.5 $S=1.3^*$ (Charged mode; $m^0 - m^\pm = 8 \pm 3$)	50.3 ± 1.1 $S=1.3^*$	0.797 $\pm .045$	$K\pi$ $K\pi\pi$	≈ 100 0.2	288 216
☞							
Q	$\left\{ \begin{array}{l} K_A(1240) \text{ or } C \\ K_A(1280) \text{ to } 1400 \end{array} \right.$	$\left\{ \begin{array}{l} 1242 \pm 10 \\ 1280 \text{ to } 1400 \end{array} \right.$	$\left\{ \begin{array}{l} 127 \pm 25 \\ \text{See note (t)} \end{array} \right.$	$\left\{ \begin{array}{l} 1.54 \pm .16 \end{array} \right.$	Resonance interpretation unclear; only 270 MeV above $K^* \pi$ threshold; see notes (r) and (t)		
					$\left\{ \begin{array}{l} K\pi\pi \\ +[K^* \pi] \\ +[K\rho] \end{array} \right.$	Only mode seen Large Seen	
K_N (1420)	$1/2(2^+)$	1408 $\pm 10^S$	107 ± 15 $S=2.6^*$	1.98 $\pm .15$	$K\pi$ $K^* \pi$ $K\rho$ $K\omega$ $K\eta$	56.9 ± 4.0 $S = 1.2^*$ 27.4 ± 3.2 $S = 1.1^*$ 9.2 ± 3.5 $S = 1.2^*$ 4.5 ± 1.8 2.0 ± 1.8	609 406 311 291 474
See note (s). $J^P = 3^-$ still possible							
☞							
L(1770)	$1/2(A)$	1770 $\pm 10^S$	50-140 (t)	3.13	$K\pi\pi$ [$K^*(1420)\pi$] $K\pi\pi\pi$	Dominant Large; in $K^+ p$ expts. dominant Possibly seen	791 300 760
Resonance interpretation unclear; only 215 MeV above $K_N(1420)\pi$ threshold; see note (t).							
☞							

Mesons

⇒ The following bumps, excluded above, are listed among the data cards: $\sigma(410)$; $M(953)$; $H(990)$; $\eta_N(1080)$; $A_1(1170)$; $\rho\rho(1410)$; $K_S K_S(1440)$; $R(1750)$; η or $\rho(1830) \rightarrow 4\pi$; ϕ or $\pi(1830) \rightarrow \omega\pi\pi$; $\rho(2100)$; $T(2200)$; $\rho(2275)$; $NN_{I=1}(2350)$; $NN_{I=0}(2375)$; $X^-(2500)$; $X^-(2620)$; $X^-(2800)$; $X^-(2880)$; $X^-(3030)$; $X^-(3075)$; $X^-(3145)$; $X^-(3475)$; $X^-(3535)$; $K_N(1080-1260)$; $K_{A(I=3/2)}(1175)$; $K_{A(I=3/2)}(1265)$; $K_N(1660)$; $K^*(2240) \rightarrow \bar{Y}N$. (See note (q).)

* Quoted error includes scale factor $S = \sqrt{\chi^2/(N-1)}$. See footnote to Stable Particles Table.

† Square brackets indicate a subreaction of the previous (unbracketed) decay mode.

§ This is only an educated guess; the error given is larger than the error of the average of the published values (see listings for the latter).

(a) ΓM is approximately the half-width of the resonance when plotted against M^2 .

(b) For decay modes into ≥ 3 particles, p_{\max} is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated by using the averaged central mass values, without taking into account the widths of the resonances.

(c) The values given for $M(\rho)$ and $\Gamma(\rho)$ and their errors are not average values from various experiments, but rather are intended to give the range where we believe the actual values are most likely to fall. Contrast the results tabulated in this note (references in the listings).

	$M(\text{MeV})$	$\Gamma(\text{MeV})$	
ρ^0	774 ± 5	111 ± 5	} From $e^+e^- \rightarrow \pi^+\pi^-$, fitted to Gounaris-Sakurai formula.
ρ^0	768 ± 10	140 ± 14	
ρ^-	764 ± 2	147 ± 4	} From physical region fits to $\pi N \rightarrow \pi\pi N$, using energy-dependent width.
ρ^0	775 ± 3	145 ± 9	
ρ^0	768 ± 2	132 ± 13	} From pole extrapolation in $\pi N \rightarrow \pi\pi N$
ρ^0	759 ± 7	119 ± 20	
ρ^0	760	131	

(d) The e^+e^- branching ratio is from $e^+e^- \rightarrow \pi^+\pi^-$ experiments only. The $\omega\rho$ interference is then due to $\omega\rho$ mixing only, and is expected to be small. See note in listings. The $\mu^+\mu^-$ branching ratio is compiled from 3 experiments; each possibly with substantial $\omega\rho$ interference. The error reflects this uncertainty; see notes in listings. If $e\mu$ universality holds, $\Gamma(\rho^0 \rightarrow \mu^+\mu^-) = \Gamma(\rho^0 \rightarrow e^+e^-) \times$ phase space correction.

(e) Empirical limits on fractions for other decay modes of $\rho(765)$ are $\pi^+\gamma < 0.5\%$, $\pi^+\eta < 0.8\%$, $\pi^+\pi^+\pi^-\pi^- < 0.15\%$, $\pi^+\pi^+\pi^-\pi^0 < 0.2\%$.

(f) Depends mainly on one experiment, $\gamma C \rightarrow \pi^+\pi^-C$. Uncertainties about ρ shape and production mechanism could increase error.

(g) Empirical limits on fractions for other decay modes of $\omega(784)$ are $\pi^+\pi^-\gamma < 5\%$, $\pi^0\pi^0\gamma < 1\%$, $\eta + \text{neutral}(s) < 1.5\%$, $\mu^+\mu^- < 0.02\%$, $\pi^0\mu^+\mu^- < 0.2\%$.

(h) This $\eta' \rightarrow \gamma\gamma$ value is from a constrained fit under the assumption that $\eta\pi\pi$, $\rho^0\gamma$, and $\gamma\gamma$ are the only existing decay modes. Note that direct measurement of the $\eta' \rightarrow \gamma\gamma$ branching fraction gave the slightly different result of $(9.9^{+4.4}_{-3.9})\%$.

(i) Empirical limits on fractions for other decay modes of $\eta'(958)$: $\pi^+\pi^- < 2\%$, $\pi^+\pi^-\pi^0 < 5\%$; $\pi^+\pi^+\pi^-\pi^- < 1\%$, $\pi^+\pi^+\pi^-\pi^-\pi^0 < 1\%$, $6\pi < 1\%$, $\pi^+\pi^-e^+e^- < 0.6\%$, $\pi^0e^+e^- < 1.3\%$, $\eta e^+e^- < 1.1\%$, $\pi^0\rho^0 < 4\%$, $\pi^0\omega < 8\%$.

(j) Empirical limits on fractions for other decay modes of $\phi(1019)$ are $\pi^+\pi^- < 5\%$, $\eta\gamma < 8\%$, $\eta + \text{neutrals} < 13\%$, $\pi^+\pi^-\gamma < 4\%$, $\omega\gamma < 5\%$, $\rho\gamma < 2\%$, $\pi^0\gamma < 0.35\%$.

(k) Width of $\eta_{0+}(1070) \rightarrow K_S K_S$ uncertain due to proximity of threshold. The data also allow a fit with scattering length and effective range. See listings and W. Beusch's review, Philadelphia Conference 1970, p. 185.

(l) $\rho\pi$ fraction of 3π mode difficult to distinguish because ρ bands cover most of the Dalitz plot.

(m) Empirical limits on fractions for decay modes of $B(1235)$: $\pi\pi < 30\%$, $K\bar{K} < 2\%$, $4\pi < 50\%$, $\phi\pi < 1.5\%$, $\eta\pi < 25\%$, $(\bar{K}K)^{\pm}\pi^0 < 8\%$, $K_S K_S \pi^{\pm} < 2\%$, $K_S K_L \pi^{\pm} < 6\%$.

(n) For an up-to-date discussion see K. W. J. Barnham and G. Goldhaber, UCRL-20293. They adjust the absolute mass scale of different expts. by ~ 10 MeV, and then find that a broad continuum of parameters will fit the data; the values above are typical. Unsplit and unresolved A_2 peaks have the following m and Γ : $\rho\pi$ and $\eta\pi$ modes-- $m = (1300 \pm 4)$ MeV, $\Gamma = (84 \pm 4)$ MeV; $\bar{K}K$ mode-- $m = (1316 \pm 3)$ MeV, $\Gamma = (60 \pm 15)$ MeV.

(o) There is only a weak indication for a $K^*\bar{K} + \bar{K}^*K$ mode of the $f'(1514)$. If this mode does not exist, the $K\bar{K}$ branching fraction will have to be reported as $80 \pm 13\%$ (rather than $72 \pm 12\%$ as given in the table), and $\eta\pi\pi$ as $20 \pm 13\%$.

Mesons

- (p) The possible $\omega\pi\pi$ decay mode of the A3 has mass 1690 MeV and width 45 MeV, in agreement with R2.
 (q) We tabulate here $Y = 0$ bumps with $M \geq 1700$ MeV, for which no satisfactory grouping into particles is yet possible. See listings.

Name	I^G	J^P	M(MeV)	Γ (MeV)	Decay modes observed	Tentative grouping
$K\bar{K}$ (1740)	1		1740	≈ 120	$K^0 K^{\pm}$	} R(1750)
R3(1750)	1,2		1748 ± 15	≤ 38	$(MM)^- \rightarrow 1/3 / >3$ charg. part. $>14 / <80 / 15$	
$\pi\pi$ (1764)			1764 ± 15	$87 + \frac{14}{20}$	$\pi^+ \pi^-$	
$K\bar{K}\pi$ (1820)	0,1,2		1820 ± 12	50 ± 23	$K_S^0 K^0 \pi^0$	} 1830 region
R4(1830)	1,2		1830 ± 15	≤ 30	$(MM)^-$	
η/ρ (1830)	+		1832 ± 6	42 ± 11	$\pi^+ \pi^- \pi^+ \pi^-$	
ϕ/π (1830)	-		1848 ± 11	67 ± 27	$\omega \pi^+ \pi^-$, possibly $\omega\rho^0$	
$X^-(2086)$	1,2		2086 ± 38	≈ 150	$(MM)^-$ backward	} ρ (2100)
ρ (2120)	1^+	3^- (?)	2120	< 249	$\pi^+ \pi^-$, $\bar{p}p$	
$\pi\pi$ (2157)	1^+	(odd)-	2157 ± 10	68 ± 22	$\pi^+ \pi^0$	} T region Seems to require >1 resonance
$N\bar{N}$ (2190)	1^-		2190	20-80	$\rho^0 \rho^0 \pi^0$, $\bar{p}p$	
	1		2190 ± 10	≈ 85	Structure in $N\bar{N}$ total σ	
T(2195)	1,2		2195 ± 15	≤ 13	$(MM)^- \rightarrow 3$ charged particles $\approx 94\%$	
3π (2207)	$\leq 3^-$		2207 ± 13	62 ± 52	$\pi^+ \pi^- \pi^0$	
4π (2200)	$1^+, 2^+, 3^+$		2207 ± 22	≈ 130	$\rho^+ \pi^+ \pi^-$, ω and η antiselected	
$KK\omega$ (2176)	$0^-, 1^+$		2176 ± 5	$20 + \frac{16}{2}$	$K_S^0 K_S^0 \omega$	
$X^-(2260)$	1,2		2260 ± 18	≤ 25	$(MM)^-$ backward	} ρ (2275)
ρ (2290)	1^+	5^- (?)	2290	< 165	$\pi^+ \pi^-$, $\bar{p}p$	
$N\bar{N}$ (2350)	1		2350 ± 10	≈ 140	Structure in $N\bar{N}$ total σ	
$N\bar{N}$ (2375)	0		2375 ± 10	≈ 190	Structure in $N\bar{N}$ total σ	
$X^-(2500)$	1,2		2500 ± 32	≈ 87	$(MM)^-$ backward	
$X^-(2620)$	1,2		2620 ± 20	85 ± 30	$(MM)^-$	} 2650 region
4π (2676)	$(1, 2, 3)^+$		2676 ± 27	≈ 150	$\rho^- \pi^+ \pi^-$, ω and η antiselected	
$X^-(2800)$	1,2		2800 ± 20	46 ± 10	$(MM)^-$	
$X^-(2880)$	1,2		2880 ± 20	≤ 15	$(MM)^-$	
$X^-(3025)$	1,2		3025 ± 20	≈ 25	$(MM)^-$	} 3030 region
$N\bar{N}$ (3035)	+		3035 ± 25	200 ± 60	4π , 6π	
$X^-(3075)$	1,2		3075 ± 20	≈ 25	$(MM)^-$	
$X^-(3145)$	1,2		3145 ± 20	≤ 10	$(MM)^-$	
$X^-(3475)$	1,2		3475 ± 20	≈ 30	$(MM)^-$	
$X^-(3535)$	1,2		3535 ± 20	≈ 30	$(MM)^-$	

- (r) See Q-region note in listings. Some investigators see a broad enhancement in mass ($K\pi\pi$) from 1250-1400 MeV (the Q region), and others see structure. Only the $K_A(1240)$ or C seems well established, whereas possible structures from 1280 to 1400 MeV cannot be disentangled. For the whole Q region the decay rate into $K^*(892)\pi$ is large, and a $K\rho$ decay is seen. The $K\eta$, $K\omega$, and $K\pi$ are less than a few percent.
 (s) The mass comes only from charged $K_N(1420) \rightarrow K\pi$ measurements, since it is not well determined in the $K\pi\pi$ system with Q effect present; average of neutral K_N^* mass is 1423, but see typed note under $K^*(892)$ mass in listings.
 (t) Four bumps (A1, A3, Q, L) appear, in mass spectra, as broad enhancements (up to 200-400 MeV wide) centred about 200 MeV above threshold; they can be interpreted as either resonances or kinematic effects. Widths and branching ratios are published, and compiled in the listing, but in some cases we do not average them because they are very sensitive to difficult background subtractions. For more discussion, see notes in data card listings.

Mixing angles for SU(3) nonets: the small table which appeared here in earlier editions has been moved to Appendix II.

Baryons

April 1974

[See notes on N's and Δ's, on possible Z*'s, and on Y*'s at the beginning of those sections in the data listings; also see notes on individual resonances in the listings.]

Particle ^a	I (J ^P) — estab.	π or K Beam T(GeV) p(GeV/c) σ = 4π*2 (mb)	Mass M ^b (MeV)	Width Γ ^b (MeV)	M ² ± ΓM ^c (GeV ²)	Partial decay mode		
						Mode	Fraction %	p or p _{max} ^d (MeV/c)
p	<u>1/2(1/2⁺)</u>		938.3		0.880	See Stable Particles		
n			939.6		0.883			
N ⁱ (1470)	<u>1/2(1/2⁺)</u> P ₁₁	T=0.53πp p=0.66 σ=27.8	1435 to 1505	165 to 400	2.16 ±0.34	Nπ Nππ	60 40	420 368
N ⁱ (1520)	<u>1/2(3/2⁻)</u> D ₁₃	T=0.61 p=0.74 σ=23.5	1510 to 1540	105 to 150	2.31 ±0.18	Nπ Nππ [Δ(1236)π] ^e Nη	50 50 [dominant] ^e ~0.6	456 410 224
N ⁱ (1535)	<u>1/2(1/2⁻)</u> S ₁₁	T=0.64 p=0.76 σ=22.5	1500 to 1600	50 to 160	2.36 ±0.18	Nπ Nη ⁰ Nππ ⁰	35 55 ~10	467 182 422
N(1670) ⁱ	<u>1/2(5/2⁻)</u> D ₁₅	T=0.87 p=1.00 σ=15.6	1655 to 1680	105 to 175	2.79 ±0.24	Nπ Nππ [Δ(1236)π] ^e ΔK Nη	40 60 [44] ^e < .3 < 1 ^j	560 525 357 200 368
N(1688) ⁱ	<u>1/2(5/2⁺)</u> F ₁₅	T=0.90 p=1.03 σ=14.9	1680 to 1692	105 to 180	2.85 ±0.21	Nπ Nππ [Δ(1236)+π] ^e ΔK Nη	60 40 [26] ^e < .2 < .5 ^j	572 538 371 231 388
N ⁱⁱ (1700) ⁱ	<u>1/2(1/2⁻)</u> S ₁₁	T=0.92 p=1.05 σ=14.3	1665 to 1765	100 to 400	2.89 ±0.42	Nπ ΔK Nη	-65 5	580 250 340
N ⁱⁱ (1780) ⁱ	<u>1/2(1/2⁺)</u> P ₁₁	T=1.07 p=1.20 σ=12.2	1650 to 1860	50 to 450	3.17 ±0.51	Nπ ΔK Nη	30 ~7 ~10 ^j	633 353 476
N(1860)	<u>1/2(3/2⁺)</u> P ₁₃	T=1.22 p=1.36 σ=10.4	1770 to 1900	180 to 330	3.46 ±0.57	Nπ Nππ ΔK Nη	25 ~5 ^j ~4 ^j	685 657 437 545
N(2190)	<u>1/2(7/2⁻)</u> G ₁₇	T=1.94 p=2.07 σ=6.21	2000 to 2260	270 to 325	4.80 ±0.67	Nπ Nππ	25	888 868
N(2220)	<u>1/2(9/2⁺)</u> H ₁₉	T=2.00 p=2.14 σ=5.97	2200 to 2245	260 to 330	4.93 ±0.65	Nπ Nππ	15	905 887
N(2650)	<u>1/2(?⁻)</u>	T=3.12 p=3.26 σ=3.67	2650	360	7.02 ±0.95	Nπ Nππ	(J+1/2) _x =0.45 ^f	1154 1140
N(3030)	<u>1/2(?)</u>	T=4.27 p=4.41 σ=2.62	3030	400	9.18 ±1.21	Nπ Nππ	(J+1/2) _x =0.05 ^f	1366 1354
Δ(1236)	<u>3/2(3/2⁺)</u> P ₃₃	T=0.195(++) p=0.304 σ=91.8	1230 to 1236	110 to 122	1.53 ±0.14	Nπ Nπ ⁺ π ⁻ Nγ	99.4 0 ~0.6	231 89 262
Δ(1650)	<u>3/2(1/2⁻)</u> S ₃₁	T=0.83 p=0.96 σ=16.4	1615 to 1695	130 to 200	2.72 ±0.28	Nπ Nππ	28 72	547 511

Baryons

Particle ^a	I (J ^P) ↳ estab.	π or K Beam T(GeV) p(GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass M ^b (MeV)	Width Γ^b (MeV)	M ² $\pm \Gamma M^c$ (GeV ²)	Partial decay mode		p or P _{max} ^d (MeV/c)
						Mode	Fraction %	
Δ (1670)	$3/2(3/2^-)$ D ₃₃	T=0.87 p=1.00 $\sigma=15.6$	1650 to 1720	175 to 300	2.79 ± 0.40	N π N $\pi\pi$	15	560 525
Δ (1890)	$3/2(5/2^+)$ F ₃₅	T=1.28 p=1.42 $\sigma=9.88$	1840 to 1920	135 to 350	3.57 ± 0.49	N π N $\pi\pi$	17	704 677
Δ (1910)	$3/2(1/2^+)$ P ₃₁	T=1.33 p=1.46 $\sigma=9.54$	1780 to 1935	230 to 420	3.65 ± 0.62	N π N $\pi\pi$	25	716 691
Δ (1950)	$3/2(7/2^+)$ F ₃₇	T=1.41 p=1.54 $\sigma=8.90$	1930 to 1980	140 to 220	3.80 ± 0.39	N π Δ (1236) π Σ K Σ (1385)K	45 ≈ 50 ~ 2 1.4	741 571 460 232
Δ (2420)	$3/2(11/2^+)$	T=2.50 p=2.64 $\sigma=4.68$	2320 to 2450	270 to 350	5.86 ± 0.75	N π N $\pi\pi$	11 >20	1023 1006
Δ (2850)	$3/2(?^+)$	T=3.71 p=3.85 $\sigma=3.05$	2850	400	8.12 ± 1.14	N π N $\pi\pi$	(J+1/2) _x =0.25 f	1266 1254
Δ (3230)	$3/2(?)$	T=4.94 p=5.08 $\sigma=2.25$	3230	440	10.4 ± 1.4	N π N $\pi\pi$	(J+1/2) _x =0.05 f	1475 1464
Z*								
Z* Evidence for states with hypercharge 2 is controversial. See listings for discussion and display of data.								
Λ	$0(1/2^+)$		1115.6		1.24	See Stable Particles		
Λ (1405)	$0(1/2^-)$ S ₀₁	p < 0 K ⁻ p	1405 $\pm 5^g$	40 $\pm 10^g$	1.97 ± 0.06	$\Sigma\pi$	100	142
Λ' (1520)	$0(3/2^-)$ D ₀₃	p=0.389 $\sigma=84.5$	1518 $\pm 2^g$	16 $\pm 2^g$	2.30 ± 0.02	N \bar{K} $\Sigma\pi$ $\Lambda\pi\pi$ [Σ (1385) π] ^e	46 ± 1 41 ± 1 9.6 ± 7 [39 ± 10] ^e	237 260 252 ⁿ 75 ⁿ
Λ' (1670)	$0(1/2^-)$ S ₀₁	p=0.74 $\sigma=28.5$	1670	15 to 38	2.79 ± 0.04	N \bar{K} $\Lambda\eta$ $\Sigma\pi$	~ 20 ~ 35 ~ 45	410 66 393
Λ'' (1690)	$0(3/2^-)$ D ₀₃	p=0.78 $\sigma=26.1$	1690	27 to 85	2.86 0.09	N \bar{K} $\Sigma\pi$ $\Lambda\pi\pi$ $\Sigma\pi\pi$	$\sim 30^h$ ~ 40 ~ 20 ~ 10	429 409 415 352
Λ (1815)	$0(5/2^+)$ F ₀₅	p=1.05 $\sigma=16.7$	1820 $\pm 5^g$	64 to 100	3.30 ± 0.15	N \bar{K} $\Sigma\pi$ Σ (1385) π	62 11 17	537 504 358
Λ (1830)	$0(5/2^-)$ D ₀₅	p=1.09 $\sigma=15.8$	1835	74 to 150	3.37 ± 0.20	N \bar{K} $\Sigma\pi$	~ 10 ~ 30	550 515
Λ (2100)	$0(7/2^-)$ G ₀₇	p=1.68 $\sigma=8.68$	2100	60 to 140	4.41 ± 0.22	N \bar{K} $\Sigma\pi$ $\Lambda\eta$ Ξ K $\Lambda\omega$	25 ~ 5 < 3 < 10	748 699 617 483 443
Λ (2350)	$0(?)$	p=2.29 $\sigma=5.85$	2350	140 to 324	5.52 ± 0.55	N \bar{K}	(J+1/2) _x =0.7 f	913

Baryons

Particle ^a	I (J ^P) — estab.	π or K Beam T(GeV) p(GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass M ^b (MeV)	Width Γ^b (MeV)	M ² $\pm \Gamma M^c$ (GeV ²)	Partial decay mode			
						Mode	Fraction %	p or P _{max} ^d (MeV/c)	
Σ	$1(1/2^+)$		(+)1189.4 (0)1192.5 (-)1197.4		1.41 1.42 1.43	See Stable Particles			
$\Sigma(1385)$	$1(3/2^+) P_{13}$	p < 0 K ⁻ p	(+)1383±1 S=1.3* (-)1386±2 S=2.2*	(+)36±3 S=1.9* (-)36±6 S=3.5*!	1.92 ±0.05	$\Lambda\pi$ $\Sigma\pi$	90±3 10±3 S=1.4*	208 117	
$\Sigma(1670)^k$	$1(3/2^-) D'_{13}$	p=0.74 $\sigma=28.5$	1670	50	2.79 ±0.08	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi$ $\Sigma\pi\pi$ [$\Lambda(1405)\pi$] ^e $\Lambda\pi\pi$	~8	410 387 447 326 207 397	
$\Sigma(1750)$	$1(1/2^-) S'_{11}$	p=0.91 $\sigma=20.7$	1750	50 to 80	3.06 ±0.11	$N\bar{K}$ $\Lambda\pi$ $\Sigma\eta$	~15 seen seen	483 507 55	
$\Sigma(1765)$	$1(5/2^-) D_{15}$	p=0.94 $\sigma=19.6$	1765 ±5 ^g	~120	3.12 ±0.21	$N\bar{K}$ $\Lambda\pi$ $\Lambda(1520)\pi$ $\Sigma(1385)\pi$ $\Sigma\pi$	~44 ~15 ~14 ~13 ~1	496 518 187 315 461	
$\Sigma(1915)^i$	$1(5/2^+) F'_{15}$	p=1.25 $\sigma=13.0$	1910	70	3.65 ±0.13	$N\bar{K}$ $\Lambda\pi$ $\Sigma\pi$	~11	613 619 568	
Formation and production experiments do not agree on $\Sigma\pi/\Lambda\pi$ ratio.									
$\Sigma(2030)$	$1(7/2^+) F_{17}$	p=1.52 $\sigma=9.93$	2030	100 to 170	4.12 ±0.27	$N\bar{K}$ $\Lambda\pi$ $\Sigma\pi$ ΞK	10 to 27 14 to 38 2 to 5 < 2	700 700 652 412	
$\Sigma(2250)$	1(?) —	p=2.04 $\sigma=6.76$	2250	100 to 230	5.06 ±0.37	$N\bar{K}$ $\Lambda\pi$ $\Lambda\pi$	(J+1/2) _x =0.3 ^f	849 842 799	
$\Sigma(2455)$	1(?) —	p=2.57 $\sigma=5.09$	2455	~120	6.03 ±0.29	$N\bar{K}$	(J+1/2) _x =0.2 ^f	979	
$\Sigma(2620)$	1(?) —	p=2.95 $\sigma=4.30$	2620	~175	6.86 ±0.46	$N\bar{K}$	(J+1/2) _x =0.3 ^f	1064	
Ξ^{ℓ}	$1/2(1/2^+)$		(0)1314.8 (-)1321.3		1.73 1.75	See Stable Particles			
$\Xi(1530)^{\ell}$	$1/2(3/2^+)$	p-wave	(0)1528.9±1.1 (-)1533.8±1.9	7.3 ±1.7	2.34 ±0.01	$\Xi\pi$	100	144	
$\Xi(1820)^{\ell}$	$1/2(?)$		1795 to 1870	12 to 99	3.31 ±0.10	$\Lambda\bar{K}$ $\Xi\pi$ $\Xi(1530)\pi$ $\Sigma\bar{K}$		396 413 234 306	
$\Xi(1940)^{\ell}$	$1/2(?)$		1894 to 1961	42 to 140	3.72 ±0.18	$\Xi\pi$ $\Xi(1530)\pi$		499 336	
Seen in both final states; not clear if one, or more, states present.									
Ω^-	$0(3/2^+)$		1672.5		2.80	See Stable Particles			



Baryons

- * Quoted error includes an S(scale) factor. See footnote to Stable Particles Table.
- An arrow at the left of the Table indicates a candidate that has been omitted because the evidence for the existence of the effect and (or) for its interpretation as a resonance is open to considerable question. See listings for information on the following: $N(1700) D''_{13}$, $N(1990) F_{17}$, $N(2040) D''_{13}$, $N(3245)$, $N(3690)$, $N(3755)$, $\Delta(1690) P''_{33}$, $\Delta(1960) D_{35}$, $\Delta(2160) P''_{33}$, $Z_0(1780)$, $Z_0(1865)$, $Z_1(1900)$, $\Lambda(1330)$, $\Lambda(1750) P_{01}$, $\Lambda(1860) P_{03}$, $\Lambda(1870) S''_{01}$, $\Lambda(2010) D''_{03}$, $\Lambda(2020) F_{07}$, $\Lambda(2100) F''_{05}$, $\Lambda(2585)$, $\Sigma(1440)$, $\Sigma(1480)$, $\Sigma(1620) S'_{11}$, $\Sigma(1620) P'_{11}$, $\Sigma(1660) D'_{13}$, $\Sigma(1690)$, $\Sigma(1880) P''_{11}$, $\Sigma(1940) D''_{13}$, $\Sigma(2070) F''_{15}$, $\Sigma(2080) P_{13}$, $\Sigma(2100) G_{17}$, $\Sigma(3000)$, $\Xi(1630)$, $\Xi(2030)$, $\Xi(2250)$, $\Xi(2500)$.
- a. For the baryon states, the name [such as $N(1470)$] contains the mass, which may be different for each new analysis. The value chosen is the rounded average from Table II of the note on N's and Δ 's in the baryon listings. For Y^* 's and Ξ^* 's, the mass is an educated guess obtained by looking at the reported values. The convention for using primes in the names is as follows: when there is more than one resonance on a given Argand diagram, the first has been designated with a prime, the second with a double prime, etc. The name (col. 1) is the same as can be found in large print in the listings.
 - b. See note on N's and Δ 's in baryon listings. For M and Γ of most baryons we report here an interval instead of an average. Averages are appropriate if each result is based on independent measurements, but inappropriate here where the spread in parameters arises because different models or procedures have been applied to a common set of data. Where only one value is given it is either because only one experiment reports that state or because the various experiments agree. An error is quoted only when the various experiments averaged have taken into account the systematic errors.
 - c. For this column M is the rounded average which also appears in the name column. For the N's and Δ 's, Γ is the average quoted on Table II of the N's and Δ 's note in the baryon listings; for the Y^* 's and Ξ^* 's, Γ is taken as the center of the interval given in the column labeled " Γ ".
 - d. For decay modes into >3 particles p_{\max} is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated using the averaged central mass values, without taking into account the widths of the resonances.
 - e. Square brackets indicate a sub-reaction of the previous unbracketed decay mode.
 - f. This state has been seen only in total cross sections. J is not known; x is Γ_{e1}/Γ .
 - g. This is only an educated guess; the error given is larger than the error of the average of the published values (see listings for the latter).
 - h. Reported values of elasticity range from .18 to .34, each with a small error; $x=.30$ is only a guess. All the other branching fractions are dependent upon this choice. The reported values of $x \cdot x_e$ for the other channels favor the large value of x. An $x=.25$ or lower would violate unitarity.
 - i. Only information coming from partial-wave analyses has been used here. For the production experiments results see the listings.
 - j. Value obtained in an energy-dependent partial-wave analysis which uses a t-channel-poles-plus-resonance parametrization. The values of the couplings obtained for the resonances may be affected by double counting.
 - k. In this energy region the situation is still confused. Formation experiments suggest two states: $P_{11}(1620)$ decaying mainly into $\Sigma\pi$, and $D_{13}(1670)$ with branching fractions $\Sigma\pi(40\%)$, $\Lambda\pi(10\%)$, $\Sigma\pi\pi(<14\%)$. Production experiments report four states: $\Sigma(1620)$ seen only in the $\Lambda\pi$ mode, $\Sigma_1(1660)$ with appreciable $\Lambda\pi$ and $\Sigma\pi$ modes, $\Sigma_2(1660)$ with main decay mode $\Lambda(1405)+\pi$ (that is, $\Sigma\pi\pi$), and $\Sigma(1690)$ seen in the $\Lambda\pi$ mode. Of these four, Σ_1 and Σ_2 seem to be on firmer ground than the other two and both seem to have $J^P = 3/2^-$ like the $D_{13}(1670)$ seen in formation experiments. Two resonances of the same spin and parity have been hypothesized as the origin of much of the complexity observed in production experiments. With the addition of the $P_{11}(1620)$, there are three candidates that eventually might be required to clarify the situation.
 - l. Only $\Xi(1530)$ is firmly established; information on the other states comes from experiments that have poor statistics due to the fact that the cross sections for $S=-2$ states are very low. For Ξ states, because of the meager statistics, we lower our standards and tabulate resonant effects if they have at least a four-standard-deviation statistical significance and if they are seen by more than one group. So $\Xi(2030)$, with main decay mode $\Sigma\bar{K}$, reported as a 3.5-standard-deviation effect, is not tabulated. See the listings for the other states.
 - m. See note on $\Delta(1236)$ in the baryon listings. Values of mass and width are dependent upon resonance shape used to fit the data.
 - n. This p value is the average value obtained by folding over the two Breit-Wigner shapes.
 - o. The preliminary results of DIEM 70 quoted in the listings have been revised so that they are now in agreement with the values quoted in the present table (G. Smadja, private communication).

PHYSICAL AND NUMERICAL CONSTANTS*
PHYSICAL CONSTANTS

N	=	6.022169(40)×10 ²³ mole ⁻¹ (based on A _{C12} = 12)
c	=	2.9979250(10)×10 ¹⁰ cm sec ⁻¹
e	=	4.803250(21)×10 ⁻¹⁰ esu = 1.6021917(70)×10 ⁻¹⁹ coulomb
1 MeV	=	1.6021917(70)×10 ⁻⁶ erg
ħ	=	6.582183(22)×10 ⁻²² MeV sec
	=	1.0545949(80)×10 ⁻²⁷ erg sec
ħc	=	1.9732891(66)×10 ⁻¹¹ MeV cm = 197.32891(66) MeV fermi
	=	0.6240088(24) GeV mb ^{1/2}
α	=	e ² /ħc = 1/137.03602(21)
k _{Boltzmann}	=	1.380622(59)×10 ⁻¹⁶ erg K ⁻¹
	=	8.61708(37)×10 ⁻¹¹ MeV K ⁻¹ = 1 eV/11604.85(49)K
m _e	=	0.5110041(16) MeV = 9.109558(54)×10 ⁻³¹ kg
m _p	=	938.2592(52) MeV = 1836.109(41) m _e = 6.72211(63)m _{π±}
	=	1.00727661(8)m ₁ (where m ₁ = 1 amu = $\frac{1}{12}$ m _{C12} = 931.4812(52)MeV)
r _e	=	e ² /m _e c ² = 2.817939(13) fermi (1 fermi = 10 ⁻¹³ cm)
λ _e	=	ħ/m _e c = r _e α ⁻¹ = 3.861592(12)×10 ⁻¹¹ cm
a _{∞ Bohr}	=	ħ ² /m _e e ² = r _e α ⁻² = 0.52917715(81)A (1A = 10 ⁻⁸ cm)
σ _{Thomson}	=	$\frac{8}{3}\pi r_e^2$ = 0.6652453(61)×10 ⁻²⁴ cm ² = 0.6652453(64) barns
μ _{Bohr}	=	eħ/2m _e c = 0.5788381(18)×10 ⁻¹⁴ MeV gauss ⁻¹
μ _{nucleon}	=	eħ/2m _p c = 3.152526(21)×10 ⁻¹⁸ MeV gauss ⁻¹
$\frac{1}{2}\omega_e^{\text{cyclotron}}$	=	e/2m _e c = 8.794014(27)×10 ⁶ rad sec ⁻¹ gauss ⁻¹
$\frac{1}{2}\omega_p^{\text{cyclotron}}$	=	e/2m _p c = 4.789484(27)×10 ³ rad sec ⁻¹ gauss ⁻¹

Hydrogen-like atom (nonrelativistic, μ = reduced mass):

$$\frac{v}{c} \text{rms} = \frac{ze^2}{n\hbar c}; E_n = \frac{\mu}{2} v^2 = \frac{\mu z^2 e^4}{2(n\hbar)^2}; a_n = \frac{n^2 \hbar^2}{\mu z e^2}$$

$$R_\infty = m_e e^4 / 2\hbar^2 = m_e c^2 \alpha^2 / 2 = 13.605826(45) \text{ eV (Rydberg)}$$

pc = 0.3 Hp (MeV, kilogauss, cm); 0.3 (which is 10⁻¹¹c) enters because there are ≈ 300 "volts"/esu volt.

1 year (sidereal)	=	365.256 days = 3.1558×10 ⁷ sec (≈ π×10 ⁷ sec)
density of dry air	=	1.205 mg cm ⁻³ (at 20°C, 760 mm)
acceleration by gravity	=	980.62 cm sec ⁻² (sea level, 45°)
gravitational constant	=	6.6732(31)×10 ⁻⁸ cm ³ g ⁻¹ sec ⁻²
1 calorie (thermochemical)	=	4.184 joules
1 atmosphere	=	1033.2275 g cm ⁻²
1 eV per particle	=	11604.85(49)°K (from E = kT)

NUMERICAL CONSTANTS

π	=	3.1415927	1 rad	=	57.2957795 deg	√π	=	1.7724539
e	=	2.7182818	1/e	=	0.3678794	√2	=	1.4142136
ln 2	=	0.6931472	ln 10	=	2.3025851	√3	=	1.7320508
log ₁₀ 2	=	0.3010300	log ₁₀ e	=	0.4342945	√10	=	3.1622777

*Compiled by Stanley J. Brodsky, based mainly on the adjustment of the fundamental physical constants by B. N. Taylor, W. H. Parker, and D. N. Langenberg, *Rev. Mod. Phys.* **41**, 375 (1969). The figures in parentheses correspond to the 1 standard deviation uncertainty in the last digits of the main number.

CLEBSCH-GORDAN COEFFICIENTS AND SPHERICAL HARMONICS

Note: A $\sqrt{\quad}$ is to be understood over every coefficient; e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation: $\begin{matrix} J & J & \dots \\ M & M & \dots \end{matrix}$

$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$

$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$

$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$

$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$

$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$

$Y_l^{-m} = (-1)^m Y_l^m^*$

$\langle j_1 j_2 m_1 m_2 | j_1 j_2 J M \rangle$
 $= (-1)^{J-j_1-j_2} \langle j_2 j_1 m_2 m_1 | j_2 j_1 J M \rangle$

SU(3) CONVENTIONS

for Isoscalar Factor Table on next page

Since January 1970 we have used the convention that the first particle shall be a baryon, the second a meson (R. Levi Setti, Proceedings of Lund Conference, 1969, p. 339 and Table II). Note, for comparison, that the de Swart table of 8×8 is merely labeled with symbols like $(I_1 = 1/2, Y_1 = 1, I_2 = 1, Y_2 = 0)$, which can be read either as $(N\pi)$ or $(K\Sigma)$. Since there are no decuplet mesons, however, his 8×10 table is unambiguous; it must be read with the meson first.

The de Swart convention violates the other convention that the $N, N\pi$ coupling shall be $D + F$ (as opposed to $-D + F$). To get $D + F$ one must use the first line of the "N" table, which reads. . . $3\sqrt{5}/10 |8_D\rangle + 1/2 |8_F\rangle$ as opposed to. . . $-3\sqrt{5}/10 |8_D\rangle + 1/2 |8_F\rangle$. The first line must then be labeled $N\pi$ rather than $K\Sigma$, i.e., with the baryon first.

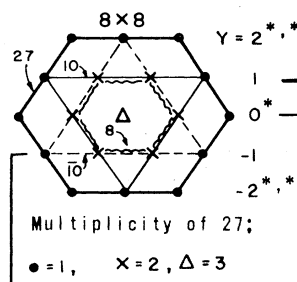
Levi Setti further advocates the convention of writing the baryon first for $SU(2)$ as well as $SU(3)$. For example, the sign of the amplitudes as plotted on his and our Argand plots comes from using our $SU(2)$ Clebsch-Gordan coefficients (Condon Shortley notation) and writing the baryon first. To make it easier to abide by this universal convention we have changed de Swart's 8×10 ($SU(3)$) table to 10×8 , with the help of his Eq. (14.3):

$$\langle \mu_1 \mu_2 | \mu \rangle = \xi_1 (-1)^{I_1 + I_2 - I} \langle \mu_2 \mu_1 | \mu \rangle .$$

SU(3) ISOSCALAR FACTORS

Adapted from J. J. de Swart, *Rev. Mod. Phys.* **35**, 916 (1963)
 (See note on previous page concerning conventions)

$$\{8\} \otimes \{8\} = \{27\} \oplus \{10\} \oplus \{10^*\} \oplus \{8\}_1 \oplus \{8\}_2 \oplus \{1\}.$$



* Five single-coefficient tables are omitted. The one involving a $\{10^*\}$ has a negative coefficient, i.e. $(NK|10^*) = -1$. The others, involving $\{27\}$ and $\{10\}$, are all +1.

		$Y=1 \ I=1/2 \ N$				$Y=1 \ I=3/2 \ \Delta$			
$\xi_1 \rightarrow$		27	8 _D	8 _F	10*	$\xi_1 \rightarrow$		27	10
Nπ		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	-1/2	Nπ		$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	-1/2	ΣK		$\sqrt{2}/2$	$\sqrt{2}/2$
Nη		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	1/2				
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	-1/2				

		$Y=0 \ I=0 \ \Lambda$				$Y=0 \ I=1 \ \Sigma$						
$\xi_1 \rightarrow$		27	8 _D	1	8 _F	$\xi_1 \rightarrow$		27	8 _D	8 _F	10	10*
NK		$\sqrt{15}/10$	$-\sqrt{10}/10$	1/2	$\sqrt{2}/2$	NK		$\sqrt{5}/5$	$-\sqrt{30}/10$	$-\sqrt{6}/6$	$-\sqrt{6}/6$	$-\sqrt{6}/6$
ΣK		$-\sqrt{15}/10$	$-\sqrt{10}/10$	-1/2	$\sqrt{2}/2$	ΣK		$\sqrt{5}/5$	$-\sqrt{30}/10$	$-\sqrt{6}/6$	$\sqrt{6}/6$	$-\sqrt{6}/6$
Σπ		$-\sqrt{10}/20$	$-\sqrt{15}/5$	$\sqrt{6}/4$	0	Σπ		0	0	$\sqrt{6}/3$	$\sqrt{6}/6$	$-\sqrt{6}/6$
Δη		$3\sqrt{30}/20$	$-\sqrt{5}/5$	$-\sqrt{2}/4$	0	Ση		$\sqrt{30}/10$	$\sqrt{5}/5$	0	1/2	1/2
						Δπ		$\sqrt{30}/10$	$\sqrt{5}/5$	0	-1/2	-1/2

		$Y=-1 \ I=1/2 \ \Xi$				$Y=-1 \ I=3/2$			
$\xi_1 \rightarrow$		27	8 _D	8 _F	10	$\xi_1 \rightarrow$		27	10*
Ξπ		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	1/2	Ξπ		$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	1/2	ΣK		$\sqrt{2}/2$	$\sqrt{2}/2$
Ξη		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	1/2				
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	-1/2				

The phase factor $\xi_1 = \pm 1$, from de Swart's Table I, enters in his symmetry formula (14.3):

$$\langle \mu_1 \mu_2 | \mu \rangle = \xi_1 (-1)^{I_1 + I_2 - I} (\mu_2 \mu_1 | \mu).$$

This factor is irrelevant if you are doing your own self-consistent calculations; it enters when you try to check someone else who chose $\mu_2 \otimes \mu_1$ instead of $\mu_1 \otimes \mu_2$.

$$\{10\} \otimes \{8\} = \{35\} \oplus \{27\} \oplus \{10\} \oplus \{8\}.$$

* Four single coefficient tables are omitted; only the $\{27\}$ is -1; the three with $\{35\}$ are +1.

Multiplicity of 35:
 • = 1, X = 2

		$Y=1 \ I=1/2 \ N$				$Y=1 \ I=3/2 \ \Delta$				
$\xi_1 \rightarrow$		27	8			$\xi_1 \rightarrow$		35	27	10
Δπ		$-\sqrt{5}/5$	$-2\sqrt{5}/5$			Δπ		1/4	$-\sqrt{5}/4$	$\sqrt{10}/4$
ΣK		$-2\sqrt{5}/5$	$\sqrt{5}/5$			Δη		$\sqrt{5}/4$	3/4	$\sqrt{2}/4$
						ΣK		$\sqrt{10}/4$	$-\sqrt{2}/4$	-1/2

		$Y=0 \ I=0 \ \Lambda$				$Y=0 \ I=1 \ \Sigma$				$Y=0 \ I=2$	
$\xi_1 \rightarrow$		27	8			$\xi_1 \rightarrow$		35	27	10	8
Σπ		$-\sqrt{10}/5$	$-\sqrt{15}/5$			Σπ		$\sqrt{3}/6$	$-3\sqrt{5}/10$	$\sqrt{3}/3$	$-\sqrt{30}/15$
ΞK		$-\sqrt{15}/5$	$\sqrt{10}/5$			Ση		$\sqrt{2}/2$	$-\sqrt{30}/10$	0	$-\sqrt{5}/5$
						ΞK		$\sqrt{3}/3$	$-\sqrt{5}/5$	$-\sqrt{3}/3$	$\sqrt{30}/15$
						ΔK		$\sqrt{3}/6$	$\sqrt{5}/10$	$\sqrt{3}/3$	$2\sqrt{30}/15$

		$Y=-1 \ I=1/2 \ \Xi$				$Y=-1 \ I=3/2$			
$\xi_1 \rightarrow$		35	27	10	8	$\xi_1 \rightarrow$		35	27
Ξπ		1/4	$-7\sqrt{5}/20$	$\sqrt{2}/4$	$-\sqrt{5}/5$	Ξπ		$\sqrt{2}/2$	$-\sqrt{2}/2$
Ξη		3/4	$3\sqrt{5}/20$	$-\sqrt{2}/4$	$-\sqrt{5}/5$	Ξη		$\sqrt{2}/2$	$-\sqrt{2}/2$
ΩK		$\sqrt{2}/4$	$-3\sqrt{10}/20$	-1/2	$\sqrt{10}/5$	ΣK		$\sqrt{2}/2$	$\sqrt{2}/2$
ΣK		1/2	$\sqrt{5}/10$	$\sqrt{2}/2$	$\sqrt{5}/5$				

		$Y=-2 \ I=0 \ \Omega^-$		$Y=-2 \ I=1$			
$\xi_1 \rightarrow$		35	10	$\xi_1 \rightarrow$		35	27
Ωη		$\sqrt{2}/2$	$-\sqrt{2}/2$	Ωπ		1/2	$-\sqrt{3}/2$
ΞK		$\sqrt{2}/2$	$\sqrt{2}/2$	ΞK		$\sqrt{3}/2$	1/2

SPECIAL RELATIVITY, PHASE SPACE, AND CROSS SECTIONS

Notation. 4-vector in c.m. $p = (w, \vec{p})$; in lab $P = (W, \vec{P})$, $T = W - m$.
 Solid-angle element $d\omega = 2\pi d\cos\theta$; $d\Omega = 2\pi d\cos\Theta$.
 $p^2 = w^2 - \vec{p}^2 = m^2$ is invariant. Cross section σ is invariant.

Lorentz Transformation

$$\begin{pmatrix} w \\ p_x \\ p_y \\ p_z \end{pmatrix} = \begin{pmatrix} \bar{w} & -\bar{\eta} & 0 & 0 \\ -\bar{\eta} & \bar{w} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} W \\ P_x \\ P_y \\ P_z \end{pmatrix} \quad \text{If } \theta \text{ and } \Theta \text{ are measured with respect} \\ \text{to the transformation axis } x, \quad (1)$$

$$\frac{p_x}{P_x} = \tan\theta = \frac{|\vec{P}| \sin\Theta}{-\bar{\eta}W + \bar{w}|\vec{P}| \cos\Theta}$$

If particle 1 is beam, 2 is target, then $(W_2, \vec{P}_2) = (m_2, \vec{0})$ and
 $\bar{w} = (W_1 + m_2)/\sqrt{s}$, $\bar{\eta} = \vec{v}\beta = |\vec{P}_1|/\sqrt{s}$, $|\vec{P}_1| = |\vec{P}_2| = \bar{\eta}m_2 = |\vec{P}_1| m_2/\sqrt{s}$. (2)
 For $m_1 = m_2$, $\bar{v}^2 = 1 + T_1/2m_1$.

General Lorentz Transformation [characterized by $\vec{\beta}$, with $\bar{w} = (1 - \beta^2)^{-1/2}$
 and $\bar{\eta} = \vec{v}\beta$: $w = \bar{w}W - \bar{\eta}\vec{P}$; $\vec{p} = \vec{P} - \bar{\eta}\frac{W + w}{\bar{w} + 1}$.

A Useful Transformation: Consider two 4-vectors $Q = (E, \vec{Q})$ and $q = (e, \vec{q})$.
 In the rest frame of Q [$Q' = (M, \vec{0})$], q becomes (q, \vec{q}')

$e' = Q \cdot q/M$ and $\vec{q}' = \vec{q} - f\vec{Q}$,
 where $Q^2 = M^2$ and $f = (e + e')/(E + M)$. These equations follow from example
 (b), p. 34 of Hagedorn.* They are particularly useful when Q is a sum of
 four-vectors that correspond to a resonant state.

Invariants. Notation: $1 + 2 \rightarrow 1' + 2'$.

$$s = (p_1 + p_2)^2 = m_1^2 + m_2^2 + 2(w_1 w_2 - \vec{p}_1 \cdot \vec{p}_2), \quad (3)$$

$$t = (p_1' - p_1)^2 = m_1^2 + m_1'^2 - 2(w_1 w_1' - \vec{p}_1 \cdot \vec{p}_1'), \quad (i = 1, 2), \quad (4)$$

$$u = (p_1' - p_2)^2 = (p_2' - p_1)^2 \quad [\text{use (6), below}]. \quad (5)$$

General relation: $s + t + u = m_1^2 + m_1'^2 + m_2^2 + m_2'^2$. (6)

In lab system $P_2 = (m_2, \vec{0})$, and writing $W = m + T$,

$$s = m_1^2 + m_2^2 + 2W_1 m_2 = (m_1 + m_2)^2 + 2T_1 m_2, \quad (3, \text{lab})$$

$$t = m_2^2 + m_1'^2 - 2W_1' m_2 = (m_2 - m_1')^2 - 2T_1' m_2. \quad (4, \text{lab})$$

In c.m. system $dt = +2|\vec{p}_1| |\vec{p}_1'| d\cos\theta$. (4, cm)

For elastic scattering ($m_1 = m_1'$, $m_2 = m_2'$), (4) and (5) in c.m. become

$$t = -2p^2(1 - \cos\theta) = -4p^2 \sin^2\theta/2, \quad (4, \text{el})$$

$u = (m_1^2 - m_2^2)^2/s - 2p^2(1 + \cos\theta) = (m_1^2 - m_2^2)^2/s - 4p^2 \cos^2\theta/2$. (5, el)

For elastic scattering, using (4, lab), (4, el), and (2),

$$T_2 = \frac{2\vec{P}_1^2 m_2}{s} \sin^2(\theta/2) \quad (\text{useful for calculating } \delta\text{-ray energies}). \quad (7)$$

Two-Body States. Energies and momenta in c.m.

$$w_1 = \frac{s + m_1^2 - m_2^2}{2\sqrt{s}}, \quad p_1^2 = p_2^2 = \frac{1}{4s} [s - (m_1 + m_2)^2] [s - (m_1 - m_2)^2]. \quad (8)$$

3- and 4-Body States. Let $m_{ij}^2 = (p_i + p_j)^2$, etc.; then

$$\sum_{i < j} m_{ij}^2 = \sum m_1^2 + m_{123}^2 = \text{const.} \quad (i, j = 1, 2, 3) \quad [\text{follows from (6)}] \quad (9)$$

$$= 2\sum m_1^2 + m_{1234}^2 = \text{const.} \quad (10)$$

$$\sum_{i < j < k} m_{ijk}^2 = \sum m_1^2 + 2m_{1234}^2 = \text{const.} \quad (i, j, k = 1, 2, 3, 4.)$$

 R_n , Invariant Volume in n-Body Momentum Space

A useful invariant is $\int d^4p \delta(p^2 - m^2) = \int \frac{d^3\vec{p}}{2w} = \int \frac{p^2 d|\vec{p}| d\omega}{2w} = \frac{1}{2} \int |\vec{p}| dwd\omega$.
 $R_2 = \pi |\vec{p}_1|/\sqrt{s}$, $R_3 = \pi^2 \int dw_1 dw_2 = (\pi^2/4s) \int dm_{12}^2 dm_{23}^2$.

Recurrence Relation for Factoring R_n (see e.g., Hagedorn, p. 93*)

Write $N \rightarrow 1, 2, \dots, k, k+1, \dots, n$ (R_n),

as $N \rightarrow K, k+1, \dots, n$ (R_{n-k+1}),
 $\left\{ \begin{array}{l} \leftarrow 1, 2, \dots, k \\ (R_k) \end{array} \right\}$ then $R_n = \int d(m_k^2) R_k R_{n-k+1}$,

or as $N \rightarrow K, L$
 $\left\{ \begin{array}{l} \leftarrow k+1, \dots, n \\ (R_k) \end{array} \right\}$ then $R_n = \int d(m_k^2) d(m_L^2) R_k R_{n-k+1} \frac{\pi^p(KL)}{\sqrt{s}}$

Cross Sections and Decay Rates†

For a system of n particles with overall four-momentum p and final momenta q_1, \dots, q_n [$q_i = (e_i, \vec{q}_i)$], define **Lorentz Invariant Phase Space**

$$d\text{LIPS}(s; q_1, \dots, q_n) = (2\pi)^4 \delta^4(p - \sum q_i) \frac{1}{(2\pi)^{3n}} \prod_{i=1}^n \frac{d^3\vec{q}_i}{2e_i}. \quad (11)$$

Note that $R_n = (2\pi)^{3n-4} \int d\text{LIPS}$.

For $1 + 2 \rightarrow n$ particles or $1 \rightarrow n$ particles, in general $|i\rangle \rightarrow |f\rangle$,

$$\sigma_{if} = \frac{1}{4F} \int |T_{if}|^2 d\text{LIPS}(s; q_1, \dots, q_n), \quad (12)$$

or

$$\Gamma_{if} = \frac{1}{2m_1} \int |T_{if}|^2 d\text{LIPS}(m_1^2; q_1, \dots, q_n), \quad (13)$$

where T_{if} is an invariant matrix element. F is Møller's invariant flux factor, $F^2 = (p_1 \cdot p_2)^2 - p_1^2 p_2^2$. In every system where \vec{p}_1 and \vec{p}_2 are collinear, $F = w_1 w_2 |\vec{v}_1 - \vec{v}_2|$ ($\vec{v} = \vec{p}/w$). If 1 is beam, 2, target ($p_2 = 0$), then $F = |\vec{p}_1| m_2 = |\vec{p}_1| \sqrt{s}$.

For elastic scattering in c.m., $\frac{d\text{LIPS}}{d\Omega} = \frac{1}{(4\pi)^2} \frac{|\vec{p}_1|}{\sqrt{s}}$, and (12) yields

$$\frac{d\sigma}{d\Omega} = \frac{|T|^2}{(8\pi)^2 s} \quad \text{or} \quad \frac{d\sigma}{dt} = \frac{|T|^2}{64\pi |\vec{p}_1|^2 s}. \quad (14)$$

The normalization is such that the optical theorem reads

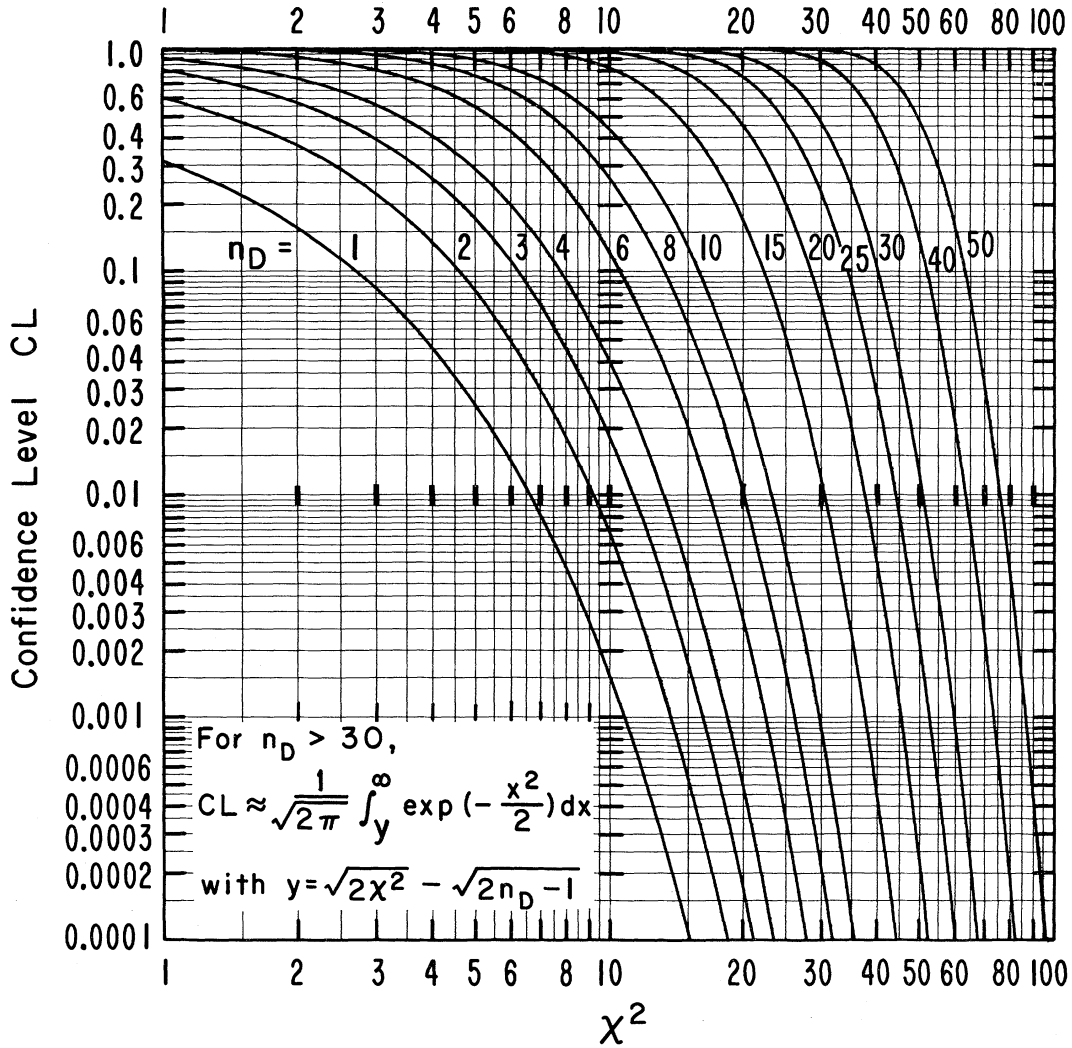
$$\text{Im } T|_{t=0} = 2|\vec{p}_1| \sqrt{s} \sigma_{\text{tot}}. \quad (15)$$

The choice of Eq. (14) implies a particular normalization of any spinors that may occur in T .† The advantage of this normalization is that it greatly simplifies the structure of T by putting factors such as $\frac{1}{(2\pi)^3} \frac{1}{2E}$ into the phase space where they really belong. In addition, the labels, i, f , refer to specific spin (helicity) states, so that the usual "average and sum" rule is implicit.

* R. Hagedorn, *Relativistic Kinematics*, W. A. Benjamin, New York, 1964.

† See, for example, Chaps. 1 and 2 of H. Pilkuhn, *The Interactions of Hadrons*, John Wiley & Sons, New York, 1967.

CONFIDENCE LEVEL VS. χ^2 FOR n_D DEGREES OF FREEDOM



For any n_D , $\langle \chi^2 \rangle = n_D$, $\delta(\chi^2) = \sqrt{2n_D}$. For large n_D , χ^2 becomes normally distributed about n_D . Thus in the notation of the box in the figure,

$$y_1 = (\chi^2 - n_D) / \sqrt{2n_D} \text{ has unit s. d.}$$

A better approximation, due to Fisher,* is that χ , not χ^2 , is normally distributed, specifically

$$y_2 = \sqrt{2\chi^2} - \sqrt{2n_D - 1} \text{ has unit s. d.}$$

One sees then that y_1 underestimates small C.L.'s. Thus for $n = 50$ and $\chi^2 = 80$, $y_1 = 3.0$ and C.L. = 0.13% vs. $y_2 = 2.7$, C.L. = 0.35%.

*R. A. Fisher, Statistical Methods for Research Workers, Oliver and Boyd, Edinburgh.

GAUSSIANLIKE DISTRIBUTIONS

The distribution

$$P_{2n+1}(x) = \frac{1}{2^n n! \sigma^{2n+2}} x^{2n+1} \exp\left[-\frac{x^2}{2\sigma^2}\right]$$

is normalized so that $\int_0^\infty P_{2n+1}(x) dx = 1$; the normalization is valid for $n > -1$ and not necessarily integral ($(\frac{1}{2})! = \sqrt{\pi}/2$). For $n = -1/2$ it reduces to the Gaussian distribution. Through a change of variables it yields the χ^2 distribution for n_D degrees of freedom:

$$P_{n_D}(\chi^2) = \frac{1}{2^{n_D/2} \left(\frac{n_D}{2} - 1!\right)} (\chi^2)^{n_D/2-1} \exp\left[-\chi^2/2\right].$$

Relation between standard deviation σ and mean deviation α :

$$2\sigma^2 = \pi\alpha^2; \sigma = 1.4826 \text{ probable error.}$$

Odds against exceeding one standard deviation = 2.15:1; two, 21:1; three, 370:1; four, 16,000:1; five, 1,700,000:1.

ATOMIC AND NUCLEAR PROPERTIES OF MATERIALS

Material	Z	A	Cross Section	Collision Length		Minimum		Radiation Length		Density
			σ^a barns	L_{coll}^b g cm ⁻²	L_{coll}^b cm	$-dE/dx^c$ MeV g ⁻¹ cm ²	$-dE/dx^c$ MeV cm ⁻¹	L_{rad}^d g cm ⁻²	L_{rad}^d cm	
H ₂	1	1.01	0.063	26.5	374 ^e	4.13	0.292 ^e	58.0	819 ^e	0.0708 ^e
D ₂	1	2.01	0.100	33.4	202 ^e	2.07	0.342 ^e	116	703 ^e	0.165 ^e
He	2	4.00	0.16	42.0	336 ^e	1.94	0.242 ^e	85.4	683 ^e	0.125 ^e
Li	3	6.94	0.23	50.4	94.3	1.69	0.902	78.7	148	0.534
Be	4	9.01	0.28	55.0	29.9	1.60	2.96	63.7	34.7	1.848
C	6	12.01	0.33	60.4	f	1.78	f	42.4	f	≈1.55 ^f
N ₂	7	14.01	0.36	63.6	78.9 ^e	1.81	1.46 ^e	37.8	46.7 ^e	0.808 ^e
Ne	10	20.18	0.465	72.1	60.1 ^e	1.73	2.08 ^e	29.1 ⁱ	24.2 ^{e, i}	1.200 ^{e, k}
Al	13	26.98	0.57	79.2	29.3	1.62	4.37	24.0	8.9	2.70
Fe	26	55.85	0.92	101.2	12.8	1.48	11.6	13.9	1.8	7.87
Cu	29	63.54	1.00	105.4	11.8	1.44	12.9	12.0	1.34	8.96
Sn	50	118.69	1.55	129.7	17.8	1.28	9.4	8.89	1.22	7.31
W	74	183.85	2.02	150.8	7.80	1.17	22.6	6.89	0.36	19.3
Pb	82	207.19	2.20	156.2	13.8	1.13	12.8	6.52	0.58	11.35
U	92	238.03	2.42	163.6	≈8.63	1.09	≈20.6	6.13	≈0.32	≈18.95
Air				64.6	53620 ^g	1.81	0.0022 ^g	36.5	30290 ^g	0.001205 ^g
Freon (CF ₃ Br)				87.1	≈58.0	1.52	≈2.3	16.6	≈11	≈1.5
H ₂ (bubble chamber, 27°K)				26.5	442 ^h	4.13	0.248 ^h	58.0	97 ^h	≈0.060 ^h
H-Ne mixture (bubble chamber) ^j				67.3	96.1	1.83	1.28	29.8 ⁱ	42.5 ⁱ	.70
H ₂ O				57.2	57.2	2.03	2.03	35.7	35.7	1.00
Ilford Emulsion				103.0	27.0	20.9	5.49	11.2	2.91	3.815
LiF				63.8	24.2	1.69	4.46	39.0	14.8	2.64
Mylar (C ₅ H ₄ O ₂)				59.1	42.8	1.91	2.64	39.6	28.7	1.38
NaI				119.0	32.4	1.32	4.84	9.58	2.61	3.67
Polyethylene (CH ₂)				51.0	≈55.5	2.09	≈1.92	44.1	≈48	≈0.92
Polystyrene (CH) ^k				54.9	≈52.3	2.03	≈2.14	43.4	≈41.3	≈1.05
Propane (C ₃ H ₈ , bubble chamber)				48.9	119.3	2.28	0.935	44.6	109	0.41

a. $\sigma = \sigma_{\text{natural}} = \pi (\hbar/m_p c)^2 \times A^{2/3} = 62.8 \text{ mb} \times A^{2/3}$
 b. $L_{\text{coll}} = A/(N\sigma_{\text{natural}}) = 26.5 \text{ g cm}^{-2} \times A^{1/3}$
 c. From W. H. Barkas and M. J. Berger, *Tables of Energy Losses and Ranges of Heavy Charged Particles*, NASA SP-3013 (1964).
 d. Mainly from *High Energy and Nuclear Physics Data Handbook*, W. Galbraith and W. S. C. Williams, Ed. (N. I. R. N. S., Rutherford Lab., Chilton, Didcot, Berks.) 1964.
 e. For liquid phase at 1 atm. and boiling temperature.
 f. Density variable.
 g. At 20°C.
 h. May vary by about ±3%, depending on operating conditions.
 i. From F. R. Huson, *Ionization Loss, Range, Straggling and Multiple Scattering*, BNL 11386 (1967).
 j. 53.7 atomic percent Ne.
 k. Density of gas at STP = $0.900 \times 10^{-3} \text{ g cm}^{-3}$, i. e., 0.75×10^{-3} times the density (1.200) of the boiling liquid.
 l. Typical scintillator; e. g., PILOT B has an atomic ratio H/C = 1.1.

MULTIPLE COULOMB SCATTERING*

The rms projected angle θ due to multiple Coulomb scattering (only) of a particle of charge z (in units of electron charge), momentum p (in MeV/c), and velocity v (in units of c) is

$$\theta_{\text{proj}} = z \frac{15}{\beta v} \sqrt{\frac{L}{L_{\text{rad}}}} (1 + \epsilon) \text{ radians;}$$

where L = length in scatterer.

For $L \geq 1/10 L_{\text{rad}}$, ϵ is generally $< 1/10$. The distribution of θ is not truly Gaussian.†

The rms projected displacement y on traversing an absorber of thickness L is

$$y_{\text{rms}} = L \theta_{\text{proj}} \sqrt{3}.$$

*Mainly from G. Z. Molière, *Naturforsch.* 3 (a), 78 (1948).

†See, for example, the experimental work of A. D. Hansen, L. H. Lanzl, E. M. Lyman, and M. B. Scott, *Phys. Rev.* 84, 634 (1951).

RADIOACTIVITY AND RADIATION PROTECTION

Unit of activity = Curie:

$$1 \text{ Ci} = 3.7 \times 10^{10} \text{ disintegrations/sec}$$

Unit of exposure dose for x and γ radiation = Roentgen:

$$1 \text{ R} = 1 \text{ esu/cm}^2 = 87.8 \text{ erg/g} (5.49 \times 10^7 \text{ MeV/g}) \text{ of air}$$

Unit of absorbed dose = rad:

$$1 \text{ rad} = 100 \text{ erg/g} (6.25 \times 10^7 \text{ MeV/g}) \text{ in any material}$$

Unit of dose equivalent (for protection) = rem:

$$1 \text{ rem (Roentgen equivalent for man)} = 1 \text{ rad} \times \text{QF},$$

where QF (quality factor) depends upon the type of radiation and other factors. For γ rays and HE protons, $\text{QF} \approx 1$; for thermal neutrons, $\text{QF} \approx 3$; for fast neutrons, QF ranges up to 10; and for α particles and heavy ions, QF ranges up to 20. Maximum permissible occupational dose for the whole body:

$$5 \text{ rem/year (or } \approx 100 \text{ millirem/week)}$$

Fluxes (per cm²) to liberate 1R in carbon:

$$3 \times 10^7 \text{ minimum ionizing singly charged particles}$$

$$0.9 \times 10^9 \text{ protons of 1 MeV energy}$$

(These fluxes are correct to within a factor of 2 for all materials.)

Natural background: 120 to 130 millirem/year

cosmic radiation (charged particles + neutrons) ~ 25

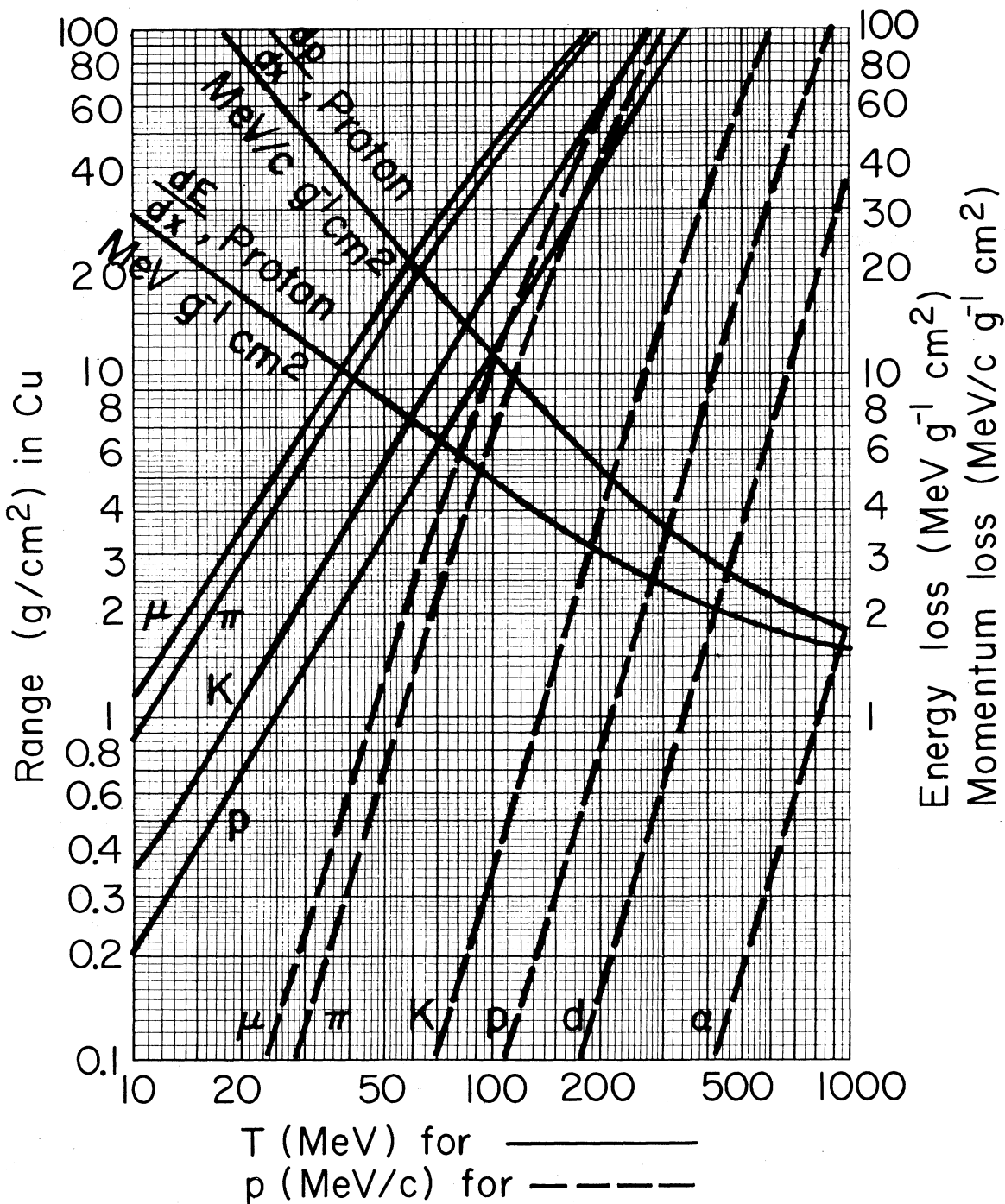
cosmic radiation (γ rays) ~ 25

radiation from rocks and air (γ rays) ~ 73

Cosmic ray background in counters: $\sim 1/\text{sec/cm}^2/\text{ster}$

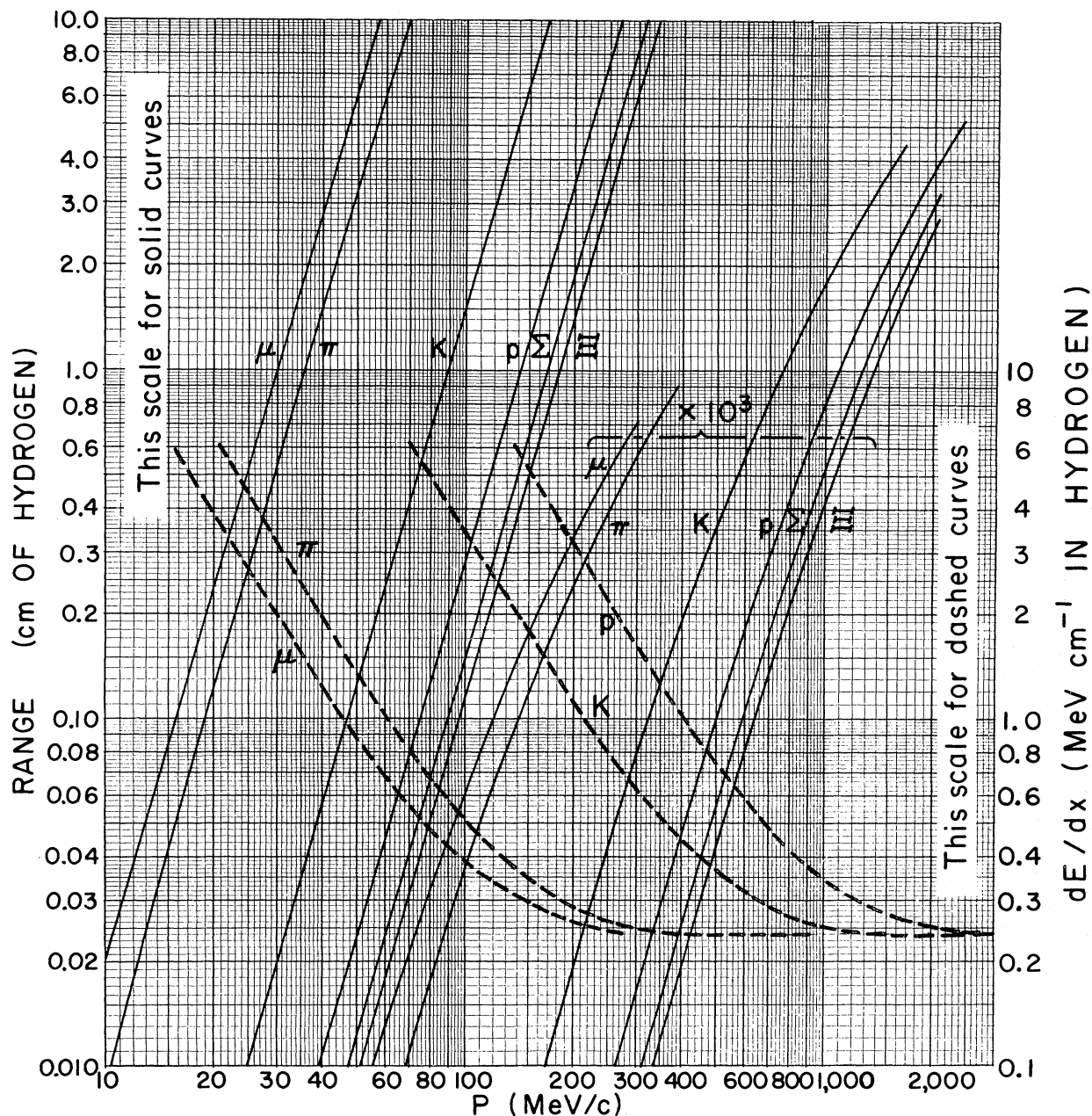
millirem/yr

RANGE AND ENERGY LOSS IN COPPER



Range and energy/momentum loss in copper, based on a calculation assuming a nominal mean excitation potential of 310 eV. (Calculation by W. A. Aron, UCRL-1325, 1951). The abscissa is to be read as kinetic energy T for the solid curves and momentum p for the dashed curves.

RANGE AND ENERGY LOSS IN LIQUID HYDROGEN



Range and energy loss in liquid hydrogen bubble chamber, determined by a μ^+ range of 1.103 ± 0.003 cm from the $\pi^+ \rightarrow \mu^+ \nu$ decay. Liquid hydrogen conditions: $T = 27.6 \pm 0.1^\circ K$; $P = 48 \pm 5$ psia; $\rho = (5.86 \pm 0.06) 10^{-2} g/cm^3$. (Data by Clark and Diehl, UCRL-3789, 1957.) Bubble chamber physicists: note that the number of bubbles per cm is proportional to $1/\beta^2$, not to dE/dx .

DATA CARD LISTINGS

Illustrated Key

Name of particle as it appears in table. **XX(1200)**

Arrow indicates this particle omitted from table. **74 XX MESON (1200, J^{PC}= -) I=1**

Quantity tabulated below. **74 XX(1200) MASS (MEV)**

M	L	1216.	11.	MERRILL	66 HRC	0 3.2 K-P	7/66
M	L	(1192.)	(16.)	LYNCH	67 HRC	+ 2.7 PI-P	6/67
M	L	1198.	10.	PIERCE	68 ASPK	+ 2.1 K-P	9/68
M	S	(1208.)	7.	FENNER	69 HRC	0 4.2 PI+P	9/69
M	S	80 1210.	8.	SMITH	70 MMS	- 3.5 PI-P	2/71*
M	S	SUPERSEDES EARLIER RESULT					
M	AVG	1206.9	5.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			

Code for quantity tabulated (M=mass, W=width, etc.)

Symbols used to key together data card and related comments.

Number of events above background.

Measured values (parentheses indicate value not used in average).

± Error in measured value (- field blank if error symmetric; parentheses on error only indicate data not used in average due to problems with error estimation).

Average value (and error) of quantity measured.

Vertical bar indicates average; width of horizontal bar on top is error (scaled) in average.

Value and error for each experiment.

Particle name, and quantum numbers (if known).

Particle code (for internal use only).

General comments on particle.

Abbreviated reference for this result; full reference given below.

Measurement technique (see abbreviations on next page.)

Charge(s) of particle detected.

Reaction producing particle, or comments.

Date this result punched (asterisk indicates result added or changed since previous edition).

Scale factor > 1 indicates inconsistent data.

Ideogram to display inconsistent data; curve is sum of Gaussians, one for each experiment (area of Gaussian = 1/error; width of Gaussian = ± error).

Contribution of experiment to χ^2 (if no entry present, experiment not used in calculating χ^2 or scale factor because of large error).

74 XX(1200) WIDTH (MEV)

W	35.	5.	MERRILL	66 HRC	0/3.2 K-P	7/66	
W	50.	10.	PIERCE	68 ASPK	+ 2.1 K-P	9/68	
W	70.	40.	FENNER	69 HRC	0 4.2 PI+P	9/69	
W	(60.)	OR LESS	SMITH	70 MMS	- 3.5 PI-P	2/71*	
W	AVG	38.4	6.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)			

WEIGHTED AVERAGE = 38.4 ± 6.0
ERROR SCALED BY 1.3

74 XX(1200) PARTIAL DECAY MODES

P1	XX(1200) INTO 3PI	DECAY MASSES
P2	XX(1200) INTO K KBAR	139+ 139+ 139
		493+ 493

74 XX(1200) BRANCHING RATIOS

R1	XX(1200) INTO 3PI/TOTAL	(P1)/TOTAL
R1	.66 .02	MERRILL 66 HRC 0 3.2 K-P 7/66
R1	(.68) (.03)	LYNCH 67 HRC + 2.7 PI-P 6/67
R1	LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION	
R1	FIT 0.675 0.012	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)
R2	XX(1200) INTO K KBAR/TOTAL	(P2)/TOTAL
R2	.35 .05	PIERCE 68 ASPK + 2.1 K-P 9/68
R2	FIT 0.325 0.012	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)
R3	XX(1200) INTO K KBAR/3PI	(P2)/(P1)
R3	.50 .03	FENNER 69 HRC 0 4.2 PI+P 9/69
R3	.41 .04	SMITH 70 MMS - 3.5 PI-P 2/71*
R3	AVG 0.468 0.026	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)
R3	FIT 0.480 0.026	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

Partial decay mode (labeled by P_i).

Branching ratio (labeled by R_j).

Value (and error) of quantity measured, as determined from constrained fit (using all measured branching ratios for this particle).

Branching ratio R_j in terms of partial decay modes P_i above (X_i sometimes used for P_i/total).

REFERENCES FOR XX(1200)

MERRILL 66/PRL 16 143	A. MERRILL	(SACLAY+CERN)
LYNCH 67 PR 155 610	B. LYNCH	(BNL)
PIERCE 68 PL 278 230	N. PIERCE	(LRL)
FENNER 69 NC 619 372	D. FENNER, R. BEANE	(NYSE+AMEX)
SMITH 70/PRL 24 147	J. SMITH	(SLAC)

References listed by year, then author.

Abbreviated reference form used on data cards above.

Journal, report, preprint, etc. (see abbreviations on next page).

Author(s)

Institution(s) of author(s) (see abbreviations on next page).

Stable Particles

(Immune to Strong Decay)

For notation, see illustrated key at beginning of data card listings.

CODE EVENTS QUANTITY ERROR* FRDR-- REFERENCE YR TECN SIGV COMMENTS DATE ABOVE BACKGROUND PUNCHED

γ 0 GAMMA (0,J=1)
0 GAMMA MASS (IN UNITS OF 10**21 MEV)

M	(6.)	OR LESS	PATEL	65	SATELLITE DATA	10/69
M	(6.)	OR LESS	GINTSBURG	64	SATELLITE DATA	10/69
M	(2,3)	OR LESS	GOLDHABER	68	SATELLITE DATA	10/69

REFERENCES

0 GAMMA
GINTSBURG 64 SOV. ASTR. AJ7 536 M. A. GINTSBURG (ACAD SCI, USSR)
PATEL 65 PL 14 105 V. L. PATEL (DURHAM)
GOLDHABER 68 PRL 21 567 A. GOLDHABER, M. NIETO (STONY BROOK)

ν_e 1 F-NEUTRINO (0,J=1/2)
1 F-NEUTRINO MASS (KEV)

M	(0,25)	OR LESS	LANGFUR	52	CNTR	
M	(0,15)	OR LESS	HAMILTON	53	CNTR	
M	(0,55)	OR LESS	MOR-0,2R	58	CNTR	
M	0,06	OR LESS	CL=90	BERGKVIST	69	CNTR

EL. STATIC. MAG. SP 11/69

REFERENCES

1 E-NEUTRINO (0,J=1/2)
LANGER 52 PR 88 689 L M LANGER, R J D MOFFAT (INDIANA)
HAMILTON 53 PR 92 1521 D HAMILTON, W P ALFORD, L GROSS (PRINCETON)
FRIEDMAN 58 PR 109 2214 LEWIS FRIEDMAN, LINCOLN G SMITH (BNL)
BERGKVIST 69 CERN 69-7 91 KARL-ERIK BERGKVIST (UNIV STOCKHOLM)

ν_μ 2 MU-NEUTRINO (0,J=1/2)
2 MU-NEUTRINO MASS (MEV)

M	(3,5)	OR LESS	BARKAS	56	EMUL	
M	(4,0)	OR LESS	DUIZIAK	59	CNTR	
M	(3,6)	OR LESS	FEINBERG	63	RVUE	7/66
M	(3,0)	OR LESS	ALLCOCK	65	RVUE	7/66
M	(2,5)	OR LESS	BARNDON	65	ASPK	
M	(2,1)	OR LESS	SHAFFER	65	CNTR	7/66
M	1,6	OR LESS	CL=90	BOOTH	67	CNTR
M	2,2	OR LESS	CL=90	HYMAN	67	HERC
M	(0,46)	(0,64)	(0,46)	FRANK	68	CNTR

CONF LEV = 68PCT 3/68
PRELIMINARY 9/68

REFERENCES

2 MU-NEUTRINO (0,J=1/2)
BARKAS 56 PR 101 778 W H BARKAS, W BIRNBAUM, F M SMITH (LRL)
DUIZIAK 59 PR 114 336 W F DUIZIAK, R SAGANE, J VEDDER (LRL)
FEINBERG 63 ARNS 13 631 G FEINBERG, L M LEDERMAN (COLUMBIA)
ALLCOCK 65 PPSL 85 875 G R ALLCOCK (LIVERPOOL)
BARNDON 65 PRL 14 449 BARNDON, NORTON, PEOPLES + (COLUM+STONY BROOK)
SHAFFER 65 PRL 14 923 R E SHAFFER, CROWE, JENKINS (LRL)
BOOTH 67 PL 26A 39 BOOTH, JOHNSON, WILLIAMS, WORMALD (LIVERPOOL)
HYMAN 67 PL 25 B 376 *LOKEN, PEWITT, MCKENZIE, KEYES + (ARG+CORN+NWU)
FRANK 68 VIENNA AHS. 36* FRANK, GAMET, LAKIN (SHAM+LIVP+STAN)

e 3 ELECTRON (0,5,J=1/2)
3 ELECTRON MASS (MEV)

M	(.511006) ± 0.000021	COHEN	65	RVUE	
M	.5110041 ± 0.000016	TAYLOR	69	RVUE	USING NEW E/H

7/70

3 ELECTRON LIFETIME (UNITS 10**21 YR)

T	OVER	2.0	MOE	65	CNTR	6/66
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3 ELECTRON MAGNETIC MOMENT (E/2ME)

MM	(1,0011609)	++(24)*10**7	SCHUPP	61	CNTR	-
MM	(1,001159622)	++(27)*10**9	WILKINSON	63	CNTR	-
MM	(1,00116168)	++(22)*10**6	RICH	66	CNTR	+ POSITRON
MM R	(1,001159557)	++(30)*10**9	RICH	68	CNTR	-
MM R	RICH 68 IS REEVALUATION OF WILKINSON 63					
MM	(1,0011596389)	++(31)*10**9	TAYLOR	70	RVUE	2/71*
MM	1,001159644	++(7)*10**9	WESLEY	70	CNTR	- NEW MEAS.

6/66
9/66
6/68
2/71*

REFERENCES

3 ELECTRON (0,5,J=1/2)
SCHUPP 61 PR 121 1 A A SCHUPP, R W PIDD, H R CRANE (MICHIGAN)
WILKINSON 63 PR 130 852 D T WILKINSON, H R CRANE (MICHIGAN)
COHEN 65 RMP 37 537 E R COHEN, J W M DUMOND (NAASC+CALTECH)
MOE 65 PR 140 B 992 M K MOE, F REINES (CASE INST TECHNOLOGY)
RICH 66 PRL 17 271 A RICH, H R CRANE (MICHIGAN)
RICH 68 PRL 20 967 A RICH (MICHIGAN)
TAYLOR 69 RMP 41 375 *PARKER, LANGENBERG (PRIN+UC+PENN)
WESLEY 70 PRL 24 1320 J. C. WESLEY, A. RICH (ANNA)

μ 4 MUON (106,J=1/2)
4 MUON MASS (MEV)

M	(105,659)	(0,002)	FEINBERG	63	RVUE	
M	105,6599	.0014	TAYLOR	69	RVUE	USING NEW E/H

7/70

4 MUON LIFETIME (UNITS 10**6)

T	2,198	0,001	0,001	FARLEY	62	CNTR
T	2,203	0,004		LUNDY	62	CNTR
T	2,202	0,003	0,003	ECKHAUSE	63	CNTR
T	2,197	0,002	0,002	MEYER	63	CNTR
T	2,198	0,002	0,002	MEYER	63	CNTR
T	AVG	2,1983	0,0008	0,0008	AVERAGE	(ERROR INCL. SCALE FACTOR OF 1.0)

11/67
7/66

4 RATIO OF LIFETIME OF MU+ TO MU-

DT	1,000	0,001	MEYER	63	CNTR	LIFETIME MU+/MU-
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7/66

4 MUON ANOMALOUS MAGN. MOMENT (10**6E/(2*MUON MASS))

MM	1162,0	5,0	CHARPAK	62	CNTR	+
MM R	(1165,75)	(0,71)	BAILEY	68	CNTR	+ STOR. RINGS
MM B	(1166,25)	(0,24)	BAILEY	68	CNTR	- STOR. RINGS
MM B	1166,16	0,31	BAILEY	68	CNTR	+ STOR. RINGS
MM	1166,16	0,31	BAILEY	68	CNTR	+ STOR. RINGS
MM	1166,14	0,31	BAILEY	68	CNTR	+ STOR. RINGS

5/69
5/69
5/69
5/69
5/69

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

4 MUON PARTIAL DECAY MODES

P1	MUON INTO E (E-NEU) (MU-NEU)	.5*	0*	0
P2	MUON INTO E 2GAMMA	.5*	0*	0
P3	MUON INTO 2ELECTRONS	.5*	.5*	.5
P4	MUON INTO E GAMMA	.5*	0	0

DECAY MASSES

4 MUON BRANCHING RATIOS

R1	MUON INTO E+2GAMMA (IN UNITS OF 10**5)	(P2)/(P1)
R1	(1,610R LESS C.L.=.90	FRANKEL 63 OSPK
R2	MUON INTO 3E (IN UNITS OF 10**7)	(P3)/(P1)
R2 F	(5,010R LESS C.L.=.90	PARKER 62 CNTR
R2 F	(1,310R LESS C.L.=.90	ALIKHANOV 62 OSPK
R2 F	(1,510R LESS C.L.=.90	FRANKEL 63 CNTR
R2 F	(1,2510R LESS C.L.=.90	BABAIEV 63 OSPK
R2 F	FOUR ABOVE EXPERIMENTS EVALUATED UPPER LIMITS ASSUMING A SECOND ORDER V-4 NEUTRINO LOOP DIAGRAM. LIMITS NOT SIGNIFICANTLY CHANGED BY ASSUMING A CONSTANT MATRIX ELEMENT.	
R3	MUON INTO E+GAMMA (IN UNITS OF 10**8)	(P4)/(P1)
R3	4,3 OR LESS C.L.=.90	FRANKEL 63 OSPK
R3	2,2 OR LESS C.L.=.90	PARKER 64 OSPK
R3	2,9 OR LESS C.L.=.90	KORENCHEN 70 OSPK

2/71*

Stable Particles

For notation, see illustrated key at beginning of data card listings.

Table with 5 columns: RHO/RHO C/RHP/RHO D/RHO E/RHO F/RHO G/RHO H/RHO I/RHO J/RHO K/RHO L/RHO M/RHO N/RHO O/RHO P/RHO Q/RHO R/RHO S/RHO T/RHO U/RHO V/RHO W/RHO X/RHO Y/RHO Z/RHO AA/RHO AB/RHO AC/RHO AD/RHO AE/RHO AF/RHO AG/RHO AH/RHO AI/RHO AJ/RHO AK/RHO AL/RHO AM/RHO AN/RHO AO/RHO AP/RHO AQ/RHO AR/RHO AS/RHO AT/RHO AU/RHO AV/RHO AW/RHO AX/RHO AY/RHO AZ/RHO BA/RHO BB/RHO BC/RHO BD/RHO BE/RHO BF/RHO BG/RHO BH/RHO BI/RHO BJ/RHO BK/RHO BL/RHO BM/RHO BN/RHO BO/RHO BP/RHO BQ/RHO BR/RHO BS/RHO BT/RHO BU/RHO BV/RHO BW/RHO BX/RHO BY/RHO BZ/RHO CA/RHO CB/RHO CC/RHO CD/RHO CE/RHO CF/RHO CG/RHO CH/RHO CI/RHO CJ/RHO CK/RHO CL/RHO CM/RHO CN/RHO CO/RHO CP/RHO CQ/RHO CR/RHO CS/RHO CT/RHO CU/RHO CV/RHO CW/RHO CX/RHO CY/RHO CZ/RHO DA/RHO DB/RHO DC/RHO DD/RHO DE/RHO DF/RHO DG/RHO DH/RHO DI/RHO DJ/RHO DK/RHO DL/RHO DM/RHO DN/RHO DO/RHO DP/RHO DQ/RHO DR/RHO DS/RHO DT/RHO DU/RHO DV/RHO DW/RHO DX/RHO DY/RHO DZ/RHO EA/RHO EB/RHO EC/RHO ED/RHO EE/RHO EF/RHO EG/RHO EH/RHO EI/RHO EJ/RHO EK/RHO EL/RHO EM/RHO EN/RHO EO/RHO EP/RHO EQ/RHO ER/RHO ES/RHO ET/RHO EU/RHO EV/RHO EW/RHO EX/RHO EY/RHO EZ/RHO FA/RHO FB/RHO FC/RHO FD/RHO FE/RHO FF/RHO FG/RHO FH/RHO FI/RHO FJ/RHO FK/RHO FL/RHO FM/RHO FN/RHO FO/RHO FP/RHO FQ/RHO FR/RHO FS/RHO FT/RHO FU/RHO FV/RHO FW/RHO FX/RHO FY/RHO FZ/RHO GA/RHO GB/RHO GC/RHO GD/RHO GE/RHO GF/RHO GG/RHO GH/RHO GI/RHO GJ/RHO GK/RHO GL/RHO GM/RHO GN/RHO GO/RHO GP/RHO GQ/RHO GR/RHO GS/RHO GT/RHO GU/RHO GV/RHO GW/RHO GX/RHO GY/RHO GZ/RHO HA/RHO HB/RHO HC/RHO HD/RHO HE/RHO HF/RHO HG/RHO HH/RHO HI/RHO HJ/RHO HK/RHO HL/RHO HM/RHO HN/RHO HO/RHO HP/RHO HQ/RHO HR/RHO HS/RHO HT/RHO HU/RHO HV/RHO HW/RHO HX/RHO HY/RHO HZ/RHO IA/RHO IB/RHO IC/RHO ID/RHO IE/RHO IF/RHO IG/RHO IH/RHO II/RHO IJ/RHO IK/RHO IL/RHO IM/RHO IN/RHO IO/RHO IP/RHO IQ/RHO IR/RHO IS/RHO IT/RHO IU/RHO IV/RHO IW/RHO IX/RHO IY/RHO IZ/RHO JA/RHO JB/RHO JC/RHO JD/RHO JE/RHO JF/RHO JG/RHO JH/RHO JI/RHO JJ/RHO JK/RHO JL/RHO JM/RHO JN/RHO JO/RHO JP/RHO JQ/RHO JR/RHO JS/RHO JT/RHO JU/RHO JV/RHO JW/RHO JX/RHO JY/RHO JZ/RHO KA/RHO KB/RHO KC/RHO KD/RHO KE/RHO KF/RHO KG/RHO KH/RHO KI/RHO KJ/RHO KK/RHO KL/RHO KM/RHO KN/RHO KO/RHO KP/RHO KQ/RHO KR/RHO KS/RHO KT/RHO KU/RHO KV/RHO KW/RHO KX/RHO KY/RHO KZ/RHO LA/RHO LB/RHO LC/RHO LD/RHO LE/RHO LF/RHO LG/RHO LH/RHO LI/RHO LJ/RHO LK/RHO LL/RHO LM/RHO LN/RHO LO/RHO LP/RHO LQ/RHO LR/RHO LS/RHO LT/RHO LU/RHO LV/RHO LW/RHO LX/RHO LY/RHO LZ/RHO MA/RHO MB/RHO MC/RHO MD/RHO ME/RHO MF/RHO MG/RHO MH/RHO MI/RHO MJ/RHO MK/RHO ML/RHO MN/RHO MO/RHO MP/RHO MQ/RHO MR/RHO MS/RHO MT/RHO MU/RHO MV/RHO MW/RHO MX/RHO MY/RHO MZ/RHO NA/RHO NB/RHO NC/RHO ND/RHO NE/RHO NF/RHO NG/RHO NH/RHO NI/RHO NJ/RHO NK/RHO NL/RHO NM/RHO NO/RHO NP/RHO NQ/RHO NR/RHO NS/RHO NT/RHO NU/RHO NV/RHO NW/RHO NX/RHO NY/RHO NZ/RHO OA/RHO OB/RHO OC/RHO OD/RHO OE/RHO OF/RHO OG/RHO OH/RHO OI/RHO OJ/RHO OK/RHO OL/RHO OM/RHO ON/RHO OO/RHO OP/RHO OQ/RHO OR/RHO OS/RHO OT/RHO OU/RHO OV/RHO OW/RHO OX/RHO OY/RHO OZ/RHO PA/RHO PB/RHO PC/RHO PD/RHO PE/RHO PF/RHO PG/RHO PH/RHO PI/RHO PJ/RHO PK/RHO PL/RHO PM/RHO PN/RHO PO/RHO PP/RHO PQ/RHO PR/RHO PS/RHO PT/RHO PU/RHO PV/RHO PW/RHO PX/RHO PY/RHO PZ/RHO QA/RHO QB/RHO QC/RHO QD/RHO QE/RHO QF/RHO QG/RHO QH/RHO QI/RHO QJ/RHO QK/RHO QL/RHO QM/RHO QN/RHO QO/RHO QP/RHO QQ/RHO QR/RHO QS/RHO QT/RHO QU/RHO QV/RHO QW/RHO QX/RHO QY/RHO QZ/RHO RA/RHO RB/RHO RC/RHO RD/RHO RE/RHO RF/RHO RG/RHO RH/RHO RI/RHO RJ/RHO RK/RHO RL/RHO RM/RHO RN/RHO RO/RHO RP/RHO RQ/RHO RR/RHO RS/RHO RT/RHO RU/RHO RV/RHO RW/RHO RX/RHO RY/RHO RZ/RHO SA/RHO SB/RHO SC/RHO SD/RHO SE/RHO SF/RHO SG/RHO SH/RHO SI/RHO SJ/RHO SK/RHO SL/RHO SM/RHO SN/RHO SO/RHO SP/RHO SQ/RHO SR/RHO SS/RHO ST/RHO SU/RHO SV/RHO SW/RHO SX/RHO SY/RHO SZ/RHO TA/RHO TB/RHO TC/RHO TD/RHO TE/RHO TF/RHO TG/RHO TH/RHO TI/RHO TJ/RHO TK/RHO TL/RHO TM/RHO TN/RHO TO/RHO TP/RHO TQ/RHO TR/RHO TS/RHO TT/RHO TU/RHO TV/RHO TW/RHO TX/RHO TY/RHO TZ/RHO UA/RHO UB/RHO UC/RHO UD/RHO UE/RHO UF/RHO UG/RHO UH/RHO UI/RHO UJ/RHO UK/RHO UL/RHO UM/RHO UN/RHO UO/RHO UP/RHO UQ/RHO UR/RHO US/RHO UT/RHO UV/RHO UW/RHO UX/RHO UY/RHO UZ/RHO VA/RHO VB/RHO VC/RHO VD/RHO VE/RHO VF/RHO VG/RHO VH/RHO VI/RHO VJ/RHO VK/RHO VL/RHO VM/RHO VN/RHO VO/RHO VP/RHO VQ/RHO VR/RHO VS/RHO VT/RHO VU/RHO VV/RHO VW/RHO VX/RHO VY/RHO VZ/RHO WA/RHO WB/RHO WC/RHO WD/RHO WE/RHO WF/RHO WG/RHO WH/RHO WI/RHO WJ/RHO WK/RHO WL/RHO WM/RHO WN/RHO WO/RHO WP/RHO WQ/RHO WR/RHO WS/RHO WT/RHO WU/RHO WV/RHO WW/RHO WX/RHO WY/RHO WZ/RHO XA/RHO XB/RHO XC/RHO XD/RHO XE/RHO XF/RHO XG/RHO XH/RHO XI/RHO XJ/RHO XK/RHO XL/RHO XM/RHO XN/RHO XO/RHO XP/RHO XQ/RHO XR/RHO XS/RHO XT/RHO XU/RHO XV/RHO XW/RHO XX/RHO XY/RHO XZ/RHO YA/RHO YB/RHO YC/RHO YD/RHO YE/RHO YF/RHO YG/RHO YH/RHO YI/RHO YJ/RHO YK/RHO YL/RHO YM/RHO YN/RHO YO/RHO YP/RHO YQ/RHO YR/RHO YS/RHO YT/RHO YU/RHO YV/RHO YW/RHO YX/RHO YY/RHO YZ/RHO ZA/RHO ZB/RHO ZC/RHO ZD/RHO ZE/RHO ZF/RHO ZG/RHO ZH/RHO ZI/RHO ZJ/RHO ZK/RHO ZL/RHO ZM/RHO ZN/RHO ZO/RHO ZP/RHO ZQ/RHO ZR/RHO ZS/RHO ZT/RHO ZU/RHO ZV/RHO ZW/RHO ZX/RHO ZY/RHO ZZ/RHO

Table with 5 columns: FISHER 59 PRL 3 349, ASTURRY 60 RCHC CONF 60 542, DEVONS 60 PRL 5 330, LATHROP 60 NC 17 109, LATHROP 60 NC 17 114, REITER 60 PRL 5 22, TELEDGI 60 RCHC CONF 60 713, CHARPAK 61 PRL 6 128, HUTCHINS 61 PRL 7 129, SHAPIRO 62 PR 125 1022, FAIRLEY 66 NC 45A 281, FISHER, LENTIC, LUNDY, MEUNIER, STROTT (CFRN), ASTURRY, HATTERSLEY, HUSSAIN + (LIVERPOOL), DEVONS, GDAL, LEDERMAN, SHAPIRO (COLUMBIA), J LATHROP, R A LUNDY, V L TELEDGI + (EFINS), J LATHROP, R A LUNDY, S PENMAN + (EFINS), REITER, ROMANOWSKI, SUTTON + (CARNEGIE), V L TELEDGI (CFRN), CHARPAK, FARLEY, GARWIN, MULLER, SPNS + (CFRN), D P HUTCHINSON, J MENES + (COLUMBIA), G SHAPIRO, L M LEDERMAN (COLUMBIA), FAIRLEY, BAILEY, BROWN, GIESCH + (CFRN), 8 CHARGED PION (140, JPG=0--1) I=1, 8 CHARGED PI MASS (MEV), M 139.37 0.20 CROWE 54 CNTR - 6/68, M 139.68 0.15 BARKAS 56 EMUL + 2/71*, M 139.577 0.013 SHAFER 67 CNTR + MESONIC ATOMS 2/71*, M M (139.550) (0.008) MALSBRUG 71 PRELIMINARY 2/71*, M M THIS VALUE DIFFERS CONSIDERABLY FROM PREVIOUS VALUE. THIS DISCREPANCY IS BEING INVESTIGATED. VALUE MOMENTARELY NOT USED. 2/71*, M AVG 139.577 0.013 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/71*, M FIT 139.576 0.011 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/71*, 8 P1+ MU+ MASS DIFFERENCE (MEV), D 34.00 0.076 BARKAS 56 EMUL 8/67, D 33.89 0.076 BARKAS 56 EMUL 8/67, D 145 33.881 0.035 HYMAN 67 HEFC + K-HE 2/71*, D 33.925 0.025 RHOOT 70 CNTR + MAGNETIC SPECT. 2/71*, D AVG 33.916 0.019 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/71*, D FIT 33.916 0.011 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/71*, 8 ((MASS P1+) - (MASS P1-)) / AVERAGE (PERCENT), DM 0.02 0.05 GREENBERG 69 CNTR 5/70, 8 CHAR. PI LIFETIME (UNITS 10**--9), T 25.6 0.5 0.5 CROWE 57 RVUE 6/68, T 25.6 0.8 0.8 ANDERSON 40 CNTR 6/68, T 8000 25.66 0.32 0.32 ASHKIN 60 CNTR + 9/66, T 26.02 0.04 ECKHAUS 65 CNTR + 6/68, T 25.6 0.3 BARDON 66 CNTR 6/68, T 25.9 0.3 DUNAITSEV 66 CNTR 6/68, T N (26.40) (0.08) KINSEY 66 CNTR + 5/66, T N SYSTEMATIC ERRORS IN CALIB. IN THIS EXP. DISCUSSED BY NORDBERG 67 8/67, T 26.67 0.24 LORKOWICZ 66 CNTR 9/66, T 26.04 0.05 NORDBERG 67 CNTR + 8/67, T 26.02 0.04 GREENBERG 69 CNTR 5/70, T AVG 26.024 0.024 0.024 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.0) 8 MEANLIFE DIFFERENCE, (+) - (-) / AVGE. (PERCENT), DT N THIS QUANTITY IS A MEASURE OF CPT INVARIANCE IN N. I., DT 0.23 0.40 LORKOWICZ 66 CNTR SEE NOTE L 9/66, DT L ABOVE IS THE MOST CONSERVATIVE VALUE QUOTED BY AUTHORS 9/66, DT 0.4 0.7 BARDON 66 CNTR 7/66, DT -0.14 0.25 PETRUKHIN 68 CNTR 8/68, DT 0.055 0.055 0.071 AYRES 69 CNTR NEW EXPERIMENT 10/69, DT AVG 0.053 0.061 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0), 8 CHARGED PION PARTIAL DECAY MODES, P1 CHAR. PION INTO MU (MU-NEU) 105+ 0, P2 CHAR. PION INTO E (E-NEU) .5+ 0, P3 CHAR. PION INTO MU (MU-NEU) GAMMA 105+ 0+ 0, P4 CHAR. PION INTO P10 E (E-NEU) 134+ .5+ 0, P5 CHAR. PION INTO E NEU GAMMA .5+ 0+ 0, P6 CHAR. PION INTO E NEU E+ E- .5+ 0+ .5+ .5, 8 CHARGED PION BRANCHING RATIOS, R1 CHAR. PION INTO MU NEU GAMMA (UNITS 10**--6) (P3)/(P1) 26 1.24 0.25 CASTAGNOL 58 EMUL E (MU) LT. 3.38 MV, R2 CHAR. PION INTO E NEU (UNITS 10**--6) (P2)/(P1) 1.21 0.07 ANDERSON 60 CNTR, R2 1.247 0.028 DI CAPUA 64 CNTR, R2 AVG 1.242 0.026 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0), R3 CHAR. PION INTO P10 E (E-NEU) (P4)/(P1) 36 0.97 0.20 BARTLETT 64 OSPK, R3 38 1.07 0.21 BACASTOW 65 OSPK + 6/66, R3 1.10 0.26 BERTMAN 65 OSPK 7/66, R3 43 1.1 0.2 DUNAITSEV 65 CNTR 6/68, R3 332 1.00 0.08 0.10 DEPOMMER 68 CNTR 3/68, R3 AVG 1.023 0.069 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0), R4 CHAR. PION INTO E NEU GAMMA (UNITS 10**--8) (P5)/(P1) 143 3.0 0.5 DEPOMMER 63 CNTR GAM KE 50-90 MEV 6/66, R5 CHAR. PION INTO E NEU E+ E- (UNITS 10**--8) (P6)/(P1) (3.4) OR LESS C.L. = .90 KORENCHEN 70 OSPK + 2/71*

Stable Particles

For notation, see illustrated key at beginning of data card listings.

REFERENCES
9 CHARGED PION [140, JPG=0--] I=1

CROWE 54 PR 96 470	K M CROWE, R H PHILLIPS (LRL)
BARKAS 56 PR 101 778	W H BARKAS, W BIRNBAUM, F M SMITH (LRL)
CROWE 57 NC 541	K M CROWE (STANFORD HEPL)
CASTAGNO 58 PR 112 1779	C CASTAGNO, I M MUCHNIK (ROME I F)
ANDERSON 60 PR 119 2050	H L ANDERSON, T FUJII, R H MILLER + (EPFINS)
ASHKIN 60 NC 16 490	ASHKIN, FAZZINI, FIDECARD, LIPMAN + (CERN)
DEPOMMIE 63 PL 7 285 1451	P DEPOMMIE, HEINTZE, RIBBI, SBERGEL (CERN)
BARTLETT 64 PR 1368 1452	BARTLETT, DEVONS, MEYER, ROSEN (COLUMBIA)
DI CAPUA 64 PR 1338 1333	DI CAPUA, GARLAND, PINDROM, STRELZOFF (COLUM)
RACASTON 65 PR 139 8407	+GHESQUIERE, NIEGAND, LARSEN (LRL+SLAC)
BERTRAM 65 PR 139 A 617	BERTRAM, MEYER, CARR IGAN + (MICH+CARNEGIE)
DUNAITSEV 65 JETP 20 58	DUNAITSEV, PETRUKHIN, PROKOSHIN + (DUBNA)
ECKHAUSE 65 PL 19 348	ECKHAUSE, HARRIS, SHULER + (WILLIAM AND MARY)
BARDON 66 PRL 16 775	BARDON, DORF, DORFAN, KRIEGER + (COLUMBIA)
DUNAITSEV 66 PL 23 283	*KUTYIN, PROKOSHIN, RASUVAEV, SIMONOV (DUBNA)
KINSEY 66 PR 144 1132	KINSEY, LORKWICZ, NORDBERG (ROCHESTER UNIV)
LOBKOWICZ 66 PRL 17 548	LOBKOWICZ, MELISSINOS, NAGASHIMA + (ROCH+BNL)
HYMAN 67 PL 258 376	+LOKEN, PEWITT, DEARICK + (ANL+CARN+NWES)
NORDBERG 67 PL 248 594	NORDBERG, LOBKOWICZ, BURMAN (ROCHESTER UNIV)
SHAFFER 67 PR 163 1451	SHAFFER, E. SHAFFER (LRL)
SEF ALSO PRL 14 923	SHAFFER, CROWE, JENKINS (LRL)
DEPOMMIE 68 NUC PHYS 84 189	DEPOMMIE, DUCLOS, HEINTZE, KLEINKNECHT + (CERN)
PETRUKHIN 68 JINP 91-3862	PETRUKHIN, RYKALIN, KHAZINS, CISEK (DUBNA)
AYRES 69 UCRL-18369	DAVID S AYRES (THESSIS) (LRL)
ALSO 67 PR 157 1288	AYRES, CALDWELL, GREENBERG, KENNEY, KURZ + (LRL)
ALSO 68 PRL 21 261	AYRES, CORMACK, GREENBERG, KENNEY + (LRL+UCSB)
GREENBERG 65 PR 125 1267	+AYRES, CORMACK, KENNEY, CALDWELL + (LRL+UCSB)
ROOTH 70 PL 328 723	+JOHNSON, WILLIAMS, WORMALD (LIVP)
KORENCHENKO 70 JINR P1-5250	KORENCHENKO, KOSTIN, MICHELMACHER + (JINR)
MALSBURG 71 THESSIS	C. VON DER MALSBURG (HEIDELBERG)

PAPERS NOT REFERRED TO IN DATA CARDS

MERRISON 62 ADVP 11 1	A W MERRISON (LIVERPOOL)
SHAPIRO 62 PR 125 1022	G SHAPIRO, L M LEORFMAN (COLUMBIA)
CZIRR 63 PR 130 341	JOHN B CZIRR (LRL)

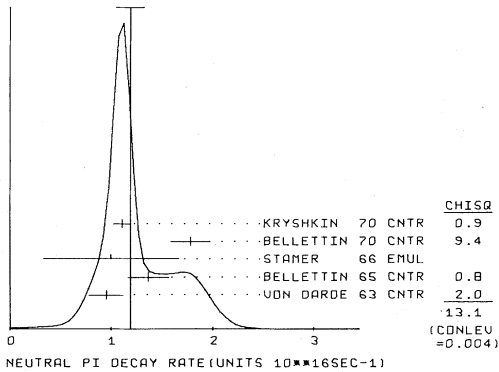
π^0
9 NEUTRAL PION [135, JPG=0--] I=1

D (5.37) (1.0)	PANOFSKY 51 CNTR -	
D 4.40 0.31	CHINDENSKY 54 CNTR -	
D 4.62 0.05	HADDUCK 59 CNTR -	
D 4.60 0.04	HILLMAN 59 CNTR -	
D 4.55 0.07	CASSELLS 59 CNTR -	
D 4.6056 0.0055	CZIRR 63 CNTR -	
D 4.59 0.03	PETRUKHIN 63 CNTR -	
D 4.6034 0.0052	VASILEVSKY 66 CNTR -	9/66
D AVG 4.6041 0.0037	AVFRAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

9 PION LIFETIME (UNITS 10**=-16)

T N 76 (1.9) (0.5) (0.5)	GLASSER 61 EMUL	
T N 45 (2.3) (1.1) (1.0)	TJETGE 62 EMUL	
T N 88 (2.8) (0.9) (0.9)	KOLLER 63 EMUL	SEE STAMER 66
T 1.05 0.18	VON DARDE 63 CNTR	
T N 75 (1.7) (0.5) (0.5)	SHWE 64 EMUL	
T 0.730 0.105	BELLETTINI 65 CNTR	6/66
T N 67 (1.6) (0.6) (0.5)	EVANS 65 EMUL	6/66
T K 232 1.0 0.5	STAMER 66 EMUL	9/67
T 0.9 0.06	BELLETTINI 70 CNTR	PRIM.EFF. ON NUC 7/70
T 0.9 0.068	KRYSHKIN 70 CNTR	PRIMAKOFF EFFECT 12/70*
T N OLD EMULSION MEASUREMENTS NOT USED BECAUSE OF POSSIBLE SYSTEMATIC		
T N SHIFT TO LARGER LIFETIME VALUES		
T K INCLUDES EVENTS OF KOLLER 63		8/67
T 0.839 0.103 0.092	AVFRAGE (ERROR INCL. SCALE FACTOR OF 2.1)	
T	(SEE IDEOGRAM BELOW)	

WEIGHTED AVERAGE = 1.19 ± 0.14
ERROR SCALED BY 2.1



9 NEUTRAL PION PARTIAL DECAY MODES

P1	PIO INTO 2GAMMA	DECAY MASSES
P2	PIO INTO E+ E- GAMMA	0+ 0
P3	PIO INTO 4ELECTRONS	.5+ .5+ 0
P4	PIO INTO 3 GAMMA	.5+ .5+ .5+ .5
		0+ 0+ 0

9 NEUTRAL PION BRANCHING RATIOS

R1	PIO INTO (GAMMA E+ E-)/(2GAMMA)	(P2)/(P1)	
R1	(0.0119) THEORFT. CALC. JOSEPH 60 HBC	QUANTUM FLECT.	9/66
R2	0 (5.0) OR LESS	DUCLOS 65 CNTR	CL=90 PERCENT 6/66
R1 S	3071 0.01166 0.00047	SAMIOS 61 HBC	PI-P TO PIO N
R1	SAMIOS VALUE USES PANOFSKY RATIO = 1.62		
R1	AVG 0.0117 0.0004	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R2	PIO INTO (3 GAMMA)/(2 GAMMA) (UNITS 10**=-6)	(P4)/(P1)	
R2	0 (5.0) OR LESS	KUTIN 65 CNTR	90 PERCENT C.L. 3/68
R3	PIO INTO (E+E+E-)/(2 GAMMA) (UNITS 10**=-5)	(P3)/(P1)	
R3	(3.47) THEORETICAL CAL. KROLL 55	QUANTUM FLECT.	9/66
R3 N	146 3.18 0.30	SAMIOS 62 HBC	SEE NOTE N BELOW 6/66
R3 N	ABOVE VALUE USES PANOFSKY RATIO=1.62		

REFERENCES
9 NEUTRAL PION [135, JPG=0--] I=1

PANOFSKY 51 PR 91 565	W K H PANOFSKY, R L AAMDT, J HADLEY (LRL)
CHINDENSKY 54 PR 93 586	W CHINDENSKY, J STEINBERGER (COLUMBIA)
KROLL 55 PR 98 1355	N KROLL, W WADA (COLUMBIA+BNL)
CASSELLS 59 PPS 74 92	CASSELLS, JONES, MURPHY, O. NEILL (LIVERPOOL)
HADDUCK 59 PRL 3 478	HADDUCK, ABASHIAN, CROWE, CZIRR (LRL)
HILLMAN 59 NC 14 887	HILLMAN, MIDDELKOOP, YAMAGATA, ZAVATTINI (CERN)
BUDAGOV 60 JETP 11 755	BUDAGOV, VIKTOR, DZHELEPOV, ERMOLOV + (JINR)
JOSEPH 60 NC 16 997	D W JOSEPH (FEI)
GLASSER 61 PR 123 1014	R G GLASSER, N SEEMAN, B STILLER (BNL)
SAMIOS 61 PR 121 275	N P SAMIOS (COLUMBIA+BNL)
SAMIOS 62 PR 126 1844	SAMIOS, PLANO, PRODELL + (COLUMBIA+BNL)
TJETGE 62 PR 127 1324	J TJETGE, W PUESCHEL (MAX PLANCK INST)
CZIRR 63 PR 130 341	JOHN B CZIRR (LRL)
KOLLER 63 NC 27 1405	F L KOLLER, S TAYLOR, T HUETTER (STEVENS)
KOLLER 63	SEE ALSO STAMER 66
PETRUKHIN 63 SIENA CONF 208	V I PETRUKHIN, YU O PROKOSHIN (JINR)
VON DARD 63 PL 4 51	VON DARDEL, DEKKERS, MERMOD, VAN PUTTEN + (CERN)
SHWE 64 PR 1369 1439	H SHWE, F N SMITH, W H BARKAS (LRL)
BELLETTINI 65 NC 40 A 1139	BELLETTINI, REMORA, BRACCINI + (PISA+FRNZE)
DUCLOS 65 PL 19 253	DUCLOS, FREYTAG, HEINTZE + (CERN+HEIDELBERG)
EVANS 65 PR 139 B 982	D A EVANS (OXFORD)
KUTIN 65 JETP LETT 2 243	KUTIN, PETRUKHIN, PROKOSHIN (JINR)
STAMER 66 PR 151 1108	STAMER, TAYLOR, KOLLER, HUETTER + (STEVENS)
VASILEVSKY 66 PL 23 281	VASILEVSKY, VISHNYAKOV, DUNAITSEV + (DUBNA)
BELLETTINI 70 NC 66A 243	BELLETTINI, REMORA, LUBELSMYER + (PISA+BNL)
KRYSHKIN 70 JETP 30 1037	*STERLIGOV, USNV (TUMSK)

K^+
10 CHARGED K (494, JP=0-) I=1/2

M 493.9 0.2	CHEN 57 RVUE +	
M 493.7 0.3	BARKAS 63 EMUL -	
M 493.78 0.17	GREINER 65 EMUL +	VIA TAU DECAY 7/66
M 493.81 0.12	AVFRAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
M FIT 493.84 0.11	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	2/71*

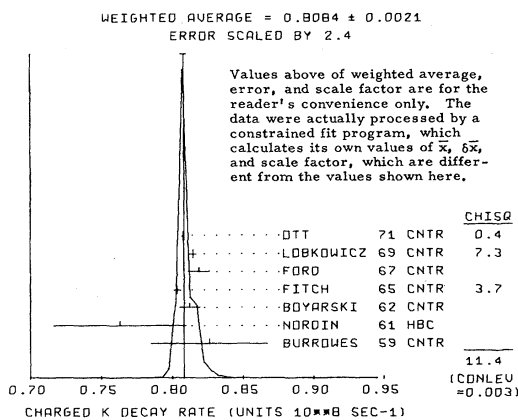
CHARGED K CONSTRAINED FIT
OVERALL FIT OF LIFETIME, WIDTHS AND BRANCHING RATIOS USES 46 DATA POINTS TO DETERMINE SEVEN QUANTITIES. OVERALL FIT HAS $\chi^2/S = 73.9$. MAIN CONTRIBUTION (13.5) COMES FROM R10 OF HA017. I WOULD SEE NO REASON TO REJECT THIS EXPERIMENT AT THIS TIME.

10 CHAR. K LIFETIME (UNITS 10**=-8)

T 0 (0.95) (0.36) (0.25)	ILOFF 56 EMUL	
T 0 52 (1.40) (0.31) (0.31)	FISERBERG 58 EMUL	
T 1.21 0.06 0.06	BURKOWES 59 CNTR	
T 0 33 (1.38) (0.24) (0.24)	FREDEN 60 EMUL	
T 0 (1.25) (0.22) (0.17)	BARKAS 61 EMUL	
T 0 51 (1.27) (0.36) (0.23)	BROWNIK 61 EMUL	
T 293 1.31 0.08 0.08	NORDIN 61 HBC -	
T (1.24) (0.07)	NORDIN 61 RVUE -	
T 1.231 0.011 0.011	BYARSKI 62 CNTR +	
T 1.249 0.0098	FITCH 65 CNTR +	K AT REST 8/67
T 1.221 0.011	FORD 67 CNTR +	
T 1.2272 0.0036	LOBKOWICZ 69 CNTR +	K IN FLIGHT 9/66
T 34 1.2380 0.0016	OTT 71 CNTR	STOPPING K 2/71*
T 0 OLD EXPERIMENTS WITH LARGE ERRORS EXCLUDED FROM AVERAGING		
T AVG 1.2370 0.0032 0.0032	AVERAGE (ERROR INCL. SCALE FACTOR OF 2.4)	
T FIT 1.2371 0.0026	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.9)	(SEE IDEOGRAM BELOW)

Stable Particles

For notation, see illustrated key at beginning of data card listings.



10 LIFETIME DIFFERENCE, (+) - (-) / AVG. (PERCENT)

DT N	THIS QUANTITY IS A MEASURE OF CPT INVARIANCE IN W.1.			
DT	0.47	0.30	FORD	67 CNTR
DT	0.090	0.098	LOBKOWICZ	69 CNTR
DT	0.114	0.093	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.21)	

10 DECAY RATES DIFF. (+) - (-) / AV. (PERCENT)

D1	DIFFERENCE IN K MU2 RATES ((W1+)-(W1-))/W1		
D1	-0.54	0.41	FORD 67 CNTR
D2	DIFFERENCE IN TAU RATES ((W2+)-(W2-))/W2		
D2	-0.50	0.40	FLETCHER 67 DSPK
D2 F	(-0.04)	(0.21)	FORD 67 CNTR SFE NOTE F
D2 F 2M	(0.10)	(0.14)	FORD 70 ASPK SEE NOTE F
D2	0.0R	0.12	FORD 70 ASPK SEE NOTE F
D2 F	THE LAST IS THE COMBINED RESULT OF FORD 67 AND FORD 70		
D2	0.07	0.12	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
D3	DIFFERENCE IN TAU PRIME RATES ((W4+)-(W4-))/AVERAGE		
D3	1.82	1.8	HERZO 69 DSPK
D4	DIFFERENCE IN K PI2 RATES ((W2+)-(W2-))/AVERAGE		
D4	0.8	1.2	HERZO 69 DSPK

10 CHARGED K PARTIAL DECAY MODES

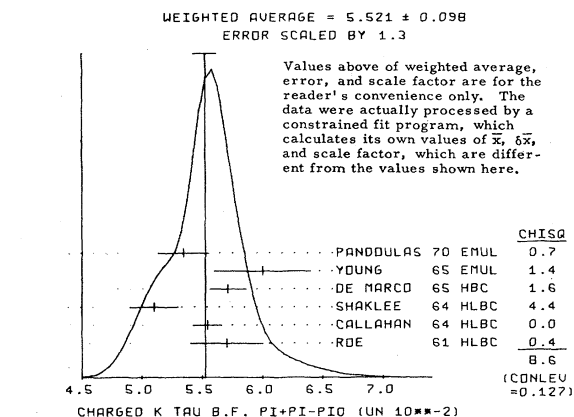
P1	CHAR. K INTO MU NEU	K MU2	105+ 0	DECAY MASSES
P2	CHAR. K INTO PI P10	K PI2	139+ 134	
P3	CHAR. K INTO PI P1+ PI-	TAU	139+ 139+ 139	
P4	CHAR. K INTO PI P2P10	TAU PRIME	139+ 134+ 134	
P5	CHAR. K INTO MU P10 NEU	K MU3	105+ 134+ 0	
P6	CHAR. K INTO E P10 NEU	K E3	134+ 0	
P7	POSIT. K INTO PI+ PI- E-NEU	K E+ 4	139+ 139+ .5+ 0	
P8	POSIT. K INTO PI+ PI+ E-NEU	K E+ 4	139+ 139+ .5+ 0	
P9	POSIT. K INTO PI+ PI+ MU- NEU	K+MU+ 4	139+ 139+ 105+ 0	
P10	CHAR. K INTO E NEU	K E 2	.5+ 0	
P11	CHAR. K INTO MU NEU GAMMA	K MU RAD	105+ 0+ 0	
P12	CHAR. K INTO PI P10 GAMMA	K PI RAD	139+ 134+ 0	
P13	CHAR. K INTO PI P1+ PI- GAMMA	TAU RAD	139+ 139+ 139+ 0	
P14	CHAR. K INTO PI E+ E- NEU	PI E E	139+ .5+ .5	
P15	CHAR. K INTO PI MU+ MU-	PI MU MU	139+ 105+ 105	
P16	CHAR. K INTO PI GAMMA GAMMA	PI GAM GAM	139+ 0+ 0	
P17	CHAR. K INTO PI F NEUTRINO GAMMA	PI E NEU GAM	139+ .5+ 0+ 0	
P18	NEG. K INTO PI+ E- E-	PI+ E- E-	139+ .5+ .5	
P19	CHAR. K INTO PI NEU NEU	PI NEU NEU	139+ 0+ 0	
P20	CHAR. K INTO E NEU GAMMA	K E RAD	.5+ 0+ 0	
P21	CHAR. K INTO PI GAMMA	K PI GAM	139+ 0	
P22				

10 CHARGED K DECAY RATES

W1	CHAR. K INTO MU NEU (K MU)	(UN. 10**6 SEC-1) (P1)	
W1	51.2	0.8	FORD 67 CNTR
W1	51.55	0.25	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
W2	CHARG. K INTO PI P1+ PI- (TAU) <td>(UN. 10**6 SEC-1) (P3)</td> <td></td>	(UN. 10**6 SEC-1) (P3)	
W2 F	(4.496)	(0.030)	FORD 57 CNTR -- SEE NOTE F
W2 F 2M	(4.529)	(0.032)	FORD 70 ASPK SEE NOTE F
W2	4.511	0.024	FORD 70 ASPK SEE NOTE F
W2 F	THE LAST IS THE COMBINED RESULT OF FORD 67 AND FORD 70		
W2	4.513	0.023	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
W3	CHAR. K INTO (TAU) - (TAU PRIME) <td>(UNITS 10**6 SEC-1) (P3-P4)</td> <td></td>	(UNITS 10**6 SEC-1) (P3-P4)	
W3	USED FOR DELTA I=1/2 TEST		
W3	3.158	0.039	FROM FIT
W4	CHAR. K INTO (MU P10 NEU) + (E P10 NEU) <td>(UNITS 10**6 SEC-1) (P5+P6)</td> <td></td>	(UNITS 10**6 SEC-1) (P5+P6)	
W4	USED FOR DELTA I=1/2 TEST		
W4	6.51	0.12	FROM FIT

10 CHARGED K BRANCHING RATIOS

R 0	OLD DATA EXCLUDED				
R1	CHAR. K INTO MU NEU (MU2)	(UNITS 10**2)	(P11)/TOTAL		
R1 0	(58.5)	(3.0)	BIRGE 56 EMUL +		
R1 0	(56.9)	(2.4)	ALEXANDER 57 EMUL +		
R1 0	OLD EXPERIMENTS NOT INCLUDED IN AVERAGING			1/71*	
R1	63.77	0.28	FROM FIT		
R2	CHAR. K INTO PI P10 (PI2)	(UNITS 10**2)	(P21)/TOTAL		
R2 0	(27.7)	(2.7)	BIRGE 56 EMUL +		
R2 0	(23.2)	(2.2)	ALEXANDER 57 EMUL +		
R2 0	EARLIER EXPERIMENTS NOT AVERAGED				
R2	(21.0)	(0.4)	CALLAHAN 65 HLBC	SFF R17	
R2	(21.6)	(0.6)	TRILLING 65 RVUE	6/66	
R2	20.92	0.29	FROM FIT		
R3	CHAR. K INTO PI P1+ PI-(TAU)	(UNITS 10**2)	(P31)/TOTAL		
R3 0	(5.6)	(0.4)	BIRGE 56 EMUL +		
R3 0	(6.8)	(0.4)	ALEXANDER 57 EMUL +		
R3 0	(5.2)	(0.3)	TAYLOR 59 EMUL +		
R3 0	EARLIER EXPERIMENTS NOT AVERAGED				
R3	5.7	0.3	ROF 61 HLBC +	9/66	
P1	2332	5.54	0.12	CALLAHAN 64 HLBC +	
R3	540	5.1	0.2	SHAKLEE 64 HLBC +	
R3	5.71	0.15	DE MARCO 65 HBC	9/66	
R3	6.0	0.4	YOUNG 65 EMUL +	6/66	
R3 P	693	5.34	0.21	PANDOLAS 70 EMUL +	10/70*
R3 P	INCLUDES EVENTS OF TAYLOR 59				
R3	5.521	0.098	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R3	5.583	0.030	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R3			(SFE IDEOGRAM BELOW)		



10 CHARGED K BRANCHING RATIOS

R4	CHAR. K INTO PI P2P10 (TAU PRIME)	(UNITS 10**2)	(P41)/TOTAL		
R4 0	(2.1)	(0.5)	BIRGE 56 EMUL +		
R4 0	(2.2)	(0.4)	ALEXANDER 57 EMUL +		
R4 0	(1.5)	(0.2)	TAYLOR 59 EMUL +		
R4 0	EARLIER EXPERIMENTS NOT AVERAGED				
R4	1.7	0.2	ROF 61 HLBC +		
R4	1.8	0.7	SHAKLEE 64 HLBC +		
R4 P	1.98	1.53	0.11	PANDOLAS 70 EMUL +	
R4 P	INCLUDES EVENTS OF TAYLOR 59			10/70*	
R4	1.413	0.087	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R4	1.677	0.044	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R5	CHAR. K INTO MU P10 NEU (MU3)	(UNITS 10**2)	(P51)/TOTAL		
R5 0	(2.8)	(1.0)	BIRGE 56 EMUL +		
R5 0	(5.9)	(1.3)	ALEXANDER 57 EMUL +		
R5 0	(2.8)	(0.4)	TAYLOR 59 EMUL +		
R5 0	EARLIER EXPERIMENTS NOT AVERAGED				
R5	3.20	0.11	FROM FIT		
R6	CHAR. K INTO E P10 NEU (E3)	(UNITS 10**2)	(P61)/TOTAL		
R6 0	(3.2)	(1.3)	BIRGE 56 EMUL +		
R6 0	(5.1)	(1.3)	ALEXANDER 57 EMUL +		
R6 0	EARLIER EXPERIMENTS NOT AVERAGED				
R6	5.0	0.5	ROF 61 HLBC +		
R6	4.29	4.7	0.3	SHAKLEE 64 HLBC +	
R6	4.78	0.26	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R6	4.855	0.070	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R7	CHAR. K INTO (PI2 + MU3)/(TOTAL)	(P2+P3)/TOTAL			
R7	WE COMBINE THESE TWO MODES FOR EXPTS MEASURING THEM IN XENON RC				
R7	BECAUSE OF DIFFICULTIES OF SEPARATING THEM THERE				
R7	23.4	1.1	ROF 61 HLBC +	11/67	
R7	25.4	0.9	SHAKLEE 64 HLBC +	11/67	
R7	24.60	0.98	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)		
R7	24.12	0.29	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)		
R8	POSIT. K INTO PI+ PI+ E- NEU	(UNITS 10**7)	(P81)/TOTAL		
R8	(20.)	OR LESS	BIRGE 65 FRC + 95 PER CT CONF	9/66	
R8	(18.9)	OR LESS	ELY 69 HLBC + 95 PER CT CONF	10/69	
R9	POSIT. K INTO PI+ PI- MU+ NEU	(UNITS 10**5)	(P91)/TOTAL		
R9	0.77	0.54	0.50	CLINE 65 FRC +	8/66
R10	POSIT. K INTO PI+ PI+ MU- NEU	(UNITS 10**6)	(P101)/TOTAL		
R10	(3.010R	LESS	BIRGE 65 FRC + 95 PER CT CONF	8/66	

Stable Particles

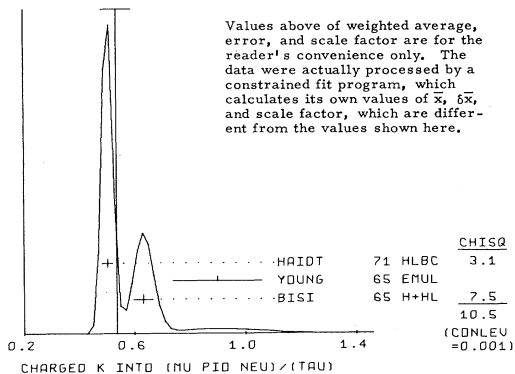
For notation, see illustrated key at beginning of data card listings.

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R11 CHAR. K INTO E NEU (UNITS 10**-5) (P111)/TOTAL
R11 (16.0) OR LESS HORREANI 66 HRC + CONLEV=0.95 11/67
R11 4 2.1 1.8 1.3 BOWEN 67 OSPK + 8/67
R11 BOWEN RESULT SHOULD BE CORRECTED TO 1.91+1.7--1.21 BECAUSE OF
R11 K+ TO E+ NEU GAMMA DECAYS BEFORE COMPARING WITH BOTTERILL 67 R2R
R12 CHAR. K INTO PI GAMMA GAMMA/TOTAL (UNITS 10**-4) (P171)/TOTAL
R12 (1.1) OR LESS CHEN 68 OSPK + T(P1)60 TO 90MEV 5/68
R12 (0.4) OR LESS KLEMS2 70 OSPK + T(P1)1GT 117 MEV 12/70*
R13 CHAR. K INTO PI PTD GAMMA (UNITS 10**-4) (P131)/TOTAL
R13 18 (2.2) (0.7) CLINE 66 FRC + P1+ KE 55-80 MEV 8/66
R13 E 0 (1.9) OR LESS ENMERSON 69 OSPK P1+ 55-80 MEV 10/69
R13 F 90 PER CENT CONFIDENCE
R14 CHAR. K INTO PI P1+ P1- GAMMA (UNITS 10**-4) (P141)/TOTAL
R14 1.0 0.4 STAMER 65 EMUL + 8/66
R15 CHAR. K INTO PI E+ F- (UNITS 10**-6) (P151)/TOTAL
R15 1 (1.1) OR LESS CAMERINI 64 FRC + 8/66
R15 (0.4) OR LESS CLINE 67 FRC + 11/67
R15 (4.4) OR LESS BISI 67 DBC + 90 PER CT CONF 11/67
R16 CHAR. K INTO PI MU+ MU- (UNITS 10**-6) (P161)/TOTAL
R16 (3.0) OR LESS CAMERINI 65 FRC + 90 PER CT CONF 8/66
R16 (2.4) OR LESS BISI 67 DBC + 90 PER CT CONF 11/67
R17 CHAR. K INTO (PI P10)/TAU (P21)/(P3)
R17 134 3.24 0.34 YOUNG 65 EMUL + 8/66
R17 1045 3.96 0.15 CALLAHAN 66 FBC + 9/66
R17 AVG 3.84 0.27 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)
R17 FIT 3.746 0.058 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
R18 CHAR. K INTO (PI 2P10)/TAU (P41)/(P3)
R18 2027 0.303 0.009 BISI 65 H+HL + 8/66
R18 17 0.393 0.099 YOUNG 65 EMUL + 8/66
R18 AVG 0.3037 0.0090 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R18 FIT 0.3004 0.0078 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
R19 CHAR. K INTO (MU P10 NEU)/TAU (P51)/(P3)
R19 2175 0.632 0.035 BISI 65 H+HL + 8/66
R19 33 0.90 0.16 YOUNG 65 EMUL + 8/66
R19 H 1505 (0.510) (0.017) EICHTEN 68 HLRC + 11/68
R19 H 1505 0.503 0.019 HAIDT 71 HLRC + 12/70*
R19 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68
R19 AVG 0.536 0.054 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.2)
R19 FIT 0.573 0.019 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
(SEE IDEOGRAM BELOW)

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WEIGHTED AVERAGE = 0.536 ± 0.054
ERROR SCALED BY 3.2



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R25 CHAR. K INTO (E P10 NEU)/(MU NEU) (P61)/(P1)
R25 472 0.0797 0.0054 AUERBACH 67 OSPK + 8/67
R25 THE VALUE .0785+-0.0025 GIVEN IN THE ABOVE REF IS AN AVERAGE OF
R25 AUERBACH 67 R25 AND CESTER 66 R23.
R25 960 0.0775 0.0033 ROTTERILL 68 ASPK + 5/68
R25 561 0.0069 0.006 GARLAND 68 OSPK + 4/68
R25 350 0.0069 0.006 ZELLER 69 ASPK + 10/69
R25 AVG 0.0783 0.0025 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
R25 FIT 0.0761 0.0011 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R26 CHAR. K INTO (MU P10 NEU)/(MU NEU) (P51)/(P1)
R26 310 0.0602 0.0046 AUERBACH 67 OSPK + 8/67
R26 426 0.0055 0.006 GARLAND 68 OSPK + 4/68
R26 240 0.0054 0.009 ZELLER 69 ASPK + 10/69
R26 AVG 0.0569 0.0029 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R26 FIT 0.0502 0.0018 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
R27 CHAR. K INTO (MU NEU)/(TAU) (P11)/(P3)
R27 R 427 (10.38) (0.82) YOUNG 65 EMUL + 9/66
R27 R DELETED FROM OVERALL FIT BECAUSE YOUNG 65 CONSTRAINS HIS RESULTS
R27 R TO ADD UP TO 1. ONLY YOUNG MEASURED MU2 DIRECTLY.
R27 FIT 11.421 0.085 FROM FIT
R28 CHAR. K INTO (E NEU)/(MU NEU) (UNITS 10**-5) (P111)/(P1)
R28 10 1.9 0.7 0.5 ROTTERILL 67 ASPK + 11/67
R28 8 1.8 0.8 0.6 MACEK 69 ASPK + 4/69
R28 66 2.15 0.35 CLARK 70 OSPK 11/70*
R28 AVG 2.04 0.28 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R29 CHAR. K INTO (MU P10 NEU)/(E P10 NEU) (P51)/(P6)
R29 C1509 0.703 0.056 CALLAHAN 66 HLRC 6/68
R29 5601 0.667 0.017 ROTTERILL 68 ASPK + 6/68
R29 A 1398 (0.604) (0.022) EICHTEN 68 HLRC 10/68
R29 A (0.596) (0.025) HAIDT 71 HLRC + 12/70*
R29 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68
R29 A ONLY INDIVIDUAL RATIOS INCLUDED IN FIT--SEE R19 AND R20--
R29 C FROM THIS EXPERIMENT WE USE ONLY THE MU3/F3 RATIO AND DO NOT
R29 C INCLUDE IN THE FIT THE RATIOS MU3/TAU AND E3/TAU SINCE THEY
R29 C SHOW LARGE DISAGREEMENTS WITH THE REST OF THE DATA.
R29 AVG 0.670 0.016 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R29 FIT 0.659 0.022 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
R30 CHAR. K INTO PI E NEU GAMMA/PI E NEU (P131)/(P6)
R30 8 0.012 0.008 BELLOTTI 67 HLRC + 11/67
R30 8 GAMMA ENERGY GREATER THAN 30 MEV
R31 K- INTO PI+ E- E-/TOTAL (UNITS 10**-5) (P191)/TOTAL
R31 TEST OF LEPTON NUMBER CONSERVATION
R31 (1.5) OR LESS CHANG 68 HRC - CL= 3/68
R32 CHAR. K INTO PI NEU NEU/ TOTAL (UNITS 10**-6) (P201)/TOTAL
R32 (100.0) OR LESS CL=90 CAMERINI 69 HLRC + TEST NEUTR. CURR. 5/70
R32 K 1.2 OR LESS THE RATIOS MU3/TAU AND E3/TAU SINCE THEY
R32 K ASSUMES PI+ SPECTRUM SAME AS P10 SPECTRUM IN KE3 DECAY
R33 CHAR. K INTO E NEU GAMMA / TOTAL (UNITS 10**-5) (P211)/TOTAL
R33 M (7.1) OR LESS CL=90 MACEK 70 OSPK + (P1) 234 TO 247 12/70*
R33 M ABOVE IS MEASUREMENT OF STRUCTURE-DEPENDENT DECAY ONLY
R34 CHAR. K INTO PI GAMMA / TOTAL (UNITS 10**-6) (P221)/TOTAL
R34 4. OR LESS, CL=90 KLEMS 70 OSPK + 1/71*
R35 CHAR. K INTO (TAU)/(TAU PRIME) (P3)/(P4)
R35 USED FOR DELTA I=1/2 TEST
R35 FIT 3.329 0.086 FROM FIT

```

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{\langle \delta P_i \delta P_i \rangle}$, while the off-diagonal elements are the normalized correlation coefficients $\langle \delta P_i \delta P_j \rangle / (\delta P_i \cdot \delta P_j)$. For the definitions of the individual P_i see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

P 1	P 2	P 3	P 4	P 5	P 6
P 1 .638+-0.003					
P 2 -.850 .209+-0.003					
P 3 -.138 -.062 .056+-0.000					
P 4 -.127 -.051 .145 .017+-0.000					
P 5 -.257 -.188 .083 .002 .032+-0.001					
P 6 -.230 -.198 .168 .009 .481 .049+-0.001					

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., $G_i \equiv \Gamma_i = \Gamma_{total} P_i$, in appropriate units. In analogy to the matrix above, the diagonal elements are $G_i \pm \delta G_i$, where $\delta G_i = \sqrt{\langle \delta G_i \delta G_i \rangle}$, while the off-diagonal elements are the normalized correlation coefficients $\langle \delta G_i \delta G_j \rangle / (\delta G_i \cdot \delta G_j)$. Note that, because of the error in Γ_{total} , the errors and correlations here are not directly derivable from those above.

G 1	G 2	G 3	G 4	G 5	G 6
G 1 .515+-0.003					
G 2 -.669 .169+-0.002					
G 3 -.106 -.038 .045+-0.000					
G 4 -.099 -.041 .139 .014+-0.000					
G 5 -.213 -.175 .083 .002 .026+-0.001					
G 6 -.166 -.173 .165 .010 .482 .039+-0.001					

```

R20 CHAR. K INTO (E P10 NEU)/TAU (P61)/(P3)
R20 230 0.90 0.06 HORREANI 66 HRC + 8/66
R20 37 0.90 0.16 YOUNG 65 EMUL + 8/66
R20 854 0.94 0.09 BELLOTTI 67 HLRC 11/67
R20 H 4385 (0.866) (0.021) EICHTEN 68 HLRC + 11/68
R20 H 4385 0.850 0.019 HAIDT 71 HLRC + 12/70*
R20 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68
R20 AVG 0.858 0.018 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R20 FIT 0.870 0.013 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R21 POSIT. K INTO (PI+ P1- E+ NEU)/TAU (UNITS 10**-4) (P71)/(P3)
R21 69 6.7 1.5 BIRGE 65 FBC + 8/66
R21 269 5.83 0.63 ELY 69 HLRC 11/68
R21 AVG 5.96 0.58 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R22 POSIT. K INTO (PI+ P1- MU+ NEU)/TAU (UNITS 10**-4) (P91)/(P3)
R22 1 (2.5) APPROX GREINER 64 EMUL + 8/66
R22 7 2.57 1.55 BISI 67 DBC + 11/67
R23 CHAR. K INTO (E P10 NEU)/(MU+ P12) (UNITS 10**-2) (P61)/(P1+P2)
R23 1479 6.89 0.21 CESTER 68 OSPK + 8/67
R23 5110 6.16 0.22 ESCHSTRUT 68 OSPK + 3/68
R23 AVG 6.02 0.15 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R23 FIT 5.733 0.090 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R24 CHAR. K INTO (PI P10)/(MU NEU) (P21)/(P1)
R24 0.3253 0.0065 AUERBACH 67 OSPK + 8/67
R24 1600 0.305 0.019 ZELLER 69 ASPK + 10/69
R24 AVG 0.3230 0.0065 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
R24 FIT 0.3280 0.0059 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

```


Stable Particles

Slope Parameter for $K \rightarrow 3\pi$ Decays

For the 3π decays of the K mesons we list the slope parameter "g" which is defined by

$$|M|^2 \propto 1 + g \frac{(s_3 - s_0)}{m_{\pi^+}^2} \quad (1)$$

where

$$s_i = (\underline{p}_K - \underline{p}_i)^2 = (m_K - m_i)^2 - 2m_K T_i \quad (2)$$

$$s_0 = \frac{1}{3} \sum s_i = \frac{1}{3}(m_K^2 + m_1^2 + m_2^2 + m_3^2) \quad (3)$$

$\underline{p}_K, \underline{p}_i$ are the four-vectors for the K and the i^{th} pion, and the index 3 refers to the odd pion. (4)

We refer to the three possible charged decays as τ, τ', τ^0

$$\begin{aligned} \tau^\pm & K^\pm \rightarrow \pi^+ \pi^- \pi^\pm \\ \tau'^\pm & K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm \\ \tau^0 & K^0 \rightarrow \pi^+ \pi^- \pi^0 \end{aligned}$$

where the odd pion is the third one.

There is no strong evidence so far that a second order term is needed in Eq. (1). Most of the experimenters have fitted their data with a second order term, which turned out to be consistent with zero. We use the value of g obtained when the second order term was dropped from the fit. HEUSSE 70 have studied the $K_L^0 \rightarrow \pi^0 \pi^0 \pi^0$ decay where only a second order term could explain deviation from uniformity of the Dalitz plot. They also get results consistent with a zero coefficient.

In the literature other definitions of slope parameters have appeared. We have converted to the definition of g in Eq. (1) whatever experimental quantity has been reported. We give the conversion to the definition (1) for two of the most widely used parameterizations and tabulate the conversion factors for the reader's convenience.

(a) For analysis of charged K's the expression often used is:

$$|M|^2 = 1 + a \frac{Y}{y}$$

with

$$Y = \frac{3T_3 - Q}{Q}, \quad Q = m_K - \sum m_i$$

The relevant formulae are:

$$Y = -\frac{3}{2} \frac{s_3 - s_0}{m_K Q} + \Delta, \quad \text{with } \Delta = \frac{m_{12} - m_3}{Q} \left(2 - \frac{m_3 + m_{12}}{m_K} \right)$$

and

$$g = \frac{-c_y a_y}{1 + a_y \Delta}, \quad \text{with } c_y = \frac{3}{2} \frac{m_{\pi^+}^2}{m_K Q}$$

(b) For the analysis of K^0 decay the expression often used is:

$$|M|^2 = 1 + 2a_t \frac{m_K}{m_{\pi^+}^2} (2T_3 - T_{3 \text{ max}})$$

with

$$T_{3 \text{ max}} = \frac{m_K^2 + m_3^2 - 4m_{12}^2}{2m_K} - m_3$$

The relevant transformations are

$$T_3 = -\frac{s_3 - s_0}{2m_K} + \frac{Q}{3} (1 + \Delta)$$

and

$$g = \frac{-2a_t}{1 + a_t c_t}, \quad \text{with } c_t = \frac{2m_K}{m_{\pi^+}^2} \left[\frac{2}{3} Q(1 + \Delta) - T_{3 \text{ max}} \right]$$

For the reader's convenience we give a table of numerical values for Q, $T_{3 \text{ max}}$, Δ , c_y and c_t , obtained using the masses from our August 1970 edition.

	Q	$T_{3 \text{ max}}$	Δ	c_y	c_t
τ^\pm	74.96	48.15	0	0.7894	0.0924
τ'^\pm	84.24	53.27	-0.0789	0.7025	-0.0778
τ^0	83.54	53.92	0.0798	0.7028	0.3176

Some K^0 authors use the above form of matrix element:

$$|M|^2 = 1 + 2a_u \frac{m_K}{m_{\pi^+}^2} (2T_3 - T_{3 \text{ max}}),$$

but define

$$T_{3 \text{ max}} = \frac{2}{3} Q.$$

The relevant transformation is then

$$g = \frac{-2a_u}{1 + a_u c_u}, \quad \text{with } c_u = \frac{2m_K}{m_{\pi^+}^2} \Delta = 0.2272.$$

Older K^0 analyses were done using

$$|M|^2 = 1 + a_v \frac{T_3}{m_K}$$

The relevant transformation is then

$$g = \frac{-c_v a_v}{1 + d_v a_v}$$

Stable Particles

For notation, see illustrated key at beginning of data card listings.

with

$$c_v = \frac{m_{\pi^+}^2}{2m_K^2} = 0.0393$$

$$\text{and } d_v = \frac{Q}{3m_K} (1 + \Delta) = 0.0604.$$

10 CHARGED K ENERGY DEPENDENCE OF DALITZ PLOT
MATRIX ELEMENT SQUARED = I + G (S3-S01)/(MPI**2)

```

GT* LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS CHARGED K INTO PI PI*PI-
    THESE EXPTS FIT **2*1*AY*. WE LIST G IN THE MAIN LISTING AND
    GIVE AY AT RIGHT. G=-1.5*AY*(MPI**2)/(MK*Q). SEE NOTE ABOVE.
GT+Z 5428 -0.22 0.024 ZINCHENKO 67 HBC + AY=0.28+-03 10/69
GT+ 9994 -0.218 0.016 RUTLER 68 HBC + AY=0.277+-020 10/69
GT+ 24 (-0.2231 0.0061) FORD 70 ASPK PRELIMINARY 1/71*
GT+ G17898 (-0.1961 0.012) GRAUMAN 70 HLRC + AY=0.228+-030 8/70*
GT+ 39819 -0.201 0.008 HOFFMASTE 70 HLRC + INCLUDES GRAUMAN 1/71*
GT+ G EMULS. DATA ADDED - ALL EVENTS INCLUDED BY HOFFMASTE 70 1/71*
GT+ Z ALSO INCLUDES DMC EVENTS
GT+ AVG * * * * * -0.2057 * * * * * AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

GT- LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS CHARGED K INTO PI PI*PI-
    FOR DEFINITION OF AY SEE NOTE UNDER S10GT+
GT- F 1347 (-0.2201 0.0351) FERRO-LUZ 61 HBC - AY=0.29+-045 10/69
GT-M 5778 -0.190 0.023 MOCSDSU 68 HBC - AY=0.242+-029 10/69
GT- 50919 -0.194 0.007 MAST 69 HBC - AY=0.247+-009 10/69
GT- F NO RADIATIVE CORRECTIONS INCLUDED
GT- M ALSO INCLUDES DMC EVENTS
GT- AVG * * * * * -0.1937 * * * * * AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

DG ((GT+)-(GT-))/((GT+)+(GT-)) IN PERCENT
DG 24 -0.70 0.53 FORD 70 ASPK 11/70*

GTP LINEAR ENERGY DEPENDENCE (G) FOR TAU PRIME DECAY CHA.K INTO PI PIP0
GTP 1792 0.48 0.04 KALMUS 64 HLBC + 10/69
GTP 1874 0.586 0.098 BISI 65 HLRC + ALSO HBC 10/69
GTP 4048 0.516 0.020 DAVISON 69 HLRC + ALSO EMUL 10/69
GTP 1157 0.602 0.049 AUBERT 70 HLRC + USES PI* SPECT. 1/71*
GTP 1527 0.602 0.061 AUBERT 70 HLRC + USES GAMMAS ALSO. 1/71*
GTP 198 0.516 0.074 PANDOLAS 70 EMUL + 10/70*
GTP AVG * * * * * -0.527 * * * * * AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.11
    
```

Note on K⁺ and K⁰ Form Factors

The definition of all the variables listed in this section can be found in the text. As in the past we keep the values of ξ as obtained in μ polarization measurements (ξ_B) separated from the values obtained from branching ratios and spectra (ξ_A), but combine them in order to display in an ideogram the discrepancies in $\text{Re } \xi$.

Many approximations are usually made to extract these parameters from the experimental data:

- 1) Scalar and tensor currents. There is no evidence for scalar or tensor currents so pure vector current is usually assumed.
- 2) $\text{Im } \xi$ so far is consistent with 0, and this is usually assumed in most of the experiments.
- 3) The q^2 dependence of the f_{\pm} form factors is assumed to be linear, as in Eq. (38) in the main text.
- 4) Radiative corrections are not serious; they change λ_+ by about 0.005 (GINSBERG 67 and 70).
- 5) Momentum transfer dependence of ξ . From Eq. (38) and Eq. (39) in the main text, it is clear that in the first approximation one has

$$\xi(q^2) = \xi(0)[1 + (\lambda_- - \lambda_+)t],$$

where $t = q^2/m_{\pi}^2$. This shows that the three parameters $\xi(0)$, λ_- , and λ_+ can be highly correlated. Many of the

old experiments have assumed $\lambda_+ = \lambda_- = 0$; however, recently it has become clear that in order to eliminate the discrepancies between ξ_A and ξ_B , λ_+ and λ_- should be included in the fits. We discuss each method of determining ξ separately.

a) Branching ratio analysis. From Eq. (40) in the main text it is clear that $\frac{d\xi}{d\lambda_+} \approx -\frac{1.4}{0.125} \approx -11$. Therefore for $\lambda_+ = 0.05$, $\Delta\xi = -0.55$. The old hypothesis $\lambda_+ = 0$ certainly is not very good.

b) Dalitz plot analysis. From Eq. (37) in the main text, it appears that the Dalitz plot distribution of the K_{e3} decay is essentially dependent only upon λ_+ and independent from ξ ; therefore, the values of λ_+ obtained from K_{e3} data can be directly averaged.

On the other hand for the $K_{\mu 3}$ Dalitz plot, all three parameters are involved and they are correlated. For example, CHIEN 70 (in a neutral K experiment) find that λ_+ is not sensitive to ξ and λ_- variations, whereas ξ and λ_- are highly correlated. HAIDT 71 (in a K^+ experiment) find that the uncorrelated parameters are $\xi(0)$ and $\xi(6.8)$. Because of these correlations, we list the values of λ_+ from $K_{\mu 3}$ decays separately from those obtained from K_{e3} decays and do not list λ_- at all. We refer the reader to the original papers for discussion of the λ_- parameter.

c) Muon polarization analysis. The μ polarization depends only on $\xi(q^2)$; the λ_+ and λ_- parameters do not appear explicitly in Eq. (41) in the main text. For large statistics one could determine $\xi(q^2)$ and then derive $\xi(0)$, λ_+ , and λ_- . This has not been possible with the statistics available so far.

In addition to the most common correlation among ξ , λ_+ , and λ_- , there are many experimental difficulties. For example, in the Dalitz plot both ξ and λ_+ are most sensitive to the low E_{π} region, which happens to be the least populated area; hence, the parameters are determined with a very small fraction of the data. Many experiments are sensitive only to part of the Dalitz plot. Others fit only one projection of it, which may yield more than one solution.

In conclusion, due to various experimental difficulties and inhomogeneous treatment of the data by the various experiments, in this edition we continue our past policy of not quoting average values, and limit our review to a listing of experimental data. Again, as in the past, the table does not include any values for λ and ξ .

For a recent review on $K_{\mu 3}$ form factors, see Gaillard and Chounet.¹

Stable Particles

For notation, see illustrated key at beginning of data card listings.

HERZO 69 PR 186 1403 +BANNER, REIER, BERTRAM, EDWARDS + (ILL)
 LORKOWICZ 69 PR 185 1676 +MELISSINOS, NAGASHIMA, TEWKSBURY+ (ROCH, BNL)
 ALSO 66 PRL 17 568 LORKOWICZ + (I)
 MACEK 69 PRL 22 32 MACEK, MANN, MC FARLANE, ROBERTS+ (PENN, TEMPLE)
 MAST 69 PR 183 1200 +GERSHIN, ALSTON-GARJOST, BANGERTER+ (LRL)
 ZFLLER 69 PR 182 1420 ZELLER, HADDOCK, HELLAND, PAHL+ (UCLA, LRL)

AUFERT 70 LAL-123R (KFEV) +HEUSSE, PASCAU, VIALLE, VILAIN+ (ORSA+BRUX)
 BOITTFEL 70 PL 318 325 +BRDWN, CLEGG, CORBETT, CULLIGAN+ (OXF)
 CLARK 70 UCL 20079 +CORK, ELIOFF, KERTH, MCKEYNOlds+ (LRL)
 FORD 70 PRL 25 1370 +PIROU, REMMEL, SMITH, SQUDER (PRIN)
 GRAUMAN 70 PR D1 1277 +KOLLER, TAYLOR, PANDOULAS+ (STEV, SETON, LEHI)
 ALSO 69 PRL 23 737 +KOLLER, TAYLOR, PANDOULAS+ (STEV, SETON, LEHI)

HOFFMAST 70 SIT-P25A HOFFMASTER, KOLLER, TAYLOR+ (STEV+SHU+LEHI)
 KLEMS 70 PRL 24 1086 +HILDFRAND, STIENING (LRL, CHIC)
 KLEMS2 70 PRL 25 473 KLEMS, HILDFRAND, STIENING (LRL, CHIC)
 MACEK 70 PR D1 1249 +MANN, MCFARLANE, ROBERTS (PENN)
 PANDOUILA 70 PR D2 1205 +TAYLOR, KOLLER, GRAUMAN + (STEV, SETON)

HAINT 71 PR D3 10 AACHEN+BARI+CERN+EP+NIJMEGEN+ORSAY+PADVA+
 ALSO 69 PL 298 691 +AACH, BARI, CERN, EPOL, NIJM, ORSAY, PADO, TORI
 OTT 71 PR D3 52 OTT, PRITCHARD (CMCL)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS

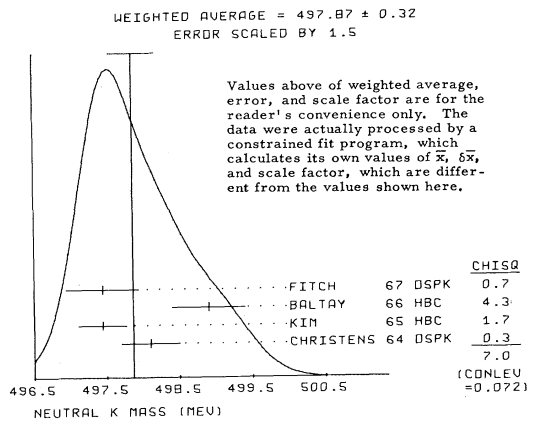
BLOCK 62 CERN CONF 371 BLOCK, LENDINARA, MONARI (INU+BOLOGNA)
 PAPERS NOT REFERRED TO IN DATA CARDS

BRENE 61 NP 22 553 BRNE, EGARDT, QVIST (NORD)
 BIRGE 63 PRL 11 35 BIRGE, ELY, GIDAL, CAMERINI + (LRL+WIS+BARL)
 ADAIR 64 PL 12 67 ADAIR, LEIPUNER (YALE, BNL)
 CARIBBO 64 PL 9 352 CARIBBO, MAKSYMOWICZ (CERN)
 ALSO 64 PL 11 340 CARIBBO, MAKSYMOWICZ (CERN)
 ALSO 65 PL 14 72 CARIBBO, MAKSYMOWICZ (CERN)
 CARIBBO 66 BERKELEY CONF 33 CARIBBO (CERN)
 GINSBERG 67 PR 162 1570 EDWARD S GINSBERG (U. MASS BOSTON)
 WILLIS 67 HEIDELBERG 273 W J WILLIS +RAPPORTEUR TALK (YALE)
 CRONIN 58 VIENNA CONF 241 RAPPORTEUR TALK (PRINCETON)
 HAINT 2 69 PL 298 696 +STEIN+AACH+BAR+CERN+EPOL+NIJM+ORS+PA+TO
 FEARING 70 PR D3 542 +FISCHBACK, SMITH (SUNY+ROHR)
 GINSBERG 70 PR D1 1229 E S GINSBERG (MIT HAIFA)

K⁰
 11 NEUTRAL K (JP=0-) I=1/2
 11 KO MASS (MEV)

M	498.1	0.4	CHRISTENS 64 DSPK		6/66
M	2223	497.44	0.33	KIM 65 HRC	KO FROM PBAR P 6/66
M	4500	498.9	0.5	BALTAY 56 HBC	KO FROM PBAR P 6/66
M		497.44	0.50	FITCH 67 DSPK	11/67
M	AVG	497.87	0.32	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)	
M	FIT	497.79	0.15	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

(SEE IDEOGRAM BELOW)



11 KO-K CH. MASS DIFFERENCE (MEV)

D	3.9	0.6	ROSENFELD 59 HBC	-	
D	6.4	1.1	CRAWFORD 59 HBC	+	
D	9	0.90	BURNSTEIN 65 HRC	-	
D	7	3.71	0.35	KIM 65 HBC	- K+ P TO KO N 4/68
D	417	3.95	0.21	HILL 68 DBC	+ K+D TO KOPP 3/68
D	AVG	3.92	0.14	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
D	FIT	3.95	0.13	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

REFERENCES
 11 NEUTRAL K (JP=0-) I=1/2

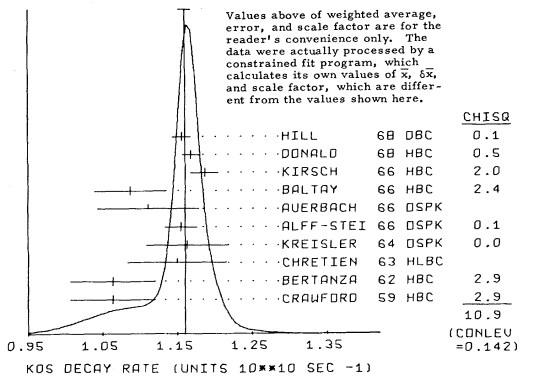
CRAWFORD 59 PRL 2 112 CRAWFORD, CPESTI, GOOD, STEVENSON, TICH0 (LRL)
 ROSENFEL 59 PRL 2 110 A H ROSENFELD, F. SOLMITZ, R D TRIPP (LRL)
 CHRISTEN 64 PRL 13 138 CHRISTENSON, CRONIN, FITCH, TURLAY (PRINCETON)
 BURNSTEIN 65 PR 138 B 895 R A BURNSTEIN, H A RUBIN (MARYLAND)
 KIM 65 PR 140 B 1334 J K KIM, L KIRSCHAD MILLER (COLUMBIA)
 BALTAY 66 PR 142 932 BALTAY, SANDWEISS, STONEHILL + (YALE+BNL)
 FITCH 67 PR 164 1711 FITCH, ROTH, RUSS, VERNON (PRINCETON)
 HILL 68 PR 168 1734 HILL, ROBINSON, SAKITT, CANTER (BNL, CARNEGIE)

K⁰
 12 SHORT-LIVED NEUTRAL K (498, JP=0-) I=1/2
 12 KOS LIFETIME (UNITS 10**10 SEC)

T	0	90	(1.07)	(0.13)	(0.13)	HOLDT 58 CC	
T	512	0.94	0.05	0.05	CRAWFORD 59 HBC		
T	0	63	(1.09)	(0.18)	(0.15)	BOWEN 50 CC	
T	0	OLD	EXPTS WITH LOW STATISTICS NOT INCLUDED IN AVERAGE.				6/68
T	378	0.94	0.05	0.05	BERTANZA 62 HRC		
T	503	0.97	0.05		CHRETIEN 63 HLBC		
T	545	0.86	0.04		KREISLER 64 DSPK		
T		0.866	0.016		ALFF-STEI 66 DSPK	9/66	
T	572	0.90	0.06	0.05	AUERBACH 66 DSPK	8/67	
T	4500	0.92	0.04		BALTAY 66 HBC	5/66	
T	B	(0.904)	(0.024)		BOTT-BODE 66 DSPK	9/66	
T	5000	0.843	0.013		KIRSCH 66 HBC	6/66	
T	19994	0.856	0.009		DONALD 68 HBC	6/68	
T	20000	0.865	0.009		HILL 68 DBC	6/68	
T	B	KOS LIFETIME NOT THE PRIMARY QUANTITY MEASURED IN THIS EXPT.				6/68	
T	H	HILL 68 GIVES A DETAILED DISCUSSION OF SYSTEMATICS ENCOUNTERED					
T	H	IN THIS TYPE OF EXPERIMENT.					
T	AVG	0.8619	0.0062	0.0062	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.2)		
T	FIT	0.8619	0.0058	0.0058	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)		

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 1.1602 ± 0.0084
 ERROR SCALED BY 1.2

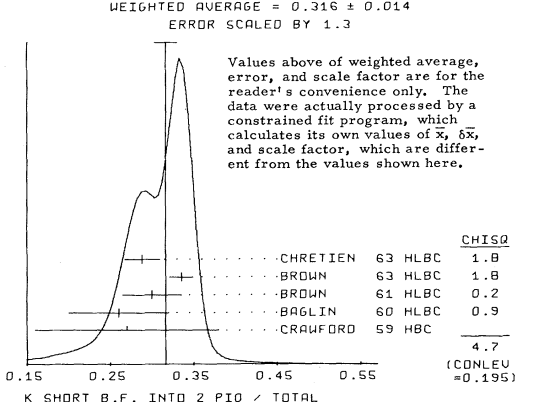


12 KOS PARTIAL DECAY MODES

P1	KOS INTO PI+ PI-	139+ 139
P2	KOS INTO PI0 PI0	134+ 134
P3	KOS INTO MU+ MU-	105+ 105
P4	KOS INTO E+ E-	+5 +5
P5	KOS INTO PI+ PI- GAMMA	139+ 139+ 0

12 KOS BRANCHING RATIOS

R1	KOS INTO (PI+ PI-)/TOTAL	(PI)/TOTAL		
R1	0.68	0.04	CRAWFORD 59 HBC	
R1	0.70	0.08	COLUMBIA 60 HBC	
R1 U	(0.7401)	(0.024)	ANDERSON 62 HBC	
R1 U 1648	0.684	0.011	DOYLE 69 HBC	PI-P TO LAM, KO 2/71*
R1 U	ANDERSON RESULT NOT PUBLISHED, EVENTS ADDED TO DOYLE SAMPLE 2/71*			
R1	AVG	0.684	0.011	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R1	FIT	0.6865	0.0045	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)
R2	KOS INTO (PI0 PI0)/TOTAL	(PI0)/TOTAL		
R2	0.27	0.11	CRAWFORD 59 HBC	
R2	0.26	0.06	BAGLIN 60 HLBC	
R2	0.30	0.035	BRODWN 61 HLBC	
R2	1066	0.335	0.014	BRDWN 63 HLBC
R2	198	0.288	0.021	CHRETIEN 63 HLBC
R2	AVG	0.316	0.014	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R2	FIT	0.3135	0.0045	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)



Stable Particles

For notation, see illustrated key at beginning of data card listings.

Table of particle data listings including fields for KOL INTO, UNITS, (P6)/(P2+P3+P4), and various numerical values for different particle types and experiments.

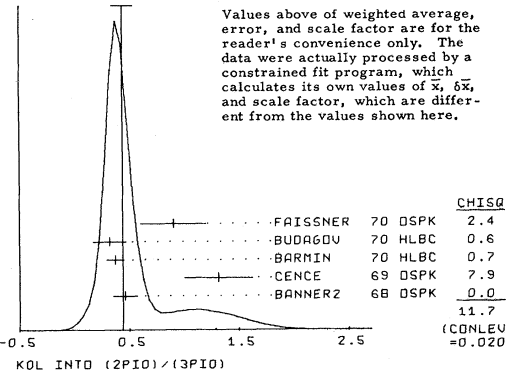


Table listing CHISO values for experiments: FAISSNER 70 DSPK 2.4, BUDAGOV 70 HLBC 0.6, BARMIN 70 HLBC 0.7, CENCE 69 DSPK 7.9, BANNERZ 68 DSPK 0.0.

Continuation of particle data listings with fields for UNITS 10**=3, (P5)/(P3+P4), and numerical values for experiments like DEBOUARD and FITCH.

Table of particle data listings for experiments like FOETH and CLARK, including fields for UNITS 10**=5 and (P7)/(P5).

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± delta P_i, where delta P_i = sqrt(delta P_i delta P_i), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of branching fractions and correlation coefficients for P1 through P11.

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., G_i = Gamma_i / Gamma_total P_i, in appropriate units. In analogy to the matrix above, the diagonal elements are G_i ± delta G_i, where delta G_i = sqrt(delta G_i delta G_i), while the off-diagonal elements are the normalized correlation coefficients (delta G_i delta G_j) / (delta G_i delta G_j). Note that, because of the error in Gamma_total, the errors and correlations here are not directly derivable from those above.

Matrix of decay mode rates and correlation coefficients for G1 through G11.

13 NEUTRAL K ENERGY DEPENDENCE OF DALITZ PLOT

Table listing MATRICELEMENT SQUARED = 1 + G (S3-S0)/(MPI**2) and GTO LINEAR ENERGY DEPENDENCE (G1 FOR TAU DECAYS) with various experimental data points.

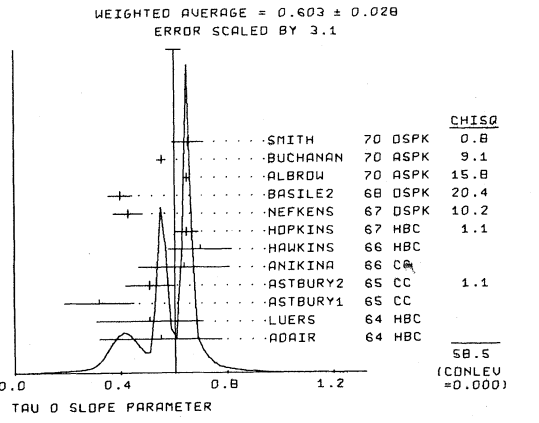
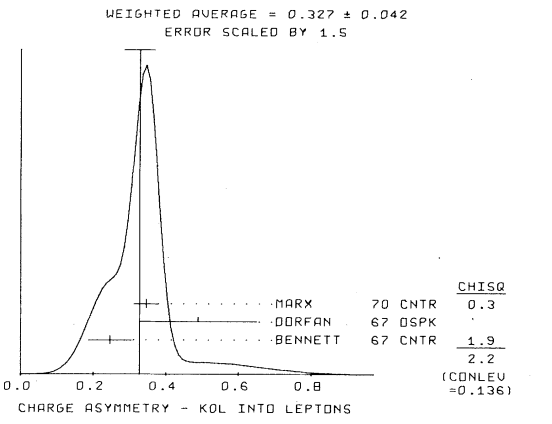


Table listing CHISO values for experiments: SMITH 70 DSPK 0.8, BUCHANAN 70 ASPK 9.1, ALBROW 70 ASPK 15.8, BASILEZ 68 DSPK 20.4, NEFKENS 67 DSPK 10.2, HOPKINS 67 HBC 1.1, HAWKINS 66 HBC, ANIKINA 66 CC, ASTBURY2 65 CC 1.1, ASTBURY1 65 CC, LUERS 64 HBC, ADAIR 64 HBC.

Stable Particles

For notation, see illustrated key at beginning of data card listings.

13 CP VIOLATION PARAMETERS IN KOL DECAYS
 SEE TEXT FOR DEFINITIONS OF CP VIOLATION PARAMETERS
 -----13 CHARGE ASYMMETRY IN TAU DECAYS-----
 TEXT SECTION IV E.2
 SFF SCRIBANO 70 FOR DEFINITION (THIS SIGMA+-1, A=1 FOR MAX ASYMMETRY
 (M)**2 = 1* SIG+- (2/SQRT(3)) * ((T+)-(T-))/ TMAX) AS SCRIBANO 70
 A DECAY ASYMMETRY PARAMETER FOR PI+ PI- PI0 (UNITS 10**(-2))
 A .3M 1.8 OR LESS CL=.90 BLANPIED 68 CNTR 0 CAL. BY SCRIBANO. 1/71*
 A .3M 0.6 OR LESS CL=.90 SCRIBANO 70 CNTR 0 1/71*
 A 4400 8.2 OR LESS CL=.90 SMITH 70 DSPK 0 CALCULATED BY US. 1/71*
 -----13 CHARGE ASYMMETRY IN LEPTONIC DECAYS (PERCENT)-----
 TEXT SECTION IV F.2
 SUCH ASYMMETRY VIOLATES CP . IT IS RELATED TO REAL(EPSILON).
 A1 KOL INTO (MU+PI-NU)-(MU-PI+NU)/(MU+PI-NU)+(MU-PI+NU) 11/67
 A1 D 1M (0.403) (0.134) DORFAN 67 DSPK DERIVED FROM R16 11/67
 A1 D 1M 0.490 0.160 DORFAN 67 DSPK 10/70*
 A1 D THE 2ND VALUE LISTED FOR DORFAN 67 RESULTS FROM A REANALYSIS OF THAT
 A1 EXPT. SFE H.A.PACIOTTA, THESIS, U C BERKELEY.
 A2 KOL INTO (E+PI-NU)-(E-PI+NU)/(E+PI-NU)+(E-PI+NU)
 A2 B 10M (0.224) (0.034) BENNETT 67 CNTR SFF NOTE B 11/67
 A2 B 10M 0.246 0.053 BENNETT 67 CNTR 10/70*
 A2 10M 0.346 0.033 MARX 70 CNTR 10/70*
 A2 B THE 2ND VALUE LISTED FOR BENNETT 67 RESULTS FROM A REANALYSIS OF
 A2 THAT EXPT. SEE H.SAAL, THESIS, COLUMBIA UNIV.
 A2 AVG 0.327 0.043 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
 AL KOL INTO ((L+)-(L-))/((L+)+(L-)) (COMBINED A1 AND A2)
 AL 10M 0.246 0.059 BENNETT 67 CNTR 2/71*
 AL D 1M 0.49 0.16 DORFAN 67 DSPK 2/71*
 AL 10M 0.346 0.033 MARX 70 CNTR 2/71*
 AL AVG 0.327 0.042 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
 (SEE IDOGRAM BELOW)



-----13 PARAMETERS FOR KOL INTO 2PI DECAY-----
 TEXT SECTION IV E.2
 ETA+- = A(KL TO PI+PI-)/(AKS TO PI+PI-)
 ETA00 = A(KL TO PI0PI0)/(AKS TO PI0PI0)
 THE MAGNITUDES OF ETA+- AND OF ETA00 ARE DERIVED FROM BR. RATIOS.
 FOR THE QUANTITIES MEASURED BY THE INDIVIDUAL EXPERIMENTS SEE LISTINGS
 OF S13R9 AND S13R2 (ETA+-) AND OF S13R17 AND S13R19 (ETA00).
 FOR THE READER'S CONVENIENCE WE LIST HERE THE DERIVED QUANTITIES ETA+-
 (CALLED E+- BELOW) AND (ETA00)**2 (CALLED EOS BELOW)
 EOS (ETA00)**2 = (A(KL TO 2PI0)/(AKS TO 2PI0))**2 (UNITS 10**(-6))
 EOS 58 5.06 1.4 BANNER 68 DSPK 10/69
 EOS 9 -2. 7.0 BARTLETT 68 DSPK 10/69
 EOS F 190 (13.) (4.) GAILLARD 69 DSPK 10/69
 EOS 133 14.1 3.4 GENCE 69 DSPK 10/69
 EOS 29 4.09 0.9 BARMAN 70 HLBC 12/70*
 EOS 30 3.61 1.9 BUDAGOV 70 HLBC 10/70*
 EOS 11.03 4.3 CHOLLET 70 DSPK CU REG.+4 GAMMAS 10/70*
 EOS F 172 9.9 3.4 FAISSNER 70 DSPK 12/70*
 EOS F FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69
 EOS AVG 4.0 1.0 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
 EOS FIT 5.0 1.0 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.5) 2/71*
 (SEE IDOGRAM BELOW)
 E+- ETA+- = A(KL TO PI+PI-)/(AKS TO PI+PI-) UNITS 10**(-3)
 E+- 45 (1.94) CHRISTENS 64 DSPK 10/69
 E+- 4 (2.02) GALBRAITH 65 DSPK 10/69
 E+- (1.66) BALE 66 DSPK 10/69
 E+- (1.935) BOTT-BODE 66 DSPK 10/69
 E+- 525 1.91 .06 FITCH 67 DSPK 10/69
 E+- FIT 1.95 0.03 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/71*
 ER RATIO OF ETA00 OVER ETA+-
 ER 1.09 0.09 RUBBIA 70 CNTR PRELIMINARY 2/71*
 F00 PHASE OF ETA 00 (DEGREES)
 F00 FIRST QUADRANT PREFERRED GORBI 69 DSPK 11/69
 F00 51. 30. CHOLLET 70 DSPK CU REG.+4 GAMMAS 10/70*

WEIGHTED AVERAGE = 4.9 ± 1.0
 ERROR SCALED BY 1.6

 PHASE OF ETA +- (DEGREES)
 DM IS (KOL-KOS MASS DIFFERENCE / HBAR) IN UNITS OF 10**10 SEC-1
 SEE SECTION D OF KOL LISTINGS FOR LATEST VALUE
 WE HAVE ADDED THE MASS DEPENDENCE AND PROPAGATED THE ERROR IN DM
 USING DM=0.5398+-0.0033 FOR BENNETT 69, BOHM 69, FAISSNER 69, AND
 JENSEN 70.
 F+- 45.0 50.0 FITCH 65 DSPK BE REGEN 11/67
 F+- 30.0 45.0 FIRESTONE 66 HBC 11/67
 F+- 70.0 21.0 BOTT-BODE 67 DSPK C REGEN 11/67
 F+- 25.0 35.0 MISCHKE 67 DSPK CU REGEN 7/68
 F+- N (51.0) (11.0) BENNETT 68 CNTR CU REG. USES 7/68
 F+- C 34.5 10.0 BENNETT 69 CNTR CU REGEN 2/71*
 F+- B 47.6 12.1 BOHM 69 DSPK VACUUM REGEN 2/71*
 F+- F 46.2 7.4 FAISSNER 69 ASPK CU REGEN 2/71*
 F+- J 43.2 4.4 JENSEN 70 ASPK VACUUM REGEN 2/71*
 F+- AVG 43.6 3.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 COMMENTS
 F+- C BENNETT 69 USES MEASUREMENT OF (F+-)-(PHI) OF ALF-STFI 66 2/71*
 F+- C BENNETT 69 F+- = 36.9+-10.0, NOT INCLUDING ERROR IN DM 2/71*
 F+- C DM DEPENDENCE OF BENNETT 69 IS 69*(DM-0.545) DEG 2/71*
 F+- B 30HM 69 F+-=41+-12, NOT INCLUDING ERROR IN DM 2/71*
 F+- B DM DEPENDENCE OF BOHM 69 IS 479*(DM-0.526) DEG 2/71*
 F+- F FAISSNER 69 ERROR ENLARGED TO INCLUDE ERROR IN REGENERATOR PHASE 11/69
 F+- F FAISSNER 69 F+-=49.3+-7.4, NOT INCLUDING ERROR IN DM 2/71*
 F+- F DM DEPENDENCE OF FAISSNER 69 IS 205*(DM-0.555) DEG 2/71*
 F+- J JENSEN 70 F+-=42.2+-4.0, NOT INCLUDING ERROR IN DM 2/71*
 F+- J DM DEPENDENCE OF JENSEN 70 IS 576*(DM-0.538) DEG 2/71*
 F+- N BENNETT 69 IS A REEVALUATION OF BENNETT 68 11/69

Superweak Model Predictions

The superweak model of Wolfenstein, Phys. Letters 13, 562 (1964) predicts that

$$\phi_{+-} = \phi_{00} = \tan^{-1} \left(\frac{2\Delta m \tau_S}{\hbar} \right)$$

and

$$\text{Re } \epsilon = |\eta_{+-}| \left[1 + \left(\frac{2\Delta m \tau_S}{\hbar} \right)^2 \right]^{-1/2}$$

The $K_L^0 - K_S^0$ mass difference, the K_S lifetime, and $|\eta_{+-}|$ given in these tables result in the predictions that

$$\phi_{+-} = \phi_{00} = (42.94 \pm 0.26)^\circ$$

and

$$\text{Re } \epsilon = (1.428 \pm 0.023) \times 10^{-3}$$

These can be compared with the experimental values

$$\phi_{+-} = (43.6 \pm 3.3)^\circ$$

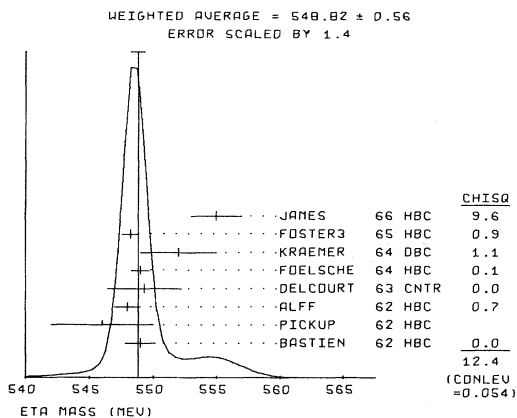
$$\phi_{00} = (51 \pm 30)^\circ$$

$$\text{Re } \epsilon = (1.64 \pm 0.24) \times 10^{-3}$$

where ϵ has been computed from δ , the charge asymmetry parameter for leptonic K_L^0 decays, and $(\text{Re } \epsilon, \text{Im } \epsilon)$, the $\Delta S = -\Delta Q$ amplitude, using Eq. (29) of the text.

Stable Particles

For notation, see illustrated key at beginning of data card listings.



14 ETA PARTIAL DECAY MODES

Mode	Decay Masses	CHISQ	
P1	ETA INTO 2GAMMA	0+ 0	
P2	ETA INTO 3PIO	134+ 134+ 134	
P3	ETA INTO PI+ PI- PIO	139+ 139+ 134	
P4	ETA INTO PI+ PI- GAMMA	139+ 139+ 0	
P5	ETA INTO E+E-PIO	134+ .5+ .5	
P6	ETA INTO E+E-PI+PI-	139+ 139+ .5+ .5	
P7	ETA INTO PIO 2GAMMA	134+ 0+ 0	
P8	ETA INTO E+GAMMA	.5+ .5+ 0	
P9	ETA INTO 2PIO GAMMA	134+ 134+ 0	
P10	ETA INTO PI+PI-PIO GAMMA	139+ 139+ 134+ 0	
P11	ETA INTO PI+PI- 2GAMMA	139+ 139+ 0+ 0	
P12	ETA INTO MU+MU-	105+ 105	
P13	ETA INTO MU+MU-GAMMA	105+ 105+ 0	
P14	ETA INTO MU+MU-PIO	105+ 105+ 134	

14 ETA DECAY RATES

Mode	Units (keV)	Rate	Notes
W1	ETA INTO 2GAMMA	10.931	(0.2)
W1	ETA INTO 2GAMMA	10.931	(0.2)
BEMPRAD	67 CNTR		
PRINAKOFF EFFECT	11/67		

The above value for $\Gamma_{\gamma\gamma}$ assumes that $\Gamma_{\gamma\gamma}/\Gamma_{\text{total}} = 31.4\%$. However, the results of that experiment may be stated more generally than is given in the paper, as

$$\Gamma_{\gamma\gamma} \times \frac{\Gamma_{\gamma\gamma}}{\Gamma_{\text{total}}} = 0.380 \pm 0.083 \text{ keV}$$

(private communication from C. Bemporad). Thus our new value of

$$\Gamma_{\gamma\gamma}/\Gamma_{\text{total}} = 38.6 \pm 1.1\%$$

would give

$$\Gamma_{\gamma\gamma} = 0.99 \pm 0.22 \text{ keV}$$

and

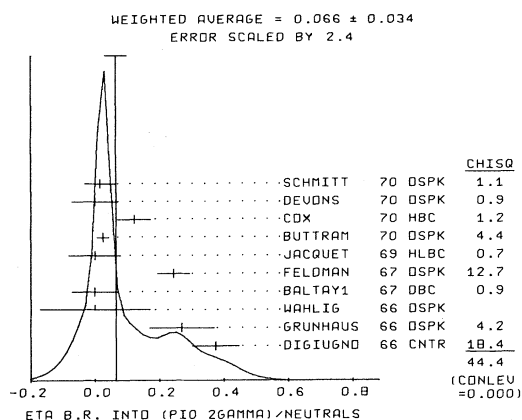
$$\Gamma_{\text{total}} = 2.63 \pm 0.59 \text{ keV.}$$

ETA DECAY INTO NEUTRALS

As is well known, there are great inconsistencies among the various experiments that report etas decay into neutrals. The controversy is over whether the mode $\eta \rightarrow \pi^0\gamma\gamma$ is ≈ 0 (as some experiments indicate) or $\geq 20\%$ (as other experiments indicate).

The discrepancies are displayed in the ideogram below, in which all ten relevant experiments have been converted to a common ratio, $\pi^0\gamma\gamma/\text{neutrals}$. Also upper limits, $< x$, have been converted to $0 \pm x$. The confidence level for consistency of all ten is $< 10^{-5}$!

In previous editions we were able to point out that it was the older experiments which gave the $> 20\%$ values, whereas the more recent experiments had been giving $\leq 1\%$. However, the recent experiment of COX 70 gives about $8 \pm 3\%$, so the controversy is not yet settled.



We feel that we should consider all ten experiments on an a priori equal basis, and then follow the prescription of deleting large χ^2 experiments until the confidence level rises to some reasonable value. If we remove the Feldman and DiGiugno experiments, χ^2 decreases from 45 (for all ten) to about 9 (for the remaining 8). Accordingly we have removed those 2 experiments and used the remaining eight experiments in our overall fit.

14 ETA BRANCHING RATIOS

Mode	Value	Experiment	Notes
R1 N	10 (2.5)	(1.0)	PICKUP 62 HBC
R1 N	53 (3.20)	(1.26)	BASTIEN 62 HBC
R1 N	(2.7)	(0.8)	SHAFER 62 HBC
R1	2.6	.9	BUSCHBECK 63 HBC
R1 N	290 (4.5)	(1.0)	JAMES 66 HBC
THESE EXPERIMENTS HAVE NOT BEEN USED IN COMPUTING THE AVERAGES AS THEY WERE UNABLE TO CLEARLY SEPARATE PARTIAL MODES (3) AND (4) FROM EACH OTHER. THE REPORTED VALUES THUS PROBABLY CONTAIN SOME (UNKNOWN) FRACTION OF MODE (4).			
R1	2.64	0.23	BALTAY2 67 DBC
R1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)		
R1	2.64	0.22	
R1	2.59	0.15	
R2	ETA INTO 2GAMMA/CHARGED	(P11)/(P3+P4)	
R2	0.99	0.48	CRAWFORD 63 HBC
R2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1) FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R2	1.387	0.084	

Stable Particles

For notation, see illustrated key at beginning of data card listings.

R3 ETA INTO (P10 2GAMMA)/NEUTRALS (P7)/(P1+P2+P7)
R3 S (0.375) (0.072) DIGIUGNO 66 CNTR ERROR DOUBLED 6/66
R3 THE ERRORS OF DIGIUGNO+ 66 HAVE BEEN INCREASED BY A FACTOR
R3 OF TWO, TO TAKE INTO ACCOUNT POSSIBLE SYSTEMATIC ERRORS, AS
R3 SUGGESTED BY THE AUTHORS.
R3 R .27 .10 GRUNHAUS 66 OSPK 8/67
R3 S (.244) (.05) FELDMAN 67 OSPK 8/67
R3 S SFE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.
R3 .026 .019 RUITRAM 70 OSPK 12/70*
R3 .122 .052 .044 COX 70 HBC 6/70
R3 (.07) OR LESS DEVONS 70 OSPK .9 CONF.LEVEL 12/70*
R3 R 16 .016 .047 SCHMITT 70 OSPK 12/70*
R3 P SCHMITT 70 IS A REANALYSIS BUNIA TOV 67 12/70*
R3 F (0.11) (0.03) STRUGALSK 70 HLBC 2/71*
R3 E THIS MEASUREMENT HAS BEEN EXCLUDED BECAUSE THE ERROR APPEARS
R3 E TO BE SERIOUSLY UNDERESTIMATED 2/71*
R3 AVG .042 0.023 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
R3 FIT 0.045 0.017 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

R4 ETA INTO (P1+ P1- GAMMA)/(P1+ P1- P10) (P4)/(P3)
R4 0.14 0.08 FOELSCH 64 HBC 7/66
R4 M 24 (0.73) (0.25) PAULI 64 DBC 7/66
R4 M THIS EXPERIMENT HAS NOT BEEN INCLUDED IN THE AVERAGES SINCE
R4 M IT IS NOT CLEAR THAT THEIR CLASS R EVENTS ARE ACTUALLY FROM ETAS.
R4 0.30 0.06 CRAWFORD 66 HBC 6/66
R4 .10 .10 KRAEMER 64 DBC 7/66
R4 .196 .041 FOSTER3 65 HBC 7/66
R4 .25 .035 LITCHFIEL 67 DBC 8/67
R4 0.28 0.04 BALTAY2 67 DBC 11/67
R4 7250 .201 .006 GORMLEY 70 ASPK 6/70
R4
R4 AVG 0.2041 0.0079 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
R4 FIT 0.2040 0.0076 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.2041 ± 0.0079
ERRR SCALED BY 1.4

Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{x} , $\delta\bar{x}$, and scale factor, which are different from the values shown here.

R5 ETA INTO (3P10+ 2/3(P10 2GAMMA)/ P1+P1-P10) (P2+2/3P7)/P3
R5 0.83 0.32 CRAWFORD 63 HBC 7/66
R5 2.0 1.0 FOELSCH 64 HBC 7/66
R5 0.90 0.24 FOSTER1 65 HBC 7/66
R5
R5 AVG 0.91 0.19 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R5 FIT 1.406 0.091 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

R6 ETA INTO 3P10/2GAMMA (P2)/(P1)
R6 1.901 OR MORE CHRETIEN 62 PBC 11/67
R6 0.88 0.16 BALTAY1 67 DBC 1/68
R6 1.1 0.2 GENCE 67 OSPK 12/70*
R6 0.75 0.09 DEVONS 70 OSPK 12/70*
R6
R6 AVG 0.824 0.085 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)
R6 FIT 0.785 0.040 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R7 ETA INTO 2GAMMA/(P1+ P1- P0) (P11)/(P3)
R7 1.61 0.39 FOSTER1 65 HBC 7/69
R7 401 1.72 .25 BAGLIN 69 HLBC 7/69
R7
R7 AVG 1.69 0.21 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R7 FIT 1.67 0.10 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R8 ETA INTO NEUTRAL/(P1+ P1- P10) (P1+P2+P7)/(P3)
R8 280 3.6 0.8 KRAEMER 64 DBC 7/66
R8 3.8 1.1 PAULI 64 DBC 9/66
R8 2.89 0.56 ALFF-STEI 66 HBC 1/68
R8 244 3.6 0.6 FLATTEZ 67 HBC 1/68
R8
R8 AVG 3.35 0.35 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R8 FIT 3.12 0.18 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R9 ETA INTO (E+E-P10)/(P1+P1-P10) (UNITS 10**=-2) (P5)/(P3)
R9 (1.1) OR LESS PRICE 65 HBC 8/67
R9 0 (0.77) OR LESS FOSTER2 65 HBC .9 CONF.LEVEL 11/67
R9 (1.42) OR LESS BAGLIN1 67 HLBC .9 CONF.LEVEL 11/67
R9 0 (1.6) OR LESS BILLING 67 HLBC .9 CONF.LEVEL 11/67

R10 ETA INTO (E+E-P1+P1-1)/TOTAL (UNITS 10**=-2) (P6)/TOTAL
R10 (0.7) OR LESS RITTENBER 65 HBC 6/66

R11 ETA INTO (E+E-P1+P1-1)/(P1+P1-GAMMA) (P6)/(P4)
R11 1 0.026 0.026 GROSSMAN 66 HBC 6/66

R12 ETA INTO 2 GAMMA/NEUTRALS (P11)/(P1+P2+P7)
R12 S (0.416) (0.044) DIGIUGNO 66 CNTR ERROR DOUBLED 6/66
R12 .44 .07 GRUNHAUS 66 OSPK 8/67
R12 S (.579) (.052) FELDMAN 67 OSPK 8/67
R12 S SFE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.
R12 T THIS RESULT FROM COMBINING CROSS-SECTIONS FROM TWO DIFFERENT EXPTS.
R12 .59 .033 BUNIA TOV 67 OSPK 11/67
R12 .535 .018 RUITRAM 70 OSPK 12/70*
R12 .486 .036 COX 70 HBC 6/70
R12 0.57 0.09 STRUGALSK 70 HLBC 2/71*
R12
R12 AVG 0.535 0.018 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R12 FIT 0.535 0.013 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R13 ETA INTO 3P10/NEUTRALS (P21)/(P1+P2+P7)
R13 S (0.209) (0.084) DIGIUGNO 66 CNTR ERROR DOUBLED 6/66
R13 S (1.29) (.10) GRUNHAUS 66 OSPK 8/67
R13 S (.177) (.035) FELDMAN 67 OSPK 8/67
R13 S SEE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.
R13 .41 .033 BUNIA TOV 67 OSPK 11/67
R13 R REDUNDANT INFORMATION FROM THIS EXPERIMENT
R13 R (.439) (.024) RUITRAM 70 OSPK 12/70*
R13 .392 .042 COX 70 HBC 6/70
R13 0.32 0.09 STRUGALSK 70 HLBC 2/71*
R13
R13 AVG 0.397 0.025 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R13 FIT 0.420 0.016 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R14 ETA INTO P10 (2GAMMA)/2GAMMA (P71)/(P1)
R14 (.5) OR LESS WAHLIG 66 SPRK .9 CONF LEVEL 7/66
R14 0.0 0.14 BALTAY1 67 DBC 11/67
R14 P (0.05) (0.04) BONAMY 67 SPRK PRELIMINARY RESULT 11/67
R14
R14 FIT 0.085 0.033 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

R15 ETA INTO (E+E-P10)/TOTAL (UNITS 10**=-2) (P5)/TOTAL
R15 (0.7) OR LESS RITTENBER 65 HBC 6/66
R15 (0.084) OR LESS RAZIN 58 DBC .9 CONF LEVEL 6/68

R16 ETA INTO 2GAMMA/(3P10 + P10 2GAMMA) (P11)/(P2+P7)
R16 0.80 .25 BACCI 63 CNTR 7/66
R16
R16 FIT 1.149 0.062 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R17 ETA INTO (P1+P1-P10 GAMMA)/(P1+P1-P10) (P10)/(P3)
R17 (1.07) OR LESS FLATTE 67 HBC 8/67
R17 (.009) OR LESS PRICE 67 HBC 8/67
R17 (0.016) OR LESS BALTAY2 67 DBC .95 CONF LEVEL 11/67
R17 (0.017) OR LESS ARNOLD 68 HLBC .9 CONF LEVEL 9/68

R18 ETA INTO (P1+P1- 2GAMMA)/(P1+P1-P10) (P11)/(P3)
R18 (.009) OR LESS PRICE 67 HBC 8/67
R18 (0.016) OR LESS BALTAY2 67 DBC .95 CONF LEVEL 11/67

R19 ETA INTO 3P10/(P1+ P1- P10) (P21)/(P3)
R19 1.3 .4 BAGLIN2 67 HLBC 8/67
R19 1.47 0.20 0.17 BULLOCK 68 HLBC 9/68
R19 199 1.50 .15 .29 BAGLIN 69 HLBC 7/69
R19
R19 AVG 1.46 0.13 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R19 FIT 1.312 0.083 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R20 ETA INTO 2GAMMA/(1/3P10+2/3(P10 2GAMMA)) (P11)/(P2+2/3P7)
R20 1.10 0.5 MULLER 63 DBC 7/66
R20
R20 FIT 1.188 0.059 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R21 ETA INTO NEUTRALS/TOTAL (P1+P2+P7)/TOTAL
R21 .79 .08 BUNIA TOV 67 OSPK 11/67
R21
R21 FIT 0.722 0.012 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)

R22 ETA INTO (P120 2GAMMA)/TOTAL (P7)/TOTAL
R22 (1.2) OR LESS JACQUET 69 HLBC .95 CONF.LEVEL 6/70
R22
R22 FIT 0.033 0.012 FROM FIT

R23 ETA INTO MU+MU-/TOTAL (UNITS 10**=-5) (P12)/TOTAL
R23 0 (2.1) OR LESS WEHMANN 68 OSPK .95 CONF.LEVEL 4/69

R24 ETA INTO MU+MU-P10/TOTAL (UNITS 10**=-4) (P14)/TOTAL
R24 (5.1) OR LESS WEHMANN 68 OSPK 4/68

R25 ETA INTO MU+MU-/2GAMMA (UNITS 10**=-5) P(12)/(P1)
R25 5.9 2.2 HYAMS 69 OSPK 7/69

R26 ETA INTO (P10 2GAMMA)/(3P10 + P10 2GAMMA) (P7)/(P2+P7)
R26 N 0.1 0.3 KANDFSKY 70 OSPK 2/71*
R26 N WE HAVE CHANGED THE ERROR ON THIS EXPERIMENT FROM +0.3,-0.1
R26 N TO THE ABOVE +0.3,-0.3 SINCE IT IS CLEAR FROM FIGURE 7 IN
R26 N THE ARTICLE THAT A CENTRAL VALUE OF 0.0 IS ABOUT AS PROBABLE
R26 N AS THE QUOTED VALUE OF 0.1 2/71*
R26
R26 FIT 0.097 0.035 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i^2)}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \cdot \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 7
P 1	.384+-.011				
P 2	-.200	.303+-.011			
P 3	-.444	-.234	.231+-.010		
P 4	-.378	-.202	.700	.047+-.002	
P 7	-.351	-.554	-.179	-.135	.033+-.011

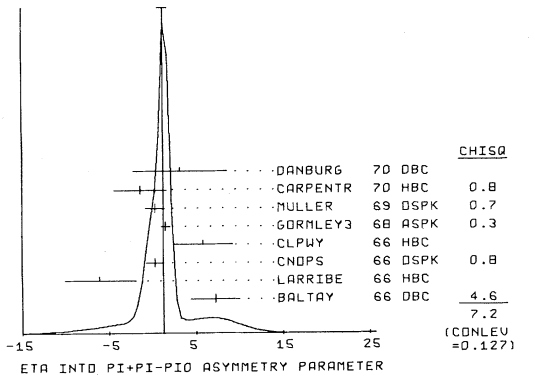
Stable Particles

For notation, see illustrated key at beginning of data card listings.

14 ETA C-NONCONSERVING DECAY PARAMETER
A DECAY ASYMMETRY PARAMETER FOR PI+ PI- P10 (UNITS 10**-2)
A 1351 7.2 2.8 BALTAY 66 DBC 8/66

A AVG 1.21 0.54 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 1.21 +/- 0.54
ERROR SCALED BY 1.3



H. Yuta and S. Okubo [Phys. Rev. Letters 21, 781 (1968)] have pointed out that an asymmetry in the decay eta -> pi+ pi- pi0 of about 2% need not imply a breakdown of C invariance, since an asymmetry of this amount could be caused by an interference between the eta and the 3pi background. Gormley et al. [Phys. Rev. Letters 22, 198 (1969)], however, believe that this effect can account for only <= 0.23% in their experiment (above). Also see: A. Frenkel and G. Vesztegombi, "C-Violation in eta-Decay," Nucl. Phys. B15, 429 (1970) and K. Taggart, "Asymmetry and Background in eta -> 3pi," Phys. Rev. D 2, 1960 (1970).

B DECAY ASYMMETRY PARAMETER FOR PI+ PI- GAMMA (UNITS 10**-2)
B 33 -2. 17. CRAWFORD 66 HBC 11/66
B 4 -4. 8. LITCHFIELD 67 DBC 8/67

REFERENCES
14 ETA(549,JPG=0-+1)=0
PEVSNR 61 PRL 7 421 PEVSNR,KRAEMER,NUSSBAUM,RICHARDSON + (JHU)
ALFF 62 PRL 9 322 ALFF,BERLEY,COLLEY,BRUGGER + (COL+RUTGERS)
BASTIEN 62 PRL 8 114 BASTIEN,BERGE,DAHL,FERRI-LUZZI + (LRL)
CHRSTIEN 62 PRL 9 127 CHRSTIEN + (BRAND+BROWN+HARVARD+MIT+ADVO)
PICKUP 62 PRL 8 329 F PICKUP,ROBINSON,SALANT (NRCC+BNL)
SHAFFER 62 CERN CONF 307 J SHAFFER,FERRI-LUZZI,MURRAY + (UC+LRL)

ALFF-STE 66 PR 145 1072 ALFF-STEINBERGER,BERLEY + (COLUMBIA+RUTGERS)
BALTAY 66 PRL 16 1224 BALTAY,FRANZINI,KIM,KIRSCH+(COLUMBIA+STONY BROOK)
CRAWFORD 66 PRL 16 333 F.S.CRAWFORD,L.R.PRICE (LRL)
DIGIUGNO 66 PRL 16 767 DIGIUGNO+GORMLEY,SILVESTRI+(SNAP+FRST+FRSC)
GROSSMAN 66 PR 146 993 R.GROSSMAN,L.PRICE,F.CRAWFORD (LRL)
GRUNHAUS 66 THESIS J.GRUNHAUS (COLUMBIA)
JAMES 66 PR 142 896 F.E.JAMES,H.L.KRAYBILL (YALE+BNL)
JONES 66 PL 23 597 JONES,BENNIE,BOJANE,HORSEY,MASON,+ (ICL+RUTH)
WAHLIG 66 PRL 17 221 WAHLIG,SHIRATA,MANNELLI (MIT+PIISA)

AGLIN1 67 PL 248 637 BAGLIN,BEZUGUET,DEGRANGE,+ (E.POLY+UC)
AGLIN2 67 BAPS 12 567 BAGLIN,BEZUGUET,DEGRANGE,+ (E.POLY+UC)
BALTAY 67 PRL 19 1495 BALTAY,FRANZINI,KIM,NEWMAN+(COLUMB+BRAND)
BALTAY2 67 PRL 19 1498 BALTAY,FRANZINI,KIM,NEWMAN+(COLUM+STONY BK)
BEMPORAD 67 PL 258 380 BEMPORAD,BRACCINI,FOA,LUBEL,SMEY+(PIISA,BONN)
AND PRIVATE COMMUNICATION
BILLIG 67 PL 258 435 BILLIG,BULLOCK,ESTEN,GOVAN,+ (UCL,OXF)
BONAMY 67 HEIDELBERG CONF. BONAMY,SONDEREGGER (SACLAY)
BUNIATOV 67 PL 258 560 BUNIATOV,ZAVATTINI,DEINET,+ (CERN,KARLS)
GENCE 67 PRL 19 1393 GENCE,PETERSON,STENGER,CHIU+(THAW1+LRL)
FELDMAN 67 PRL 18 968 FELDMAN,FERRI-LITCHFIELD,HALPERN,+ (PENN)
FLATTE 67 PRL 18 976 S.M.FLATTE (LRL)
FLATTE2 67 PR 163 1441 S.M.FLATTE AND C.G.WOHL (LRL)
LITCHFIELD 67 PL 248 486 LITCHFIELD,RANGAN,SEGAR,SMITH+(RUTH+SACLAY)
PRICE 67 PRL 18 1207 L.R.PRICE,F.S.CRAWFORD (LRL)

ARNOLD 68 PL 278 466 +PATY,BAGLIN,BINGHAM+(STRB+MADR+EPDL+BERK)
BAZIN 68 PRL 20 895 BAZIN,GOSHAM,ZACHER,+ (PRINCETON,QUEENS)
BULLOCK 68 PL 278 402 +ESTEN,FLEITING,GOVAN,HENDERSON,OWEN+(LOUC)
WEHMANN 68 PRL 20 748 WEHMANN,ENGELS+(HARV+CASE+SAC+COR+MGILL)

AGLIN 69 PL 298 445 BAGLIN,BEZUGUET,+ (EPDL,BERK,MADR,STRB)
ALSD 70 NP 822 66 +BEZUGUET,DEGRANGE,MUSSET+(EPOL,MADR,STRB)
HYAMS 69 PL 298 123 HYAMS,KOCH,POTTER,VON LINDERN,+ (CERN,PMI)
JACQUET 69 NC 58 743 JACQUET,NGUYEN-KHAC,HAATUF+(EPDL,BERG)

BUTTRAM 70 PRL 25 1358 +KREISLER,MISCHKE (PRIN)
COX 70 PRL 24 534 COX,FORTNEY,GOLESON (DUKE)
DEVONS 70 PR(D) 1 1936 +GRUNHAUS,KOZLOWSKI,NEMETHY+(COLU+SYR)
GORMLEY 70 PR 20 501 GORMLEY,HYMAN,LEE,NASH,PEPPLES+(COLU+BNL)
AND ALSO REVIS 18(L)THESIS MICHAEL GORMLEY (COLU)
KANOFSKY 70 NC 68 413 A. KANOFSKY (LEHI)
SCHMITT 70 PL 328 638 +BUNIATOV,ZAVATTINI,DEINET+(CERN,KARL)
STRUGALS 70 JINR-EL-5256 STRUGALSKI,CHUUILO,GEMESY,+ (JINR)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS
BASTIEN 62 PRL 8 114 BASTIEN,BERGE,DAHL,FERRI-LUZZI,MILLER+(LRL)
ARMONY 62 PRL 8 117 D CARMONY, A ROSENFELD, VAN DE WALLE (LRL)
ROSENFEL 62 PRL 8 293 A ROSENFELD, D CARMONY, VAN DE WALLE (LRL)

REFERENCES ON ETA ASYMMETRY PARAMETERS

BALTAY 66 PRL 16 1224 BALTAY,FRANZINI,KIM,KIRSCH+(COLUM+STONY BK)
CNOPS 66 PL 22 546 CNOPS,FINOCCHIARI,SASSALLE+(CERN+SACL)
CRAWFORD 66 PRL 16 333 F.S.CRAWFORD,L.R.PRICE (LRL)
LARRIBE 66 PL 23 606 LARRIBE,LEVEQUE,MULLER,PAULI,+ (SACL+RUTH)
CLPHY 66 PR 149 1044 COLUMBIA,LRL,PURDUE,WISCONSIN,YALE

16 PROTON (938,J=1/2) I=1/2
16 PROTON MASS (MEV)
M (938.256) (0.005) COHEN 65 RVUE 7/66
M (938.2592) .0052 TAYLOR 69 RVUE USING NEW E/H 7/70

16 PROTON LIFETIME (UNITS 10**26 YR)
Y 10**20 YRS OR MORE GOLDHABE 54 TH 232 FISS.MODE INDEPEN
FLEROV (0.002) OR MORE FLEROV 57 TH 232 FISS.MODE INDEPEN
T B (1.5) OR MORE BACKENSTO 60 CNTR
T B (60.0) OR MORE KROPP 65 CNTR
Y (200.0) OR MORE GURR 67 CNTR DEP. ON DECAY MODE
T B KROPP AND BACKENSTOSS SENSITIVE TO PARTICULAR DECAY MODES OF PROT

16 PROTON MAGNET. MOMENT(E/2MP)
ARMONY (2.792763(0.000030) COHEN 65 RVUE
MM 2.792782 .000017 TAYLOR 69 RVUE USING NEW E/H 7/70

16 PROTON ELECTRIC DIPOLE MOMENT (IN UNITS OF 10**-23 E CM)
NONZERO VALUE IMPLIES VIOLATION OF T AND P IN EM INTERACTION
EDM 16 700. 900. HARRISON 49 MBR 10/69

REFERENCES
16 PROTON (938,J=1/2) I=1/2
GOLDHABE 54 PR 96 1157 FN072 GOLDHABER,F.REINES+(LDS ALAMOS,BNL)
FLEROV 57 SDV PHYS DOK 3 78 FLEROV,KLOCHKOV,SKOBKIN,TEPETEVE (USSR)
BACKENSTO 60 NP 16 749 BACKENSTOSS,FRAUENFELDER,HYMAN+(CERN)
COHEN 65 RMP 37 537 E R COHEN, J W M DUMOND (NAASC+CALTECH)
KRAPP 65 PR 137 B 740 W R KRAPP,F REINES (CASE INST TECHNOLOGY)
GURR 67 PR 158 1321 GURR,KRAPP,REINES,MEYER (CASE,JOHANNESBURG)
HARRISON 69 PRL 22 1263 HARRISON,SANDERS,WRIGHT (CLARENDON OXFORD)
TAYLOR 69 RMP 41 375 +PARKER,LANGENBERG (PRIN+UCI+PENN)

n
17 NEUTRON (939,J=1/2) I=1/2
17 NEUTRON-PROTON MASS DIF.(MEV)
D M 1.29344 0.00007 MATTAUCH 65 RVUE 3/71*
D M WE HAVE CONVERTED MATTAUCH NEUTRON-HYDROGEN MASS DIFFERENCE TO 3/71*
D M NEUTRON-PROTON MASS DIFFERENCE USING CURRENT VALUE OF ELECTRON MASS 3/71*
D M AND A HYDROGEN BINDING ENERGY OF 13.6 EV 3/71*

Stable Particles

For notation, see illustrated key at beginning of data card listings.

17 NEUTRON MAGNETIC MOMENT (MAGNETONS, 938.2 MEV)				
MM	-1.913148	0.000066	COHEN	56 RVUE 7/66

17 NEUTRON ELECTRIC DIPOLE MOMENT (IN UNITS OF 10**+23 E CM) TEST OF C VIOLATION IN THE EN INTERACTION				
FDH	(5.1)	OR LESS	BAIRD	69 MBR 10/69

17 NEUTRON LIFETIME (UNITS 10**+3 SEC)				
THE MEASUREMENT OF THE NEUTRON LIFETIME BY SOSNOVSKII 59 HAS BEEN DISCARDED SINCE 1. IT DISAGREES WITH THE BETTER AND MORE RECENT RESULT OF CHRISTENSEN 67. 2. THE VALUE OF GA/GV DERIVED FROM THE NEW VALUE OF THE LIFETIME AGREES WELL WITH THE GA/GV VALUE OBTAINED FROM THE FREE NEUTRON DATA.				
T	(1.012)	(0.021)	SOSNOVSKI 59 PILE	SEE NOTE E 7/68
E ERROR CHANGED BECAUSE ERROR IN CROSS SECTION FOR NEUTRON ABSORPTION IN GOLD HAS BEEN REDUCED				
T	0.932	0.014	CHRISTENSEN 67 PILE	3/68

17 BETA DECAY COUPLING CONSTANTS				
AV	GA/GV (SEE TEXT FOR SIGN CONVENTION)			
AV B	(-1.18)	(0.02)	BHALLA	66 RVUF 11/67
AV	-1.250	0.064	CONFORTO	67 RVUE SEE NOTE C BELOW
AV	-1.23	0.01	CHRISTENSEN 67 CNTR	SEE NOTE D BELOW
AV B	THIS VALUE NOT USED SINCE CONFORTO'S LIFETIME HAS BEEN DISCARDED			
AV C	CONFORTO VALUE COMBINES ALL FREE NEUTRON DECAY DATA			
AV D	CHRISTENSEN MEASUREMENT NOT SENSITIVE TO SIGN OF GA/GV			
AV	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
AV AVG	-1.2310	0.0098		
F	PHASE ANGLE OF GA RELATIVE TO GV (DEGREES)			
F C	176.1	6.4	CONFORTO	67 RVUE 11/68
F S	181.3	1.3	ERDZOLMSKI TO CNTR POLAR. NEUTRON	10/69
F C	VALUE DERIVED FROM FREE NEUTRON DECAY ONLY			
F S	ONLY STATISTICAL ERROR QUOTED			
F	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
F AVG	181.1	1.3		

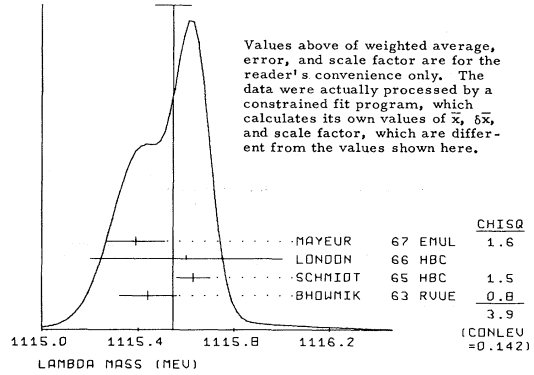
REFERENCES				
17 NEUTRON (939, J=1/2) I=1/2				
COHEN 56 PR 104 283 V W COHEN, CORNGOLD, RAMSEY (ANL+HARVARD)				
SOSNOVSK 59 JETP 9 717 SOSNOVSKII, SPIVAK, PROKDEEV + (IAE MOSCOW)				
YATTAUCH 65 NP 67 1 +THIELE, HAPSTRA (M. PLANCK INST. CHEM., +)				
BHALLA 66 PL 19 691 C P BHALLA (ALABAMA)				
CHRISTEN 67 PL 268 11 +NIELSEN, DAHSEN, BROWN, RUSTADT (DENMARK)				
CONFORTO 67 APAP 22 15 G. CONFORTO (CERN)				
BAIRD 69 PR 179 1285 +MILLER, DRESS, RAMSEY (ORNL, HARV)				
TAYLOR 69 RMP 41 375 +PARKER, LANGENBERG (PRINCETON+PENN)				
ERDZOLM 70 SJNP 11 583 ERDZOLMSKI, RONDARENKO, + (KURC MOSCOW)				
ALSO 68 PL 278 597 ERDZOLMSKI, RONDARENKO + (KURC IN MOSCOW)				
PAPERS NOT REFERRED TO IN DATA CARDS				
JACKSON 57 PR 106 517 JACKSON, TREIMAN, WYLD (PRINCETON)				
COHEN 65 RMP 37 537 E R COHEN, DUDDOND (NAAASC+CAL INST TECH)				

18 LAMBDA (1115, J=1/2+) I=0				
18 LAMBDA MASS (MEV)				
M N SINCE OUR FINAL VALUES FOR THE SIGMA AND LAMBDA MASSES COME FROM DOING AN OVERALL FIT TO ALL MEASURED MASSES AND MASS DIFFERENCES, WE HAVE USED THE UNCORRELATED MEASUREMENTS FROM SCHMIDT 65 RATHER THAN THE ONES COMING FROM THE OVERALL FIT REPORTED IN THAT PAPER. SINCE THERE SEEMS TO BE NO CONVINCING ARGUMENT AS TO WHY ONE SHOULD IGNORE DATA USING RANGE MEASUREMENTS, WE HAVE INCLUDED HERE VALUES DEPENDING ON PION AND PION RANGES.				
M	1115.44	0.12	BHOWMIK	63 RVUE + SEE NOTE L BELOW
L	ABOVE LAMBDA MASS HAS BEEN RAISED 35 KEV TO ACCOUNT FOR 46 KEV INCREASE IN PION MASS AND 11 KEV DECREASE IN CHARGED PION MASS.			
M	635(1115.86)	(0.09)	BALTAY	65 HBC ERROR IS STATIS. 6/66
M	488 1115.63	0.07	SCHMIDT	65 HBC SEE NOTE N 6/68
M	5 1147(1115.74)	(0.04)	CHIEN	66 HBC 6.9 PBAR P 9/67
M	5 972(1115.69)	(0.05)	CHIEN	66 HBC 6.9 PBAR PANTIL 9/67
M	1115.6	0.4	LONDON	66 HBC 6/66
M	(1116.0)	(0.2)	BADIER	67 HBC 2.4 PBAR P, LLBAR 8/67
M	105 1115.39	0.12	MAYEUR	67 EMUL 11/67
M	ERROR PURELY STATISTICAL			
M	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)			
M AVG	1115.544	0.075		
M	FIT 1115.596 0.063 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW) 2/71*			

DM	0.05	0.06	CHIEN	66 HBC 6.9 PBAR P 9/67
DM	0.29	0.15	BADIER	67 HBC 2.4 PBAR P 8/67
DM	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)			
DM AVG	0.083	0.083		

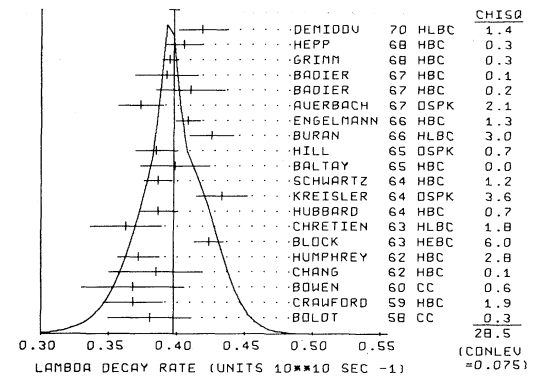
18 LAMBDA LIFETIME (UNITS 10**+10)				
T	198	2.63	0.21	0.21 BOLDT 58 CC
T	825	2.72	0.16	0.16 CRAWFORD 59 HBC
T	140	2.72	0.29	0.27 POWEN 60 CC
T	196	2.60	0.28	0.20 CHANG 62 HBC
T	799	2.69	0.11	0.11 HUMPHREY 62 HBC
T	2239	2.36	0.06	0.06 BLOCK 63 HBC
T	706	2.76	0.20	CHRISTEN 63 HBC
T	794	2.59	0.09	HUBBARD 64 HBC
T	2260	2.31	0.10	KREISLER 64 OSPK
T	1378	2.59	0.07	SCHWARTZ 64 HBC
T	435	2.51	0.16	BALTAY 65 HBC 6/66
T	2534	2.4	0.1	HILL 65 OSPK 6/66
T	916	2.35	0.09	BURAN 66 HBC

WEIGHTED AVERAGE = 1115.544 ± 0.075
ERROR SCALED BY 1.4



T	5 1147	(2.50)	(0.14)	CHIEN	66 HBC	6.9 PBAR P	9/67
T	5 972	(2.70)	(0.20)	CHIEN	66 HBC	6.9 PBAR P, ANTI	9/67
T	2213	2.452	0.066	0.054 ENGELMANN	66 HBC		9/66
T	585	2.68	0.13	0.11 AUERBACH	67 OSPK		8/67
T		2.44	0.15	BADIER	67 HBC	2.4 PBAR P	6/68
T		2.55	0.15	BADIER	67 HBC	2.4 PBAR P, ANTI	6/68
T	8342	2.535	0.035	GRIMM	68 HBC		6/68
T	2600	2.47	0.08	HEPP	68 HBC		8/68
T	1059	2.39	0.10	DEMIDOV	70 HLBC	PI-P, 3.86 GEV/C	12/70*
T	ERROR PURELY STATISTICAL						
T	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.2) (SEE IDEOGRAM BELOW)						
AVG	2.517	0.024	0.023				

WEIGHTED AVERAGE = 0.3974 ± 0.0037
ERROR SCALED BY 1.2



DT	0.064	0.085	BADIER	57 HBC	2.4 PBAR P	8/67
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18 LAMBDA MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

MM	-1.5	0.5	COOL	67 OSPK
MM	0.0	0.6	KERNAN	63 CC
MM	8553	-1.39	0.72	ANDERSON 64 HBC
MM	151	-0.5	0.28	CHARRIERE 65 EMUL
MM	900	-0.70	0.27	HILL 65 OSPK
MM		-0.70	0.07	DAML-JENS 70 EMUL PULSED MAGNET
MM	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
MM AVG	-0.704	0.065		

18 LAMBDA PARTIAL DECAY MODES

P1	LAMBDA INTO PROTON PI-	DECAY MASSES
P2	LAMBDA INTO NEUTRON PI0	938+ 139
P3	LAMBDA INTO PROTON MU- NEUTRINO	938+ 134
P4	LAMBDA INTO PROTON E- NEUTRINO	938+ 105+ 0
		938+ .5+ 0

Stable Particles

For notation, see illustrated key at beginning of data card listings.

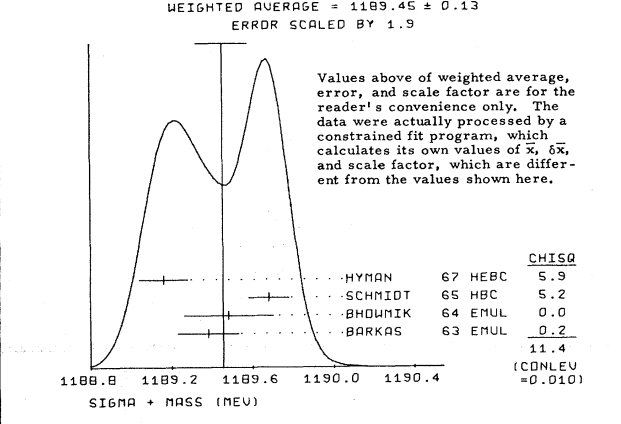
Table with columns: Particle name, Measurement type, Value, Error, Reference. Rows include R1, R2, R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, R14, R15, R16, R17, R18, R19, R20, R21, R22, R23, R24, R25, R26, R27, R28, R29, R30, R31, R32, R33, R34, R35, R36, R37, R38, R39, R40, R41, R42, R43, R44, R45, R46, R47, R48, R49, R50.

Table with columns: Particle name, Measurement type, Value, Error, Reference. Rows include A-, A+, A0, A1, A2, A3, A4, A5, A6, A7, A8, A9, A10, A11, A12, A13, A14, A15, A16, A17, A18, A19, A20, A21, A22, A23, A24, A25, A26, A27, A28, A29, A30, A31, A32, A33, A34, A35, A36, A37, A38, A39, A40, A41, A42, A43, A44, A45, A46, A47, A48, A49, A50.

Table with columns: Particle name, Measurement type, Value, Error, Reference. Rows include AV, AV2, AV3, AV4, AV5, AV6, AV7, AV8, AV9, AV10, AV11, AV12, AV13, AV14, AV15, AV16, AV17, AV18, AV19, AV20, AV21, AV22, AV23, AV24, AV25, AV26, AV27, AV28, AV29, AV30, AV31, AV32, AV33, AV34, AV35, AV36, AV37, AV38, AV39, AV40, AV41, AV42, AV43, AV44, AV45, AV46, AV47, AV48, AV49, AV50.

Table with columns: Particle name, Measurement type, Value, Error, Reference. Rows include ANDERSON, BARLW, BERTER, BURAN, CHEN, ENGLMANN, LONDON, ALJBRACH, BAUER, ENGLMANN, LONDON, ANDERSSON, GRIMM, HEPP, MERRILL, DAUBER, DOYLE, MALONEY, BAGGETT, CHU, DAHL, DEMIDOV, OLSEN.

Table with columns: Particle name, Measurement type, Value, Error, Reference. Rows include ARMENTER, BALTAY, BERGE, SIGMA+, MASS, N, SEE NOTE PRECEDING LAMBDA MASS LISTINGS, BAPKAS, BHOWMIR, KREFT, SCHMIDT, HYMAN.



Note on sigma Lifetime Errors

When combining lifetimes, we first convert mean lives tau to decay rates Gamma, since for small numbers of events the distribution of the decay rates is more nearly Gaussian. However, in checking input data it is useful to bear in mind the theoretical minimum statistical error delta_min(tau) in the mean life itself. This is

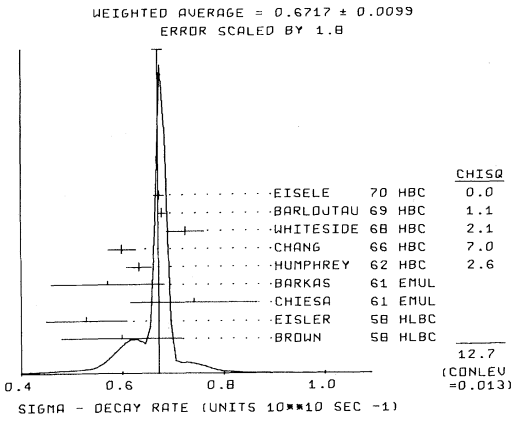
delta_min(tau) = 1 / (sqrt(N))

Stable Particles

For notation, see illustrated key at beginning of data card listings.

Table of particle data listings including references, sigma values, and mass listings. Includes sections for '20 SIGMA- LIFETIME', '20 SIGMA- PARTIAL DECAY MODES', '20 SIGMA- BRANCHING RATIOS', and '20 SIGMA- DECAY PARAMETERS'.

Table of particle data listings including references, sigma values, and mass listings. Includes sections for '20 SIGMA- LIFETIME', '20 SIGMA- PARTIAL DECAY MODES', '20 SIGMA- BRANCHING RATIOS', and '20 SIGMA- DECAY PARAMETERS'.



Stable Particles

For notation, see illustrated key at beginning of data card listings.

AV GV/GA FOR SIGMA TO LAMBDA BETA DECAY (SEE TEXT FOR SIGN CONVENTION) ... AV AVG 0.35 0.18 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

AV1 GA/GV FOR SIGMA TO NEUTRON BETA DECAY (SEE TEXT FOR SIGN CONVENTION) ... AV1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES

BROWN 58 CERN CONF 270 FISLER 58 NC SER10 10 150 ... ANG 1 69 ZPHY 223 103 ANG 2 69 ZPHY 228 151 ... BROWN 57 PR LOR 1036

***** PAPERS NOT REFERRED TO IN DATA CARDS *****

21 SIGMA 0 (1193, JP=1/2) I=1 D1 N SEE NOTE PRECEDING LAMBDA MASS LISTINGS. ... D1 AVG 4.849 0.069 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

21 (SIGMA 0) - (LAMBDA) MASS DIFFERENCE (MEV) DL N SEE NOTE PRECEDING LAMBDA MASS LISTINGS. ... DL FIT 76.024 0.089 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

21 SIGMA 0 LIFETIME (UNITS 10**14) T 1.0 OR LESS DAVIS 62 EMUL

21 SIGMA 0 PARTIAL DECAY MODES P1 SIGMA 0 INTO LAMBDA GAMMA 1115+ 0 P2 SIGMA 0 INTO LAMBDA F+ E- 1115+ .5+ .5

21 SIGMA 0 BRANCHING RATIOS R1 SIGMA 0 INTO (LAMBDA E+ F-)/TOTAL (0.0054) THEORET. CAL. FEINBERG 5R (P2)/(P1+P2) QUANTUM ELECT. 9/66

REFERENCES

FEINBERG 5R PR 109 1019 G.FEINBERG ... COURANT 63 PRL 10 409 ... ALFF 65 PR 137 B1105 ALFF, GELFAND, NAUENBERG

22 XI- (1321, JP=1/2) I=1/2 22 XI- MASS (MEV) M H 11(1317.0) (2.2) WANG 61 HLRC ... M AVG 1321.32 0.17 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

22 XI+BAR MASS (MEV) M1 S 1(1322.0) (1.3) BROWN 62 HBC ... M1 S THE ERROR IS STATISTICAL ONLY

22 MASS DIFFERENCE, (XI-) - (ANTI-XI-) IN MEV DM 1.0 1.1 CHIEN 66 HRC 6.9 PRAR P 9/67

22 XI- MAGNETIC MOMENT (MAGNETONS, 939.26 MEV) MM 2724 -0.1 2.1 BINGHAM 70 DSPK - K-P AT 1.8 GEV/C 2/71*

22 XI- LIFETIME (UNITS 10**10) T H 11 (3.5) (3.4) (1.23) WANG 61 HLRC ... T AVG 1.660 0.037 0.035 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.1)

22 XI+BAR LIFETIME (UNITS 10**10) T1 S 5 (1.5) (0.55) CHIEN 66 HRC ... T1 S 12 (1.9) (0.7) (0.5) SHEN 67 HBC

22 XI- PARTIAL DECAY MODES P1 XI- INTO LAMBDA PI- 1115+ 139 P2 XI- INTO LAMBDA E- NEUTRINO 111+ .5+ 0

22 XI- BRANCHING RATIOS R1 XI- INTO (LAMBDA E- NEU)/(LAMBDA PI-) (UNITS 10**3) (P2)/(P1) ... R1 1 (155) EFFECTIVE DENOM. CARMONY 63 HBC



Stable Particles

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle name, experiment name, units, and total count. Includes entries for XI- INTO (NEUTRON PI-), XI- INTO (LAMBDA MU- NEUTRINO), and XI- INTO (SIGMA E- NEUTRINO).

Table titled '22 XI- DECAY PARAMETERS' listing various decay parameters like ALPHA XI-, PHASE ANGLE, and branching ratios for different decay modes.

Table titled '22 XI- (1321,JP=1/2) I=1/2' listing references for various experiments and authors such as Fowler, Wang, and Cerny.

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS



Table with columns for particle name, experiment name, units, and total count. Includes entries for XI 0 (1314,JP=1/2) I=1/2 and XI 0 MASS (MEV).

Table titled '23 XI 0 LIFETIME (UNITS 10**=-10)' listing lifetime measurements from various experiments like Jauneau, Cerny, and Hubbard.

Table titled '23 XI 0 PARTIAL DECAY MODES' listing partial decay widths for various decay channels like XI 0 INTO LAMBDA PI 0 and XI 0 INTO PROTON PI-.

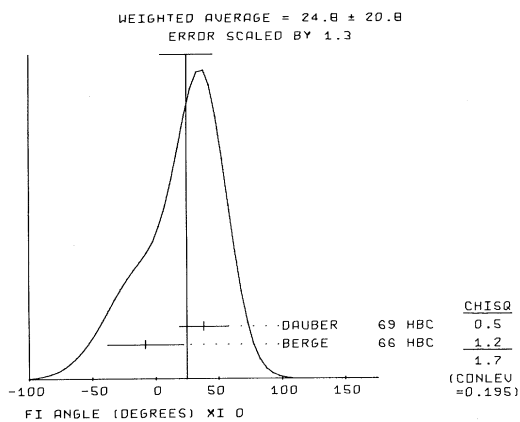
Table titled '23 XI 0 BRANCHING RATIOS' listing branching ratios for various decay modes like XI 0 INTO (PROTON PI-)/(LAMBDA PI 0) and XI 0 INTO (SIGMA+ E- NEU)/(LAMBDA PI 0).

Table titled '23 XI 0 DECAY PARAMETER' listing decay parameters for various experiments and authors like Jauneau, Cerny, and Hubbard.

Table titled '23 XI 0 DECAY PARAMETER' listing decay parameters for various experiments and authors like Jauneau, Cerny, and Hubbard.

For notation, see illustrated key at beginning of data card listings.

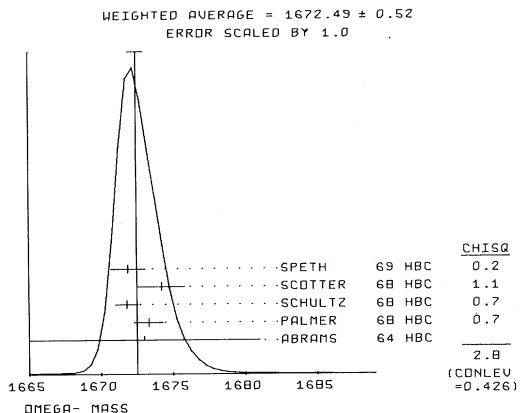
Stable Particles



 REFERENCES
 23 XI 0 (I314, JP=1/2) I=1/2
 ALVAREZ 59 PRL 2 215 ALVAREZ, EBERHARD, GOOD, GRAZIANO, TICHO+ (LRL)
 JAUNEAU 63 SIENA CONF 1 1 JAUNEAU+ (PARIS+CERN+LOND+RUTH+BERGEN)
 ALSO 63 PL 4 49 JAUNEAU+ (PARIS+CERN+LOND+RUTH+BERGEN)
 TICHO 63 BNL CONF 410 HAROLD K TICHO (UCLA)
 CARMONY 64 PRL 12 482 CARMONY, PJERRON, SCHLEIN, SLATER, STORK+ (UCLA)
 HUBBARD 64 PR 135 B 183 HUBBARD, BERGE, KALBFLEISCH, SHAFER+ (LRL)
 PJERRON 65 PRL 16 275 + SCHLEIN, SLATER, SMITH, STORK, TICHO (UCLA)
 PJERRON 65 THESIS G M PJERRON (UCLA)
 BERGE 66 PR 147 945 BERGE, EBERHARD, HUBBARD, MERRILL+ (LRL)
 HUBBARD 66 UCL 11510 J RICHARD HUBBARD (THESIS, BERKELEY) (LRL)
 LONDON 66 PR 143 1034 LONDON, RAU, GOLDBERG, LICHTMAN+ (BNL+SYRACUS) (LRL)
 MERRILL 66 BERKELEY CONF MERRILL, SHAFER, BERGE (LRL)
 ALSO 66 UCL 16455 DEANE MERRILL (THESIS, BERKELEY) (LRL)
 PALMER 64 PL 268 323 PALMER, RADOJICIC, RAU, RICHARDSON+ (BNL, SYR)
 DAUBER 69 PP 179 1262 +BERGE, HUBBARD, MERRILL, MILLER (LRL)

Ω⁻
 24 Ω⁻ (1675, JP=3/2+) I=0
 QUANTUM NUMBERS ASSIGNED FROM SU3
 THERE ARE 28 REPORTED Ω⁻ EVENTS
 SEE PREVIOUS EDITION (RMP 41 109) FOR MORE DETAILS
 24 Ω⁻ MASS (MEV)

M	1(1620.0)	(25.0)	(10.0)	EISENBERG	54 EMUL
M	1 1673.0	8.0		ABRAMS	64 HBC INTO XI- P10
M	3 1673.3	1.0		PALMER	68 HBC K-P 4.6 ± .5 GEV/C 11/69
M	3 1671.9	0.2		SCHULTZ	68 HBC K-P 5.5 GEV/C 11/69
M	5 1674.2	1.6		SCOTTER	68 HBC K-P 6. GEV/C 11/69
M	6 1671.9	1.2		SPETH	69 HBC K-P 10. GEV/C 11/69
M	AVG	1672.49	0.52	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) (SEE IDEOGRAM BELOW)	



24 Ω⁻ Ω⁻ LIFETIME (UNITS 10⁻¹⁰ SEC.)

T A	1	(1.63)	ABRAMS	64 HBC	7/66	
T A	1	(0.71)	BARNES 1	64 HBC	7/66	
T A	1	(1.4)	BARNES 2	64 HBC	7/66	
T A	1	(1.85)	COLLEY	65 HBC	7/66	
T A	1	(1.5)	RICHARDSON	65 HBC	7/66	
T A	1	(0.93)	ARCLV COL	68 HBC	11/67	
T A	1	(2.6)	ARCLV COL	68 HBC	11/67	
T A	1	(1.6)	ARCLV COL	68 HBC	11/67	
T A	1	(0.21)	ARCLV COL	68 HBC	11/67	
T A	1	(1.20)	SCHULTZ	68 HBC	11/67	
T A	1	(0.06)	SCHULTZ	68 HBC	11/67	
T A	1	(0.63)	SCHULTZ	68 HBC	11/67	
T A	1	(0.25)	SCOTTER	68 HBC	6/68	
T A	1	(0.30)	SCOTTER	68 HBC	6/68	
T A	1	(0.71)	SCOTTER	68 HBC	6/68	
T A	1	(0.08)	SCOTTER	68 HBC	6/68	
T A	1	(1.04)	SCOTTER	68 HBC	6/68	
T A	1	(2.31)	SCOTTER	68 HBC	6/68	
T A	ALLISON INCLUDES ALL ABOVE + 3 MORE BNL EVENTS, UNPUBLISHED.				6/68	
T	21	1.31	0.37	0.24	ALLISON 68 RVUE	6/68
T	1	(2.3)			SPETH 69 HBC	10/69
T	1	(0.31)			SPETH 69 HBC	10/69

24 Ω⁻ Ω⁻ PARTIAL DECAY MODES

P1	Ω ⁻ INTO Λ ⁰ K-	DECAY MASSES
P2	Ω ⁻ INTO XI ⁰ P1-	1115+ 493
P3	Ω ⁻ INTO XI ⁰ P1 ⁰	1314+ 139
		1321+ 134

24 Ω⁻ Ω⁻ BRANCHING RATIOS

27 EXAMPLES OF Ω⁻ DECAYS HAVE BEEN REPORTED. 16 HAVE DECAYED INTO Λ⁰ K-, 9 INTO XI⁰ P1-, 2 INTO XI⁰ P1⁰, AND ONE IS AMBIGUOUS BETWEEN Λ⁰ K- AND XI⁰ P1-. 1 BNL EVENT HAS NOT BEEN DESCRIBED.

R1	Ω ⁻ INTO Λ ⁰ K-	P1
R1	2 EVENTS	PALMER 68 HBC 11/69
R1	3 EVENTS	SCHULTZ 68 HBC 11/69
R1	5 EVENTS	SCOTTER 68 HBC 11/69
R1	6 EVENTS	SPETH 69 HBC 11/69
R2	Ω ⁻ INTO XI ⁰ P1-	P2
R2	1 EVENTS	ABRAMS 64 HBC 11/69
R2	4 EVENTS	PALMER 68 HBC 11/69
R2	3 EVENTS	SCOTTER 68 HBC 11/69
R2	1 EVENT	SPETH 69 HBC 11/69
R3	Ω ⁻ INTO XI ⁰ P1 ⁰	P3
R3	1 EVENT	PALMER 68 HBC 11/69
R3	1 EVENT	SCOTTER 68 HBC 11/69

 REFERENCES
 24 Ω⁻ (1675, JP=3/2+) I=0
 EISENBERG 54 PR 96 541 Y EISENBERG (CORNELL)
 ABRAMS 64 PRL 13 670 + BURNSTEIN, GLASSER+ (MARYLAND+USNRL)
 BARNES 1 64 PRL 12 204 V E BARNES, CONNOLLY, CRENNELL, GULWICK+ (BNL)
 BARNES 2 64 PL 12 134 V E BARNES, CONNOLLY, CRENNELL, GULWICK+ (BNL)
 COLLEY 65 PL 19 152 COLLEY+ODD+ (BIRM+GLASGOW+MUN+OXF+RHUL)
 RICHARDS 65 BAPS 10 115 RICHARDSON, BARNES, CRENNELL+ (BNL+SYRACUS)
 SAMIOS 65 ARGONNE CONF 189 N P SAMIOS (RVUE) BNL
 ARCLV CO 68 NUC PHYS B4 326 AACHEN+BERLIN+CERN+LONDON IMP. COLL.+VIENNA
 ALLISON 68 PRIV. COMM. JOHN ALLISON (LANCASTER)
 PALMER 68 PL 268 323 PALMER, RADOJICIC, RAU, RICHARDSON+ (BNL, SYR)
 SCHULTZ 68 PR 168 1509 SCHULTZ+ (ILL, ARGONNE, NORTHWESTERN, WISC)
 SCOTTER 68 PL 268 474 SCOTTER+ (BIRM, GLASGOW, IC LONDON, MUNICH, OXF)
 SPETH 69 PL 298 252 SPETH+ (AACHEN, BERLIN, CERN, LONDON, VIENNA)

Mesons

For notation, see illustrated key at beginning of data card listings.

CODE EVENTS QUANTITY ERROR+ ERROR- REFERENCE YR TECN SIGN COMMENTS DATE ABOVE PUNCHED BACKGROUND

π

PI MESON (JPG=0--) I=1
SEE LISTING OF STABLE PARTICLES

$\sigma(410)$

7 OLD SIGMA MESON (410, JPG=0++) I=0
NO EVIDENCE FOR RESONANCE OMITTED FROM TABLE.
SEE NOTE ON ETA 0+(700-1000)
SEE REVIEW OF MORGAN AND PISUT 70.

REFERENCES ON SIGMA

SAMIOS 62 PRL 9 139	+BACHMAN, LEA+ (BNL+GCNY+CO+KY)
RL'KHINT 63 JETP 17 80	RL'KHINTSEVA, GREIBINNIK, ZHUKOV + (MURINA)
ROTH 63 PR 132 2314	+ABSHIAN (LRL)
KIRZ 63 PR 130 2481	+SCHWARTZ + TRIPP (LRL)
BARISH 64 PR 135 B 416	BARISH, KURZ, PEREZ-MENDEZ, SOLOMON (LRL)
CRAWFORD 64 PRL 13 421	+GROSSMAN+LLOYD+PRICE+FOMLER (LRL)
DEL FABR 64 PRL 12 674	DEL FABRO, DE PRETIS, JONFS+ (FRASCATI)
KALMUS 64 PRL 13 99	+KERNAN, PU, POWELL, DOWD (LRL+MICHIGAN)
BIRGE 65 PR 130 9 1600	+FLY+GIDAL+KALMUS+CAMERINI+ (LRL+WISC)
BROWN 65 CORAL GABLES 219	BROWN+FAIER (NORTHWESTERN)
JACOBS 66 PRL 16 669	+SELOVE (LRL)
KITPELMAN 66 PRL 22 119	+ALLEN+GODDEN, MARSHALL + (COLORADO+EDW)
LOVELACE 66 PRL 22 332	LOVELACE, HEINZ, DONNACHIE (CERN)
ANDERSON 67 PRL 18 89	+FUKUI+KESSLER+ (CHIC+ANL+OTT+MCGILL+QMC)
CORNETT 67 PR 156 1451	+DIMERELL+MIDDLEMAS+NEWTON (DXF+RUTHERF)
MALAMUD 67 PRL 19 1056	F. MALAMUD + P. E. SCHLEIN (UCLA)
WALKER 67 PRL 18 630	+CARROLL, GARFINKEL, OH (WISCONSIN)
BANDER 68 PR 168 1679	M. BANDER, G.L. SHAN, J.R. FULCO (UCI+UCSB)
BISWAS 68 PR 27 B 513	+CASON, JOHNSON, KENNEY, POIRIER+ (NOTRE DAME)
EISENHAN 68 PRL 20 758	EISENHAN, MISTRY, MOSTEK + (CORNELL)
FOSTER 68 NP B 6 107	+GAVILLET+LABROSSE+MONTANET+ (CERN+PARIS)
JONES 68 PR 166 1405	+CALDWELL+ZACHAROV+HARTING+HEULER+ (CERN)
MARATECK 68 PRL 21 1613	+HAGPITIAN+ (PENN+LRL+GOLD+PURD+TNTD+WISC)
DAVISON 69 PR 180 1333	+RACASTON+BARKAS+ (RIVS+BERK)
FLY 69 NP 180 1319	+GIDAL, HAGPITIAN, + (BERK+LUC+WISC)
GUTAY 69 NP B 12 31	+CARMONY, CASON, LOEFFLER, MEIERE (PURDUE)
HALL 69 NP B 12 573	+MURRAY, RIDDIFORD (BIRMINGHAM)
RODFORTS 69 PRL 29 B 368	R. G. ROBERTS, F. WAGNER (CERN)
RAJJOY 70 PRL 24 948	+GRIVES, VANBERG, MAGLIG+ (PENN+RUTG+UPN+ANL)
MAUNG 70 PRL 33 B 521	+MASEK, MILLER, RUDERMAN, VERNON, + (UCSD+LRL)
MORGAN 70 SPRINGER TRACTS	MOD. PHYS., VOL. 55, P. 1. MORGAN, PISUT (RHEL+CERN)

η

ETA (549, JPG=0++) I=0
SEE LISTINGS OF STABLE PARTICLES

$\eta_0+(700-1000)$
OR ϵ

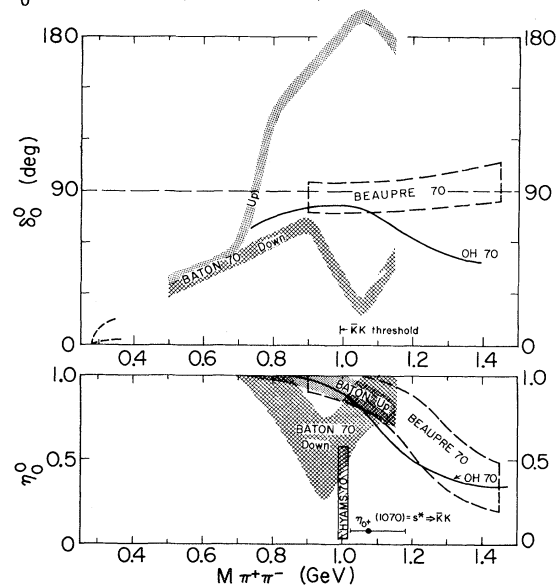
14 ETA 0+(700-1000, JPG=0++) I=0
ALSO CALLED EPSILON (720)

Information about the $\pi\pi$ system in the I = 0 S wave comes mainly from phase-shift analyses of the reaction $\pi^+p \rightarrow \pi^+\pi^+n$. Although no method used is free from serious objections, all analysts in the past (BATON 65, BATON 67, GUTAY 67, JOHNSON 67, WALKER 67, BISWAS 68, MARATECK 68, MALAMUD 69, SCHARENQUIVEL 69) have agreed that the S wave is near the unitarity limit in the

region 650-900 MeV. However there was an ambiguity between two sets of phase shifts which crossed at about 700 MeV, leading to four possible tracks, called "down-up," "down-down," The "down-up" solution, preferred by MARATECK 68 and SCHARENQUIVEL 69, leads to an " $\epsilon(720)$ " resonance, with Γ about 150 MeV. Experiments on the $\pi^0\pi^0$ system (BROWN 68, SMITH 69, DEINET 69, SONDEREGGER 69) have not yet removed this ambiguity. For a review, see MORGAN 70.

In the figure we have collected recent results on the S-wave phase δ_0^0 from analyses where the S_0 wave is permitted to be inelastic, as required by direct measurements of $\bar{K}K$ production (HYAMS 70). BATON 70 now rule out the low-energy "up" solution, and their "down-down" solution is mildly preferred because it has a rather big inelasticity compatible with that reported by HYAMS 70 at 1 GeV; moreover, above 1 GeV their "down-down" solution more-or-less agrees with that of OH 70 and BEAUPRE 70.

In the threshold region, MAUNG 70 have determined the scattering length to be $a_0 = 0.28 \pm 0.21$, in agreement with the previous most probable value, $a_0 = 0.16 \pm 0.04$ (MORGAN 70).



Mesons

For notation, see illustrated key at beginning of data card listings.

Table of references for mesons, including authors like J.P. Baton, G. Laurens, J. Regnier, and various institutions like SLAC and CERN.

ρ(765)

9 RHO (765, JPG = 1-+) I=1

THESE ARE WIDE FLUCTUATIONS IN THE MEASURED VALUES FOR MASS AND WIDTH OF THE RHO DUE TO DIFFERENCES IN PRODUCTION MECHANISM, BACKGROUND, METHOD OF ANALYSIS AND PARAMETRIZATION UNCERTAINTIES IN THEORY GIVE RISE TO SYSTEMATIC ERRORS OF ABOUT 20 MEV IN MASS AND WIDTH.

THE FOLLOWING ENTRIES ARE THE MOST SIGNIFICANT ONES. THEY ILLUSTRATE THE DISCREPANCIES, AND ARE ALSO REPEATED IN FOOTNOTE (C) OF THE MESON TABLE.

D1 BATON 67,70 (RHO-0 FROM POLE EXTRAPOLATION)
D2 HYAMS 68 (RHO-0 FROM PHYSICAL REGION FITS)
D3 MARATECK 68 (RHO-0 FROM POLE EXTRAPOLATION)
D4 PISUT 68 (RHO-0 FROM PHYSICAL REGION FITS)
D5 AUGUSTIN 2 69 (RHO-0 INTERFERING WITH OMEGA, FROM E+E- COLL BEAMS)
D6 AUSLENDER 69 (RHO-0 FROM E+E- COLLIDING BEAMS)
D7 MALAMUD 69 (RHO-0 FROM PHYSICAL REGION FITS)
D8 RODS 69 (COMPILE OF RHO-0 FROM E+ E- COLLIDING BEAMS)
D9 SCHARENQUIVEL 69 (RHO-0 FROM POLE EXTRAPOLATION)

CHARGE PLUS ONLY
M R (760.0) (9.0) CARMONY 64 HRC + 3.5 P1+P, CUT 4
M R 760.0 10.0 ARMENISE 65 HRC + 2.8 P1+P
M R (765.0) (15.0) ALFF-STEI 66 HRC + 2.3 P1+P 6/66
M R (783.0) (6.0) JAMES 66 HRC + 2.1 P1+P 6/66
M R (758.0) (10.0) JAMES 66 HRC + 2.1 P1+P, CUT 2.5 8/66
M R 777.0 7.0 ARC COLL. 68 HRC + 8 P1+P TO P+3P1 5/68

CHARGE PLUS AND MINUS
M S (750.0) (3.0) BALTAY 66 HRC +- 0.0 PBAR P
M R 755.0 10.0 ALLES-BDR 67 HRC +- 5.7 PBAR P 12/66
M R 730.0 11.0 BARLOW 67 HRC +- 1.2 PBAR P 11/66
M R 782.0 5.0 FOSTER 68 HRC +- PBAR P AT REST 6/68

CHARGE MINUS ONLY
M R (748.0) KENNEY 62 HRC - 1.2 P1-P
M R 130 (775.0) (2.0) GUIRAGOSS 63 HRC - 3.3 P1-P
M R (768.0) (5.0) BLIEDEN 65 MMS - 3.5 P1-P 6/66
M R 772.0 19.0 FIDECARD 66 OSPK - 2.5 P1-P, CUT 8 11/66
M R (760.0) (5.0) HAGOPIAN 66 HRC - 3.0 P1-P 6/66
M R (765.0) (5.0) HAGOPIAN 66 HRC - 2.14 P1-P, CUT 12 9/67
M R 6014 (757.6) (6.6) JACOBS 66 HRC - 2-3P1- 6/68
M R 2775 (753.5) (10.5) JACOBS 66 HRC - 2-3P1- CUT 20 6/68
M R (749.0) (2.0) EISENER 67 HRC - 2.1 P1-P 10/66
M R 752.0 14.0 BANNER 67 MMS - 1.8 P1-P, PNM 9/67
M C 7556 (755.0) (5.0) DI RATON 67 HRC - 2.8 P1-P 10/67
M R 751.0 5.0 CLEAR 67 HRC - 3 P1-P 7/67
M R (777.0) (6.0) MILLER 67 HRC - 2.7 P1-P, CUT 5 9/66
M R (775.0) (5.0) MILLER 67 HRC - 2.7 P1-P, CUT 10 9/66
M R (764.0) (5.0) MILLER 67 HRC - 2.7 P1-P, CUT 20 9/66
M R (752.0) (2.0) EISENER 67 HRC - 6.2 P1-P 9/67
M R 12773 764.3 1.9 1.9 P4 PISUT 68 RVUE - 1.7-3.2 P1-P, CUT 10 6/68
M A 12773 (764.3) (19.2) (3.3) P4 PISUT 68 RVUE - 1.7-3.2 P1-P, CUT 10 6/68
M A ERRORS ARE 2 STD AND INCLUDE SYSTEMATIC UNCERTAINTIES FROM THEORY
M R 776.0 4.0 REYNOLDS 69 HRC - 2.26 P1-P 5/70

Main table of meson data entries, including authors like SAMIOS, ABOLINS, GUIRAGOSS, and various meson types like RHO, OMEGA, and their production methods.

MIXED CHARGES
M 240 (752.0) ALITTI 63 HRC - 0 1.6 P1-P
M 290 (765.0) CHADWICK 63 HRC +- 0.0 PBAR P
M 150.0 9.0 FRENCH 67 HRC +- 3.4 PBAR P 6/67
M 775.0 2.0 JOHNSON 68 HRC - 0 3.7-4.2 P1-P 7/69

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

9 RHO(0) - RHO(+-) MASS DIFFERENCE (MEV)
D 2.4 2.1 PISUT 68 RVUE P1 N TO RHO N 5/68

9 RHO WIDTH (MEV)
SEE NOTE ON RHO MASS ABOVE

CHARGE PLUS ONLY
M R 90.0 10.0 SAGLAY 63 HRC + 2.8 P1+P
M R (77.0) (20.0) CARMONY 64 HRC + 3.5 P1+P, CUT 4
M R 150.0 10.0 ARMENISE 65 HRC + 2.8 P1+P 5/70
M R (100.0) ALFF-STEI 66 HRC + 2.3 P1+P 6/66
M R (177.0) (15.0) JAMES 66 HRC + 2.1 P1+P 7/66
M R (147.0) (19.0) JAMES 66 HRC + 2.1 P1+P, CUT 2.5 8/66
M R 149.0 22.0 ARC COLL. 68 HRC + 8 P1+P TO P+3P1 5/68

CHARGE PLUS AND MINUS
M S (150.0) (30.0) BALTAY 66 HRC +- 0.0 PBAR P 6/66
M S (150.0) (30.0) BALTAY 66 HRC +- 0.0 PBAR P 6/66
M R 146.0 31.0 ALLES-BDR 67 HRC +- 5.7 PBAR P 12/66
M R 130.0 25.0 BARLOW 67 HRC +- 1.2 PBAR P 11/66
M R 145.0 10.0 FOSTER 68 HRC +- PBAR P AT REST 6/68

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

Mesons

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle name, mass, width, and other properties. Includes entries for OMEGA INTO (ZPI0 GAMMA)/(PI0 GAMMA), OMEGA INTO (ETA GAMMA)/(PI0 GAMMA), etc.

WEIGHTED AVERAGE = 0.66 ± 0.17
ERROR SCALED BY 1.4

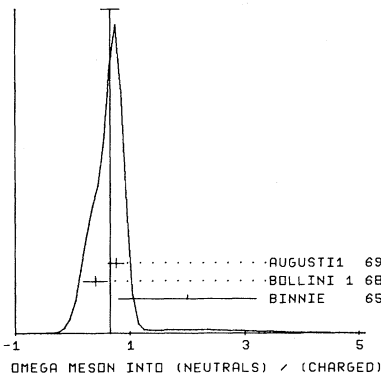


Table with columns for particle name, mass, width, and other properties. Includes entries for OMEGA INTO NEUTRALS / TOTAL, OMEGA INTO (PI PI)/(TOTAL), etc.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± delta P_i, where delta P_i = sqrt(delta P_i^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of correlation coefficients for P1, P2, and P3.

REFERENCES FOR OMEGA

Table listing references for the OMEGA particle, including authors, journals, and years. Includes entries like MAGLIC, ALVAREZ, ROSENFELD, STEVENSON, etc.

M(953)

59 M(953) G=+1
WHILE MASS AND WIDTH ARE CONSISTENT WITH THOSE OF THE ETA PRIME(958), THE (PI+ PI- GAMMA) DECAY DOES NOT SHOW A RHOD SIGNAL, UNLIKE THE ETA PRIME. THIS IS TAKEN AS EVIDENCE FOR A NEW PARTICLE, WHILE THIS DIFFERENCE IN DALITZ PLOT DISTRIBUTION APPEARS SIGNIFICANT, IT STILL NEEDS FURTHER CONFIRMATION TO BE REGARDED AS WELL ESTABLISHED.

Table with columns for particle name, mass, width, and other properties. Includes entries for M(953) G=+1, M(953) G=+1, etc.

For notation, see illustrated key at beginning of data card listings.

REFERENCES FOR η

AGUILAR 70 PRL 25 1635 AGUILAR-BENITEZ, RASSANO, SAMIOS, BARNES (BNL)

$\eta'(958)$

2 ETA PRIME (958, JPC=0-+) I=0

KNOWN ALSO AS X0
(JPC = 2- NOT YET EXCLUDED.)
(SEE NOTE ON QUANTUM NUMBERS AT END OF ETA PRIME LISTINGS)

2 ETA PRIME MASS (MEV)

M	85	(957.0)		DAUBER	64 HBC	1.95 K-P	
M	K	(958.0)	(1.0)	KALBFLEISCH	64 HBC	2.7 K-P	6/66
M	K			KALBFLEISCH	64 SUPERSEDED BY RITTENBERG 69		
M		957.0	3.0	BADIER	65 HBC	3.0 K-P	
M		960.0	2.0	TRILLING	65 HBC	3.65 P+ P	9/66
M		956.0	10.0	CONN	66 HBC	3.3 P+D	6/66
M		950.0	3.0	LONDON	66 HBC	2.2 K-P	6/66
M		960.0	9.0	MOTT	69 HBC	4.1-5.5 K-P	7/69
M		957.	1.	RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
M		956.0	2.0	AGUILAR	70 HBC	3.9-4.6 K-P	1/71*
M	AVG	957.46	0.75			AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

2 ETA PRIME WIDTH (MEV)

M	85	(4.0) OR LESS		DAUBER	64 HBC	1.95 K-P	
M	K	(7.0) OR LESS		KALBFLEISCH	64 HBC	2.7 K-P	6/66
M	K			KALBFLEISCH	64 SUPERSEDED BY RITTENBERG 69		
M		(30.0) OR LESS		BADIER	65 HBC	3.0 K-P	
M		(15.0) OR LESS		LONDON	66 HBC	2.2 K-P	6/66
M		(10.0) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
M		(20.0) OR LESS		AGUILAR	70 HBC	3.9-4.6 K-P	1/71*

2 ETA PRIME PARTIAL DECAY MODES

P1	ETA PRIME INTO PI+ PI- ETA	DECAY MASSES	139+ 139+ 548
P1	P1(N) ETAS DECAY INTO ALL NEUTRALS		
P1	P1(C) ETAS DECAY CHARGED		
P2	ETA PRIME INTO P0 P0 ETA	134+ 134+ 548	
P2	P2(N) ETAS DECAY INTO ALL NEUTRALS		
P2	P2(C) ETAS DECAY CHARGED		
P3	ETA PRIME INTO P1+ PI- GAMMA	139+ 139+ 0	
P3	(INCLUDING RHO GAMMA)		
P4	ETA PRIME INTO GAMMA GAMMA	0+ 0	
P4	ETA PRIME INTO RHO GAMMA	0+ 765	
P10	ETA PRIME INTO P1+ PI- E+ E-	139+ 139+ .5+ .5	
P11	ETA PRIME INTO 2 PI	139+ 139	
P12	ETA PRIME INTO 3 PI	139+ 139+ 134	
P13	ETA PRIME INTO 4 PI	139+ 139+ 139+ 139	
P14	ETA PRIME INTO 5 PI	139+ 139+ 139+ 139	
P15	ETA PRIME INTO P0 GAMMA GAMMA	134+ 0+ 0	
P16	ETA PRIME INTO P0 E+ E- (VIOLATES C IN BORN APPROX.)	134+ .5+ .5	
P17	ETA PRIME INTO ETA E+ E- (VIOLATES C IN BORN APPROX.)	548+ .5+ .5	
P18	ETA PRIME INTO P0 RHO 0 (VIOLATES C)	134+ 765	
P19	ETA PRIME INTO P0 OMEGA (VIOLATES C)	134+ 783	

2 ETA PRIME BRANCHING RATIOS

Note on $\eta'(958)$ branching ratios

In our calculation of the constrained branching fractions of the $\eta'(958)$ we assume the following decay modes:

- a) $\eta\pi\pi$ (including $\eta\pi^0\pi^0$, 72% of the η' s have neutral decays),
- b) $\rho^0\gamma$,
- c) $\gamma\gamma$.

Note that the average of the directly observed $\gamma\gamma$ values (9.9 \pm 4.4% (BOLLINI 68, BENSINGER 70) is slightly different from the result of the overall fit, (6.6 \pm 3.7)%, because of independent measurements of ($\eta' \rightarrow$ all neutrals)/($\eta' \rightarrow$ total). In the fit we do not use the constraint

$$R = \Gamma(\eta' \rightarrow \eta\pi^+\pi^-) / \Gamma(\eta' \rightarrow \eta\pi^0\pi^0) = 2$$

from I-spin conservation. The result of the fit is in agreement with it, $R = 2.2 \pm 0.5$.

R1 ETA PRIME INTO (PI+ PI- ETA (NEUTRAL DEC.))/TOTAL

R1 K 68 (0.36) (0.05) KALBFLEISCH 64 HBC 2.7 K-P 10/66

R1 K 281 (0.34) (0.026) SUPERSEDED BY RITTENBERG 69 1.7-2.7 K-P 9/69

R1 FIT 0.317 0.017 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R2 ETA PRIME INTO (PI+ PI- NEUTRALS) / TOTAL

R2 33 0.35 0.06 BADIER 65 HBC 3.0 K-P 10/66

R2 39 0.4 0.1 LONDON 66 HBC 2.2 K-P 10/66

R2 AVG 0.353 0.051 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

R2 FIT 0.373 0.017 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R3 ETA PRIME INTO (PI+ PI- ETA (CHRG. DECAY))/TOTAL

R3 K 44 (0.12) (0.02) KALBFLEISCH 64 HBC 2.7 K-P 10/66

R3 K KALBFLEISCH 64 SUPERSEDED BY RITTENBERG 69

R3 7 0.07 0.04 BADIER 65 HBC 3.0 K-P 10/66

R3 10 0.1 0.04 LONDON 66 HBC 2.2 K-P 10/66

R3 107 0.123 0.014 RITTENBERG 69 HBC 1.7-2.7 K-P 9/69

R3 AVG 0.116 0.013 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

R3 FIT 0.1220 0.0066 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R4 ETA PRIME INTO (PI+ PI- NEUTRALS (EXCLUDING PI+ PI- ETA (NEUTR. DEC.))) / TOTAL

R4 K 10 (0.05) (0.04) KALBFLEISCH 64 HBC 2.7 K-P 10/66

R4 K 42 (0.045) (0.029) SUPERSEDED BY RITTENBERG 69 1.7-2.7 K-P 9/69

R4 FIT 0.056 0.014 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R5 ETA PRIME INTO (NEUTRALS) / TOTAL

R5 K 54 (0.25) (0.05) KALBFLEISCH 64 HBC 2.7 K-P 10/66

R5 K KALBFLEISCH 64 SUPERSEDED BY RITTENBERG 69

R5 16 0.26 0.17 BADIER 65 HBC 3.0 K-P 10/66

R5 32 0.3 0.1 LONDON 66 HBC 2.2 K-P 10/66

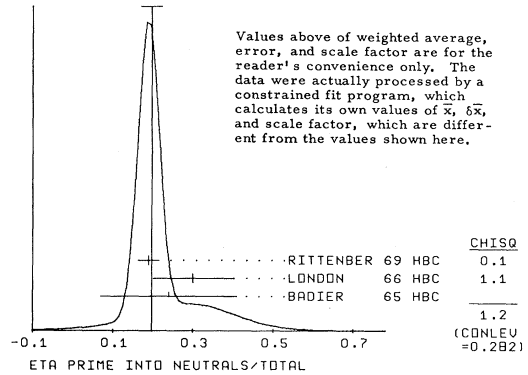
R5 123 0.189 0.026 RITTENBERG 69 HBC 1.7-2.7 K-P 9/69

R5 AVG 0.197 0.027 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)

R5 FIT 0.211 0.021 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.197 \pm 0.027
ERROR SCALED BY 1.1



R5 ETA PRIME INTO (PI+ PI- GAMMA (INCLUDING RHO GAMMA))/TOTAL

R5 K 42 (0.27) (0.04) KALBFLEISCH 64 HBC 2.7 K-P 10/66

R6 K KALBFLEISCH 64 SUPERSEDED BY RITTENBERG 69

R6 B 35 (0.34) (0.09) BADIER 65 HBC 3.0 K-P 10/66

R6 B CONTRADICTORY BACKGROUND SUBTRACTION

R6 20 0.2 0.1 LONDON 66 HBC 2.2 K-P 10/66

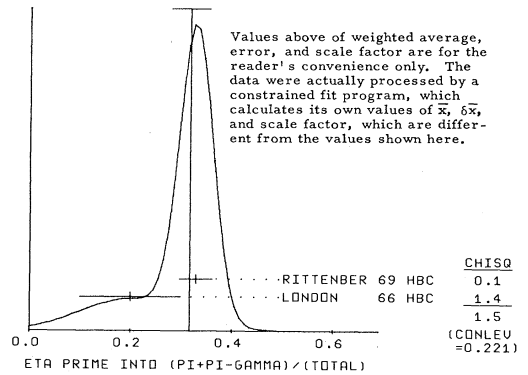
R6 298 0.329 0.033 RITTENBERG 69 HBC 1.7-2.7 K-P 9/69

R6 AVG 0.316 0.038 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)

R6 FIT 0.294 0.027 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.316 \pm 0.038
ERROR SCALED BY 1.2



Mesons

For notation, see illustrated key at beginning of data card listings.

R7	ETA PRIME INTO (PI+ PI- GAMMA (INCLUDING RHO GAMMA))/(PI PI ETA)	0.25	0.14	DAUBER	64 HBC	1.95 K-P	10/66
R7	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)	0.460	0.066				
R8	ETA PRIME INTO (PI0 E+ E-)/TOTAL	(0.013) OR LESS		RITTENBERG	65 HBC	2.7 K-P	10/66
R9	ETA PRIME INTO (ETA E+ E-)/TOTAL	(0.011) OR LESS		RITTENBERG	65 HBC	2.7 K-P	10/66
R10	ETA PRIME INTO (PI0 RHO0)/TOTAL	(0.04) OR LESS		RITTENBERG	65 HBC	2.7 K-P	10/66
R11	ETA PRIME INTO (PI0 OMEGA)/TOTAL	(0.08) OR LESS		RITTENBERG	65 HBC	2.7 K-P	10/66
R12	ETA PRIME INTO (PI+ PI- E+ E-)/TOTAL	(0.006) OR LESS		RITTENBERG	65 HBC	2.7 K-P	10/66
R13	ETA PRIME INTO (2 PI)/TOTAL	(0.07) OR LESS	COMP.BY LONDON	66 HBC			10/66
R14	ETA PRIME INTO (3 PI)/TOTAL	(0.07) OR LESS	COMP.BY LONDON	66 HBC			10/66
R15	ETA PRIME INTO (4 PI)/TOTAL	(0.01) OR LESS	COMP.BY LONDON	66 HBC			10/66
R16	ETA PRIME INTO (6 PI)/TOTAL	(0.01) OR LESS	COMP.BY LONDON	66 HBC			10/66
R18	ETA PRIME INTO (RHO0 GAMMA)/(PI PI ETA)	0.31	0.15	DAVIS	68 HBC	5.5 K-P	9/68
R18	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)	0.460	0.066				
R19	ETA PRIME INTO (2 GAMMA)/TOTAL	(0.055) (0.035) (0.030)	BOLLINI	68 CNTR	1.9 PI-P		9/68
R19	CHANGED BY US BECAUSE OF NEW CROSS SECTION, BENSINGER 70	0.085	0.055	0.045	BOLLINI	68 CNTR	1.9 PI-P
R19		0.126	0.075		BENSINGER	70 DRC	2.2 PI+ D
R19	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	0.098	0.042				
R19	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	0.066	0.037				
R20	ETA PRIME INTO (PI+PI-)/TOTAL	(0.02) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
R21	ETA PRIME INTO (PI+PI-PI0)/TOTAL	(0.05) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
R22	ETA PRIME INTO (PI+PI+PI-PI-)/TOTAL	(0.01) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
R23	ETA PRIME INTO (PI+PI+PI-PI-)/TOTAL	(0.01) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
R24	ETA PRIME INTO (PI+PI+PI-PI- NEUTRALS)/TOTAL	(0.01) OR LESS		RITTENBERG	69 HBC	1.7-2.7 K-P	9/69
R25	ETA PRIME INTO (RHO0 GAMMA)/(ALL PI+ PI- GAMMA)	0.94	0.20	AGUILAR	70 HBC	3.9-4.6K-P	1/71*
R26	ETA PRIME INTO (PI0 PI0 ETA INTO 3 PI0)/TOTAL	0.11	0.06	BENSINGER	70 DBC	2.2 PI+ D	1/71*
R26	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	0.060	0.015				

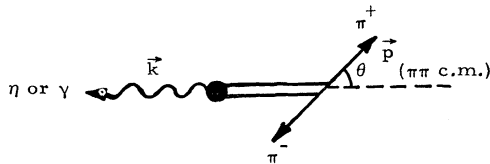
FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i^2)}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4
P 1	.439+-.024			
P 2	-.428	.201+-.049		
P 3	-.408	-.302	.279+-.027	
P 4	-.158	-.826	-.019	.066+-.035

UNCERTAINTY IN THE J^P ASSIGNMENT OF η'(958)

For the dominant (64%) $\pi\pi\eta$ decay mode of the η' , since the Dalitz plot population is rather flat (DAUBER 64, LONDON 66, RITTENBERG 69, DUFLEY 69), and in particular does not vanish at the edges of the plot, the $J^P = \text{Normal}$ series may be ruled out.



In the notation of the sketch, any Normal matrix element would have a factor $\sin\theta$ and would thus go to zero at the edge of the Dalitz plot [C. Zemach, Phys. Rev. 133, B1201 (1964)].

Since 1^+ is excluded by the observation of the $\gamma\gamma$ decay mode (BOLLINI 68, BENSINGER 70), this leaves us with $0^-, 2^-, \dots$ of the Abnormal series. Unfortunately, no information permits us at present to distinguish between 0^- and 2^- [A. Zaslavskij, V. Ogievetskij and W. Tybor, Soviet Journ. of Nucl. Phys. 9, 498 (1969)].

The $\pi\pi\eta$ decay mode (RITTENBERG 69, LONDON 66, DUFLEY 69) is fitted better with $J^P = 0^-$ than with $J^P = 2^-$, but the latter cannot be excluded, either when comparing the simplest matrix elements with the Dalitz plot density distribution or when judging from the θ -dependence or the η energy dependence of the squared element [Zaslavskij et al., op. cit., and S. Giler et al., Acta Phys. Polon. A 37, 475 (1970)].

Turning to the 29% mode $\eta' \rightarrow \pi^+\pi^-\gamma$, which RITTENBERG 69 has shown to be mainly $\rho^0\gamma$ decay, one gets a good fit with both $J^P = 0^-$ [$d\sigma/d(\cos\theta) \propto \sin^2\theta$] and $J^P = 2^-$ [$d\sigma/d(\cos\theta) \propto 6 + \sin^2\theta$ for the simplest possibility of pure M1 transition, or $d\sigma/d(\cos\theta) \propto 1.48 + \sin^2\theta$ from the vector dominance model (private communication from V. Ogievetskij)], so no distinction is possible here.

Moreover, the possible existence of a new resonance, $M(953)$, which decays into $\pi\pi\gamma$ but not into $\rho^0\gamma$, makes a study of the $\pi\pi\gamma$ final state even more ambiguous.

A longer discussion appears in our 1970 editions.

REFERENCES FOR ETA PRIME			
DAUBER	64 PRL 13 449	DAUBER, SLATER, SMITH, STORK, TICH0	(UCLA)JP
ALSO	64 DUBNA CONF 1 418	DAUBER, SLATER, L T SMITH, STORK, TICH0	(UCLA)
KALBFLEI	64 PRL 13 349	G.R. KALBFLEISCH, G. DAHL, A. RITTENBERG	(LRL)JP
BADIER	65 PL 17 337	BADIER, DEMOULIN, BARLOUTAUD + (PAR+SAC+ZEEM)	
KIENZLE	65 PL 19 438	KIENZLE, MAGLIC, LEVRAT, LEFEBVRES + (CERN)	
RITTENBERG	65 PRL 15 556	RITTENBERG, KALBFLEISCH	(LRL+ANL)
TRILLING	65 PL 19 427	+BROWN, GOLDBERG, KADYK, SCANTO	(LRL)
COHN	66 PL 21 347	COHN, MCCULLOCH, BUGG, CONDO (ORN+TENN+UNCAR)	
LONDON	66 PR 143 1034	LONDON, RAU, SAMIOS, GOLDBERG + (ANL+SYRACUSE)JP	
BOLLINI	68 NC 58 A 289	+RUHLER, DALPIAZ, MASSAM + (CERN+BGNA+STRB)	
DAVIS	68 PL 27 B 532	+AMMAR, MOTT, DAGAN, DERRICK, FIELDS (NWES+ANL)	
MOTT	69 PR 177 1966	+AMMAR, DAVIS, KR0PAc, SLATE, DAGAN+ (NWES+ANL)	
RITTENBERG	69 UCLR-18863	ALAN RITTENBERG (THESIS)	(LRL)I=0
AGUILAR	70 PRL 25 1635	AGUILAR-BENITEZ, BASSANO, SAMIOS, BARNES+ (BNL)	
BENSINGER	70 PL 33 B 505	BENSINGER, ERWIN, THOMPSON, W.D. WALKER (MISC)	
QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS			
MARTIN	66 PL 22 352	MARTIN, CRITTFN0EN, SCHR0EDER	(INDIANA)I=0
BARRARO	68 PRL 20 340	BARRARO-GALTIER, MATISON, RITTENBERG	(LRL)I=0
BARLOUTAUD	68 PL 26 B 674	BARLOUTAUD+ (SACLAY+AMSTD+BDL'G+HEI'G+E.P.)I=0	
DUFLEY	69 PL 29 B 605	+GOBBI, POUCHON, CNOPS, + (ETHZ+CERN+SACL)JP	

For notation, see illustrated key at beginning of data card listings.

$\delta(962)$ AND $\pi_N(975)$

3A DELTA MESON (962.JPG=) I = 1
AND P (975.JPG=N -) I = 1

The original, narrow $\delta^-(962)$ was seen with the CERN MMS (KIENZLE 65). Other missing-mass spectrometers (OOSTENS 66, BANNER-1 67, BANNER-2 67, ABOLINS 70) have added nothing conclusive to this evidence.

A claim for a $\delta^- \rightarrow \pi^- \pi^0$ decay at high $|t|$ in the reaction $\pi^- p \rightarrow \pi^- \pi^0 p$ (ALLEN 66) can now be discounted by virtue of the results of a compilation (ROOS 70) of the same reaction; the compilation has 7 times larger statistics in the region 640-1000 MeV and $0.5 < |t| < 1.0$ (GeV/c)², and includes the data of ALLEN 66.

Claims for 3π decay (ALLISON 67, JUHALA 68) have also been contradicted (SAMIOS 68, KRUSE 69). For a review, see SAMIOS 68, MAGLIĆ 69.

Recent experiments have, however, shown evidence for a broad resonance, also called $\pi_N(975)$. These are BARNES 69 and AMMAR 70, who see a peak in the $\eta\pi$ system, and who claim that it cannot be explained by the kinematic effect discussed below under 2a; moreover, ABOLINS 70 see a peak of width 60 MeV in a missing mass spectrum.

The following references have possible relevance to the existence of the δ .

1) The $\pi_N(1016)$ may be interpreted as a virtual bound state in the $K\bar{K}$ channel. It would then correspond to a narrow resonance at about 975^{+15}_{-10} MeV (ASTIER 67) in open channels, e. g., $\eta\pi$ or 5π .

2) Further $\eta\pi$ enhancements have been reported at masses in the 960-980 MeV region. As evidences for a resonance they are however not yet convincing, because there are two kinematic effects that can produce $\eta\pi$ peaks in that mass region:

a) In the reactions $K^- n \rightarrow \Lambda \pi^-(MM)$ and $K^- p \rightarrow \Lambda \pi^+ \pi^-(MM)$ (studied by CRENNELL 69, MILLER 69) with selection of the missing mass (MM) in the $\eta(549)$ region, a spurious δ peak can arise from contamination with $\Lambda p^- \pi^0$ final states. This has been pointed out by CRENNELL 69.

b) In final states containing many pions [e. g., $2\pi^+ 2\pi^- \pi^0$, $(3\pi)^\pm \pi^0$], and with the ω copiously produced, the constraint of at least one η combination in the $\pi^\pm \pi^+ \pi^- \pi^0$ mass "fakes" a peak in the mass region around 960 MeV, due to reflections from the ω . This remark (by NELLEN 69) may apply to the observations of DEFOIX 68 and CAMPBELL 69. OTWINOWSKI 70

has shown, however, that the δ peak in OTWINOWSKI 69 cannot be explained this way.

If we accept $\delta \rightarrow \pi\eta$ by strong decay, then $I^G = 1^-$; nonobservation of 3π decay can be explained by choosing $J^P = 0^+$, or simply by saying that 3π background is too large to permit detection. These quantum numbers, $1^-(0^+)$, are then the same as those most likely for $\pi_N(1016)$, which could be just the $\bar{K}\bar{K}$ decay mode of the δ .

An unattractive alternative is to believe that δ is really very narrow, and guess that its $\pi\eta$ decay is G-violating electromagnetic. (It is not clear whether there would be competition from $\pi\pi\eta$ decay, which is strong but has much smaller phase space.) However, in this electromagnetic (em) case, one would also expect slightly faster decay into $\pi\pi$, and we are not sure whether this mode should have been detected. To see why we expect $\pi\pi$ decay, note that these em decays into $\pi^- \pi^0$ or $\pi^- \eta$ involve emission and reabsorption of a photon, with rates proportional to e^4 (also $\pi\pi$ has slightly larger phase space than $\pi\eta$). Neutral em decays (as in the familiar $\eta^0 \rightarrow 3\pi$) have selection rules either

$$\Delta G = \text{Yes}, \quad \Delta |I| = 1,$$

or

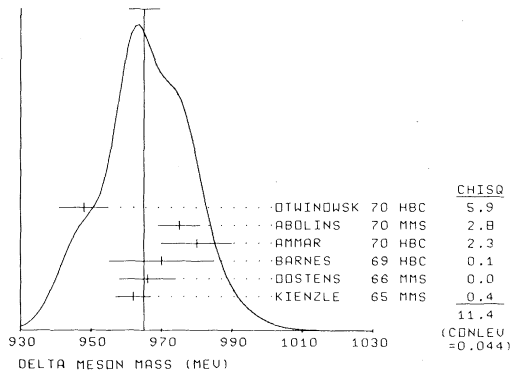
$$\Delta G = \text{No}, \quad \Delta |I| = 2,$$

but charged decays ($\delta^- \rightarrow \pi^- \pi^0$ or $\pi^- \eta$) have no such rules (except $\Delta |I| \leq 2$).

3A DELTA MESON MASS (MEV)

M	262	962.0	4.0	KIENZLE	65 MMS	-	3-5 PI- P	0/66	
M		964.0	4.0	OOSTENS	66 MMS	+	3,2 PP T3 D + MM	0/66	
M		(940.1)	APPROX.	CHUNG S	68 HRC	-	3,2 PI- P	4/70	
M		(975.01)		DEFOIX	69 HRC	+	1,2 PA P, FTA PI	3/69	
M		970.0	15.0	BARNES	69 HRC	-	4-5 K- P, PI- FTA PI	0/69	
M		(980.0)	(10.0)	MILLER	69 HRC	-	4-5 K- N, FTA PI	7/69	
M	D	(964.0)	(6.0)	OTWINOWSKI	69 HRC	+	4 PI+ P, P 70 PI	11/69	
M	D	OTWINOWSKI	69	SUPERSEDED BY OTWINOWSKI 70					
M		990.0	10.0	AMMAR	70 HRC	+	4,1,5,5K- FTA PI	5/70	
M		975.0	6.0	ABOLINS	70 MMS	+	3,2-5,3 PP- 7MM	1/71	
M		22	964.0	7.0	OTWINOWSKI	70 HRC	+	4 PI+ P, P 70 PI	11/69
M	AVG	965.0	4.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)					

WEIGHTED AVERAGE = 965.0 ± 4.4
ERROR SCALED BY 1.5



Mesons

For notation, see illustrated key at beginning of data card listings.

4 PHI PARTIAL DECAY MODES
PHI INTO K+ K-
PHI INTO K0I K0J
PHI INTO PI+ PI- (INCLUDING RHO PI)
PHI INTO PI+ PI- (VIOLATES G)
PHI INTO E+ E-
PHI INTO MU+ MU-
PHI INTO PI0 GAMMA
PHI INTO ETA GAMMA
PHI INTO PI0+PI0-GAMMA
PHI INTO OMEGA GAMMA (VIOLATES C)
PHI INTO ETA PI0 (VIOLATES C)
PHI INTO RHO GAMMA (VIOLATES C)

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS
The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i +/- delta P_i, where delta P_i = sqrt(delta P_i^2 + delta P_j^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

REFERENCES FOR PHI
BERTANZA 62 PRL 9 130
ARMENTEROS 63 SIENA CONF 2 70
GELFAND 63 PRL 11 438
GELFAND 63 DATA INCLUDED IN MILLER 65 BELOW
SCHLEIN 63 PRL 10 368
RADIER 65 PL 17 337
BERLEY 65 PP 139 B 1097
GALTIERI 65 PRL 14 279
LINDSEY 65 PRL 15 221
LINDSEY 65 DATA INCLUDED IN LINDSEY 66 BELOW
LINDSEY 65 UCRL 16526
MILLER 65 CU-237 (NEVIS 131)
GRAY, L 66 PRL 17 501
LINDSEY 66 PR 147 913
LINDSEY 66 PL 20 93
LINDSEY 66 DATA INCLUDED IN LINDSEY 66 ABOVE
LONDON, 66 PR 143 1034
ABRAMS 67 MD TECH REP 720
BARLUM 67 NC 50A 701
CHASE 67 PRL 18 710
DAHL 67 PR 143 1377
HERTZBACH 67 PR 155 1461
KHACHATU 67 PL 248 349
ABRAMS 68 PR 175 1697
ASTVACAT 68 PL 27 B 45
ALSI 67 PRL 19 869
RECKER 68 PRL 21 1504
BINNIE 68 PL 27B 106
BOLLINI 68 NC 56 A 1171
MOSTER 68 PRL 20 1057
WEHMANN 68 PRL 20 748
RALAKIN 69 IYAF 327 TRANS
BEMPORAD 69 PL 29 B 383
MOY 69 THESIS
SCOTTER 69 NC 62 A 1057
SIDOROV 69 LIVERPOOL SYMPOSIUM
BALAKIN 70 PREPRINT
BENAKSAS 70 LAL 1240
BIZOTI 70 PL 32 416
SEE ALSO REF. Y-JORBA, LIVERPOOL SYMPOSIUM
RALAKIN 70 PRIV.COMM.
EARLES 70 PRL 25 1312
HYAMS 70 PRIV.COMM.
SABRE 70 PREPRINT

NOTE ON I_G^0 = 0^+ PEAKS WITH M = 1000 - 1120 MeV [S^* -> KK-bar, eta_N(1080) -> pi pi].

Evidence for I_G^0 = 0^+ peaks in this region comes from two distinct final states: KK-bar and pi pi-. As the KK-bar peaks undoubtedly have J^P = 0^+, whereas the pi pi- peaks in some cases favor J^P = 2^+, it is not clear whether there are one or more objects in this region. For the time being we continue to treat the KK-bar and pi pi- peaks as two distinct objects, named eta_0+(1070) or S^*, and eta_N(1080), respectively.

Final understanding of these resonances, which can decay into both KK-bar and pi pi-, requires coupled-channel analyses. For a summary of the S-wave pi pi partial wave amplitude in this region, see the earlier mini-review on the eta_0+(700-1000) or epsilon, meson.

The disagreement between some of the observed widths of the KK-bar peaks is related to an ambiguity in interpretation either as a resonance (above threshold) or as a scattering length effect in the KK-bar channel. Computation of the average values of mass and width of the S^* would be useless because of the large discrepancies. In the meson table we give the values estimated in the review of BEUSCH 70.

The eta_N(1080) is seen in pi pi- -> pi pi- n predominantly at backward decay angles, cos theta < -0.75. OH 70 state that this "bump is almost certainly the result of p-d interference".

Mesons

For notation, see illustrated key at beginning of data card listings.

Note that the selection made in some HBC experiments to reduce the background under the $\eta_N(1080)$ in the reaction $\pi^- p \rightarrow \pi^+ \pi^- n$ may lead to a sample of events ambiguous with $\pi^- p \rightarrow p \pi^- \pi^0$. This is because selecting on small momentum transfer to the $\pi^+ \pi^-$ system, together with large π^- in/out scattering angle, leads to rather high lab momenta of the π^+ , so that ionization cannot be used to discriminate between the two hypotheses (BATON 70, footnote p. 525, and private communications from G. Laurens).

$\eta_0^+(1070)$ or S^*		3 ETA (1070, JPC=0+0+) I=0 NAMED S* BY CRENNELL 66. SEE NOTE ABOVE.
3 ETA (1070) MASS (MEV)		
M	(1000.0)	APPROX RINGHAM 42 HBC 6-18 PI-N
M	(1000.0)	APPROX RIGI 62 HBC 10.0 PI-P
M	(1000.0)	APPROX ERWIN 42 HBC 2.10 PI-P
M	30(1030.0)	APPROX. BALTAY 64 HBC 3.7 PBAR P
M	30(1030.0)	APPROX. BARMIN 44 HBC 2.8 PI-P
M	20 1068.0	10.0 CRENNELL 66 HBC 6.0 PI-P
M	120	SCATT. LENGTH FITS BETTER. HESS 66 HBC 1.6-4.2 PI-P
M	730 1070.0	6.0 5.0 BEUSCH 67 DSPK 5.7, 1.2 PI-P
M	54 1030.0	10.0 ALITTI 48 HBC, DBC 3.6-5.0 K-N
M	1046.0	7.0 AGUILAR- 69 HBC 0.7, 1.2 PBAR P
M	AGUILAR 69	SEES INDICATION OF D-WAVE IN ETA(1060) REGION.
M	H (1070.)	(SEE NOTE (1)) HOANG 69 DSPK 4 PI-P, KS KS N
M	H (1040.)	(SEE NOTE (2)) HOANG 69 DSPK 4 PI-P, KS KS N
M	H (1030.)	(SEE NOTE (3)) HOANG 69 DSPK 4 PI-P, KS KS N
M	H	SCATTERING LENGTH FITS EQUALLY WELL. VALUES ARE (IN F)
M	H	+(1.2 +0.3/-0.4) AND SIGNIFICANT IN PART (FOR 4 GEV/C DATA)
M	H	+(1.0 +0.3/-0.3) AND SIGNIFICANT IN PART (FOR 5 GEV/C DATA)
M	B (1112.)	(SEE NOTE (1)) BEUSCH 70 DSPK 4.6 PI-P
M	B (1053.)	(SEE NOTE (2)) BEUSCH 70 DSPK 4.6 PI-P
M	B	SCATTERING LENGTH FITS ONLY WITH EFFECTIVE RANGE. VALUE (F) IS
M	B	+(11.3 +0.35/-0.35) WITH EFF. RANGE = 1.4 F. (NOTE THAT SCATT.
M	B	LENGTH PLUS EFFECTIVE RANGE TERM ALWAYS SIMULATE RESONANCE IF
M	B	RELATIVE SIGN CHOSEN NEGATIVE.)
M	(1)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 0.
M	(2)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 1.
M	(3)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 2.
M	AVG	1061.7 10.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.9)

3 ETA (1070) WIDTH (MEV)		
W	20 80.0 15.0 19.0 CRENNELL 66 HBC 6.0 PI-P	
W	168.0 21.0 19.0 BEUSCH 67 DSPK 5.7, 1.2 PI-P	
W	BEUSCH 67 ASSUMES NO S WAVE SCATTERING LENGTH. WITH S WAVE THE WIDTH	
W	BECOMES NARROWER THAN QUOTED ABOVE.	
W	54 45.0 35.0 15.0 ALITTI 48 HBC, DBC 3.6-5.0 K-N	
W	40.0 29.0 AGUILAR- 69 HBC 0.7, 1.2 PBAR P	
W	H (1200.)	(SEE NOTE (1)) HOANG 69 DSPK 4 PI-P, KS KS N
W	H (1100.)	(SEE NOTE (2)) HOANG 69 DSPK 4 PI-P, KS KS N
W	H (1150.)	(SEE NOTE (3)) HOANG 69 DSPK 4 PI-P, KS KS N
W	H	SCATTERING LENGTH FITS EQUALLY WELL. SEE UNDER MASS ABOVE.
W	R (1361.)	(SEE NOTE (1)) BEUSCH 70 DSPK 4.6 PI-P
W	S (208.)	(SEE NOTE (2)) BEUSCH 70 DSPK 4.6 PI-P
W	B	SCATT. LENGTH WITH EFFECTIVE RANGE ALSO FITS. SEE UNDER MASS.
W	(1)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 0.
W	(2)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 1.
W	(3)	FIT ASSUMING BRANCHING RATIO (2 PI)/(K KBAR) OF 2.
W	AVG	85.8 27.4 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.9)

3 ETA (1070) PARTIAL DECAY MODES		
P1	ETA (1070) INTO KKBAR	493+ 497
P2	ETA (1070) INTO P1PI	139+ 134

3 ETA (1070) BRANCHING RATIOS			
R1	ETA (1070) INTO (PI P1)/(K KBAR)	66 HBC	90 PCT CONF LEV
R1	(2,5) OR LESS	CRENNELL	6 PI-P
R1	1.0	LAI	68 HBC

***** REFERENCES FOR ETA(1070) *****

WANG	61 JETP 13 323	WANG TSU-TSENG, VEKSLER, VRANA, + (JINR)
RIGI	62 CERN CONF 247	A RIGI, S BRANOT, R CARRARA + (CERN)
BINGHAM	62 CERN CONF 240	H H BINGHAM, M BLOCH + (PARIS-SEC POLY-CERN)
ERWIN	62 PRL 9 34	ERWIN, HOYER, MARCH, WALKER, WAGLER (WIS+BNL)
BALTAY	64 DUBNA CONF 1 409	BALTAY, LACH, CRENNELL, OREN, STUMP + (YALE+BNL)
BARMIN	64 DUBNA CONF 1 433	BARMIN, DOLGOLENKO, YEROFEEV, KRISTINI + (ITEP)
CRENNELL	66 PRL 16 1025	CRENNELL, KALBFLEISCH, LAI, SCARR, SCHU+ (BNL)
HESS	66 PRL 17 1109	+DAHL+HARDY+KIRZ+MILLER (BNL)
BARLOW	67 NC 504 701	+LILLETOL+MONTANET+(CERN+CDF+IR+LIVERPOOL)
REUSCH	67 PL 25 B 357	+FISCHER, GOBBI, ASTBURY, MICHELINI+(ETH+CERN)
DAHL	67 PR 163 1377	+HARDY+HESS+KIRZ+MILLER (LRL)

ALITTI	68 PRL 21 1705	+BARNES, CRENNELL, FLAMINIO, GOLDBERG, + (BNL)
LAI	68 PHILAD. CONF. P. 303	KWAN WU LAI (BNL)
PHELAN	68 THESIS	JAMES J. PHELAN (ANL+ST. LOUIS INTV)
ALSO	68 PRL 21 316	HOANG, EARTLY, PHELAN, ROBERTS + (ANL+CHIC+NDAM)
AGUILAR-	69 PL 29 B 241	M. AGUILAR-RENETEZ, J. BARLOW, + (CERN+CDF)
ALSO	BARLOW 67	
ALSO	69 NP B 14 195	M. AGUILAR-BENITEZ, J. BARLOW, + (CERN+CDF)
HOANG	69 NC 61 A 325	T.F. HOANG (ANL)
HOANG	69 PR 185 1363	+EARTLY, PHELAN, ROBERTS, + (ANL+ILL(ICH)+NDAM)
BADIER	70 NP B 22 512	+BONNET, DREVILLON, RAURILLIER, + (EPOL+INP)
BATON	70 PL 33 B 528	+LAURENS, REIGNIER (SACLAY)
REUSCH	70 PHILA. CONF. P. 185	M. REUSCH (ETH+CERN)
HYAMS	70 PHILA. CONF. P. 41	+KOCH, BEUSCH, + (CERN+MUNT+ETH+LOIC+HAWA)
OH	70 PR D 1 2494	+GARFINKEL, MORSE, WALKER, PRENTICE+ (ISC+TNT)

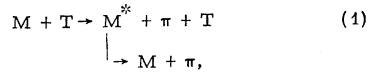
A1(1070)

10 A1 MESON (1070, JPC=1+-) I=1

The Threshold Enhancements
A1, A3, Q, and L

We present here first some general remarks about these four related peaks; second, some remarks specifically about A1.

The four peaks under study generally arise from a reaction of the form



where T is the target (proton or deuteron), M the beam particle, and M^* a resonance in the $M\pi$ system. The relevant quantities for each of the enhancements are given below:

M	$M^*(J^P)$	$M^* \pi$ enhancements mass in (MeV)	Mass threshold M^* and π (MeV)
π	$\rho(1^-)$	A1 (~ 1070)	920
π	$f^0(2^+)$	π_A (~ 1640) or A3	1400
K	$K^*_{890}(1^-)$	K_A (~ 1300) or Q	1030
K	$K^*_{1420}(2^+)$	K_A (~ 1770) or L	1560

These enhancements share the following features:

- 1) Large widths (150 to 400 MeV), with little structure except for compilations discussed in the note on the Q peak.
- 2) Mass about 200 MeV above $M^* \pi$ threshold.
- 3) They are produced strongly only in diffractive processes [reaction(1)]; exceptions are discussed in point 7 below.
- 4) Each enhancement has one predominant decay mode, described in reaction (1), although the Q, which is just above $K\rho$ threshold, can hardly avoid showing some $K\rho$ signal.
- 5) As would be required of a resonance, the enhancement seems to appear in a single partial wave. In fact for all four peaks, preliminary analyses favor a relative s-wave for the $M^* \pi$ system.
- 6) Some of these enhancements [reaction(1)] have

been "reasonably" well described by the Reggeized Deck model, which does not invoke resonance production in the $M^* \pi$ system. However, the duality hypothesis makes the meaning of these fits unclear.

7) There have been some reports of these enhancements produced in "nondiffractive" processes [i. e., other than reaction (1)] (e. g., Q in $\pi p \rightarrow \Lambda K^* \pi$, $A1$ in $K p \rightarrow K \rho \pi$, etc.) Sometimes this is taken as evidence supporting the resonance interpretation, since the production mechanisms are different. However, there are multi-Regge models that can produce low-mass enhancements just as the double-Regge models do. Hence the warnings of (6) apply. This warning was voiced by BERLINGHIERI 69 (footnote 5) as a warning against their own $A1$ peak produced "nondiffractively" by K^+ mesons. See also discussion of CASO 70 below.

8) There has been a considerable amount of confusion related to these enhancements. Usually the mass distributions from different experiments are quite similar except that the mass distribution from one experiment may be a smooth broad peak, while others may have one or more peak(s) on top of it. In the former case it is natural to attribute the whole enhancement to one effect, resulting in broad width and mainly one decay mode. In the latter case one tends to perform a background subtraction and various cuts to remove the broad enhancement, leaving narrow resonances with several decay modes. The width, cross section, and branching ratio are therefore very sensitive to the assumed background since it is high and rapidly varying. This probably explains why two groups sometimes state very different conclusions while their data are actually very similar.

We have already mentioned in the text (Sec. III) and under point 6 above that for these four peaks it is not yet possible to distinguish between two different interpretations (threshold enhancement vs resonance) which are linked by the notion of duality. Our procedure is to enter all four on the Meson Table, with a warning about the interpretation.

Note on Nondiffractive $A1$ Production

Very many authors have shown that $A1$ production according to reaction (1) above ($\pi p \rightarrow p A1$) can be described by a Reggeized Deck model. (Two recent papers are BRANDENBURG 70 and CHIEN-1 70.) Hence it is particularly interesting to see if $A1$ is also produced in other reactions that are free from this complication.

• ANDERSON 69 reported $A1$ production in $\pi^- p \rightarrow p A1^-$ in the backward direction. The $A1$ so produced has much steeper u dependence than exhibited by the other well-known resonances also produced in the same experiment, and this steep $d\sigma/du$ has no simple theoretical explanation. Hence we still accept this result with some reservation.

• DANYSZ 67 and FRIDMAN 68 reported $A1$ in $\bar{p} p \rightarrow 3\pi^+ 3\pi^- \pi^0$. However, the evidence is not overwhelming because of low statistics and high background. The further facts that $A1$ is neither observed in simpler final states (e. g., $\bar{p} p \rightarrow 2\pi^+ 2\pi^- \pi^0$) in the same experiments, nor observed in other $\bar{p} p$ experiments, make the case for $A1$ production in $\bar{p} p$ reactions quite dubious.

• ALLABY 69 measured the deuteron momentum spectra from the reaction $pp \rightarrow d + MM$ at 21.4 GeV/c. Broad peaks were observed with MM at 1100 and 1300 MeV. However, in these measurements the MM system was not observed, hence it is not clear if these peaks were indeed produced in a $\rho\pi$ system. Also SHIH 70 has shown that a double-Regge-pole model produces a broad $A1$ -like enhancement in reactions in which baryon exchange dominates and a $\rho\pi$ system exists in the final state.

• CASO 70 study the following final states in $\pi^- p$ reactions at 11.2 GeV/c:

- (1) $\pi^- p \rightarrow p \pi^- (k\pi^0)$, $k \geq 2$,
- (2) $p \pi^+ \pi^- \pi^-$,
- (3) $p \pi^+ \pi^- \pi^- \pi^0$,
- (4) $p \pi^+ \pi^- \pi^- (k\pi^0)$, $k \geq 2$,
- (5) $n \pi^+ \pi^+ \pi^- \pi^-$,
- (6) $p \pi^+ \pi^+ \pi^- \pi^-$,
- (7) $n \pi^+ \pi^+ \pi^- \pi^- (k\pi^0)$. $k \geq 1$.

Reactions (2) and much of (4) are of course the classical Deck examples which we are trying to avoid in this note. But the remaining reactions each do show an $A1$ peak, and when they are all added, the peak shows up quite well. However, we must restate here more precisely our reservation given in point 7 above. Figure 1(a) shows a double-Regge-pole diagram that contributes to reaction 2. The matrix element is called M_2 , and we know that it more-or-less explains the $A1$ peak. Figure 1(b) shows how the DRP diagram must be modified to explain reaction (6) which involves two more pions. One can simply add one more link to the chain, and put a ρ at the new vertex. The triple-Regge-pole matrix element M_3 can then be written $M_3 = M_2 M_1$, where M_2 already contains the $A1$ peak, and M_1 just involves another propagator. We do not see why the extra factor M_1 should destroy

Mesons

For notation, see illustrated key at beginning of data card listings.

the A1 peak. Then how can we say that reaction 6 is new evidence for a resonant interpretation of the A1 peak?

Another way to draw this reaction 6 is sketched in Fig. 1(c). M_2^1 might have much the same shape as M_2 .

One can look for new evidence where there is charge-exchange at the lower vertex, so that Pomeron exchange is forbidden. $M_2(n)$ should then simulate the A1 peak less well than $M_2(p)$. Unfortunately their charge-exchange reaction (7) does not show a very good A1 signal. We conclude that it is hard to find an independent "nondiffractive" example of A1 production.

In summary, reports on A1 production in "nondiffractive" reactions (i.e., other than $\pi^\pm p \rightarrow \rho^0 \pi^\pm p$) become increasingly abundant. (For details, see GARELICK 70.) But the most striking example is still that of ANDERSON 69, and even here we have stated some reservations. CASO 70 also looks quite convincing. In general, contrary to the overwhelming evidence of A1 production in $\pi^\pm p \rightarrow \rho^0 \pi^\pm p$ reactions reported in every π^\pm experiment, A1 production in "nondiffractive" reactions has been reported in very few experiments with much lower statistics. More experiments and re-examination of the "non-reporting" experiments are needed; e.g., it is important to examine other high-energy πp experiments for reactions 1 to 7 to compare with CASO 70.

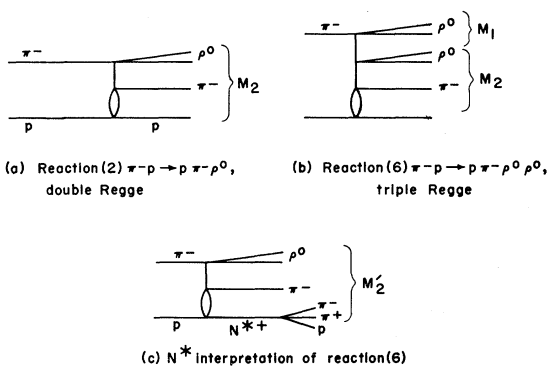


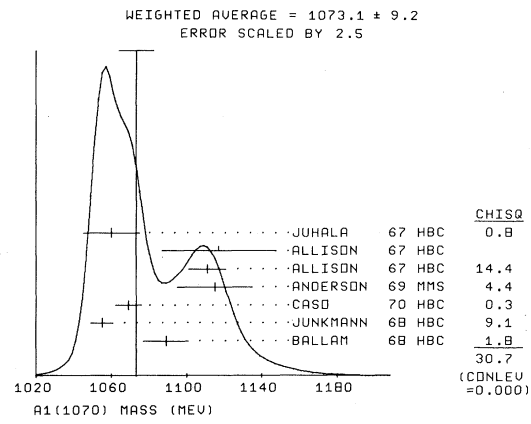
Fig. 1. Regge pole diagrams for A1 production.

10 A1 MESON MASS (MEV)

MASS AND WIDTH MIGHT HAVE LARGE SYSTEMATIC ERRORS DUE TO COMPLICATED BEHAVIOR OF BACKGROUND.

PRODUCED BY PIONS, RESONANCE INTERP. CONFUSED BY DECK EFFECT

PRODUCED BY π^+					
(1089.0)		ADERHOLZ	64 HBC	+ 4.0 $\pi^+ p$	6/68
(1080.)	APPROX.	BOESEBECK	68 HBC	+ 8 $\pi^+ p$	1/71*
(1040.0)		ARMENISE	70 HBC	+ 0.9 $\pi^+ n$	-- A1 P
PRODUCED BY π^-					
(1060.)		ASCOLI	68 HBC	- 0.5 $\pi^- p$	6/68
1089.0	12.0	BALLAM	68 HBC	- 16.0 $\pi^- p$	9/68
(1090.)	APPROX.	CHUNG	68 HBC	- 3.2+4.2 $\pi^- p$	2/67
1055.0	5.0	JUNKMANN	68 HBC	- 16. $\pi^- p$	5PI 9/69
(1119.)	(30.)	KEY	68 HBC	- 3 $\pi^- p$	9/68
S	SHOULDER ON A2 ONLY				
1069.0	7.0	CASO	70 HBC	- 11.2 $\pi^- p$	5/70
(1120.0)		CRENNELL	70 HBC	- 6. $\pi^- p$	PI-F PI 5/70
PRODUCED BY PIONS, BACKWARDS SCATT. NO DECK EFFECT.					
1115.0	20.0	ANDERSON	69 MMS	- 16 $\pi^- p$	BACKW9 8/69
PRODUCED BY PBARs, SEE TYPED NOTE.					
(1054.)	(7.)	DANYSZ	67 HBC	++ 3.3, 6 PBAR P	7/67
(1042.)	(21.)	FRIDMAN	68 HBC	++ 5.7 PBAR P	6/68
PRODUCED BY K^- , SEE TYPED NOTE.					
1111.	10.	ALLISON	67 HBC	+ 6 $K^- p$	LAM +5 PI 1/68
1117.	30.	ALLISON	67 HBC	+ 6 $K^- p$	LAM 44 PI 1/68
1060.	15.	JUHALA	67 HBC	+ 0.4, 6-5 $K^- p$	SRIDDY 1/68
PRODUCED BY K^+ , SEE TYPED NOTE.					
K^+ (1030.0)	(20.0)	ALEXANDER	69 HBC	+ 9 $K^+ p$	9/69
K^+ (1030.0)	(20.0)	BERLINGHI	69 HBC	+ 0.12, 7 $K^+ p$	9/69
K^+ FOR CONTRADICTIONARY EVIDENCE SEE		RABIN	70 AND TYPED NOTE.		
AVG	1073.1	9.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.5) (SEE IDEOGRAM BELOW)		



10 A1 MESON WIDTH (MEV)

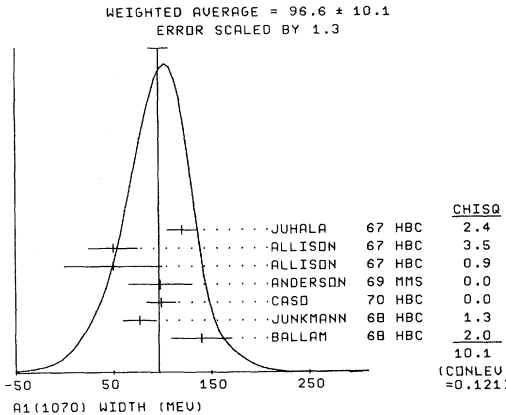
SEE NOTE UNDER A1 MESON MASS.

PRODUCED BY PIONS, RESONANCE INTERP. CONFUSED BY DECK EFFECT

PRODUCED BY π^+						
(80.0)		ADERHOLZ	64 HBC	+ 4.0 $\pi^+ p$	6/68	
(130.)	APPROX.	BOESEBECK	68 HBC	+ 8 $\pi^+ p$	1/71*	
(50.0)	OR LESS	ARMENISE	70 HBC	+ 0.9 $\pi^+ n$	-- A1 P	
PRODUCED BY π^-						
140.0	31.0	BALLAM	68 HBC	- 16.0 $\pi^- p$	9/68	
(125.)	APPROX.	CHUNG	68 HBC	- 3.2+4.2 $\pi^- p$	2/67	
77.0	17.0	JUNKMANN	68 HBC	- 16. $\pi^- p$	5PI 9/69	
K (76.)	(46.)	KEY	68 HBC	- 3.0 $\pi^- p$	11/67	
K	SHOULDER ON A2 ONLY					
99.0	15.0	CASO	70 HBC	- 11.2 $\pi^- p$	5/70	
PRODUCED BY PIONS, BACKWARDS SCATT. NO DECK EFFECT.						
98.0	45.0	20.0	ANDERSON	69 MMS	- 16 $\pi^- p$	BACKW9 8/69
PRODUCED BY PBARs, SEE TYPED NOTE.						
(33.)	(19.)	DANYSZ	67 HBC	++ 3.3, 6 PBAR P	7/67	
(130.)	APPROX.	FRIDMAN	68 HBC	++ 5.7 PBAR P	6/68	
PRODUCED BY K^- , SEE TYPED NOTE.						
50.	50.	ALLISON	67 HBC	+ 6 $K^- p$	LAM +4 PI 1/68	
50.	25.	ALLISON	67 HBC	+ 6 $K^- p$	LAM +5 PI 1/68	
120.	15.	JUHALA	67 HBC	+ 0.4, 6-5 $K^- p$	SRIDDY 1/68	
PRODUCED BY K^+ , SEE TYPED NOTE.						
(160.0)	(20.0)	ALEXANDER	69 HBC	+ 9 $K^+ p$	9/69	
B (120.0)	(30.0)	BERLINGHI	69 HBC	+ 12.7 $K^+ p$	9/69	
K^+ FOR CONTRADICTIONARY EVIDENCE SEE		RABIN	70 AND TYPED NOTE.			
(130.0)	(20.0)	BERLINGHI	69 HBC	+ 0.12, 7 $K^+ p$	9/69	
AVG	96.6	10.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)			

Mesons

For notation, see illustrated key at beginning of data card listings.



10 A1 PARTIAL DECAY MODES

Mode	Decay Masses
P1 A1 INTO RHO PI	765+139
P2 A1 INTO KBAR K	493+497
P3 A1 INTO ETA PI	548+139
P4 A1 INTO ETA PRIME PI	957+139
P5 A1 INTO 3 PI	139+139+139

10 A1 BRANCHING RATIOS

R1 A1 INTO (KBAR K)/(RHO PI)	DAHL	67	HBC	-	4.0	PI- P	.10/66
R1 A1 INTO (KBAR K)/(RHO PI)							

REFERENCES FOR A1

ADERHOLTZ 64 PL 10 226	AACH+BERL+BIRM+RONN+DESY+HAMB+IMP.COL+ MPI
DEUTSCHM 56 PL 20 82	DEUTSCHMANN, STEINBERG + (AACH+BERL+GERN)
ALLISON 67 PL 258 619	+CRUIZ+ (OX+MUN+BIRM+RUTH+GLASG+LON(IC))
DAHL 67 PR 153 1377	+HARDY+HES+KIRZ+MILLER (LRL)
DANYSZ 67 NC 51 A 801	DANYSZ+FRENCH+SIMAK (CERN)
JUHALA 67 PRL 19 1355	+LEACOCK+RHODE+KOPPELMAN+ (IOWA+COL)
ASCOLI 68 PRL 21 113	+CRAWLEY, KRUSE, MORTARA, SCHAFFER, + (ILLINOIS)
BALLAM 68 PRL 21 934	+ARROYO, CHADWICK, FRIES, GUTRAGROSSI AN+ (SLAC)JP
RHODES 68 NP B 4 501	ROESERCK, DEUTSCHMANN, + (AACH+BERL+GERN)
CASO 68 NC 54 A 983	+CRUTE+CORDS+DIAZ+ (GENOVA+HAMB+MIL+SACL)
CHUNG 68 PR 166 1491	S. U. CHUNG, O. DAHL, J. KIRZ, D. N. MILLER (LRL)
FRIDMAN 68 PR 167 1258	+MAURER, MICHALON, DUDET+ (HEIDELB+STRASBOURG)
JUNKMANN 68 NP 88 471	+COCCONI+ (AACH+BERL+GERN+WAR S)
KEY 68 PR 166 1430	+PRENTICE+COOPER+MANNER+WALKER+(TO+ANL+MIS)
ALEXANDE 69 PR 183 1168	G. ALEXANDER, A. FIRESTONE, G. GOLDBERGER (LRL)
ALLARY 69 PL 298 1398	+RIND+DODDENS+DUFET+KLOVNING+... (CERN)
ANDERSON 69 PRL 22 1390	+KOLLINS+ (BNL+CERN)
BERLINGH 69 PRL 23 42	BERLINGHI, FARBER, + (RICH)
JUHALA 69 PR 184 1461	+LEACOCK, RHODE, KOPPELMAN, LIBBY, + (AMES+COL)
ARMENISE 70 LNC 4 199	+GHIDINI, FORINO, CARTACCI, + (BARI+BGNA+FRIZ)
CASO 70 LNC 3 707	+CORDS, COSTA, OARO+ (GENOVA+DESY+HAMB+MIL+SACL)
ALSO CASO 68	
CRENNELL 70 PRL 24 781	+KARSHON, LAI, SCARR, SIMS (BNL)
RATIN 70 PRL 24 925	+GALTIERI, D'FREDDO, FLATTE, FRIEDMAN+ (LRL)
SHH 70 BNL 14059-REV	+YOUNG (BNL)

PAPERS NOT REFERRED TO IN DATA CARDS

REFERENCES ON A1.5 (1170)

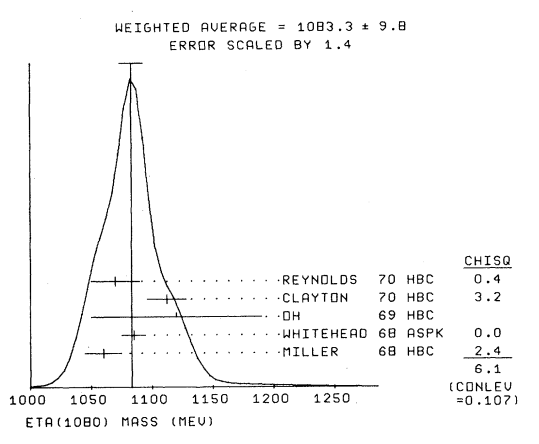
BUTTERW 67 HEIDELB. CONF. P. 28	REVIEW TALK ON MESONS AT HEIDELBERG CONF.
CASON 67 PRL 19 880	+LAMS, ATSWAS, DERADO, GROVES, + (INDTREFAME)
ASCOLI 68 PRL 21 113	+CRAWLEY, KRUSE, MORTARA, SCHAFFER, + (ILLINOIS)
DONALD 68 PL 26 B 327	+FROESE, BETTINI, + (LIVERPOOL, OSLO, PADUA)
VON KR 68 PL 278 253	+MIYASHITA, KOPPELMAN, MARSHALL, LIBBY (COL)
JUNKMANN 68 NP 88 471	+COCCONI+ (AACH+BERL+RONN+GERN+MAR S)
ARMENISE 69 LNC 2 501	+GHIDINI, FORINO, CARTACCI, + (BARI+BGNA+FRIZ)
GALLDWAY 70 PR D 1 3077	+MOTT, ALYEA, LEE, MARTIN, PRICKETT (IND)

$\eta_N(1080)$
→ $\pi\pi$

30 ETA N (1080, JPC=N+1) I=0 J GREATER THAN 1
SOME EXPERIMENTS SUGGEST J=2.
SEE NOTE ABOVE ETA(1070,0+)
OMITTED FROM TABLE

30 ETA N MASS (MEV)

M	Year	Method	Value
M	1060.0	15.0	MILLER 68 HRC
M	70 1085.0	10.0	WHITEHEAD 68 ASPK
M	1120.0	100.0	OH 69 HBC
M	1112.0	16.0	CLAYTON 70 HBC
M	(1080.0)		DIAZ 70 HBC
M	1070.0	20.0	REYNOLDS 70 HBC
M	AVG	1083.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)



30 ETA N WIDTH (MEV)

W	Year	Method	Value
W	(70.0)	OR LESS	MILLER 68 HRC
W	(25.0)	OR LESS	WHITEHEAD 68 ASPK
W	150.0	100.0	OH 69 HBC
W	(80.0)		CLAYTON 70 HBC
W	(80.0)		DIAZ 70 HBC
W	85.0	35.0	REYNOLDS 70 HRC
W	AVG	98.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

REFERENCES FOR ETA N

MILLER 68 PRL 21 1489	+GUTAY, JOHNSON, KENNEY+ (PURDUE+NDAME+SLAC)
WHITEHEAD 68 NC 53 A 817	C. WHITEHEAD + (HARWELL+STHAMP+U.C. LON)
OH 69 PRL 23 331	+WALKER, CARROLL, FIREBAUGH, + (WISC+TNT)
BATON 70 PL 33 B 528	+LAURENS, REIGNIER (SACLAY)
DIAZ 70 NP B 16 239	+SAVILLEY, LAROSSE, MONTANET+ (CERN+GDF)
CLAYTON 70 NP B 22 85	+MASON, MUIRHEAD, RIDGWOODS, + (LIV+ATFN)
OH 70 PR D 1 2494	+GARFINKEL, MORSE, WALKER, PRENTICE (WISC+TNT)
REYNOLDS 70 NP B 21 77	+ALBRIGHT, BRADLEY, + (ATHOS+FLAS+MUN+COL)

$A_{1.5}(1170)$

44 A 1.5 (1170, JPC= -) I=1
THIS ENTRY LISTS REFERENCES TO PEAKS OF LOW STATISTICAL SIGNIFICANCE IN THE 3 PI SYSTEM BETWEEN THE A1 AND THE A2. OMITTED FROM TABLE.

REFERENCES ON A1.5 (1170)

BUTTERW 67 HEIDELB. CONF. P. 28	REVIEW TALK ON MESONS AT HEIDELBERG CONF.
CASON 67 PRL 19 880	+LAMS, ATSWAS, DERADO, GROVES, + (INDTREFAME)
ASCOLI 68 PRL 21 113	+CRAWLEY, KRUSE, MORTARA, SCHAFFER, + (ILLINOIS)
DONALD 68 PL 26 B 327	+FROESE, BETTINI, + (LIVERPOOL, OSLO, PADUA)
VON KR 68 PL 278 253	+MIYASHITA, KOPPELMAN, MARSHALL, LIBBY (COL)
JUNKMANN 68 NP 88 471	+COCCONI+ (AACH+BERL+RONN+GERN+MAR S)
ARMENISE 69 LNC 2 501	+GHIDINI, FORINO, CARTACCI, + (BARI+BGNA+FRIZ)
GALLDWAY 70 PR D 1 3077	+MOTT, ALYEA, LEE, MARTIN, PRICKETT (IND)

Mesons

For notation, see illustrated key at beginning of data card listings.

Table with columns for F WIDTH (MEV), various author names (e.g., SELVJE, RONDAR, ACCENSI), and their respective HBC values and other parameters.

WEIGHTED AVERAGE = 159.5 ± 10.7
ERROR SCALED BY 2.4

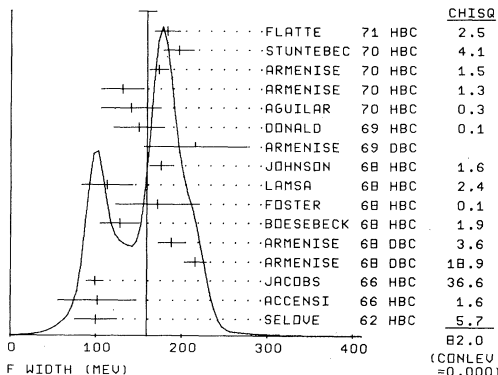


Table listing author names and HBC values corresponding to the data points in the graph above.

Table titled '5 F PARTIAL DECAY MODES' with columns for decay modes (e.g., F INTO PI+ PI-, F INTO K KBAR) and decay masses.

Table titled '5 F BRANCHING RATIOS' with columns for branching ratios and average values.

Detailed text notes regarding the data analysis, including mentions of 'ASCOLI' and 'DETERMINATION DIFFICULT BECAUSE PROXIMITY OF A2 WHICH HAS SAME NEUTRAL (K KBAR) MODES...'.

Large table of references for F mesons, listing author names, HBC values, and other parameters.

D(1285)

8 D MESON (1285; JPC= +) I=0

Table titled '8 D MESON MASS (MEV)' with columns for mass values and author names.

WEIGHTED AVERAGE = 1286.4 ± 4.3
ERROR SCALED BY 1.4

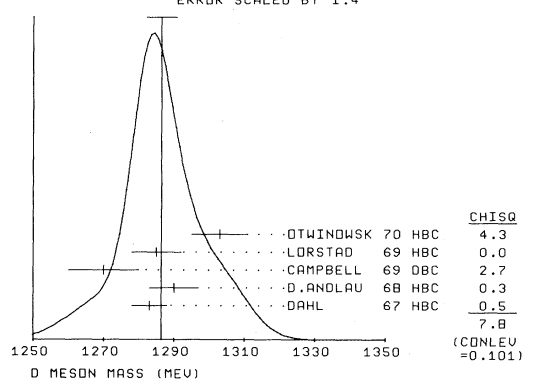


Table listing author names and HBC values corresponding to the data points in the graph above.

Mesons

For notation, see illustrated key at beginning of data card listings.

B D MESON WIDTH (MEV)

W	35.0	10.0	DAHL	67 HRC	1.6-4.2 PI- P	10/66
W	30.0	5.0	D. ANDLAU	68 HRC	1.2 PBAR P, 5-6 PFS	8/68
W	(40.0)		DEFIUX	69 HRC	1.2 PR P, 7 PI	3/69
W	30.0	15.0	CAMPBELL	69 ORC	2.7 PI+ D	8/69
W	60.0	15.0	LORSTAD	69 HRC	0.7 PB P, 4, 5-BODY	9/69
W	(52.0)	(29.0)	OTWINDWSK	70	8 PI+ P, P+6PI	9/69
W	44.0	24.0	OTWINDWSK	70 HBC	8 PI+ P, P+6PI	9/69
W	AVG	33.4	4.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		

B D MESON PARTIAL DECAY MODES

P1	D MESON INTO K KBAR PI	DECAY MASSES
P2	D MESON INTO PI PI RHO	497+ 497+ 134
P3	D MESON INTO ETA PI PI	134+ 134+ 765
P4	D MESON INTO DELTA(962) PI	58+ 134+ 134
		956+ 134

B D MESON BRANCHING RATIOS

R1	D MESON INTO (PI PI RHO) / (K KBAR PI)	CHARGED PI ONLY	10/66
R1	(4.0) OR LESS	DAHL	67 HRC
R1	THIS IS FOR (RHO0 PI+ PI-)/(K KBAR PI0)	DONALD	69 HRC
R2	D MESON INTO (K KBAR PI)/(ETA PI PI)	DEFIUX	68 HRC
R2	(0.124) (0.035)		1.2 PR P, 7 PI
R3	D MESON INTO (DELTA PI)/(ETA PI PI)	AMMAR	70 HRC
R3	POSSIBLY SEFN	OTWINDWSK	70 HRC
R3	POSSIBLY SEFN		9 PI+ P, P+6PI

REFERENCES FOR D MESON

HARLOW	67 NC 50 A 701	+MONTANET, D-ANDLAU+(CERN+CDF+IDR+LIVERPOOL)
DAHL	67 PR 163 1377	+HANDY+HESS+KIRZ+MILLER (LRL) JP
	SEE ALSO 65 PRL 14 1074	MILLER, CHUNG, DAHL, HESS, HARDY, KIRZ, (LRL+UCI)
D. ANDLAU	68 PR 8 5 693	+ASTIER, BARLOW, MONTANET+ (CDF+CERN+RAD+LIV) JP
DEFIUX	68 PL 28 B 353	+RIVET, SIAUD, CONFORTO, SHIVELY (CDF+IPE+CERN)
CAMPBELL	69 PRL 22 1204	+LICHTMAN, + (PURD)
DONALD	69 NP 8 11 551	+EDWARDS, BURAN, BETTINE, + (LIVP+NSLQ+PAD) JP
LORSTAD	69 NP 8 14 63	B. LORSTAD, D. ANDLAU, ASTIER, + (CDF+CERN) JP
OTWINDWSK	69 PL 29 B 529	S. OTWINDWSKI (WARSAW)
AMMAR	70 PR	+KROPAC, DAVIS, DERRICK+ (KANS+NWES+ANL+WISC)
OTWINDWSK	1207/VI/PH	S. OTWINDWSKI (WARSAW)

A2(1300) 12 A2 MESON (1300, JP=2+-) I=1

There are now several important experiments which disagree on the issue of whether the A2 is split or not. As can be seen from the summary table, the division between splitting and not-splitting experiments is not simply correlated to any physical property such as A2 charge, momentum transfer, beam energy, decay

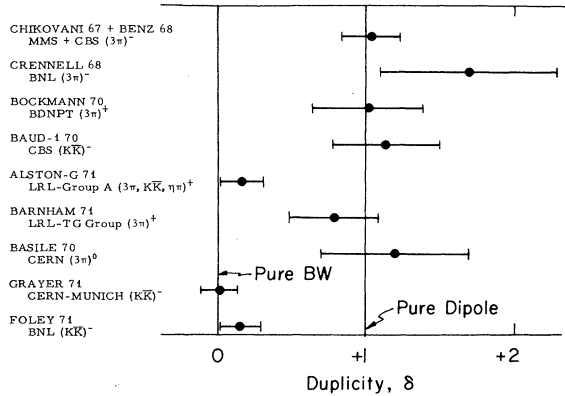
mode, etc. If indeed all experiments are correct (and we have no reason to suspect the contrary) then a symmetric dipole is excluded. It seems possible, by invoking two interfering resonances, to fit the line shape of all the observed A2 peaks. But then the mechanism responsible for the necessary energy-dependence and charge-dependence of the interference phase and/or the degree of coherence remains an unsolved problem.

The experimental evidence for and against splitting is shown in the figure, taken from G. R. Lynch, UCRL 20622. Each mass spectrum is fitted to an unphysical mixture of the form

$$dn/dm = \delta (\text{Dipole}) + (1 - \delta) (\text{BW}) + \text{Background}$$

where (Dipole) and (BW) are best fits to each hypothesis separately.

The "Duplicity" parameter δ is easily visualized and plotted because it spans the conflicting hypotheses. We repeat that it makes no physical sense to mix hypotheses.



SUMMARY OF EXPERIMENTS THAT INVESTIGATED A2 SPLITTING

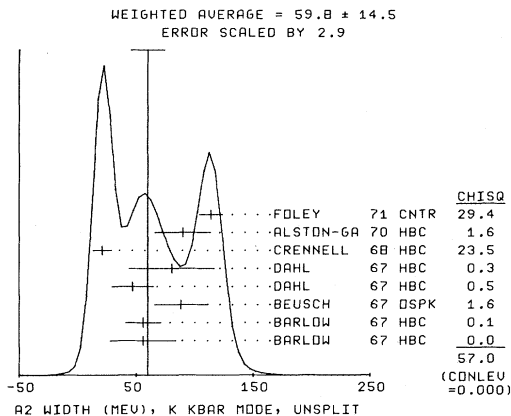
Reaction	Momentum (GeV/c)	Reference	Method	$\Gamma_r/2$ (MeV)	t range (GeV/c) ²	Events in peak	Probability of fit ^a			
							BW	Dipole	2 coh. BW	2 incoh. BW
$\pi^- p \rightarrow X^- p$	6, 7	CHIKOVANI 67	Jac. peak	8.3	0.20-0.29	1400	<<0.1%	≥40%	≥40%	<0.2%
$\rightarrow X^- p$	2.65	BENZ 68	0°	5.2	---	1100				
$\rightarrow K^- K_S^0 p$	7	BAUD-1 70	Jac. peak	<10	0.20-0.29	145	1%	>60%		
$\rightarrow \eta \pi^- p$	11	BAUD-2 70	Jac. peak	20	0.20-0.30	~90				
$\rightarrow \pi^- X^0 p$	6	CRENNELL 68	HBC	10	0.22-0.39	100				
$\rightarrow X^0 n$	3.2	BASILE 70	CNTR	7.5	0.35-0.65	2600	1%	65%		25%
$\pi^+ p \rightarrow \pi^+ \pi^+ \pi^+ p$	7	ALSTON-G 74	HBC	6.7	> 0.2	941	42%	9% } 0.2%		
$\rightarrow \eta \pi^+ p$	7	ALSTON-G 74	HBC	9.2	all	189	33%			
$\rightarrow K^+ K_S^0 p$	7	ALSTON-G 74	HBC	3.8	all	132	4%			
$\rightarrow \rho^0 \pi^+ p$	5	BOCKMANN 70	HBC	5	t' > 0.1	108	20%		63%	
$\bar{p} p \rightarrow K^+ K_1^- \pi^+$	0, 0.7, 1.2	AGUILAR-1 69	HBC	5	---	270	4%	65%		28%
$\rightarrow 2\pi^+ 2\pi^-$	1.2	DONALD 69	HBC	12	---	~100	2%			40%
$\pi^- p \rightarrow K^- K_S^0 p$	17.2	GRAYER 74	ASPK	6.7	<0.7	~950	32%	"Not a dipole"		

a. These confidence levels permit only qualitative comparisons, since they depend on the size and number of bins used and the treatment of the number of adjustable parameters; these choices differ from group to group.

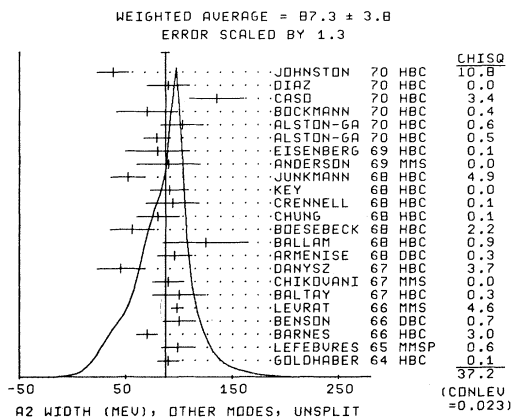
Mesons

For notation, see illustrated key at beginning of data card listings.

12 A2L MESON WIDTH (MEV)
WL A (29.) (10.) CHIKOVANI 67 MMS - 6.7 PI- P 6/68
WL A INCLUDED IN BENT 68 FIT OF 2 COHERENT SYMMETRIC POLES.
WL 22. 5. BENZ 68 MMS - 2.65 PI- P 12/68



12 A2H MESON WIDTH (MEV)
WH A (35.) (10.) CHIKOVANI 67 MMS - 6.7 PI- P 6/68
WH A INCLUDED IN BENT 68 FIT OF 2 COHERENT SYMMETRIC POLES.
WH 12.0 10.0 BENZ 68 MMS - 2.65 PI- P 12/68

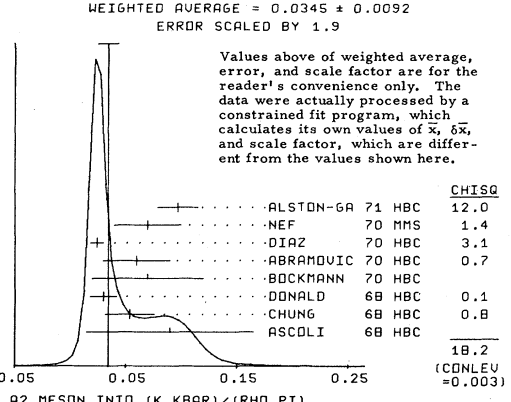


12 A2 MESON WIDTH(MEV), K KBAR, UNSPLIT, UNRESOLVED
WK 60 54.0 28.0 BARLOW 67 HRC +- 1.2 PHAR P, KK 9/67
WK 80 84. 23. 22. BEUSCH 67 OSPK 0 5-12 PI-P,KIKI 7/67

12 A2 MESON WIDTH (MEV), 3PI, ETA PI MODES, SPLITTING UNRESOLVED
W (100.0) ADEPHALZ 64 HRC 4.0 PI+P 6/66
W 90.0 10.0 GILDHARER 64 HRC +- 3.7 PI+- P 6/66

12 A2 MESON BRANCHING RATIOS
R1 A2 MESON INTO (K KBAR) / (RHO PI)
R1 (0.08) OR LESS LANDER 64 HRC + 3.5 PI+P 10/66

12 A2 MESON PARTIAL DECAY MODES
P1 A2 MESON INTO RHO PI 765+ 139
P2 A2 MESON INTO KBAR K 493+ 497



Mesons

For notation, see illustrated key at beginning of data card listings.

Table of meson data cards (R2-R10) including fields for meson name, quantum numbers, and fit parameters.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± ΔP_i, where ΔP_i = √(ΔP_i ΔP_i), while the off-diagonal elements are the normalized correlation coefficients (ΔP_i ΔP_j) / (ΔP_i ΔP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of branching fractions P1 through P4 with error bars and correlations.

REFERENCES FOR A2

References for A2 meson, listing authors, journals, and quantum numbers.

Table of meson data cards (FOSTER-FOSTER) including fields for meson name, quantum numbers, and fit parameters.

A2 I = 2 (1320)

37 A2, 2 (1320) I=2 OR GREATER SEEN AS A BUMP IN RHO- P1 SPECTRUM. EVIDENCE NOT COMPELLING. DROPPED. SEE JAN. 67 EDITION.

E(1422)

6 E MESON (1422, JP=0+) I=0 BAILLON 67 FAVOR JP=0-. DAHL 67 FAVOR 1+ BUT DO NOT EXCLUDE 2-, 0-. LORSTAD 69 FAVOR 0- OR 1+.

Table for 6 E MESON MASS (MEV) with columns for author, mass, and error.

Table for 6 E MESON WIDTH (MEV) with columns for author, width, and error.

Table for 6 E MESON PARTIAL DECAY MODES with columns for decay mode and branching fraction.

Table for 6 E MESON BRANCHING RATIOS with columns for decay mode and ratio.

Table for 6 E MESON PARTIAL DECAY MODES (continued) with columns for decay mode and branching fraction.

Mesons

For notation, see illustrated key at beginning of data card listings.

REFERENCES FOR F PRIME

BAILLON 67 NC 50A 393	+EDWARDS+D.ANDLAU+ASTIER+ (CERN+CDF+IR)
BARASH 67 PR 156 1399	BARASH,IRSCH,MILLET,TAN (COLUMBIA)
DAHL 67 PR 163 1377	+HARDY+HESS+KIRZ+MILLER (LRL II JP)
AI SD 65 PRL 14 1074	MILLER,CHUNG,DAHL,HESS,HARDY,KIRZ (LRLAUC)
FRENCH 67 NC 52A 438	+KINSON+MCDONALD+RIDDIFORD+ (CERN+BIRM)
FOSTER 68 NP B 8 174	+GAVILLET,LARROSSE,MONTANET,+ (CERN+CDEF)
BETTINI 69 NC 62 A 1038	+CRESTI,LIMENTANI,BERTAUZA,BIGI+(PADU+PISA)IC
LORSTAD 69 NP B 14 63	B.LORSTAD,D.ANDLAU,ASTIER,+ (CDF+CERN) JP

K_sK_s(1440)
ρρ(1410)

29 KSKS(1440) AND RHO RHO(1410) (JPG=V+) I GTE 0

EVIDENCE NOT YET COMPELLING. OMITTED FROM TABLE. IF RHO RHO AND K_s K_s ARE MODES OF THE SAME RESONANCE, THEN I=0.

29 KSKS AND RHO RHO MASS (MEV)

M	-----RHO RHO MODE-----	BETTINI 66 DBC 0 0. PBAR P TO 5PR	9/66
M	(1410.0)		
M	-----KS KS MODE-----	ABRAMS 67 HBC 4.25 K- P	5/67
M	POSSIBLY SEEN		
M	THE AUTHORS ASSOCIATE THE PEAK WITH THE F PRIME, BUT BACKGROUND ESTIMATION IS DIFFICULT		
M	1412. 23.	BARLOW 67 HBC 1.2 PBAR P	5/67
M	1439.0 5.0 6.0	REUSCH 67 OSPK 5.7, 1.2 P1-P	9/67
M	(1425.0)	FOLEY 71 CNTR - 20.3 P1- P,K- KS	2/71*
M	AVG 1437.5 5.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

29 KSKS AND RHO RHO WIDTH (MEV)

W	-----RHO RHO MODE-----	BETTINI 66 DBC 0 0. PBAR P TO 5PR	9/66
W	(90.0)		
W	-----KS KS MODE-----	BARLOW 67 HBC 1.2 PBAR P	5/67
W	100. 70.	REUSCH 67 OSPK 5.7, 1.2 P1-P	9/67
W	43.0 17.0 18.0	FOLEY 71 CNTR - 20.3 P1- P,K- KS	2/71*
W	(20.0) OR LESS		
W	AVG 46.4 17.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

REFERENCES FOR KSKS(1440) AND RHO RHO(1410)

BETTINI 66 NC 42A 695	+CRESTI,LIMENTANI,LORIA,PERUZZO+(PAD+PISA)
ABRAMS 67 PRL 18 620	+KEHOE,GLASSER,SECHI-ZORN,WOLSKY (MARYLAND)
BARLOW 67 NC 50 A 701	+MONTANET,D.ANDLAU+(CERN+CDF+TOR+IVERPOOL)
REUSCH 67 PL 25 B 357	+TSCHER,GORBI,ASTBURY,MICHELINI+(ETH+CERN)
DONALD 69 NP B 11 551	+EDWARDS,BURAN,BETTINI,+ (LIV+OSLO+PADU)
FOLEY 71 PRL 26 413	+LOVE,OZARI,PLATNER,LINDENBAUM,+ (BNL+CUINY)

f'(1514) 13 F PRIME (1514, JPG=2++) I=0

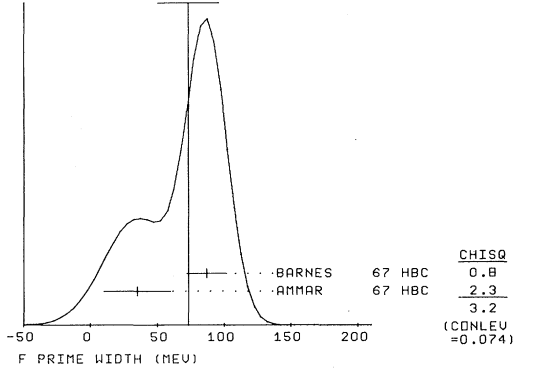
13 F PRIME(1514) MASS (MEV)

M	14(1490.0)	CRENNELL 66 HBC 6.0 P1- P	8/66
M	5(1460.) (10.)	ABRAMS 67 HBC 4.25 K- P	5/67
M	BACKGROUND ESTIMATION DIFFICULT.		
M	1515.0 7.0	AMMAR 67 HBC 5.5 K- P	5/67
M	70 1513.0 7.0	BARNES 67 HBC 4.6, 5. K- P	10/67
M	AVG 1514.0 4.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

13 F PRIME(1514) WIDTH (MEV)

W	5 (53.) (18.)	ABRAMS 67 HBC 4.25 K- P	5/67
W	BACKGROUND ESTIMATION DIFFICULT.		
W	35.0 25.0	AMMAR 67 HBC 5.5 K- P	5/67
W	70 87.0 15.0	BARNES 67 HBC 4.6, 5. K- P	10/67
W	AVG 73.2 22.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)	

WEIGHTED AVERAGE = 73.2 ± 22.9
ERROR SCALED BY 1.8



13 F PRIME PARTIAL DECAY MODES

P1	F PRIME INTO PI+ PI-	DECAY MASSES
P2	F PRIME INTO K KBAR	139+ 139
P3	F PRIME INTO K K*(890)	437+ 497
P4	F PRIME INTO ETA ETA	493+ 892
P5	F PRIME INTO PI PI ETA	548+ 548
P6	F PRIME INTO PI K KBAR	139+ 139+ 548
		139+ 497+ 497

13 F PRIME BRANCHING RATIOS

R1	F PRIME INTO (PI+ PI-)/(K KBAR)	AMMAR 67 HBC	5.5 K- P	CL=67	9/67
R1	(0.21) OR LESS	BARNES 67 HBC	4.6, 5.0 K- P		10/67
R3	F PRIME INTO (ETA ETA)/(K KBAR)	BARNES 67 HBC	4.6, 5.0 K- P		10/67
R3	(0.50) OR LESS				
R4	F PRIME INTO (PI PI ETA)/(K KBAR)	AMMAR 67 HBC	CL=0.67		10/67
R4	(0.3) OR LESS	BARNES 67 HBC	4.6, 5.0 K- P		10/67
R4	0.25 0.13				
R5	F PRIME INTO (PI K KBAR + K K*(890))/(K KBAR)	AMMAR 67 HBC	CL=0.67		10/67
R5	(0.4) OR LESS	BARNES 67 HBC	4.6, 5.0 K- P		10/67
R5	(0.14) OR LESS				

REFERENCES FOR F PRIME

CRENNELL 66 PRL 15 1025	+ KALBFLEISCH,LAI,SCARR,SCHUMANN + (BNL II)
ABRAMS 67 PRL 18 620	+KEHOE,GLASSER,SECHI-ZORN,WOLSKY (MARYLAND)
AMMAR 67 PRL 19 1071	+DAVIS,HWANG,DAGAN,DEWICK + (INU+ANL) JP
BARNES 67 PRL 19 964	+DORNAN,GOLDBERG,LEITNER + (BNL+SYRACUSE)ICJP
ALITTI 68 PRL 21 1705	+BARNES,CRENNELL,FLAMINI,GOLDBERG+ (BNL)
LORSTAD 69 NP B 14 63	B.LORSTAD,D.ANDLAU,ASTIER,+ (CDF+CERN) JP
SCOTTER 69 NC 62 A 1057	+ERSKINE,PALER,+ (RIRM+GLAS+LOIC+MPI+OXF)

π/ρ(1540)
or F₁

47 PI/RHO (1540, JPG= 1) I=1

NAMED F₁ BY AGUILAR 69

JP=2-1+ FAVORED, 0- LESS PROBABLE.

ALSO CONTAINS PHOTOPRODUCED HUMP AT 1550 MEV, I UNKNOWN

47 PI/RHO (1540) MASS (MEV)

M	10(1490.0) (20.0)	ADERHOLZ 69 HBC + 8 PI+ P,KBARPI	11/69
M	142 1540.0 5.0	AGUILAR 69 HBC 0.7PBARP,KBARPI	11/69
M	0 (1550.0) (40.0)	DAVIER 69 STRC 04.5-18 BREMS,4PI	11/69
M	0	WITH STATISTICS INCREASED FOURFOLD (MAY 70), DAVIER ET AL STILL	
M	0	FIND A BROAD FOUR PION ENHANCEMENT, BUT LITTLE EVIDENCE FOR AL+ STRUCTURE. NO INFORMATION ON ISOSPIN.	

47 PI/RHO (1540) WIDTH (MEV)

W	10 (85.0) (39.0)	ADERHOLZ 69 HBC + 8 PI+ P,KBARPI	11/69
W	142 40.0 15.0	AGUILAR 69 HBC 0.7PBARP,KBARPI	11/69
W	0 (260.0) (110.0)	DAVIER 69 STRC 04.5-18 BREMS,4PI	11/69
W	0	SEE NOTE UNDER MASS ABOVE.	

47 PI/RHO (1540) PARTIAL DECAY MODES

P1	PI/RHO (1540) INTO K KBAR PI	DECAY MASSES
P2	PI/RHO (1540) INTO K*(890) KBAR	134+ 497+ 497

REFERENCES FOR PI/RHO (1540)

ADERHOLZ 69 NP B 11 259	+BARTSCH,+ (AACH+BERL+CERN+KRAK+WARS)
AGUILAR 69 PL 29 B 379	+BARLOW,JACOBS,D.ANDLAU,ASTIER+ (CERN+CDEF)
AGUILAR 69 NP B 14 195	+BARLOW,JACOBS,D.ANDLAU,ASTIER+ (CERN+CDEF)
DAVIER 69 SLAC-PUB-666	+DERADO,FRIES,LIU,MIZLEY,ODIAN,PARK+ (SLAC)

π_A(1640)
or A3

34 PI (1640, JPG=A-) I = 1

This peak is seen in the ππ spectrum in π[±]p → π[±]π[±]p reactions. It is only 230 MeV above ππ threshold, and its interpretation is ambiguous. See note on A1, Q, A3, and L in these listings, just before the A1.

Recently CRENNELL 70 clearly demonstrated that A3 is an s-wave ππ enhancement. A recent compilation on high-energy π[±]p data (CHIEN-2 70) also shows that ππ is the only dominant mode. BRANDENBURG 70 and CHIEN-1 70 both showed that A3 production in π[±]p at 7, 13, and 20 GeV/c can be described by a Reggeized Deck model.

Mesons

For notation, see illustrated key at beginning of data card listings.

A3 production is also reported in pp and other reactions, but the statistics are relatively poor and results inconclusive.

THIS ENTRY CONTAINS G=-1 PEAKS AND THE R2 PEAK. BARTSCH 6R FIND BEST FIT WITH JP=2-. NEXT BEST 1+, 0-, 3+ NOTE THAT (OMEGA PI PI) PEAKS HAVE DIFFERENT MASS WIDTH

Table with 4 columns: M, W, A, M. Rows include mass measurements for various experiments like FORINO, VETLITSKY, BALTAY, etc. Includes a weighted average of 1640.1 ± 6.2.

WEIGHTED AVERAGE = 1640.1 ± 6.2 ERROR SCALED BY 1.1

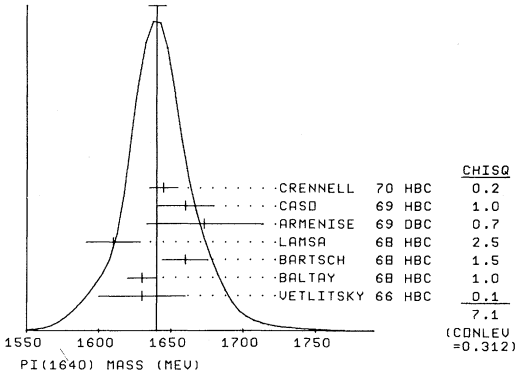


Table with 4 columns: M, W, A, M. Rows include mass measurements for various experiments like DANYSZ, DIBAL, YOST, BARNES, BEKETOV. Includes a weighted average of 1689.3 ± 7.1.

Table with 4 columns: W, W, A, W. Rows include width measurements for various experiments like VETLITSKY, BALTAY, BARTSCH, LAMSA, ARMENISE, CASO, CRENNELL, MIYASHITA, BEKETOV. Includes a weighted average of 107.7 ± 18.7.

Table with 4 columns: W, W, A, W. Rows include width measurements for various experiments like DANYSZ, YOST, BARNES. Includes a weighted average of 45.1 ± 11.5.

Table with 4 columns: P1, P2, P3, P4, P5, P6, P7, P8, P9. Rows list partial decay modes for pi(1640) into various meson pairs like 3 pi, rho pi, eta pi, etc.

34 PI (1640) BRANCHING RATIOS

Table with 4 columns: R2, R2, R2, R3, R3, R3, R3, R3, R3, R4, R4, R5, R5, R5, R5, R6, R6, R6, R7, R7, R7, R8, R8, R8, R9, R9, R9. Rows list branching ratios for various decay channels like INTO (PI+- RH0)/ (ALL PI+- PI+ PI-), INTO (PI+- F)/ (ALL PI+- PI+ PI-), etc.

REFERENCES FOR PI(1640)

Table with 4 columns: FORINO, FOCACCI, LEVRAT, LURATTI, VETLITSKY, DANYSZ, DUBAL, BALTAY, BARTSCH, CASO, FERBEL, IOFFREDDO, LAMSA, YOST, ARMENISE, BARNES, CASO, BEKETOV, BRANDENB, CRENNELL, CHIENI, CHIENI, MIYASHITA. Rows list references for various experiments.

phi_N(1650) -> rho^0 pi^0

THIS 3 PI RUMP OVERLAPS IN MASS THE PI(1640) F PI BUMP CALLED THE A3, BUT IN SOME EXPTS. ONE CAN ESTABLISH THE ENHANCEMENT IS RHO 0 PI INSTEAD OF F PI, SO THE PHI(1650) AND THE PI(1640) HAVE DIFFERENT SPIN. MATTHEWS 71 SUGGEST JP = NORMAL, A POSSIBLE RECURRENCE OF THE OMEGA MESON.

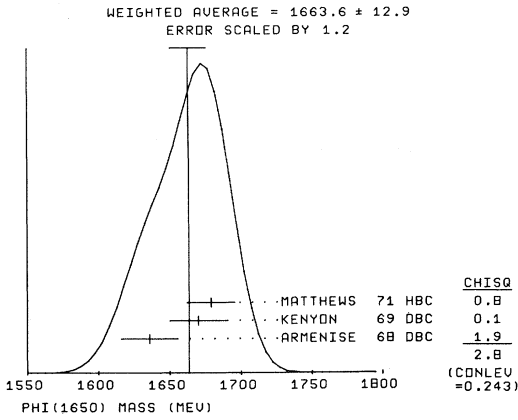
Table with 4 columns: M, W, A, M. Rows include mass measurements for various experiments like ARMENISE, KENYON, ARMENISE, GORDON, GORDON. Includes a weighted average of 1663.6 ± 12.9.

Table with 4 columns: W, W, G, W, W, W, W, W. Rows include width measurements for various experiments like ARMENISE, KENYON, GORDON, MATTHEWS. Includes a weighted average of 141.4 ± 17.1.

Table with 4 columns: P1, P2, P3. Rows list partial decay modes for phi(1650) into various meson pairs like 3 pi, rho pi, etc.

Mesons

For notation, see illustrated key at beginning of data card listings.



45 PHI (1650) BRANCHING RATIOS

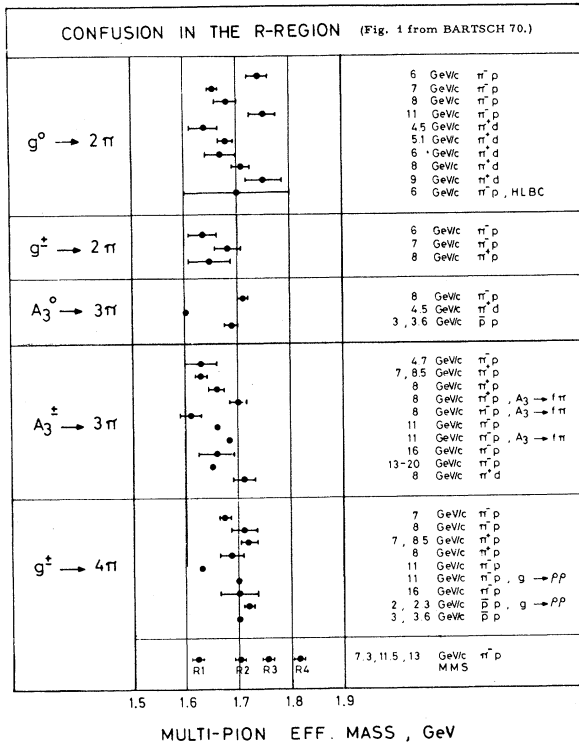
R1	PHI (1650) INTO (5 PI)/(13 PI)				
R1	0.10	0.10	KENYON 69 DBC	0 8, PI + D	8/69
R2	PHI (1650) INTO (RHO PI) / (3 PI)				
R2 G	(0.75) (0.11)		GORDON 70 DBC	0 4.2 PI + D	1/71*
R2	100 (0.75) OR GREATER		MATTHEWS 71 HBC	06.95 PI 0.2P 3PI	1/71*
R2 G	NOT CERTAIN IF PHI(1650) IS OBSERVED IN THIS EXPERIMENT				

45 PHI (1650) CROSS SECTIONS

CS	FOR A COMPILATION SEE MATTHEWS 71 HBC 06.95 PI 0.2P 3PI 1/71*				
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REFERENCES FOR PHI (1650)

ARMENISE 68 PL 268 336	+GHIDINI, FORINO* (BARI+BOLOGN+FIRENZ+ORSAY)
KENYON 69 PRL 23 146	+KINSON, SCARR, + (SNL+ORUC+ORNLL)
ARMENISE 70 LNC 4 199	+GHIDINI, FORINO, CARTACCI, + (BARI+BGNA+FRZ)
GORDON 70 CD 1195 179	THEISIS, ILLINDIS (ILL)
MATTHEWS 71 SUBM. TO PR	+PRENTICE, YOUNG, CARROLL, FIFRAUG (TNTD+MISC)



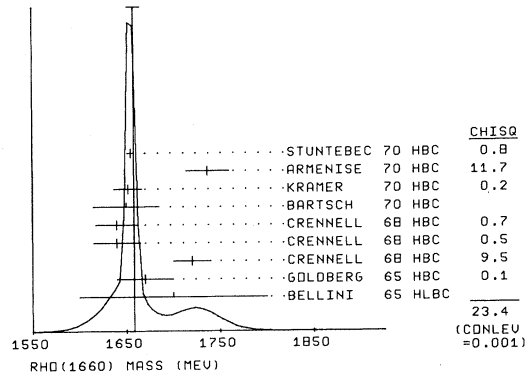
$\rho_N(1660)$ or g

15 RHO (1660, JPG+ N+ I=1)
THIS ENTRY CONTAINS THE R1, G2 PI) AND K KBAR PEAKS. FOR POSSIBLE 4 PI MODES SEE RHO(1710)
CRENNELL 68 SUGGEST JP=3- FROM THE PI PI SCATTERING ANGLE DISTR. BARTSCH 70 USES OPE MODEL, SUGGESTS ELASTICITY LESS THAN 1

15 RHO (1660) MASS (MEV)

M	1700.0	100.0	BELLINI 65 HLBC	0 6.1 PI-P	6/66
M	(1640.0)		FORINO 65 DBC	0 4.5 PI+D	6/66
M	1670.0	30.0	GOLDBERG 65 HBC	0 6 PI+D, 8 PI-P	
M R	(1630.)	(15.)	NUBAL 67 MMS	- 7, 11, 5, 12PI- P	7/67
M R	R1 PEAK FROM CERN MMS EXPT. DECAY MODES AND G PARITY UNKNOWN.				
M R	NOTE THAT THE R1 HAS MUCH SMALLER WIDTH THAN THE OTHER ENTRIES.				
M	(1683.)	(13.)	ARMENISE 68 DBC	0 5.1 PI+ D	6/68
M B	(1620.)	(20.)	ROSEBECK 68 HBC	+ 8 PI+ P	6/68
M B	ROSEBECK 68 SUPERSEDED BY BARTSCH 70				
M	1720.0	20.0	CRENNELL 68 HBC	- 6.0 PI- P	12/68
M	1640.0	25.0	CRENNELL 68 HBC	- 6.0 PI- P	12/68
M	1640.0	20.0	25.0	CRENNELL 68 HBC	+ 6.0 PI-P, KBAR K
M	(1655.0)	(10.0)	JOHNSTON 68 HBC	0 7.0 PI- P	6/68
M	131(1640.0)		ADERHOLZ 69 HBC	+ 8 PI+ P, K+K0	8/69
M	(1600.)		BARTSCH 70 HBC	- 8 PI- P	5/70
M	(1750.0)		CASO 69 HBC	0 11. PI- P, N2PI	8/69
M	1650.0	35.0	BARTSCH 70 HBC	+ 8 PI+ P, 2 PI	5/70
M	1682.0	15.0	KRAMER 70 HBC	+ 13.1 PI+ P, 2PI	11/70*
M	1737.0	23.0	ARMENISE 70 HBC	0 9 PI+ N -- GO P	1/71*
M	1655.0	6.0	STUNTEBEC 70 HBC	0 8. PI-P, PI+PI-	1/71*
M	AVG 1658.5 ± 7.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0) (SEE IDEOGRAM BELOW)				

WEIGHTED AVERAGE = 1658.5 ± 7.1
ERROR SCALED BY 2.0



15 RHO (1660) WIDTH (MEV)

W	(40.0)		FORINO 65 DBC	0 4.5 PI+D	6/66
W	180.0	40.0	GOLDBERG 65 HBC	0 6 PI+D, 8 PI-P	
W R	(21.)	OR LESS	LEVRAT 66 MMS	- 7, 12 PI- P	7/67
W R	R1 PEAK FROM CERN MMS EXPT. DECAY MODES AND G PARITY UNKNOWN.				
W R	188.	40.	ARMENISE 68 DBC	0 5.1 PI+ D	6/68
W B	(100.)	(40.)	ROSEBECK 68 HBC	+ 8 PI+ P	6/68
W B	ROSEBECK 68 SUPERSEDED BY BARTSCH 70				
W	200.0	100.0	CRENNELL 68 HBC	- 6.0 PI- P	12/68
W	200.0	100.0	CRENNELL 68 HBC	- 6.0 PI- P	12/68
W	79.0	70.0	25.0	CRENNELL 68 HBC	+ 6.0 PI-P, KBAR K
W	(180.0)	(20.0)	JOHNSTON 68 HBC	0 7.0 PI- P	6/68
W	13 (100.0)		ADERHOLZ 69 HBC	+ 8 PI+ P, K+K0	8/69
W	(200.)		BARTSCH 69 HBC	- 8 PI- P	5/70
W	171.0	65.0	CASO 69 HBC	0 11. PI- P, N2PI	8/69
W	180.0	30.0	BARTSCH 70 HBC	+ 8 PI+ P, 2 PI	5/70
W	40.0	32.0	KRAMER 70 HBC	+ 13.1 PI+ P, 2PI	11/70*
W	20.0	(8.0)	STUNTEBEC 70 HBC	0 8. PI-P, PI+PI-	1/71*
W	AVG 136.0 ± 23.9 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)				

15 RHO (1660) PARTIAL DECAY MODES

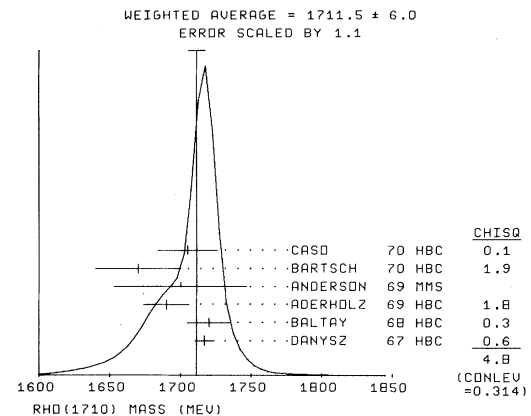
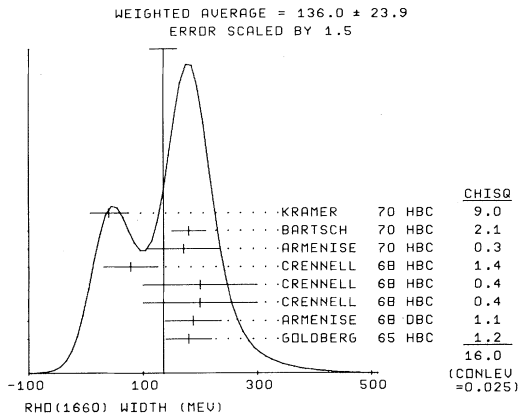
P1	RHO (1660) INTO PI PI	139+ 139	DECAY MASSES
P2	RHO (1660) INTO 4PI	139+ 139+ 139+ 139	
P5	RHO (1660) INTO K KBAR	497+ 497	

15 RHO (1660) BRANCHING RATIOS (FOR OTHER POSSIBLE MODES SEE RHO(1715) BELOW)

R1	R1 MESON FRACTION INTO ONE / THREE / FIVE OR MORE CHARGED TRACKS			
R1	(0.37) / 0.59 / 0.04	FOCACCI 66 MMS	10/66	
R3	RHO(1660) INTO (K KBAR) / (2 PI)			
R3	INDICATION SEEN	EHLICH 66 HBC	+ 0 7.9 PI- P	
R3	PROBABLY SEEN	ABRAMS 67 HBC	0 4.25 K- P	
R3	0.08	0.08	CRENNELL 68 HBC	6.0 PI- P
R3	0.08	0.03	BARTSCH 70 HBC	+ 8. PI+ P
R3	AVG 0.080 ± 0.026 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			

Mesons

For notation, see illustrated key at beginning of data card listings.



REFERENCES FOR RHO(1660)

BELLINI, DI CORATO, DU'FINO, FORINI (MILANO)	
FORINI, DI GIACCI, SARDI (BOLGNA+ORSAY+SACLAY)	
GOLDBERG (CERN+PARIS+ORSAY+MILANO+CEA+SACL)	
R. ENRICH, W. SELOVE, H. YUTA (PENNSYLVANIA)	
CERN MISSING MASS SPECTROMETER GROUP (CERN)	
CERN MISSING MASS SPECTROMETER GROUP (CERN)	
K. HEHN, G. LASPER, SECHI-ZORN, WOLSKY (MARYLAND)	
CERN MISSING MASS SPECTROMETER GROUP (CERN)	
L. DUBAL (GENEVE)	
FORINI, GARTACCI, BARI, BOLDI, FIRENZE, ORSAY	
ROSEBECK, DEUTSCHMANN, AACHEN, BERLIN, CERN	
KARSHON, LAI, SCARR, SKILLICORN (BNL)	
PRENTICE, STEENBERG, YODN (TORONTO+WISC)	
BARTSCH, AACHEN, BERL, CERN, KRAK+WARS	
SELOVE, RISHWAS, CASON, (PENN+NDAM+ROCH)	
CONTE, BENZI, (GENO+DESY+HAMB+MIL+SACL)	
GHIIDINI, FORING, CARTACCI, (PARI+BGNA+FRIZ)	
KRAUS, TSANDS, GROTE, KOTZAN, (AACH+BERL+CERN)	
BARTON, GUTAY, LICHTMAN, MILLER, (PURDUE)	
STUNTEBECK, KFNNEY, DEERY, RISHWAS, CASON, (NDAM)	

$\rho(1710)$
→ 4π

3R RHO(1710, JPC = +1 I = 1 OR 2

THIS ENTRY CONTAINS 4PI, RHO 2PI, 2RHO, OMEGA PI, AND K*BAR ENHANCEMENTS, AND THE R1 AND THE R2. NOTE THAT THE (OMEGA PI) PEAKS HAVE DIFFERENT MASS. DECAE RHO(1710) INTO 4PI MAY BE INEL. DECAYS OF G-MESON ABOVE. SEE BARTSCH 70.

3R MASS (MEV)

M	R	Ref	Mass (MeV)	Notes
M	R	RO 1717.	7.	DANYSZ 67 HBC
M	R	SEEN IN 2.5-3 PRAR P.	2PI+2PI-	-WITH 0.1, 2 PI+PI- PAIRS IN RHO BAND
M	K	(1700.)	FRNCH 67 HBC	0.3, 3.6 PBAR P
M	K	OBSERVED IN NEUTRAL (K* KBAR) MODE (G-PARITY UNKNOWN)		
M	J	(1729.)	15.	BALTAY 68 HBC + 7, 8.5 PI+ P
M	J	(1675.0)	(10.0)	JOHNSTON 68 HBC - 7.0 PI- P
M	J	NOT SEPARATED FROM 2 PI DECAY		
M	M	(1590.0)	16.0	ADERHDLZ 69 HBC + 8 PI+ P, K*BAR PI
M	M	(1700.0)	17.0	ANDERSON 69 HBC - 16 PI- P, BACKW
M	M	(1627.)	(12.)	(17.) BARNHAM 70 HBC + 10 K+ P, RHO PI PI
M	M	(1670.0)	30.0	BARTSCH 70 HBC + 8 PI+ P, 4 PI
M	M	(1705.0)	21.0	CASO 70 HBC - 11, 2PI- P, RHO 2PI
M	M	(1700.0)		MAURER 70 HBC 05.7 PRAR P, 7 PI
M	M	SEEN IN 2 RHO, NOT IN 4 PI OUTSIDE RHO BANDS.		
M	AVG		1711.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1) (SEE IDEOGRAM BELOW)

3R WIDTH (MEV)

M	R	Ref	Width (MeV)
M	R	(30.)	OR LESS
M	R	R2 PEAK FROM CERN MMS EXPT. DECAY MODES AND G-PARITY UNKNOWN.	
M	R	(40.)	(12.1)
M	R	SEEN IN 2.5-3 PRAR P. 2PI+2PI- -WITH 0.1, 2 PI+PI- PAIRS IN RHO BAND	
M	J	(100.)	35.
M	J	(40.0)	(20.0)
M	J	NOT SEPARATED FROM 2 PI DECAY	
M	M	(112.0)	60.0
M	M	(195.0)	60.0
M	M	(172.)	(25.)
M	M	(160.0)	40.0
M	M	(160.0)	
M	AVG		123.8

W OMEGA PI DECAYS AND THE R1

W	M	Ref	Decay Mode
W	M	(30.)	OR LESS
W	M	R1 PEAK FROM CERN MMS EXPT. DECAY MODES AND G-PARITY UNKNOWN.	
W	M	(130.)	73. 43.
W	M	(60.0)	

3R RHO(1710) PARTIAL DECAY MODES

PI	Decay Mode	CHISO
P1	RHO(1710) INTO 4 PI	139+ 139+ 139+ 139
P2	RHO(1710) INTO A2 PI	139+ 1300
P3	RHO(1710) INTO OMEGA PI	139+ 783
P4	RHO(1710) INTO PHI PI	1018+ 139
P5	RHO(1710) INTO 2 RHO	765+ 765
P6	RHO(1710) INTO PI PI RHO	765+ 139+ 139

3R RHO(1710) BRANCHING RATIOS

R1	R2	R3	R4
R1	R2	R3	R4
R1	R2	R3	R4
R1	R2	R3	R4

3R MASS (MEV)

M	R	Ref	Mass (MeV)
M	R	RO 1717.	7.
M	R	SEEN IN 2.5-3 PRAR P. 2PI+2PI- -WITH 0.1, 2 PI+PI- PAIRS IN RHO BAND	
M	K	(1700.)	FRNCH 67 HBC
M	K	OBSERVED IN NEUTRAL (K* KBAR) MODE (G-PARITY UNKNOWN)	
M	J	(1729.)	15.
M	J	(1675.0)	(10.0)
M	J	NOT SEPARATED FROM 2 PI DECAY	
M	M	(1590.0)	16.0
M	M	(1700.0)	17.0
M	M	(1627.)	(12.)
M	M	(1670.0)	30.0
M	M	(1705.0)	21.0
M	M	(1700.0)	
M	M	SEEN IN 2 RHO, NOT IN 4 PI OUTSIDE RHO BANDS.	
M	AVG		1711.5

REFERENCES FOR RHO(1710)

DEUTSCHMANN ET AL (AACHEN+BERLIN+CERN)	
CERN MISSING MASS SPECTROMETER GROUP (CERN)	
CERN MISSING MASS SPECTROMETER GROUP (CERN)	
FRENCH+KINSON+SIMAK (CERN+LIVERPOOL)	
FOCACCI+KIENZLE+LECHMANN+LEVRAT (GENE)	
L. DUBAL (GENEVE)	
FRENCH (GENEVE)	
KINSON+MC DONALD+IDRIFORD (CERN+BRIM)	
KUNG+VEH+FERREL (COLMBA+ROCH+RUT+YALE) I=1	
CASO (GENO)	
CONTE+CORDS+DIAZ (GENOVA+HAMB+MIL+SACL)	
PRENTICE, STEENBERG, YODN (TORONTO+WISC) I=J	
ANDERSON 69 HBC + 7, 8.5 PI+ P	
JOHNSTON 68 HBC - 7.0 PI- P	
ADERHDLZ 69 NP R 11 259	
ANDERSON 69 PRL 22 1390	
VETLITSKY 69 SUNP 9 461	
BARTSCH, (AACH+BERL+CERN+KRAK+WARS)	
GOLLINS, BLIEDEN (BNL+CERN)	
GUTHAVIN, KLIGER, KOLGANOV, LEBEDEV (ITP)	
BARNHAM 70 PRL 24 1083	
GOLDFY, JONES, KEWYON, PATHAK, PIDDIFORD (BRIM)	
BARTSCH 70 CERN/D./PHYS70-4	
CASO 70 LNC 3 707	
MAURER 70 THESIS ND, 598	
KRAUS, TSANDS, GROTE, KOTZAN (AACH+BERL+CERN)	
CONTE, TOMASINI, GORDS (GENO+HAMB+MIL+SACL)	
G. MAURER (STRASBOURG)	

Mesons

For notation, see illustrated key at beginning of data card listings.

R(1750) 39 R(1750) I=1 THIS ENTRY CONTAINS I=1 PEAKS AND THE R3 PEAK NOT A FIRMLY ESTABLISHED RESONANCE - OMITTED FROM TABLE

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for DURAL 67 MMS and STUNTEBEC 70 HBC.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for LEVRAT 66 MMS and STUNTEBEC 70 HBC.

39 R(1750) BRANCHING RATIOS R3 R3 MESON FRACTION INTO ONE / THREE / FIVE OR MORE CHARGED TRACKS

REFERENCES FOR R (1750) CERN MISSING MASS SPECTROMETER GROUP (CERN) BOESEBECK 68 HBC

ηA, I=0 (1830) → 4π, K* K

42 ETA OR RHO (1830) G=1 (JPG=+) I GTE 0 THIS ENTRY CONTAINS 4 PI AND K PI KBAR AND THE R4 MMS PEAK. R4 IS ONLY A 3 STANDARD DEVIATION EFFECT. OMITTED FROM TABLE.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for DANYSZ 67 HBC and STUNTEBEC 70 HBC.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for DANYSZ 67 HBC and STUNTEBEC 70 HBC.

Table with columns for partial decay modes and decay masses. Includes data for ETA OR RHO (1830) INTO 4 PI and INTO RHO PI PI.

REFERENCES DANYSZ 67 PL 248 309 BOESEBECK 68 HBC + 8 PI+ P, PI+ P10

φA, I=0 (1830) → 5π, K* K

43 PHI OR PI (1830) G=-1 (JPG=-) I GTE 0 THIS ENTRY CONTAINS OMEGA PI PI AND K PI KBAR AND THE R4 MMS PEAK. R4 IS ONLY A 3 STANDARD DEVIATION DEVIATION EFFECT. I=1 IF (OMEGA RHO) MODE EXISTS. OMITTED FROM TABLE.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for DANYSZ 67 HBC and STUNTEBEC 70 HBC.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for DANYSZ 67 HBC and STUNTEBEC 70 HBC.

Table with columns for partial decay modes and decay masses. Includes data for PHI (1830) INTO 5 PI and INTO OMEGA PI PI.

REFERENCES DANYSZ 67 NC 51A 801 BOESEBECK 68 HBC + 8 PI+ P, PI+ P10

S(1930) REGION 31 S(1930, JPC=) I=1 OR 2 THIS ENTRY CONTAINS, BESIDES THE S(1930) SEEN BY CHIKOVANI 66 WITH A MMS, AND CONFIRMED BY CLINE 70 IN PBAR P BACKWARD ELASTIC SCATTERING, VARIOUS OTHER PEAKS NEARBY. FOR REVIEWS, SEE ASTIER 70, MONTANET 69.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for CHIKOVANI 66 MMSP and BOESEBECK 68 HBC.

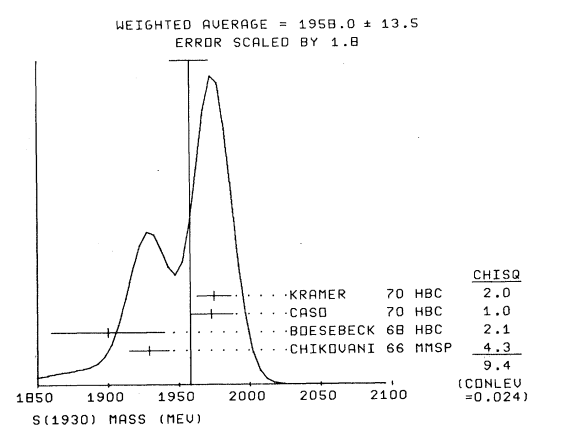


Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes data for CHIKOVANI 66 MMSP and BOESEBECK 68 HBC.

Table with columns for partial decay modes and decay masses. Includes data for S INTO PI+ PI- and INTO PBAR P.

Mesons

For notation, see illustrated key at beginning of data card listings.

31 D(SIGMA)/D(T) (MICROBARN/ (GEV/C)**2)

CS	35.0	12.0	FOCACCI	66 MMS	.22 LTE T LTE	.36	9/66
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REFERENCES FOR S(1930)

CHIKOVANI	66 PL 22 233	CFRN MISSING MASS SPECTROMETER GROUP (CERN)
FOCACCI	66 PRL 17 890	CFRN MISSING MASS SPECTROMETER GROUP (CERN)
BOESEBEC	68 NP B 4 501	BOESEBEC, DEUTSCHMANN, + (AACHEN+BERLIN+CERN)
CLINE	68 PRL 21 1268	+ENGLISH, REEDER, TERRELL, THWITY (WISCONSIN)
MONTANET	69 LUND CONF. P. 189	L. MONTANET, RAPPORTEUR (CERN)
ASTIER	70 KIEV CONF.	RAPP. TALK ON BOSON RESONANCES (CDF)
CASO	70 INC 3 707	+CORDS, COSTA, DARD, + (GENO, DESY, HAM, MILA, SACL)
CLINE	70 PREPRINT, DEC 70	D. CLINE, J. ENGLISH, D. D. REEDER (WTSC) J
	PRIV. COMM., FEB 71	SUPERSEDES DEC 70 PREPRINT
KRAMER	70 PRL 25 396	+BARTON, GUTAY, LICHTMAN, MILLER, + (PURDUE)

ρ (~2100) REGION

51 RHO (2100, JPG= +) I=1

NICHOLSON 69 SUGGEST IG=1+, JP=3- FROM ANALYSIS OF DIFFERENTIAL CROSS-SECTIONS FOR PBAR PI -- ZPT. OMITTED FROM TABLE.

51 RHO (2100) MASS (MEV)

M	2096.0	38.0	ANDERSON	69 MMS	- 16 PI- P, BACKW	8/69
M	(2120.1)		NICHOLSON	69 CNTR	0 .7-2.4 PB P, 2PI	9/69

51 RHO (2100) WIDTH (MEV)

W	(150.0)		ANDERSON	69 MMS	- 16 PI- P, BACKW	8/69
W	N	(249.)	NICHOLSON	69 CNTR	0 .7-2.4 PB P, 2PI	9/69

THE WIDTH INCLUDES RESOLUTION.

REFERENCES FOR RHO(2100)

ANDERSON	69 PRL 22 1390	+COLLINS, ALIEDEN+ (BNL+CERN)
NICHOLSON	69 PRL 23 603	NICHOLSON, RARISH, DELORME, + (CALT+ROCH+BNL)

T(2200) REGION

32 T(2200, JPG=) I=1 OR 2

THIS ENTRY CONTAINS, BESIDES THE T(2200) SEEN BY CHIKOVANI 66 WITH A MMS, VARIOUS OTHER PEAKS NEARBY. FOR REVIEWS, SEE MONTANET 69, KALBFLEISCH 70 OMITTED FROM TABLE.

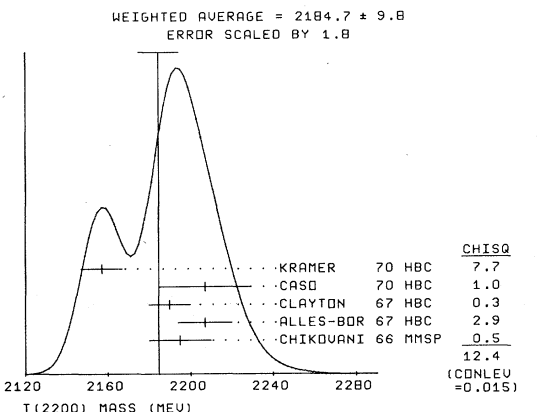
32 T(2200) MASS (MEV)

M	2195.0	15.0	CHIKOVANI	66 MMS P	- 12.0 PI-P	8/66
M	R	(2190.)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71*

32 T(2200) WIDTH (MEV)

W	(130.0)		CASO	70 HBC	- 11.2PI- P, NOTE C.	5/70
W	C	(2207.0)	SEEN IN RHO- PI+ PI-	(OMEGA AND ETA ANTISELECTED IN 4 PI SYSTEM)		5/70
W	D	2157.0	10.0	KRAMER	70 HBC + 13.1 PI+ P, 2PI	11/70*
W	D			HAS IG=1+ FROM ABSENCE OF PI+PI+ PEAK. THUS JP=(00D)-.		11/70*

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8) (SEE IDEOGRAM BELOW)



32 T(2200) WIDTH (MEV)

W	(113.0)	OR LESS	CHIKOVANI	66 MMS P	- 12.0 PI-P	8/66
W	B	(85.)	ABRAMS	67 CNTR	S CHANNEL NBAR N	7/67
W	B				SEE NOTE B UNDER T(2200) MASS ABOVE.	
W	K	BETWEEN 20 AND 80 MEV	KALBFLEISCH	69 HBC	0 5.7 PBAR P	12/66
W	K	SEEN IN PRAR P TO RHO RHO P10. IG=1-			0 5-CHANNEL PBAR P	7/69
W	(130.0)		CASO	70 HBC	- 11.2PI- P, NOTE C.	5/70
W	C		SEEN IN RHO- PI+ PI-	(OMEGA AND ETA ANTISELECTED IN 4 PI SYSTEM)		5/70
W	D	68.0 22.0	KRAMER	70 HBC + 13.1 PI+ P, 2PI		11/70*
W	D		HAS IG=1+ FROM ABSENCE OF PI+PI+ PEAK. THUS JP=(00D)-.			11/70*

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

32 D(SIGMA)/D(T) (MICROBARN/ (GEV/C)**2)

CS	29.0	10.0	FOCACCI	66 MMS	.22 LTE T LTE	.36	9/66
----	------	------	---------	--------	---------------	-----	------

32 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON

CS	(5.5)		ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71*
CS	K	(0.5)	KALBFLEISCH	69 HBC	OS CHANNEL NBAR N	7/69
CS					SEEN IN PBAR P TO RHO RHO P10. IG=1-	

REFERENCES FOR T(2200)

CHIKOVANI	66 PL 22 233	CERN MISSING MASS SPECTROMETER GROUP (CERN)
FOCACCI	66 PRL 17 890	CERN MISSING MASS SPECTROMETER GROUP (CERN)
ABRAMS	67 PRL 18 1209	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
		ALSO ABRAMS 70
ALLES-BDR	67 NC 50 A 776	ALLES-BORELLI, FRENCH, FRISK, + (CERN+BNL)G=
CLAYTON	67 HEIDBG. CONF. P. 57	+MASON, MUIRHEAD, FILIPPAS* (LIVPOLN+ATHENS)
COOPER	69 PRL 20 1059	+HYMAN, MANNFR, MUSGRAVE, VOJVODIC (ANL)
BRICMAN	69 PL 29 B 451	+FERRI-LUZZI, RIZARD, + (CERN+CAEN+SACL)
CASO	69 NC 62 A 755	+CONTE, BENZ, + (GENO+DESY+HAMB+MILA+SACL)
KALBFLEISCH	69 PL 29 B 259	G. KALBFLEISCH, R. STRAND, V. VANDERBURG (BNL)
MONTANET	69 LUND CONF. P. 189	L. MONTANET, RAPPORTEUR (CERN)
ABRAMS	70 PR D 1 1917	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
CASO	70 LNC 3 707	+CONTE, TOMASINI, CORDS+ (GENO+HAMB+MILA+SACL)
		ALSO CASO 69
KALBFLEISCH	70 PHILAD. CONF. P. 609	G. KALBFLEISCH AND D. MILLER REVUES (BNL)
KRAMER	70 PRL 25 396	+BARTON, GUTAY, LICHTMAN, MILLER, + (PURDUE)

ρ (~2275) REGION

52 RHO (2275, JPG= +) I=1

NICHOLSON 69 SUGGEST IG=1+, JP=5- FROM ANALYSIS OF DIFFERENTIAL CROSS-SECTIONS FOR PBAR PI -- ZPT. OMITTED FROM TABLE.

52 RHO (2275) MASS (MEV)

M	2260.0	18.0	ANDERSON	69 MMS	- 16 PI- P, BACKW	8/69
M	(2290.)		NICHOLSON	69 CNTR	0 .7-2.4 PB P, 2PI	9/69

52 RHO (2275) WIDTH (MEV)

W	(25.0)	OR LESS	ANDERSON	69 MMS	- 16 PI- P, BACKW	8/69
W	N	(165.)	NICHOLSON	69 CNTR	0 .7-2.4 PB P, 2PI	9/69

THE WIDTH INCLUDES RESOLUTION.

REFERENCES FOR RHO(2275)

ANDERSON	69 PRL 22 1390	+COLLINS, ALIEDEN+ (BNL+CERN)
NICHOLSON	69 PRL 23 603	NICHOLSON, RARISH, DELORME, + (CALT+ROCH+BNL)

N̄N̄_{I=1} (2350)

58 N NBAR (2350) I=1

OMITTED FROM TABLE

58 MASS

M	B	(2350.)	(10.)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71*
M	B					SEEN AS BUMP IN I=1 STATE. WIDTH MUCH LARGER THAN IN MMS EXPT.	

58 WIDTH

W	B	(140.)		ABRAMS	67 CNTR	S CHANNEL PBAR N	7/67
W	B					SEEN AS BUMP IN I=1 STATE. WIDTH MUCH LARGER THAN IN MMS EXPT.	

58 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON

CS	(3.2)		ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71*
----	-------	--	--------	---------	------------------	-------

REFERENCES FOR N NBAR (2350)

ABRAMS	67 PRL 18 1209	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
ABRAMS	70 PR D 1 1917	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)

Mesons

For notation, see illustrated key at beginning of data card listings.

U(2375) REGION

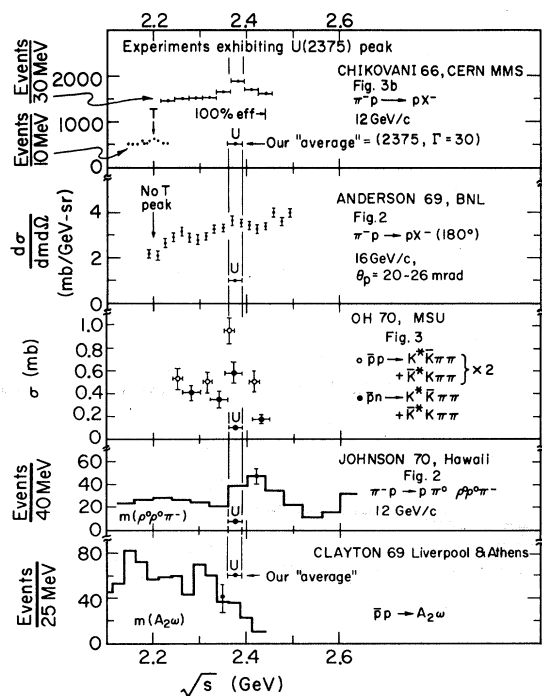
33 U(2375) I=1

Evidence for the existence of a non-strange, I=1 meson in the 2375-MeV region, having $\Gamma \lesssim 30$ MeV, originally came from the CERN Missing-Mass Spectrometer group (CHIKOVANI 66). Others have since found indications of a peak in the neighborhood; the figure shows the current situation.

Removed from our figure and data cards since the last edition is some previously reported evidence for the U in the reactions $\bar{p}p \rightarrow K\bar{K}\omega$ and $\bar{p}p \rightarrow K_S K_L +$ neutrals (RING-1 69 and RING-2 69). The groups involved now have added more data and do not see any significant peak in the U region (OH 70 and private communication).

We show the CLAYTON 69 and JOHNSON 70 mass distributions, although the peaks in these experiments do not appear to line up with those of the other experiments. There is also a peak in the I=1 $\bar{p}p$ total cross section reported by ABRAMS 67, but it is 140 MeV wide and is not shown.

Although the net evidence for the U is not very strong, we continue to include it in our Meson Table because of the concurrence of several experiments on its mass and width.



33 U(2375) MASS (MEV)										
M	C	2382.0	24.0	CHIKOVANI 66 MMS	-	12.0	P1-P		8/66	
M	C	(2324.0)	(20.0)	CLAYTON 67 HRC	+	2.5	PBAR, A2+OMEGA		11/69	
M	C	2380	10	IS MISTAKE...			PRIV. COMM. FROM MUIRHEAD.			
M	C	2370.	17.	ANDERSON 69 ASPK	-	16	P1- BKSCAT		11/69	
M	C	(2420.0)	(25.0)	JOHNSON 70 HBC	-	12.0	P1- P		1/71*	
M	C	2360.0	75.0	OH	70	HBC	-OPBAR(P,N),K*K2P1		5/70	
M	AVG	2370.7	12.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						

33 U(2375) WIDTH (MEV)									
W		(30.0)	OR LESS	CHIKOVANI 66 MMS	-	12.0	P1-P		8/66
W		(57.)		ANDERSON 69 ASPK	-	16	P1- BKSCAT		11/69
W		(80.0)	OR LESS	JOHNSON 70 HBC	-	12.0	P1- P		1/71*
W		(60.0) OR LESS		OH	70	HBC	-OPBAR(P,N),K*K2P1		5/70

33 U(2375) D(SIGMA)/D(T) (MICROBARN/(GEV/C)**2)									
CS		42.0	14.0	FOCACCI 66 MMS	.2R	LTE	LTE	.36	9/66

33 U MESON BRANCHING RATIOS									
RI		U- MESON FRACTION INTO ONE / THREE / FIVE OR MORE CHARGED TRACKS							
RI		(0.30) / 0.45 / 0.25	FOCACCI 66 MMS	-					10/66

41 N NRAR (2375) I=0									
EVIDENCE FOR RESONANCE PRELIMINARY. OMITTED FROM TABLE.									

41 MASS									
M		2375.	10.	ABRAMS 70 CNTR	S	CHANNEL	NBAR	N	1/71*

41 WIDTH									
W		(190.)		ABRAMS 70 CNTR	S	CHANNEL	NBAR	N	1/71*

41 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON									
CS		(2.5)		ABRAMS 70 CNTR					1/71*

X-(2500)									
OMITTED FROM TABLE									

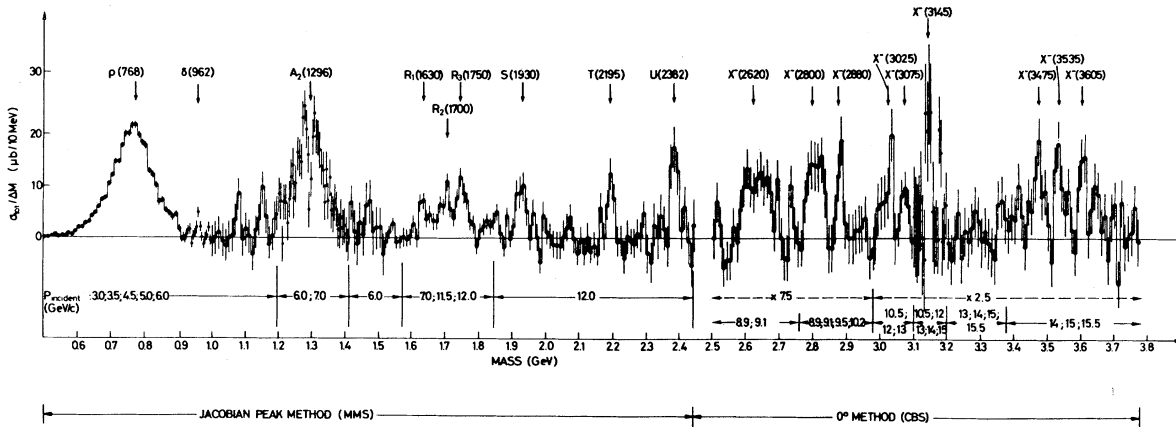
46 X-(2500) MASS (MEV)									
M		2500.0	32.0	ANDERSON 69 MMS	-	16	P1- P, BACKW9		8/69

46 X-(2500) WIDTH (MEV)									
W		(87.0)		ANDERSON 69 MMS	-	16	P1- P, BACKW9		8/69

REFERENCES FOR X-(2500)									
ANDERSON 69 PRL 22 1390				+COLLINS,+					(BNL+CARN)

For notation, see illustrated key at beginning of data card listings.

BOSON MASS-SPECTRUM IN $\pi^-p \rightarrow p(\text{BOSON})^-$, CERN 1965-70



From Baud et al. paper, presented by G. Damgaard, Proceedings 1970 Philadelphia Conference.

X⁻(2620) through X⁻(3535)

The figure on this page shows the X⁻ spectrum, as studied with the CERN Boson Spectrometer. The background is not shown but is about ten times as large as the signal. Uncertainties in where to draw this background make it extremely difficult to measure the width of these peaks, or to judge their statistical significance.

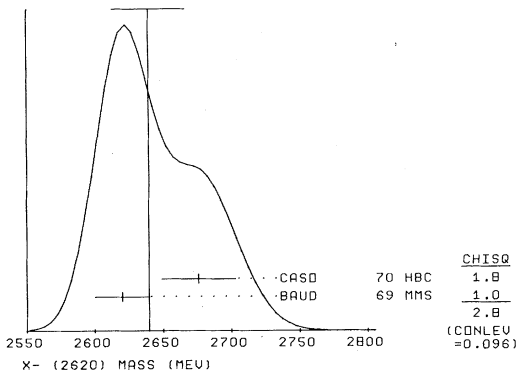
X⁻(2620)

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48 X- (2620, JPG= ) I=1 OR 2
   OMITTED FROM TABLE

-----
48 X- (2620) MASS (MEV)
W 550 2620. 20. BAUD 69 MMS - 8.-10. PI- P 9/69
W 2678.0 27.0 CASO 70 HRC - 11.2PI- P-NOTE C 5/70
W C SEEN IN RHO- PI+ PI- (OMEGA AND ETA ANTISELECTED IN 4 PI SYSTEM) 5/70
M
M AVG 2639.8 26.8 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)
      (SEE IDEOGRAM BELOW)
    
```

WEIGHTED AVERAGE = 2639.8 ± 26.8
ERRDR SCALED BY 1.7



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48 X- (2620) WIDTH (MEV)
W 550 85. 30. BAUD 69 MMS - 8.-10. PI- P 9/69
W (150.0) CASO 70 HRC - 11.2PI- P-NOTE C 5/70
W C SEEN IN RHO- PI+ PI- (OMEGA AND ETA ANTISELECTED IN 4 PI SYSTEM) 5/70
*****
REFERENCES FOR X-(2620)
BAUD 69 PL 308 129 CERN BOSON SPECTROMETER GROUP (CERN)
CASO 70 LNC 3 707 +CONTE,TOMASINI,CORDS+(GENO+HAMB+MLA+SACL)
*****

X-(2800)
49 X- (2800, JPG= ) I=1 OR 2
   OMITTED FROM TABLE

-----
49 X- (2800) MASS (MEV)
M 640 2800. 20. BAUD 69 MMS - 8.-10. PI- P 9/69

-----
49 X- (2800) WIDTH (MEV)
W 640 46. 10. BAUD 69 MMS - 8.-10. PI- P 9/69
*****
REFERENCES FOR X-(2800)
BAUD 69 PL 308 129 CERN BOSON SPECTROMETER GROUP (CERN)
*****

X-(2880)
50 X- (2880, JPG= ) I=1 OR 2
   OMITTED FROM TABLE

-----
50 X- (2880) MASS (MEV)
M 230 2880. 20. BAUD 69 MMS - 8.-10. PI- P 9/69

-----
50 X- (2880) WIDTH (MEV)
W 230 15. OR LESS BAUD 69 MMS - 8.-10. PI- P 9/69
*****
REFERENCES FOR X-(2880)
BAUD 69 PL 308 129 CERN BOSON SPECTROMETER GROUP (CERN)
*****
    
```


Mesons

For notation, see illustrated key at beginning of data card listings.

X⁻ (3030) 53 X⁻ (3030, JPG=) I=1 OR 2
OMITTED FROM TABLE

53 X⁻ (3030) MASS (MEV)

M	3035.0	25.0	ALEXANDER 70 HBC	7. PBAR P+7 PI	5/70
M	3025.0	20.0	BAUD 70 HMS	- 10.5-13 PI- P	5/70
M					
M	AVG	3028.0	15.6	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

53 X⁻ (3030) WIDTH (MEV)

W	200.0	60.0	ALEXANDER 70 HBC	7. PBAR P+7 PI	5/70
W	(25.0)	APPROX.	BAUD 70 HMS	- 10.5-13 PI- P	5/70

REFERENCES FOR X⁻(3030)

ALEXANDER 70 PRL 25 63 +BAR-NIR,DAGAN,GIDAL,GRUNHAUS+ (TEL-AVIV)
BAUD 70 PL 31 B 549 CERN BOSON SPECTROMETER GROUP (CERN)

X⁻ (3075) 54 X⁻ (3075, JPG=) I=1 OR 2
OMITTED FROM TABLE

54 X⁻ (3075) MASS (MEV)

M	3075.0	20.0	BAUD 70 HMS	- 10.5-13 PI- P	5/70
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54 X⁻ (3075) WIDTH (MEV)

W	(25.0)	APPROX.	BAUD 70 HMS	- 10.5-13 PI- P	5/70
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REFERENCES FOR X⁻(3075)

BAUD 70 PL 31 B 549 CERN BOSON SPECTROMETER GROUP (CERN)

X⁻ (3145) 55 X⁻ (3145, JPG=) I=1 OR 2
OMITTED FROM TABLE

55 X⁻ (3145) MASS (MEV)

M	3145.0	20.0	BAUD 70 HMS	- 10.5-15 PI- P	5/70
---	--------	------	-------------	-----------------	------

55 X⁻ (3145) WIDTH (MEV)

W	(10.0)	OR LESS	BAUD 70 HMS	- 10.5-15 PI- P	5/70
---	--------	---------	-------------	-----------------	------

REFERENCES FOR X⁻(3145)

BAUD 70 PL 31 B 549 CERN BOSON SPECTROMETER GROUP (CERN)

X⁻ (3475) 56 X⁻ (3475, JPG=) I=1 OR 2
OMITTED FROM TABLE

56 X⁻ (3475) MASS (MEV)

M	3475.0	20.0	BAUD 70 HMS	- 14-15.5 PI- P	5/70
---	--------	------	-------------	-----------------	------

56 X⁻ (3475) WIDTH (MEV)

W	(30.0)	APPROX.	BAUD 70 HMS	- 14-15.5 PI- P	5/70
---	--------	---------	-------------	-----------------	------

REFERENCES FOR X⁻(3475)

BAUD 70 PL 31 B 549 CERN BOSON SPECTROMETER GROUP (CERN)

X⁻ (3535) 57 X⁻ (3535, JPG=) I=1 OR 2
OMITTED FROM TABLE

57 X⁻ (3535) MASS (MEV)

M	3535.0	20.0	BAUD 70 HMS	- 14-15.5 PI- P	5/70
---	--------	------	-------------	-----------------	------

57 X⁻ (3535) WIDTH (MEV)

W	(30.0)	APPROX.	BAUD 70 HMS	- 14-15.5 PI- P	5/70
---	--------	---------	-------------	-----------------	------

REFERENCES FOR X⁻(3535)

BAUD 70 PL 31 B 549 CERN BOSON SPECTROMETER GROUP (CERN)

K K MESON (JP=0-) I=1/2
SEE LISTINGS OF STABLE PARTICLES

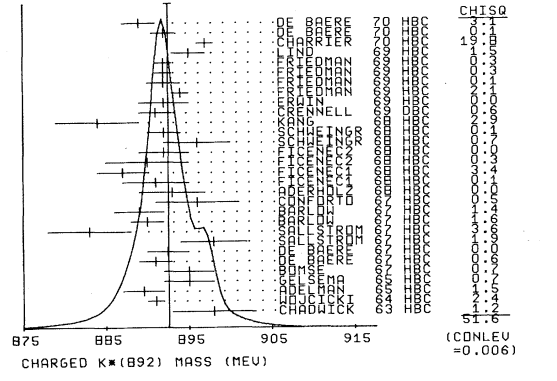
K⁻ (725) 17 KAPPA (725, JP=) I=1/2
EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.
FOR A COMPILATION, SEE APPENDIX A OF JAN 67 EDITION
(RMP 39, 1) OF THIS DATA SUMMARY.
SEE ALSO ROSENFELD, PROC. 1968 UNIV. OF PENN. CONF. ON MESON SPECTROSCOPY

K* (892) 18 K* (892, JP =1-) I=1/2

18 K* (892) MASS (MEV)

M	898.0	5.0	CHARGED ONLY. THIS IS WHAT APPEARS ON MESON TABLE			
M	3870	891.0	1.0	ONDELICK 68 HBC	+ 1.5 K+P	
M	891.5	2.5	WJACIICKI 68 HBC	- 1.7 K+P		
M	(891.0)	(3.0)	ADELMAN 68 HBC	+ 1.5 K+P	6/66	
M			FERRI-LUZ 68 HBC	+ 3.0 K+P		
M	D VALUE ABOVE SUPERSEDED BY DE BAERE TO BELOW					
M	895.0	3.0	DELSERA 65 HBC	- 1.5 K+P	1/71*	
M	895.	3.	BOMSE 67 HBC	+ 2.3 K+P	7/67	
M	891.	2.	DE BAERE 67 HBC	+ 3.5 K+P (KO PI+)	7/67	
M	892.5	2.5	DE BAERE 67 HBC	+ 3.5 K+P (K+ PI0)	7/67	
M	898.	4.	SALLSTROM 67 HBC	+ 3. K+P (KO PI+)	7/67	
M	883.	5.	SALLSTROM 67 HBC	+ 3. K+P (K+ PI0)	7/67	
M	890.	2.	BARLOW 67 HBC	+ 1.2 PBAR P	11/66	
M	889.	3.	BARLOW 67 HBC	+ 1.2 PBAR P	11/66	
M	896.0	5.0	CONFORTO 67 HBC	+ 6.6 K+P	9/67	
M	893.	4.	ADERHOLZ 68 HBC	- 10 K+P	6/68	
M	891.	4.	FICENEC1 68 HBC	- 1.3 K+P (K-PI0)	9/67	
M	887.	3.	FICENEC1 68 HBC	- 1.3 K+P (KOP1-)	9/67	
M	896.0	5.0	FICENEC2 68 HBC	- 2.7 K- (PIK-PI0)	2/69	
M	892.0	3.0	FICENEC2 68 HBC	- 2.7 K- (PIKOP1-)	2/69	
M	896.0	4.0	SCHWEINGR 68 HBC	- 4.1 K+P	9/67	
M	892.0	2.0	SCHWEINGR 68 HBC	- 5.5 K+P	9/67	
M	896.0	5.0	KANG 68 HBC	+ 6.6 K+P	7/69	
M	891.0	2.0	CRENNELL 69 DBC	+ 3.9 K+N (KOP1-)	7/69	
M	892.0	3.0	ERWIN 69 HBC	+ 3.5 K+P	9/69	
M	2886	896.	1.	FRIEDMAN 69 HBC	+ 2.1 K+P (3BDY)	9/69
M	728	892.	2.	FRIEDMAN 69 HBC	+ 2.5 K+P (3BDY)	9/69
M	3229	892.	1.	FRIEDMAN 69 HBC	+ 2.6 K+P (3BDY)	9/69
M	1027	892.	1.	FRIEDMAN 69 HBC	+ 2.7 K+P (3BDY)	9/69
M	895.	2.	LIND 69 HBC	+ 9. K+P	9/69	
M	897.0	1.0	CHARLIER 70 HBC	+ 8.25 K+P	1/71*	
M	892.0	1.5	DE BAERE 70 HBC	+ 3.5 K+P	1/71*	
M	889.0	2.0	DE BAERE 70 HBC	+ 5.0 K+P	1/71*	
M	AVG	892.55	0.45	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)		

WEIGHTED AVERAGE = 892.55 ± 0.45
ERROR SCALED BY 1.3



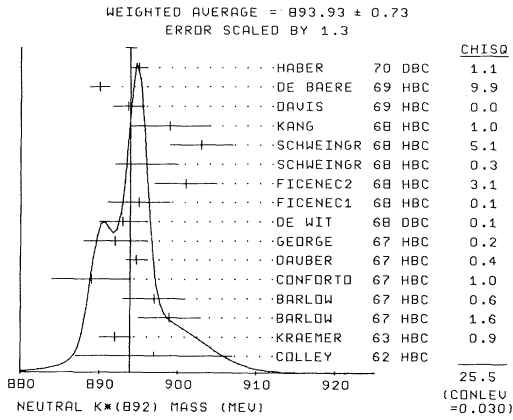
M MIXED-- CHARGED AND NEUTRAL. NOT TABULATED

M	200	(880.0)		ALEXANPER 62 HBC	+ 0 2.2 PI-P	
M	895.0	2.0		FERRI-LUZ 68 HBC	+ 0 3.0 K+P	6/66
M	(895.0)			WANGLER 65 HBC	+ 0 3.0 PI- P	6/66
M	894.	5.		FRENCH 67 HBC	+0 3-4 PBAR P	6/67
M	AVG	894.9	1.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		

M NEUTRAL ONLY. BUT WE DONT USE THIS FOR MASS DIFF. - SEE TYPED NOTE

M	70	897.0	10.0	COLLEY 62 HBC	0 2.0 PI-P		
M	200	892.0	2.0	KRAEMER 63 HBC	0 2.3 K+P		
M	150	(885.0)		SMITH 63 HBC	0 2.3 PI- P		
M	899.	4.		BARLOW 67 HBC	0 1.2 PBAR P	11/66	
M	897.	4.		BARLOW 67 HBC	0 1.2 PBAR P	11/66	
M	889.0	5.0		CONFORTO 67 HBC	0. PBAR P	9/67	
M	896.7	1.3		DAUBER 67 HBC	0 2.0 K+P	12/66	
M	892.0	4.0		GEORGE 67 HBC	0 5.0 K+P	11/67	
M	893.	3.		DE WIT 68 HBC	0 3. K+P	9/69	
M	F	895.	4.	FICENEC1 68 HBC	0 1.3 K+P (K-PI+)	11/69	
M	F	901.	6.	FICENEC2 68 HBC	0 2.7 K- (PIK-PI+)	11/69	
M	F	FICENEC ERROR RAISED		SEE TYPED NOTE			
M	M	896.0	4.0	SCHWEINGR 68 HBC	0 4.1 K+P	9/67	
M	M	903.0	4.0	SCHWEINGR 68 HBC	0 5.5 K+P	9/67	
M	M	899.0	5.0	KANG 68 HBC	0 4.4 K+P	7/69	
M	M	10700	893.7	2.0	DAVIS 69 HBC	0 12. K+P	9/69
M	M	D 2000	890.0	1.25	DE BAERE 69 HBC	0 5.0 K+P	9/69
M	M	D DE BAERE ERRORS ENLARGED BY US TO GAMMA/SUBTINI.		SEE TYPED NOTE.			
M	M	4000	895.0	1.0	HABER 70 DBC	0 3. K+N	5/70
M	AVG	893.93	0.73	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)			

For notation, see illustrated key at beginning of data card listings.



19 K*(0) - K*(±) MASS DIFF. (MEV)					
0	ALL ERRORS ENLARGED BY 1.3. SEE TYPED NOTE				
0	333 6.3 5.0	HABER 67 HBC	0 PRAR P	2/67	
0	1492 6.5 5.0	FICENEC1 68 HBC	1.3 K- P	2/69	
0	1500 9.5 5.0	FICENEC2 68 HBC	2.7 K- P	2/69	
0	AVG 7.5 3.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			

Note on K* Masses and Mass Differences

This note is divided into two discussions:

I. Basic difficulties in determining the mass difference because of interferences and biases.

II. Impossibly small errors reported by some experiments. We have increased some errors that violate the laws of statistics, and scaled up some errors that are inconsistent; but we warn that most of the errors in our data cards are inconsistent. One cannot then obtain a K* mass difference by calculating an average mass for K*⁰ and for K*[±] and just subtracting the two.

I. BASIC DIFFICULTIES

There are two difficulties in measuring a mass difference m(K*⁰) - m(K*[±]) of ~ 7 MeV when the half-width Γ/2 of the K* is 25 MeV:

- 1) Interference between the resonant amplitude and background can in general shift the peak in the mass spectrum by some fraction of Γ/2.
- 2) The two charges of K* have different topologies; this introduces differences in the measuring and fitting of the events, which can also produce mass shifts.

Some reactions (symmetric under reflection of I_z) are immune to the first difficulty. Thus compare the mass of K*⁰ produced in

$$\pi^- p \rightarrow \Lambda \pi^- K^+$$

with the mass of K*[±] in the I_z-reflected reaction

$$\pi^+ n \rightarrow \Lambda \pi^+ K^0.$$

The final-state amplitudes of each will contain not only the |K*⟩ with Ispin 1/2, but also an interfering I = 3/2 P-wave, which we can call |K*_{3/2}⟩. But I_z symmetry

forces ⟨π⁻p|ΔK*⁰⟩ to equal ⟨π⁺n|ΔK*[±]⟩; and similarly for the two K*_{3/2} amplitudes, so that the shifting of the K* peak is the same in both reactions. Nobody has published a mass difference exploiting this fact.

II. IMPOSSIBLY SMALL ERRORS

Consider a sample of N events, with their invariant masses m distributed as an S-wave Breit-Wigner resonance:

$$i.e., P(\epsilon - \epsilon_R) = \frac{1/\pi}{(\epsilon - \epsilon_R)^2 + 1} \quad (1)$$

where $\epsilon = \frac{m}{\Gamma/2}$, $\epsilon_R = \frac{m_R}{\Gamma/2}$. One can then show that the minimum possible error on the determination of the central value ϵ_R is

$$\delta_{min}(\epsilon_R) = \pm \sqrt{\frac{2}{N}}, i.e., \delta_{min}(m_R) = \pm \sqrt{\frac{2}{N}} \frac{\Gamma}{2} \quad (2)$$

This lower limit assumes no background events. In practice, with background, the error will be larger, by another factor $\alpha \approx \sqrt{2}$.

We illustrate errors with small and large backgrounds with a table summarizing the recent experiment ("Unsplit K*'s") by DAVIS 69.

Mass Errors δm of DAVIS 69

- Sample with 5% background/signal at peak
Events: K*(892), 10 700 events in resonance, $\frac{\Gamma}{2} \approx 25$ MeV.
Lower limit from Eq. (2), $\delta_{min}(m) = \sqrt{\frac{2}{N}} \frac{\Gamma}{2} = \pm 0.35$ MeV.
Their likelihood fit yields two sorts of errors:

δ₁(m). Ignore correlations, i.e., keep all parameters (background, width, etc.) fixed, vary m only:
 $\delta_1(m) = \pm 0.41$, $\delta_1(m)/\delta_{min}(m) = 1.16$.

δ₂(m). As m is varied, reoptimize other parameters.
 $\delta_2(m) = \pm 0.53$, $\delta_2(m)/\delta_{min}(m) = 1.5$.

DAVIS 69 mention $\delta_2 = 0.53$, but to hedge against systematic effects, they quote $\delta_3 = 2$ MeV. We punch 2 MeV.

- Sample with 50% background/signal at peak.
Events: K*(1420), 2200 events in resonance, $\frac{\Gamma}{2} = 50$ MeV.

$$\delta_{min}(m) = 1.6 \text{ MeV},$$

$$\delta_1(m) = \pm 2.2 \text{ MeV}, \delta_1(m)/\delta_{min}(m) = 1.4,$$

$$\delta_2(m) = \pm 2.6 \text{ MeV}, \delta_2(m)/\delta_{min}(m) = 1.6.$$

Width Errors δΓ of DAVIS 69

For width, the equivalent of Eq. (2) is $\delta_{min}(\Gamma) = \pm \sqrt{\frac{8/3}{N}} \frac{\Gamma}{2} = 1.15 \delta_{min}(m)$. For convenience we neglect the factor 1.15 and use $\delta_{min}(\Gamma) \approx \delta_{min}(m)$.

- 5% background, K*(892):
 $\delta_2(\Gamma) = \pm 1.6$ MeV, $\delta_2(\Gamma)/\delta_{min}(m) = \frac{1.6}{0.35} = 4.6$.
- 50% background, K*(1420):
 $\delta_2(\Gamma) = \pm 10$ MeV, $\delta_2(\Gamma)/\delta_{min}(m) = \frac{10}{1.6} = 6.25$.

Mesons

For notation, see illustrated key at beginning of data card listings.

We note that $\delta_2(m)/\delta_{\min}(m)$ does not change rapidly with background (1.5 at 5%, 1.6 at 50%) and hence conclude that it is hard to believe an error with $\delta_2/\delta_{\min} < 1.4 = \sqrt{2}$. We chose $\sqrt{2}$ because together with Eq. (2) it leads to the simple "realistic" result

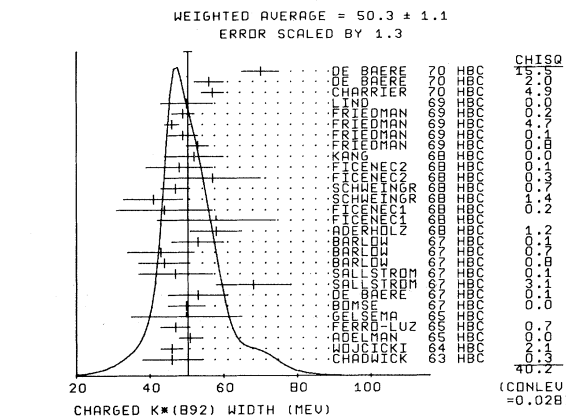
$$\delta(m) > \sqrt{2} \sqrt{\frac{2}{N}} \frac{\Gamma}{2} = \frac{\Gamma}{\sqrt{N}} \quad (3)$$

We conclude that for a sensitive subtraction like $m(K^{*0}) - m(K^{*+})$, the experiments as listed are useless, and we must either re-evaluate them all or concentrate on those two experiments that explicitly quote a mass difference. For a detailed discussion of how we have actually treated those experiments, we refer to the January 1970 edition of this note.

The table above also allows us to concoct a criterion for "realistic" errors in width $\delta(\Gamma)$. We average the 5% and 50% background results [to give $\delta(\Gamma)/\delta_{\min}(m) \approx 5$ to 6] and express the result in terms of Γ , in the style of Eq. (3). We then get the "realistic" test for widths:

$$\delta\Gamma > 4 \frac{\Gamma}{\sqrt{N}} \quad (4)$$

18 K* (892) WIDTH (MEV)	
W CHARGED ONLY. THIS IS WHAT APPEARS ON MESON TABLE	
46.0	8.0
3870 46.0	3.0
51.0	3.0
47.0	4.0
50.0	15.0
50.1	5.0
D (4.0) (4.0)	
D VALUE ABOVE SUPERSEDDED BY DE BAERE 70 BELOW	
53. 8.0	
68. 10.0	
47. 10.0	
44. 7.0	
43. 9.0	
53. 7.0	
(43.1) CONFORTO 67 HRC	+ 0.2 PBAR P 9/67
58. 16.0	ADERHOLZ 68 HRC - 10 K-P 6/68
44. 13.0	FICENEC1 68 HRC - 1.3 K-P (K-PI) 9/67
41.0	SCHWEINGR 68 HRC - 4.1 K-P 9/67
47.0	4.0
57.0	13.0
48.0	9.0
52.0	8.0
(27.0) (8.0) (16.0)	ERWIN 69 HRC + 3.5 K-P 9/69
53. 3.0	FRIEDMAN 69 HRC - 2.1 K-P (3B0Y1) 9/69
49. 4.0	FRIEDMAN 69 HRC - 2.45 K-P (3B0Y) 9/69
46. 2.0	FRIEDMAN 69 HRC - 2.6 K-P (3B0Y) 9/69
49. 3.0	FRIEDMAN 69 HRC - 2.7 K-P (3B0Y) 9/69
50. 7.0	LIND 69 HRC + 9.0 K-P 5/70
57.0	3.0
58.0	4.0
70.0	5.0
AVG	50.3 1.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)
W MIXED-- CHARGED AND NEUTRAL. NOT TABULATED	
200 60.0	5.0
51.8	3.5
(40.0) 60.0	10.0
AVG	54.9 2.8 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W NEUTRAL ONLY.	
70 60.0	10.0
200 50.0	5.0
150 (50.0)	
53. 13.0	
34. 8.0	
(43.1) CONFORTO 67 HRC	+ 0.2 PBAR P 9/67
44. 4.0	DAUBER 67 HRC 0 2.0 K-P 12/66
58. 8.0	DE WIT 68 HRC 0 3.0 K-P 9/69
52. 12.0	FICENEC1 68 HRC 0 1.3 K-P (K-PI+) 9/67
50.0	8.0
49.0	8.0
51.0	11.0
53.0	11.0
10700 53.2	1.6
D 2000 58.0	5.0
D DE BAERE ERRORS ENLARGED BY US TO 4*GAMMA/SORT(N). SEE TYPED NOTE.	
4000 54.0	3.0
AVG	52.2 1.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)



18 K* (892) PARTIAL DECAY MODES

P1	K*(892) INTO K PI	DECAY MASSES
P1	K*(892) INTO K PI	493+ 139
P2	K*(892) INTO (K PI PI)	493+ 139+ 139

18 K* (892) BRANCHING RATIOS

R1	K*(892) INTO (K PI PI)/(K PI)	BRANCHING RATIOS
R1	0 (0.00210R LESS)	WOJCICKI 64 HRC - 1.7 K-P

REFERENCES FOR K*(892)

ALSTON 61 PRL 6 300	ALSTON+ALVAREZ, EBERHARD, GOOD, GRAZIANO+(LRL)
ALEXANDE 62 PRL 8 447	ALEXANDER, KALBFLEISCH, MILLER, G SMITH (LRL)
COLLEY 62 CERN CONF 315	D COLLEY, N GELFAND + (COLUMBIA+PUTGERS)
CHADWICK 63 PL 6 309	CHADWICK, CRENNELL, DAVIES, RETTINI+(OXF+PADU)
GOLDHABE 63 ATHENS CONF 92	SULAMITH GOLDHABER (LRL)
KRAEMER 63 ATHENS CONF 130	R KRAEMER, L MADANSKY + (JOHNS HOPKINS)
SMITH 63 PRL 10 138	SMITH, SCHWARTZ, MILLER, KALBFLEISCH, HUF+(LRL)
WOJCICKI 64 PR 135 B 494	STANLEY G WOJCICKI (LRL)
ADELMAN 65 ATHENS 527	STUART LEE ADELMAN (CAVENDISH)
FERRI-LUZZI 65 NC 36 1101	FERRI-LUZZI, GEORGE, HENRI, JONGEJANS (CERN)
FERRI-LUZZI 65 NC 39 417	FERRI-LUZZI, GEORGE, GOLDSCHMIDT-CLERMONT (CERN)
GELSEMA 65 THESIS	E S GELSEMA (SEE ALSO PL 10 341) (AMSTERD)
WANGLER 65 PR 137 B 414	WANGLER, ERWIN, WALKER (WISCONSIN)
BARASH 67 PR 156 1399	BARASH, KIRSCH, MILLER, TAN (COLUMBIA)
BARLOW 67 NC 50 A 701	*MONTANET, D-ANDLU+ CERN+CDF+(OR-LIVERPOOL)
BOMSE 67 PR 158 1298	*BORENSTEIN+COLE+GILLESPIE+(JOHNS HOPKINS)
CONFORTO 67 NP 83 469	*MARECHAL, MONTANET+ (CERN+CF+IPN+LIVERPOOL)
DAUBER 67 PR 153 1403	*SCHLEIN, SLATER, TICHOU (UL)
DE BAERE 67 NC 51 A 401	*GOLDSCHMIDT-CLERMONT, HENRI+(BRUX+CERN)
FRENCH 67 NC 42A 442	*KINSON+MC DONALD+RIDDFORD+(CERN+BERM)
GEORGE 67 NC 49A 9	*GOLDSCHMIDT-CLERMONT+HENRI+(CERN+BRUX)
SALLSTROM 67 NC 49A 348	SALLSTROM+OTTER+EKSPONG (STOCKHOLM)
ADERHOLZ 68 NP 5 567	*DEUTSCHMANN+ (AACH+BERL+CERN+I.C.+VIENNA)
DE WIT 68 THESIS	S. DE WIT (AMSTERDAM)
FICENEC1 68 PR 169 1034	*HULSIER+SWANSON+TROWER (TORONTO)
FICENEC2 68 PR 175 1725	FICENEC, GORDON, TROWER (ILLINDIS)
KANG 68 PR 176 1587	Y.W.KANG (IDWA)
SCHWEIFING 68 PR 166 1317	SCHWEINGRUBER, DERRICK, FIELDS, AMMAR+(ANL+NN)
CAFENELL 69 PRL 22 487	*KARSHON-LAI, ONE ALL, SCARR (BNI)
DAVIS 69 PRL 23 1071	*DERENZO, FLATTE, ALSTON, LYNCH, SOLMITZ (LRL)
DE BAERE 69 NC 61 A 397	*GOLDSCHMIDT-CLERMONT, HENRI, + (BRUX+CERN)
ERWIN 69 NP 8 9 364	*WALKER, GOSHA, WEINBERG (WISC+PRIN+VAND)
FRIEDMAN 69 UCL-19860	J.FRIEDMAN, PH.D. THESIS (LRL)
LIND 69 NP 8 14 1	*ALEXANDER, FIRESTONE, FU, GOLDHABER (LRL) JP
JUHALA 69 PR 184 1461	*LEACOCK, RHODE, KOPPELMAN, LIBBY, + (AMES+GOLD)
ATHERTON 70 NP 8 16 416	*FRANEK, FRENCH, FRISK, REDNAR+ (CERN+PRAG)
CHARRIER 70 CERN PHYS 70-50	CHARRIERE, DRIJARD, OUNNOUDE, + (LHEB+CERN)
DE BAERE 70 CERN PHYS 70 41	*DEBAISIEUX, DE WOLF, DUFOUR, + (LHEB+CERN)
HABER 70 NP 8 17 289	*SHAPIRA, ALEXANDER+ (REHO+SACL+98NA+EPOLL)

$K_N(1080-1260)$
 $\rightarrow K\pi$

19 KN(1080-1260)
 OMITTED FROM TABLE.

For a recent review of $K\pi$ phase shifts, see ASTIER 70. However, there may be more structure than reported at Kiev, but the final results (super-seeding TRIPPE 68, SCHLEIN 69) are not yet public, so we refrain from further comments here on phase shift analyses.

Mesons

For notation, see illustrated key at beginning of data card listings.

As to the existence of peaks in the $K\pi$ mass distribution in this region, the situation is as follows (for histograms see our August 1970 edition in Phys. Letters 33B, 1):

- DODD 69, compiling $\sim 7600 K^+ p \rightarrow K_S^+ p$ events produced at 3-3.5 GeV/c, see an excess of events at $M = 1080$ MeV and a small peak at $M = 1260$ MeV, $\Gamma \approx 70$ MeV. Antiselecting for Δ_{33}^{++} , the peak is enhanced to 4.6 σ .
- CRENNELL 69 have 3044 $K^- n \rightarrow K_S^- n$ events produced at 3.9 GeV/c, and see a 5 σ peak at $M = 1160$ MeV, $\Gamma = 90$ MeV, without Δ_{33} antiselection. The effects of DODD 69 and CRENNELL 69 tend to cancel, if added.
- FIRESTONE 70 see a spike at $M = 1247 \pm 5$ MeV, $\Gamma = 20^{+9}_{-6}$ MeV in the reaction $K^+ n \rightarrow K^+ \pi^- p$ at 12 GeV/c.
- The international K^+ compilation (private communication, V. Henri, CERN) has 77,220 $K^+ p \rightarrow K^+ \pi^- p$ events produced at 2.5-12.7 GeV/c, and see no structure in the $K^+ \pi^-$ mass distribution.

Thus the combined information on narrow peaks is contradictory or cancelling, to the extent that no safe conclusions can be drawn about the existence of resonances.

REFERENCES FOR KN1080-12601	
TRIPPE 68 PL 28 B 203	+CHIEN,MALAMUD,MELLEMA,SCHLEIN,+ (UCLA)
CRENNELL 69 PRL 22 487	+KARSHON,LAI,D,NEALL,SCARR (BNL)
DODD 69 PR 177 1994	+JOLDERSMA,PALMER,SAMIOS (BNL)
GOLDSPERG 69 PL 30 B 434	SABRE COLLABOR. (SACL+AMST+BGNA+REHO+EPDL)
SCHLEIN 69 ARGONNE CONF. 446	P.SCHLEIN (UCLA)
FIRESTON 70 UCL-20091	A.FIRESTONE,G.GOLDBERGER,D.LISSAUER (LRL)
ASTIER 70 KIEV CONF.	RAPP,TALK ON BOSON RESONANCES (CDF)

$K_A, I = 3/2 (1175)$

24 KA 3/2 (1175, JP = 1 I = 3/2)
 EVIDENCE NOT COMPELLING. OMITTED FROM TABLE FOR A DISCUSSION SEE ROSENFELD 68 AND GIACOMELLI 70 WHO CONCLUDES THAT IF THIS STATE HAS WIDTH NOT LARGER THAN 100 MEV, THEN ITS PRODUCTION CROSS SECTION IS 1 OR 2 ORDERS OF MAGNITUDE SMALLER THAN THAT OF NON-EXOTIC K^* 'S.

REFERENCES ON KA3/2(1175)	
WANDLER 64 PL 9 71	T P WANDLER, A R ERWIN, M D WALKER (WISCONS)
MILLER 65 PL 15 74	MILLER,KOVACS, MCILWAIN, PALFREY + (PURDUE)
ROSENFEL 68 PHILA. CONF. P. 455	A.H.ROSENFELD (LRL)
DODD 69 PR 177 1991	+JOLDERSMA, PALMER, SAMIOS (BNL)
CHI 70 PL 32 B 409	+DERRICK, JOHNSON, MUSGRAVE, + (ANL+NWES+KANS)
GIACOMELLI 70 PL 33 B 373	G.GIACOMELLI + (BGNA+SACL+ZEEM+REHO+EPDL)

$K_A, I = 3/2 (1265)$

25 KA 3/2 (1265, JP = 1 I = 3/2)
 EVIDENCE NOT COMPELLING. OMITTED FROM TABLE FOR A DISCUSSION SEE ROSENFELD 68.

REFERENCES ON KA3/2 (1265)	
FRENCH 67 NC 52A 442	+KINSON+MC DONALD+RIDDFORD+ (CERN+BIRM)
ROSENFEL 68 PHILA. CONF. P. 455	A.H.ROSENFELD (LRL)
CHI 70 PL 32 B 409	+DERRICK, JOHNSON, MUSGRAVE, + (ANL+NWES+KANS)

Q REGION, $K\pi\pi(1240-1400)$

There is a peak in the $K^* \pi$ spectrum centered at about 1300 MeV, or just about 270 MeV above $K^* \pi$

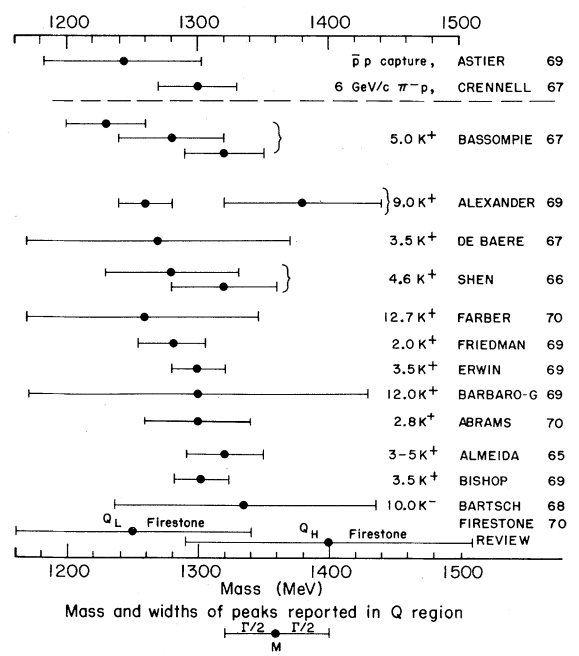
threshold. The $Q(\rightarrow K^* \pi)$ is then the hypercharge = 1 analog of the $A_1(\rightarrow \rho \pi)$, and its interpretation as resonance vs. threshold enhancement is ambiguous. (For more discussion see text Sect. III and the note on "A1, A3, Q, L", just before the A1 in these listings.)

However, in the case of the Q region, there are two additional complications:

- 1) The $K_A^*(1240)$ or "C" meson, with width only 120 MeV, seems to be produced nondiffractively in $\bar{p}p$ capture at rest, and in the reaction $\pi^- p \rightarrow \Lambda K \pi \pi$.
- 2) The Q peak in the $K\pi\pi$ mass distribution does not have a simple Breit-Wigner shape (FIRESTONE 70 and private communication from V. Henri, CERN). It can be fitted by two Breit-Wigner shapes (5-parameter fit) at all energies with reasonable χ^2 (FIRESTONE 70).

The figure shows M and Γ for all peaks reported in this region. We list first the results of the $\bar{p}p$ production of the "C", and the $\pi^- p \rightarrow \Lambda Q^0$ experiment, so as to separate them from the conventional diffractive experiments $K^+ p \rightarrow Q^+ p$. The remaining Q reports are ordered by increasing mass.

The dominant J^P assignment throughout the Q region is 1^+ . $K^* \pi$ is the dominant decay mode. While there is $K\rho$ and interference between $K^* \pi$ and $K\rho$, there is no evidence for other decay modes. From the coherent production of Q in $K^- d$ experiments, $I = 1/2$ is confirmed.



Mesons

For notation, see illustrated key at beginning of data card listings.

28 Q REGION I=1/2

28 Q REGION MASSES (MEV)

Produced by Beams Other than K Mesons	M	W	A	E	M	M	M	M
1242.0 9.0 10.0 ASTIER 69 HBC 0 PBAR P	1242.0	9.0	10.0	ASTIER	69	HBC	0	PBAR P
45(1300.)	1298.8	11.3	11.3	CRENNELL	67	HBC	0.6	PI-LK2PI

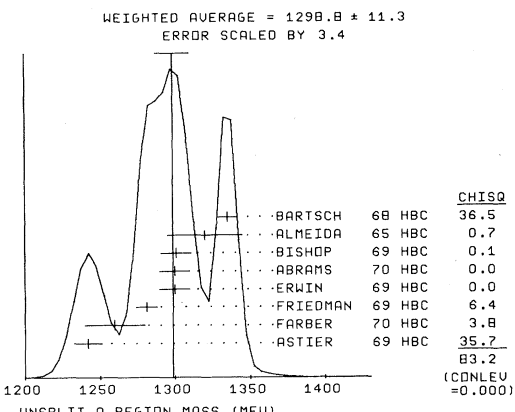
PRODUCED BY BEAMS OTHER THAN K MESONS

M	W	A	E	M	M	M	M
1230.0	11/67	1260.0	5/70	1260.0	6/70	1270.0	9/69

PRODUCED BY K BEAMS

M	W	A	E	M	M	M	M
1230.0	11/67	1260.0	5/70	1260.0	6/70	1270.0	9/69

AVG 1298.8 ± 11.3 AVERAGE ERROR INCLUDES SCALE FACTOR OF 3.4 (SEE IDEOGRAM BELOW)



28 Q REGION WIDTHS (MEV)

PRODUCED BY BEAMS OTHER THAN K MESONS

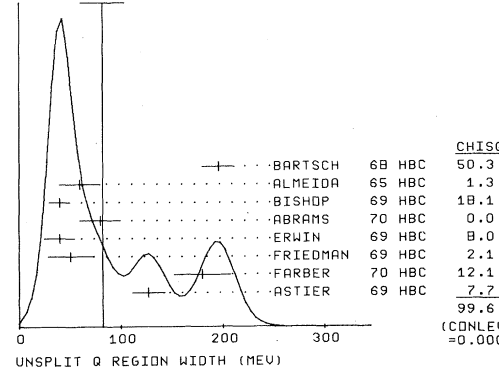
M	W	A	E	M	M	M	M	
127.0 7.0 25.0 ASTIER 69 HBC 0 PBAR P	127.0	7.0	25.0	ASTIER	69	HBC	0	PBAR P
45 (60.)	7/67	45 (60.)	7/67	CRENNELL	67	HBC	0.6	PI-P

PRODUCED BY K BEAMS

M	W	A	E	M	M	M	M
1230	11/67	1260	5/70	1260	6/70	1270	9/69

AVG 82.5 ± 21.6 AVERAGE ERROR INCLUDES SCALE FACTOR OF 3.0 (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 82.5 ± 21.6
ERROR SCALED BY 3.0



28 Q REGION PARTIAL DECAY MODES

P1	Q REGION INTO K*(RHO) PI	DECAY MASSES
P2	Q REGION INTO K RHO	493+ 139+ 139
P3	Q REGION INTO K PI	493+ 139+ 139
P4	Q REGION INTO K ETA	493+ 139+ 139
P5	Q REGION INTO K OMEGA	493+ 139+ 139
P6	Q REGION INTO K PI PI	493+ 139+ 139

28 Q REGION BRANCHING RATIOS

PRODUCED BY BEAMS OTHER THAN K MESONS

R1	Q REGION INTO (K RHO)/TOTAL (UNITS OF 10**2)	M	W	A	E	M	M	
R1	75.0 10.0 ARMENTERO 64 HRC	75.0	10.0	ARMENTERO	64	HRC	0.0	PBAR P
R2	25.0 10.0 ARMENTERO 64 HRC	25.0	10.0	ARMENTERO	64	HRC	0.0	PBAR P
R9	0.2 OR LESS (CL=.90) CRENNELL 67 HBC	0.2	OR LESS	CRENNELL	67	HBC	0.6	0.0 PI-P
R10	0.1 OR LESS (CL=.90) CRENNELL 67 HBC	0.1	OR LESS	CRENNELL	67	HBC	0.6	0.0 PI-P

PRODUCED BY K BEAMS

R1	Q REGION PEAK INTO (K PI) / (K*(RHO) PI)	M	W	A	E	M	M	
R1	0.8 OR LESS SHEN 66 HBC	0.8	OR LESS	SHEN	66	HBC	4.6	K+P, S BODY
R1	70 Q REGION PEAK INTO K*(RHO) PI AND K RHO (OVERLAPPING BANDS)	70	OR LESS	SHEN	66	HBC	4.6	K+P
R2	0.1 OR LESS SHEN 66 HBC	0.1	OR LESS	SHEN	66	HBC	4.6	K+P
R2	0.1 OR LESS SHEN 66 HBC	0.1	OR LESS	SHEN	66	HBC	4.6	K+P
R2	0.1 OR LESS SHEN 66 HBC	0.1	OR LESS	SHEN	66	HBC	4.6	K+P
R1 A	200 (1.0) BERLINGHI 67 HRC	200	(1.0)	BERLINGHI	67	HRC	12.7	K+P
R2 A	0.02 OR LESS BERLINGHI 67 HRC	0.02	OR LESS	BERLINGHI	67	HRC	12.7	K+P
R2	0.02 OR LESS, C.L.=.95 BARTSCH 68 HBC	0.02	OR LESS, C.L.=.95	BARTSCH	68	HBC	10.0	K-P
R3	0.02 OR LESS BERLINGHI 67 HRC	0.02	OR LESS	BERLINGHI	67	HRC	12.7	K+P
R4	0.02 OR LESS BERLINGHI 67 HRC	0.02	OR LESS	BERLINGHI	67	HRC	12.7	K+P
R4	12 0.01 0.005 BARTSCH 68 HBC	12	0.01 0.005	BARTSCH	68	HBC	10.0	K-P
R5	0.1 OR LESS BERLINGHI 67 HRC	0.1	OR LESS	BERLINGHI	67	HRC	12.7	K+P
R5	0.1 OR LESS BERLINGHI 67 HRC	0.1	OR LESS	BERLINGHI	67	HRC	12.7	K+P
R5	0.4 0.1 BARTSCH 68 HRC	0.4	0.1	BARTSCH	68	HRC	10.0	K-P
R5	0.47 0.18 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.9	0.47	0.18	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.9				
R6	0.21 OR LESS DE BAERE 67 HRC	0.21	OR LESS	DE BAERE	67	HRC	3.5	K+P
R7	0.22 0.08 BARTSCH 68 HRC	0.22	0.08	BARTSCH	68	HRC	10.0	K-P

REFERENCES FOR Q REGION

PRODUCED BY BEAMS OTHER THAN K MESONS

ARMENTERO 64 DUBNA CONF 1 577 ARMENTEROS, EDWARDS, D-ANDLAU + (CERN+CDF)

ALSO DURBA CONF 1 611 R ARMENTEROS (RAPORTEUR) BARASH, KIRSCH, MILLER, TAN (COLUMBIA)

ALSO 66 PRL 145 1095 KALBFLEISCH, LAI, SCARR, SCHUMANN (BNL)

CRENNELL 67 PRL 19 44 +MRECHAL, MONTANET, + (CDF+CERN+INP+LIP+IJP)

ASTIER 69 NP B 10 65 +CRESTI, LIMENTANI, FERTAUZA, RIGI+ (PARO+PISA)

BETTINI 69 NC 62 A 1038

PRODUCED BY K BEAMS

ALMEIDA 65 PL 16 184 ALMEIDA, ATHERTON, RYER, DORNAN, FORSON+ (CAMBR)

SHEN 66 PRL 17 726 +BUTTERWORTH, FU, GOLDBERS, TRILLING (LRL)

ALSO 66 (PRIVATE COMMUNICATION) GOLDBER (LRL)

BASSOMPIERRE 67 PL 26B 30 BASSOMPIERRE, GOLDSCHMIDT+ (CERN+BRUX+BIW+LIJ)

BERLINGHI 67 PRL 18 1087 BERLINGHI, FARBER, FERREL, FORMAN, (TRICH) LIJ

DE BAERE 67 NC 49A 374 +DEBAISIEUX, FAST, FILIPPAS+ (CERN+BRUX)

AND PRIVATE COMMUNICATION BY B. JONGEJANS

BARTSCH 68 NP B8 9 +CICCONI, + (AACH+BERL+CERN+DIC+VIEN)

DMSE 68 PRL 20 1519 +DRENSTEIN, CALLAHAN, GILLESPIE, + (JHHHPK) 1+

DEGRI 68 PRL 20 1194 +CALLAHAN+FTLINGR+GILLESPIE+ (JHHHPK) 1+

ALSO 70 ANTIC

Mesons

For notation, see illustrated key at beginning of data card listings.

ALFAXNDE 69 NP B 13 503
 ANDREWS 69 PRL 22 731
 BARBARO 69 PRL 22 1207
 COLLEY 69 NC A 50 519
 WERNER 69 PR 18R 2023
 BISHOP 69 NP B 9 403
 FRWIN 69 NP B 9 366
 FRIEDMAN 69 UCLL-18R60
 CHIEN 69 PL 29B 433
 CHUNG 69 PR 182 1443
 WERNER 69 PR 18R 2023

G. ALEXANDER, FIRESTONE, GOLDBERGER, + (LRL)
 +LACH, LUDLAM, SANDWEISS, BERGER, + (YALE+LRL)
 BARBARO-GALTIERI, DAVIS, FLATTE, + (LRL)
 +EASTWOOD, + (BIERM+GLAS+LIDIC+MDI+OXF+ARIEL)
 +AMMAR, DAVIS, KRUPAC, YARGER, CHU, + (NINES+ANL) 1+
 +GOSHAM, ERWIN, WALKER (WISC)
 +WALKER, GOSHAM, WEINBERG (WISC+PRIN+VAND)
 J. FRIEDMAN, PH. D. THEISIS (LRL)
 +MALAMUD, MELLEMA, RUDHICK, SCHLEIN+ (UCLA)
 +FISNER+BALL+LIJERS (BNL)
 +AMMAR, DAVIS, KRUPAC, YARGER, CHU+ NINES+ANL

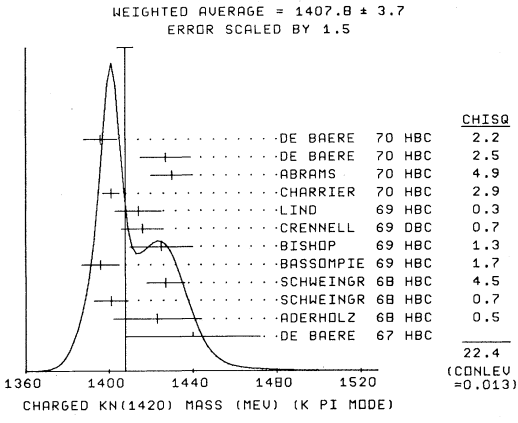
ABRAMS 70 PR D 1 2433
 ANTICH 70 NP B 20 201
 BOWLER 70 PL 31 A 318
 FARRER 70 PR D 1 78
 FIRESTON 70 PHILAD, CNF, P. 229 A. FIRESTONE REVIEW

+FISENSTEIN, KIM, MARSHALL, O-HALLORAN, + (ILL)
 +CARSON, CHIEN, COX, DENFRI, ETTLINGER, + (JHU) 1+
 M. G. BOWLER (OXFORD)
 +FERREL, SLATTERY, YUTA (ROCHA) 1+
 (LRL) 1+

K_N(1420)
 or **K****
 22 KN (1420, JP=2+) I=1/2
 JP = 3- IS UNLIKELY BUT NOT YET COMPLETELY RULED OUT.

22 KN(1420) MASS (MEV)
 M FOR DIFFICULTIES IN MEASURING MASS DIFFERENCE, SEE TYPED NOTE UNDER K*

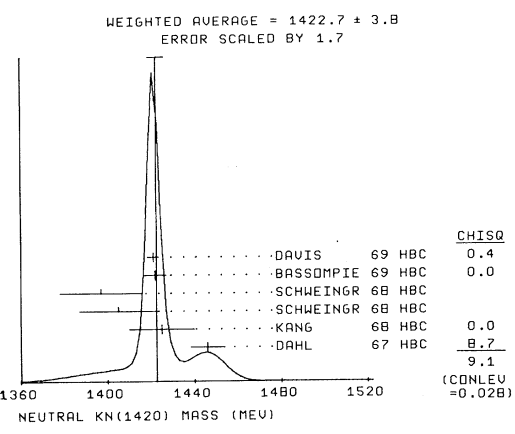
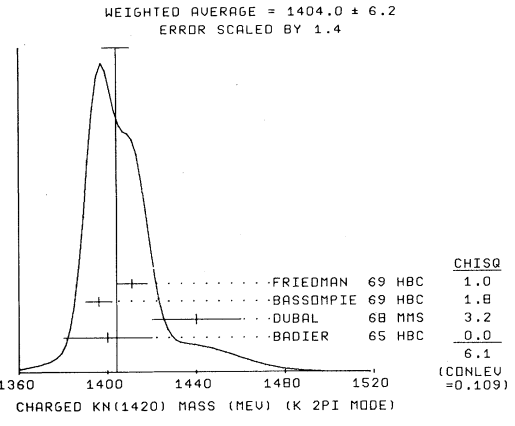
M CHARGED ONLY, WITH FINAL STATE K PI
 M O (1427.) (15.0) DE BAERE 67 HRC + 3.5 K+P (K0 PI+) 10/66
 M D VALUE ABOVE SUPERSERVED BY DE BAERE 70 BELOW 1/71*
 M 1460. 24.0 DE BAERE 67 HRC + 3.5 K+P (K+ PI0) 10/66
 M 1423. 21. ADERHOLZ 68 HRC - 10 K- P (K PI) 6/68
 M 1401.0 8.0 SCHWEINGR 68 HRC - 4.1 K- P (K PI) 9/67
 M 1427.0 9.0 SCHWEINGR 68 HRC - 5.5 K- P (K PI) 9/67
 M B 125 1396. 9. BASSOMPIE 69 HRC + 5 K+P (K PI) 11/69
 M 1425.0 15.0 BISHOP 69 HRC + 3.5 K+ P 9/69
 M 1416.0 10.0 CRENNELL 69 HRC - 3.9 K-N (K0 PI-) 7/69
 M 1414. 11. LIND 69 HRC + 9. K+ P (K0 PI+) 9/69
 M 1401.0 4.0 CHARRIER 70 HRC + 8.25 K+ P 1/71*
 M 1430. 14.0 ARAMS 70 HRC + 2.5-3.2 K+P, K2PI 11/70*
 M 1427.0 12.0 DE BAERE 70 HRC + 3.5 K+ P 1/71*
 M 1396.0 8.0 DE BAERE 70 HRC + 5.0 K+ P 1/71*
 M AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
 M AVG 1407.8 3.7 (SEE IDEOGRAM BELOW)



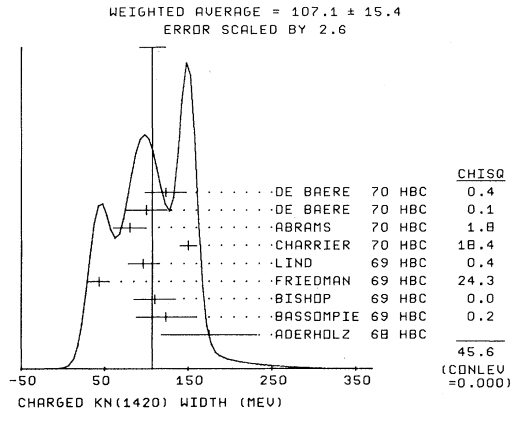
M CHARGED ONLY, WITH OTHER FINAL STATES
 M 1400.0 20.0 BADIER 65 HRC - 3. K- P (K*PI) 10/66
 M 20 1440.0 20.0 DUBAL 68 HRC - 11.5 K- P 6/68
 M R 240 1396. 6. BASSOMPIE 69 HRC + 5 K+P (K 2PI) 11/69
 M 1411. 7. FRIEDMAN 69 HRC - 2.7 K-P (K 2PI) 9/69
 M AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
 M AVG 1404.0 6.2 (SEE IDEOGRAM BELOW)

M CHARGED AND NEUTRAL
 M 1404.0 15.0 FOCARDI 65 HRC - 0.3 K- P (K PI) 10/66
 M 1390.0 30.0 SHEN 66 HRC + 0.4, 6 K+ P (K PI) 10/66
 M 1430.0 10.0 SHEN 66 HRC + 0.4, 6 K+ P (K*PI) 10/66
 M 1423.0 7.0 BASSANI 57 HRC - 0.4, 6, 5.0 K- P 10/67
 M 1420.0 10.0 GOLDBERGER 67 HRC 9.0 K+ P (K 2PI) 10/67
 M AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 M AVG 1421.2 4.7 (SEE IDEOGRAM BELOW)

M NEUTRAL ONLY
 M (1440.1) CRENNELL 67 HRC 0.6 PI- P (K 2PI) 7/67
 M 1444.0 7.9 DAHL 67 HRC 0.4 PI- P (KPI) 10/66
 M 1425.0 15.0 KANG 68 HRC 0.4, 6 K- P 7/69
 M 1405.0 14.9 SCHWEINGR 68 HRC 0.4, 1 K- P (K PI) 9/67
 M 1397.0 19.0 SCHWEINGR 68 HRC 0.4, 5 K- P (K PI) 9/67
 M B 420 1422. 5. BASSOMPIE 69 HRC 0.5 K+P (K PI) 11/69
 M BASSOMPIE ERRORS ENLARGED BY US TO GAMMA/SORTIN). SEE K* TYPED NOTE. 11/69
 M 2200 1421.1 2.6 DAVIS 69 HRC 0.12, K+ P (K+PI-) 9/69
 M AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)
 M AVG 1422.7 3.8 (SEE IDEOGRAM BELOW)



22 KN(1420) WIDTH (MEV)
 M CHARGED ONLY, WITH FINAL STATE K PI
 M O (140.0) (29.0) DE BAERE 67 HRC + 3.5 K+ P 10/66
 M D VALUE ABOVE SUPERSERVED BY DE BAERE 70 BELOW 1/71*
 M 175. 57. ADERHOLZ 68 HRC - 10 K- P (K PI) 6/68
 M B 125 123. 35. BASSOMPIE 69 HRC + 5 K+P (K PI) 11/69
 M 110.0 25.0 BISHOP 69 HRC + 3.5 K+ P 9/69
 M 43. 13. FRIEDMAN 69 HRC - 2.7 K-P (K 2PI) 9/69
 M 96. 18. LIND 69 HRC + 9. K+ P 5/70
 M 150.0 10.0 CHARRIER 70 HRC + 8.25 K+ P 1/71*
 M 80.0 20.0 ARAMS 70 HRC + 2.5-3.2 K+P, K2PI 11/70*
 M 100.0 25.0 DE BAERE 70 HRC + 3.5 K+ P 1/71*
 M 123.0 25.0 DE BAERE 70 HRC + 5.0 K+ P 1/71*
 M AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.6)
 M AVG 107.1 15.4 (SEE IDEOGRAM BELOW)



Mesons

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle type (e.g., CHARGED ONLY, NEUTRAL ONLY), mass (MEV), width (MEV), and experimental references with their respective errors and scale factors.

WEIGHTED AVERAGE = 101.2 ± 8.4
ERROR SCALED BY 1.1

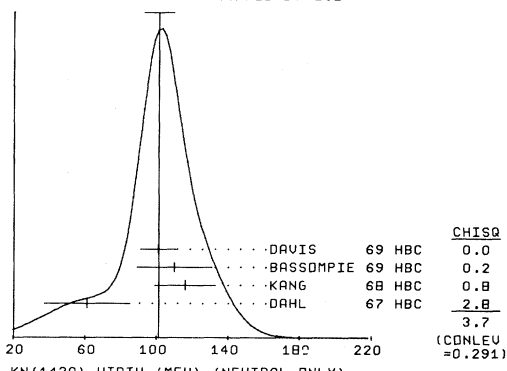


Table listing various meson decays and their experimental measurements, including decays into K Omega, K Rho, K Eta, and K Pi Pi. It includes references to specific experiments and their results.

Table titled 'FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS'. It provides a matrix of branching fractions (P1 to P5) for different decay modes. The diagonal elements represent the individual branching fractions, and the off-diagonal elements represent the normalized correlation coefficients.

Table titled 'REFERENCES FOR KN(1420)'. It lists the experimental papers and authors for the various measurements of the KN(1420) meson, including references to journals like Physical Review Letters and Nuclear Physics.

Table titled '22 KN(1420) PARTIAL DECAY MODES' and '22 KN(1420) BRANCHING RATIOS'. It provides a detailed breakdown of the branching ratios for different decay channels, such as into K Pi, K Rho, K Omega, and K Eta. It includes statistical fits and correlations between different modes.

For notation, see illustrated key at beginning of data card listings.

K_N(1660) 27 KN (1660, JP = 1) I = 1/2
EVIDENCE NOT COMPELLING, OMITTED FROM TABLE

27 KN (1660) MASS (MEV)

M	(1660.0)		CARMONY	67 HBC	- 3.8 K-P, OMEGA K	11/67
M	1660.0	10.0	JOBES	67 HBC	+ 5. K+ P	11/67
M	J		CLAIMED BY JOBES IN (K PI), (K*(890) PI), AND (K*(1420) PI)			
M	J		MODES. K PI HUMP INTERFERES MOSTLY WITH DELTA(1236).			

27 KN (1660) WIDTH (MEV)

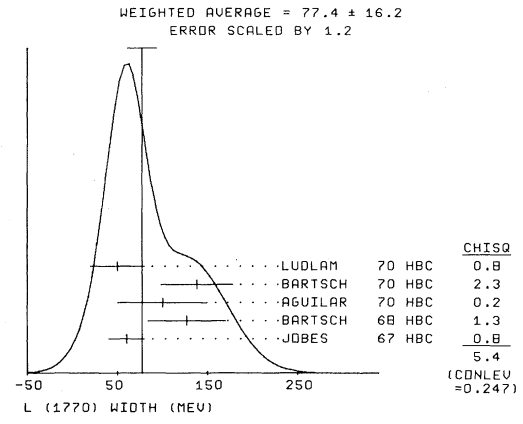
M	60.0	20.0	JOBES	67 HBC	+ 5. K+ P	11/67
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27 KN (1660) PARTIAL DECAY MODES

P1	KN*(1660)	INTO K PI	DECAY MASSES
P2	KN*(1660)	INTO K PI PI	493+ 139
P3	KN*(1660)	INTO K*(890) PI	493+ 139+ 139
P4	KN*(1660)	INTO KN*(1430) PI	892+ 139
			1409+ 139

REFERENCES FOR KN(1660)

CARMONY	67 PRL	18	615	D. CARMONY, T. HENDRICKS, L. LANDER (LA JOLLA)
JOBES	67 PL	268	49	+ BASSOMPIERRE, DE BAERE + (BIRM+CERN+BRUX)



L(1770) 23 L (1770, JP = 1) I = 1/2

For a review of the L-meson data up to May 1970, see CHIEN 70.

This peak is seen in most K⁺p and K⁻d experiments, with a large decay mode into K_N(1420)π, only ~200 MeV above that threshold. There is at present no agreement among experimenters about whether to interpret this peak as a resonance, with decay into several open channels, or as a pure K_N(1420)π threshold enhancement of possibly "kinematic" origin. We list only one branching ratio

$$L \rightarrow K_N(1420)\pi$$

$$L \rightarrow K\pi\pi$$

which has bearing on this contradiction.

23 L (1770) MASS (MEV)

M	20(1780.1)		BERLINGHI	67 HBC	+ 12.7 K+P	7/67
M	1760.0	15.0	JOBES	67 HBC	+ 5. K+ P	11/67
M	1785.0	12.0	BARTSCH	68 HBC	10.0 K- P	9/69
M	B		INCLUDED IN BARTSCH 70			
M	1745.0	20.0	AGUILAR	70 HBC	- 4.6 K- P	6/70
M	1780.0	15.0	BARTSCH	70 HBC	- 10.1 K- P	1/71*
M	(1753.0)	(10.0)	BIRMINGH.	70 HBC	+ 10. K+P, K(1420)PI	11/70*
M	(1773.0)	(20.0)	BIRMINGH.	70 HBC	+ 10. K+P, K* 2PI	11/70*
M	1760.0	15.0	LUPLAM	70 HBC	- 12.6 K- P	1/71*
M	AVG	77.4	7.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)		

23 L (1770) WIDTH (MEV)

M	20 (80.)		BERLINGHI	67 HBC	+ 12.7 K+P	7/67
M	60.0	20.0	JOBES	67 HBC	+ 5. K+ P	11/67
M	127.0	43.0	BARTSCH	68 HBC	10.0 K- P	9/69
M	B		INCLUDED IN BARTSCH 70			
M	100.0	50.0	AGUILAR	70 HBC	- 4.6 K- P	6/70
M	138.0	40.0	BARTSCH	70 HBC	- 10.1 K- P	1/71*
M	(176.0)	(23.0)	BIRMINGH.	70 HBC	+ 10. K+P, K(1420)PI	11/70*
M	(155.0)	(22.0)	BIRMINGH.	70 HBC	+ 10. K+P, K* 2PI	11/70*
M	50.0	40.0	LUPLAM	70 HBC	- 12.6 K- P	1/71*
M	AVG	77.4	16.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2) (SEE IDEOGRAM BELOW)		

23 L (1770) PARTIAL DECAY MODES

P1	L INTO K PI PI	DECAY MASSES
P2	L INTO K*(1420) PI	497+ 134+ 134
		134+1409

23 L (1770) BRANCHING RATIOS

R1	L INTO (K*(1420) PI) / (K PI PI)		
R1	LARGE	DENEGRI	68 HBC - 12.6 K- P 1/71*
R1	(1.0)	BARBARO	69 HBC + 12.0 K+ P 1/71*
R1	(0.60) OR LESS CL=0.95	AGUILAR	70 HBC - 0.6 K- P 1/71*
R1	LESS THAN 1.0	BARTSCH	70 HBC - 10.1 K- P 1/71*
R1	LARGE	BIRMINGH.	70 HBC + 10. K+P, 4-6 BODY 11/70*
R1	DOMINANT	LUPLAM	70 HBC - 12.6 K- P 1/71*

REFERENCES FOR L (1770)

BERLINGHI	67 PRL	18	1087	BERLINGHIERI+BARBER+FERBEL+FORMAN+ (ROCH)I
JOBES	67 PL	268	49	+ BASSOMPIERRE, DE BAERE + (BIRM+CERN+BRUX)
DENEGRI	68 PRL	20	1194	+ CALLAHAN+ETTLINGER+GILLIESIE+ (JOHNSHOP)
BARTSCH	68 NP	BR	9	+ CROCONI,+ (AACH+BERL+CERN+LOIC+VIEN)
ANDREWS	69 PRL	22	731	+ LACH, LUPLAM, SANDWEISS, BERGER,+ (YALE+LRL)
BARBARO	69 PRL	22	1207	BARBARO+GALTIERI, DAVIS, FLATTE,+ (LRL)
COLLEY	69 NC	A	59	519 +EASTWOOD,+ (BIRM+GLAS+LOIC+MPIM+OXF+RHEL)
AGUILAR	70 PRL	25	94	AGUILAR-BENITEZ, BARNES, PASSANO, CHUNG,+ (RNL)
BARTSCH	70 PL	33	8	186 +FEUTSCHMANN,+ (AACH+BERL+CERN+LOIC+VIEN)
BIRMINGH	70	PREPRINT		BIRMINGHAM GLASGOW COLL. (BIRM+GLAS)
CHARRIER	70	CERN PHYS	70-50	CHARRIERE, GOLDSCHMIDT-C, DEBAERE+(CERN+BRUX)
CHIEN	70	PHILA. CONF.	P.275	C.V. CHIEN (JOHNS HOPKINS)
LUPLAM	70	PR D	2	1234 +SANDWEISS, SLAUGHTER (YALE)

K*(2240) 40 K* (2240, JP = 1) I = 1/2
ENHANCEMENT SEEN IN (ANTI)HYPERON+NUCLEON) MASS.
EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

40 K* (2240) MASS (MEV)

M	15	2240.	20.	ALEXANDER	68 HBC	+ 0.9 K+P, YBAR+N+...	6/68
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40 K* (2240) WIDTH (MEV)

M	15	70.	20.	ALEXANDER	68 HBC	+ 0.9 K+P, YBAR+N+...	6/68
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REFERENCES FOR K*(2240)

ALEXANDE	68 PRL	20	755	ALEXANDER, FIRESTONE, GOLDBER, SHEN (LRL)
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Baryons

Note on Resonance Parameters from Partial-Wave Analyses

I. Speed Plots

In the discussion which follows, we use the term "speed plot" to indicate a plot showing the variation with C. M. energy m of the derivative $|dT/dm|$ of a partial-wave amplitude T . (See section IV.C of the main text.) In principle such plots are a very sensitive and useful means of searching for a resonance. A rapid increase in speed followed by a rapid decrease is certainly a good indication of the presence of a resonance. In practice these plots must be judiciously used because:

- 1) The values of dT/dm are sensitive to variations in T . It is difficult enough to determine $T(m)$; finding its derivative is necessarily more difficult.
- 2) Once the speed plot tells us that a resonance is present, the determination of precise parameters from such a plot requires additional considerations:

a) the maximum of the speed is not necessarily at the resonance mass,

b) the width cannot simply be obtained by the relation $|dT/dm|_{m=M} = 2\kappa/\Gamma$.

Consider for example the P_{33} partial-wave amplitude in π -N scattering. Since its elasticity (κ) is one, we have

$$T(m) = \frac{\Gamma(m)/2}{M-m - i\Gamma(m)/2}. \quad (1)$$

If we let $\Gamma^1(m) = d\Gamma/dm$, then we find that

$$\text{"Speed"} = \left| \frac{dT}{dm} \right| = \frac{2}{\Gamma} \frac{1+(M-m)\Gamma^1/\Gamma}{1+4(M-m)^2/\Gamma^2}. \quad (2)$$

To estimate where Eq. (2) is maximum, we let $\delta = M-m$ ($|\delta| < \Gamma/2$) and keep only linear terms in δ .

Thus

$$\left| \frac{dT}{dm} \right| \approx \frac{2}{\Gamma} \left(1 + \delta \frac{\Gamma^1}{\Gamma} \right). \quad (3)$$

Since all reasonable parametrizations of $\Gamma(m)$ agree that $\Gamma^1 \geq 0$, we may conclude that to first order the "speed" will have its maximum value at an energy less than the resonant value, $m=M$.

This effect is illustrated in Fig. 1, which is taken from UCRL-20030 π N.¹ For the P_{33} partial wave, the CERN experimental and CERN Kirsopp solutions indicate the instability of $|dT/dm|$ in the region of a resonance (the other solutions are "smooth" by the

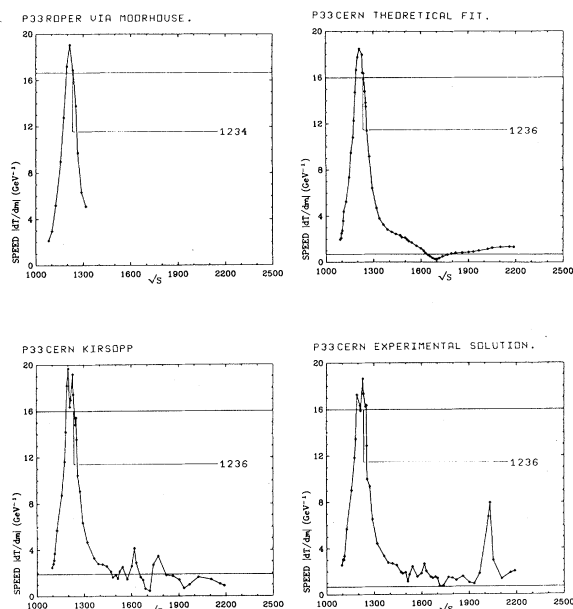


Fig. 1. Speed plots as computed from four solutions compiled in Ref. 1. ($\sqrt{s} = m =$ C. M. energy in MeV)

nature of the analysis). In addition, each of the plots, quite consistently, gives $2/\Gamma \approx 16 \text{ GeV}^{-1}$ at a resonant mass of $\sim 1236 \text{ MeV}$. This corresponds to a width at resonance of $\sim 125 \text{ MeV}$. The speed, however, peaks some 10 to 15 MeV lower in mass and at a value of $\sim 18.5 \text{ GeV}^{-1}$. Hence, were we to estimate the mass and width of the 33-resonance from the maximum speed, we would get $M \approx 1220 \text{ MeV}$ and $\Gamma = 108 \text{ MeV}$. For additional discussion on the mass and width of this resonance, see the mini-review at the beginning of the $\Delta(1236)$ listings.

II. Effect of Partial-Wave Parametrization

To appreciate the systematic differences that arise in extracting resonance parameters from the data, we have compared several recent analyses in Table I. The notation is standard: T is the scattering amplitude, $S = 1 + 2i T$, and f is the spin non-flip amplitude. The subscripts B and R refer to background and resonant contributions, respectively. The first part of the table refers to analyses of the π N system, the second part to the $\bar{K}N$.

Baryons

Table I. Representative examples of energy-dependent parametrizations used to determine resonance parameters.

	Authors	E _{CM} Range (MeV)	Data	Model	Background
πN	AYED 70	1320-2480	π [±] p, ↑, CEX, ↑	$\begin{cases} 1. T = T_B + T_R e^{i\phi_R} \\ 2. S = S_B \cdot S_R \end{cases}$	$\delta_B = \sum a_n q^n$ $\eta_B = 1 / [1 + (\sum b_m q^m)^2]$ a _n , b _m , φ _R free
	DAVIES 70	1320-1960	π [±] p, ↑, CEX, ↑	$S = e^{2i\phi} \cdot \frac{\Gamma_{el}}{M - E - i\Gamma/2}$	$\phi = a + bx + cx^2$ $B = 1 / (1 + dx + ex^2)$ $x = k^2 - (460 \text{ MeV}/c)^2$ α, a-e free
	HULL 70	2030-2560	π ⁻ p, ↑ t ≤ .6(GeV/c) ²	f = f _B + f _R	f _B = (α+i) $\frac{k\sigma}{4\pi} e^{bt/2}$ α, σ, b free
KN	BRICMAN-1 70	1780-1950	K ⁻ p, ↑, CEX	T = T _B + T _R e ^{iφ_R}	T _B = a + b(p - p ₀) a, b free, complex φ _R free, real
	GALTIERI 70	1640-2170	Δπ, Σπ	T = T _B + T _R e ^{iφ_R}	T _B = (a + bk) exp[i(c + dk)] a, b, c, d, φ _R free, real
	KIM 71	1430-1880	KN, Δπ, Σπ	K matrix, multichannel	Expand K matrix in successive intervals
	LITCHFIELD 70	1800-2170	Δπ	T = T _B + (T _R)	T _B = $\sum_{n=0}^N a_n P_n(x)$ P _n are Legendre polyns. x is linear in momentum such that x = ±1 at p _{max} p _{min} . a _n complex, free

The problem of adding background to a resonant amplitude has been extensively discussed in the literature.² It is not very clear what is meant by background. After it has been parametrized as an expansion in the incident momentum (or other similar parametrizations as shown in Table I), it may turn out to move fast enough across the Argand plot to represent a resonance. In this case the parameters obtained for the T_R part of the amplitude are very much dependent upon the parametrization of T_B, and one might have to reparametrize the amplitude with an additional resonance. This is an extreme situation, but in general any nonconstant background is very critical to the determination of resonance parameters.

We now discuss the models displayed in Table I.

A. πN System. There have been several determinations made of the parameters for nucleon resonances. We have tabulated all these parameters in Table II of the following "Note on N's and Δ's." It is useful for

extending the comparison here, which we restrict to the analyses mentioned in Table I of the present note.

1. AYED 70 have done an energy-independent phase-shift analysis and then fit their experimental Argand plots (reported in Fig. 3 of Note on N's and Δ's) in order to obtain resonance parameters. They use the two methods shown in Table I for combining background and resonances. Both essentially can be written as

$$T = T_B + \gamma T_R \quad \text{with} \quad T_B = \frac{S_B - 1}{2i} = \frac{\eta_B e^{2i\delta_B} - 1}{2i}$$

and

- Model 1: γ₁ = e^{iφ_R}, φ_R real, free parameter
- Model 2: γ₂ = η_B e^{2iδ_B}, i.e., S matrix for the background.

For only one open channel the two are identical, since unitarity requires that e^{iφ_R} = η_B e^{2iδ_B}, with η_B = 1. Model 2 is most likely invalid (see for ex-

Baryons

ample Goebel and McVoy²) when S_B becomes inelastic. This can be understood by considering the corresponding Argand amplitude:

$$T = T_B + \eta_B e^{2i\delta_B} \frac{x}{\epsilon - 1}, \quad \epsilon = \frac{2(M-m)}{\Gamma}, \quad x = \Gamma_{el}/\Gamma.$$

If, for some reason, $\eta_B \approx 0$ when $m \approx M$ and a clear resonant circle were present, than an abnormally large elasticity ($x > 1$) might be obtained. For this reason we included in the listings only those parameters from model 1. Two typical examples of how the parameters depend on the model are shown in Table II. For the rather elastic F_{15} state, the two sets of

Table II. Two examples of resonance parameters from models 1 and 2 of AYED 70.

		M	Γ	x	ϕ_R
F_{15}	Model 1	1682	109	.59	-.03
	Model 2	1694	128	.62	---
P_{13}	Model 1	1766	182	.149	-1.58
	Model 2	1508	047	.123	---

parameters are in reasonable agreement. However, there is a difference of 12 MeV in the mass, 19 MeV in the width, and .03 in the elasticity, certainly larger than the statistical errors on the data would indicate. This is an excellent example of why we do not average parameters from energy-dependent analyses when errors of ± 2 MeV are quoted. The other example is known as P_{13} (1860) in our listings. The results of the two models disagree not only with one another but also with all of the older values. In addition, model 2 seems to require a second P_{13} at ~ 1896 MeV!

2. DAVIES 70 (Table I) uses a model similar to model 1 of AYED 70. The main difference between these two analyses is in the procedure: AYED 70 do an energy-independent partial-wave analysis first and then fit the Argand plots so obtained. DAVIES 70 fit the data directly with their model; therefore they obtain smooth Argand plots (Fig. 2 of Note on N's and Δ 's). The amplitude is written in a form suggested by Goebel and McVoy,² who discuss the multichannel generalization of the Breit-Wigner approximation. In principle² α (or ϕ_R) can be constrained through unitarity from the other elements of the S matrix. The resonant parameters of DAVIES 70 are generally quite different from those of AYED 70 (see Table II in "Note on N's and Δ 's") possibly because AYED 70 include newer experimental data and use a different fitting procedure as mentioned earlier.

3. HULL 70. This is not strictly a partial-wave analysis, but it does emphasize that useful information on resonances can be obtained from the forward elastic diffraction peak. This was first noted by Damouth et al.³ and elaborated upon by Levi Setti and co-workers.⁴ Essentially, HULL 70 parametrize all background effects and the tails of resonances (outside the region considered) by the diffraction-like background f_B . Through a fit to recent forward π^-p data, supplemented by forward polarization data, they were able to establish the spin of the $H_{19}(2220)$. Their resonance parameters are in reasonable agreement with those of the more orthodox analysis of AYED 70.

B. $\bar{K}N$ system. Most of the partial-wave analyses reported in the $\bar{K}N$ system are energy dependent, therefore the resonance parameters are very much dependent upon the parametrization used to fit the experimental data. We report in Table I only a few representative examples of such analyses.

1. BRICMAN 70 uses model 1 of AYED 70 but parametrizes the background with a simple linear dependence in momentum. This simple parametrization, used extensively by the CERN-Heidelberg-Saclay collaboration, has proven quite adequate for discovering many new resonant states in the 1-GeV/c region. It was introduced as an expedient in the belief that background cannot vary too rapidly in an interval of ~ 200 MeV. In addition the initial paucity of $\bar{K}N$ data could not permit a more elaborate parametrization (AYED 70 have many background parameters).

The following studies, which we discuss, indicate several methods for extending this sort of analysis to include an ever-broader region in energy.

2. GALTIERI 70 uses a background parametrization that tends to have a circular behavior in the Argand plot. The analysis is performed by dividing the overall region into two intervals so far as the background is concerned. Continuity in $|T|$ is required and, of course, the same resonant parameters are used for both regions.

In the first region ($\lesssim 1900$ MeV), the resonant structure is somewhat better known. The inclusion of the second region in the fit is initially made with only a background contribution (with the exception of the well-known $G_{07}(2100)$ and $F_{17}(2020)$). If a resulting partial wave (background!?) shows a resonant-like behavior, the fit is repeated with the pres-

ence of a Breit-Wigner in that wave. In this manner it was found that both the D_{03} and D_{13} required a second resonant state at higher energy.

3. KIM 71 presents a multichannel K-matrix analysis of $\bar{K}N$ data from threshold to ~ 1880 MeV. In this analysis he parametrizes the S, P, D, and F_{05} waves with the K matrix. In principle this is the ideal manner in which to analyze the $\bar{K}N$ data; however, such an analysis is not immune to problems.

The first problem is the treatment of three-body states. KIM 71 handles them by introducing an additional "junk" or "quasi-two-body" channel. Such an approximation is most likely reasonable if the three-body final states constitute a small part of the total cross section. The additional channel can be to some extent constrained by the measured three-body cross section.

The next problem is the energy dependence used for the K matrix. Kim's approach is to use an effective range expansion of K^{-1} . Beyond ~ 1500 MeV, he expands K^{-1} in successive intervals of ~ 100 MeV. In each of these intervals a new set of ranges is determined by a fit, while continuity is ensured at the boundaries of the regions. This method will ensure smooth Argand plots and, very importantly, unitarity is observed. However, finding the resonance parameters by searching for the poles in the K matrix is now impractical (uses K matrix for every 100 MeV) and one has to resort to some other method (reading parameters off the Argand plot).

4. LITCHFIELD 70 first fits the data with each partial wave (except F_{17} and D_{15}) parametrized by the Legendre polynomial expansion (LPE) in Table I. He then looks for counterclockwise ("resonant like") behavior in the resulting Argand plots. Such a behavior usually requires the order N to be 4 or 5; if it exists, the fit is repeated with a lower-order LPE and the addition of a Breit-Wigner amplitude. If the new fit is just as good, he takes this as evidence for a resonant effect.

For $N=1$ the LPE corresponds to the linear background used by BRICMAN-1 70; for higher orders it can simulate the behavior of GALTIERI 70's background. Thus the LPE attempts to generalize the useful features of these analyses. As a practical point, LITCHFIELD 70 notes that the use of the LPE generally reduces the computer time for minimization. Since the a_n are relatively uncorrelated, the reduction of N to 1 and the inclusion of a Breit-Wigner would, hopefully, tend to leave a_0 and a_1 unchanged.

References

1. D. Herndon et al., " πN Partial-Wave Amplitudes, a Compilation," UCRL-20030 πN , Feb. 1970.
2. Discussion on the "addition" of background and resonances can be found, among others, in the following papers: R. H. Dalitz, *Ann. Rev. Nucl. Sci.* **13**, 339 (1963); C. J. Goebel and K. W. McVoy, *Phys. Rev.* **164**, 1932 (1967); K. T. R. Davies and M. Baranger, *Ann. Phys. (N. Y.)* **19**, 383 (1962); C. Michael, *Phys. Letters* **21**, 93 (1966); and S. Coleman, in Theory and Phenomenology in Particle Physics, edited by A. Zichichi (Academic Press, New York, 1969).
3. D. E. Damouth, L. W. Jones, and M. L. Perl, *Phys. Rev. Letters* **11**, 287 (1963).
4. R. Levi Setti, T. Lasinski, and E. Predazzi, *Phys. Rev.* **179**, 1429 (1969).

Note on N's and Δ 's

There are now complete phase-shift analyses from four different groups: The Saclay group (referred to as BAREYRE 68 and AYED 70), the Berkeley group (JOHNSON 67), the Glasgow group (DAVIES 70), and the CERN group.

The CERN group has performed two phase-shift analyses, using different methods. The CERN I solution is published as DONNACHIE-1 68 for both I spin 1/2 and 3/2. Their figures contain two sorts of results:

1. "Experimental Phase Shifts," i. e., partial-wave amplitudes at each energy at which they used experimental input. In their figures these are plotted as, η and δ at each energy, but not as Argand plots.
2. "Theoretical Fits" using smooth functions based on dispersion-relation theory. These are plotted both as smooth curves of η and δ vs energy, and as Argand plots. Brody et al.¹ have recently criticized the "Theoretical Fits" because it turns out that although the "experimental" amplitudes describe the data as well as (or better than) any other available set, the theoretical fits for some rapidly varying partial waves are too smooth. Because they are so convenient to draw and to remember, we continue to present these smooth Argand plots, having warned the reader of their limitations.

The new solution, CERN II,² covers I=3/2 only, and has been published only as Argand plots of "experimental" amplitudes.

Baryons

The most recent and extensive analysis comes from the Saclay group (AYED 70); it extends up to a mass of 2.5 GeV. They obtain values for M, Γ, x through energy-dependent fits of the Argand plots (see the preceding mini-review).

We reproduce below, in Figs. 1, 2, 3, most of the available Argand diagrams. The Berkeley diagrams, from which the authors do not yet quote resonance parameters, are reproduced here only for $I = 1/2$ partial waves.³ Table I below is a summary of all the states claimed by the various groups with our evaluation of their significance. We have included in the Baryon Table only states listed as "good" or "fair."

Spread in Values of Resonance Parameters

Values of masses, widths, and branching ratios can be obtained only from phase-shift analyses. In production experiments, in fact, it is seldom clear which of the many states at similar masses is being observed. We now have phase-shift analyses from four different groups, but we are quite far from having reliable masses and widths derived therefrom.

The problem is that the errors on the phase shifts are quite large and it is thus difficult to draw smooth curves on the Argand diagrams. In addition, except for the Glasgow solutions and the results of AYED 70, where energy-dependent fits to the data and phase shifts are done, the resonance parameters are just the result of an "eyeball" fit with the use of different methods. As a result, different authors using the same phase shifts often estimate different values of M, Γ, x . This is the case for the CERN I solution, from which three sets of parameters have been reported. Furthermore, as discussed in the preceding note, different assumptions about background and resonant shapes can often lead to quite dissimilar values for M, Γ, x even in energy-dependent analyses.

In order to make the reader aware of these problems, we present below in Table II all the different values for M, Γ, x . We also present in this table the mean values of these parameters and the "Ind. Ext. Error," or the "external error" of the individual values:

$$\langle \delta x_i \rangle = \sqrt{\frac{1}{N} \sum_1^N (x_i - \bar{x})^2}$$

The error $\delta \bar{x}$ of the mean is of course smaller by another factor $1/\sqrt{N}$ but we avoid giving it because we feel that $\bar{x}, \delta \bar{x}$ have little meaning here. In order that the reader may appreciate the spread in the values of M, Γ, x given in Table II, we discuss briefly the different methods used to determine these parameters. AYED

70 analyze their phase-shift results with an energy-dependent background and Breit-Wigner amplitudes. BAREYRE 68 uses two methods to find resonance parameters: 1 (cross section method)--the energy where the total cross section is maximum; 2 (speed method)-- the energy where the speed of variation of the amplitude in the Argand plot is maximum. CERN quotes only one method, usually where the absorption is maximum, but three different sets of values have been given. The Glasgow group (DAVIES 70) uses Breit-Wigner parametrization; A and B differ in the starting values of the minimization (CERN I solution was used for solution B). For some states no parameters have been quoted by the authors. We report in the M column of Table II our evaluation of the status of this resonance as judged on the published Argand plots. Symbols are the same as on Table I.

On the main table of Baryon Resonances we decided not to quote a value with an error, but to quote a range of masses and widths in order to point out the large indeterminacy of these parameters. So the P_{11}^+ will be $M = 1435$ to 1505 MeV, $\Gamma = 164$ to 400 MeV, etc.

The resonant parameters, of course, depend on the validity of the partial-wave analyses. Several recent experimental results have not been included in the phase-shift analyses summarized here. Regions of disagreement have been noted by Abillon et al.⁴ (in the backward direction) and by Albrow et al.⁵ (at $M = 2100$ MeV and higher). Future analyses are likely to provide different parameters, especially in the 2-GeV region.

Footnotes and References

1. A. D. Brody, D. W. G. S. Leith, B. G. Levi, B. C. Shen, D. Herndon, R. Longacre, L. Price, A. H. Rosenfeld, and P. Söding, *Phys. Rev. Letters* **22**, 1401 (1969).
2. The CERN II Argand plots have been reported by A. Donnachie, 14th International Conference on High Energy Physics, Vienna, 1968, p. 139.
3. For the complete set of Argand plots including speed versus energy, see UCRL-20030 πN (Feb. 1970) by D. J. Herndon, A. Barbaro-Galtieri, A. H. Rosenfeld. The AYED 70 solutions are not included, since they were not available at the time.
4. J. M. Abillon, A. Borg, M. Crozon, T. Leray, J. P. Mendiburu, and J. Tocqueville, *Phys. Letters* **32B**, 712 (1970).
5. M. G. Albrow, S. Andersson/Almehed, B. Bosnjakovic, C. Daum, F. C. Erne, J. P. Lagnaux, J. C. Sens, F. Udo, *Nucl. Phys.* **B25**, 9 (1970).

Baryons

PARTIAL-WAVE AMPLITUDES FOR CERN I, CERN II, AND BERKELEY SOLUTIONS

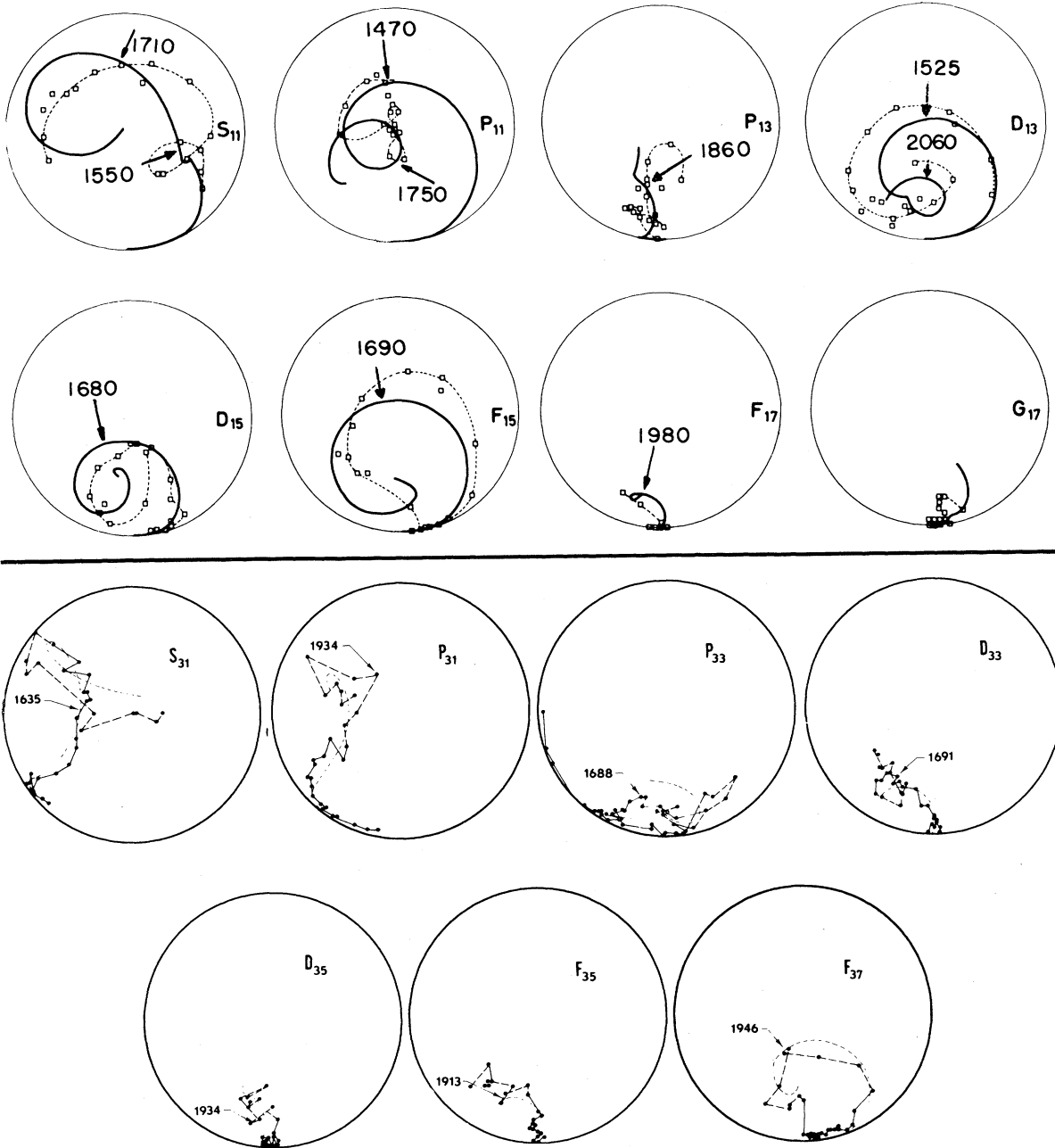


Fig. 1. Results of the phase-shift analyses of the CERN and Berkeley groups. The CERN I results are the smooth curves (dashed in the $I = 3/2$ diagrams). This analysis used dispersion relations to join and smooth the solutions found at different energies. The arrows in the $I = 1/2$ diagrams indicate approximate resonance positions; they have been drawn by us. The CERN II solution is shown (as a dot-dash line) only for the $I = 3/2$ amplitudes since the $I = 1/2$ are not available. The arrows have been drawn by the authors. The Berkeley solution is shown only for the $I = 1/2$ state (as empty squares joined by dashes).

Baryons

PARTIAL-WAVE AMPLITUDES FOR GLASGOW SOLUTIONS

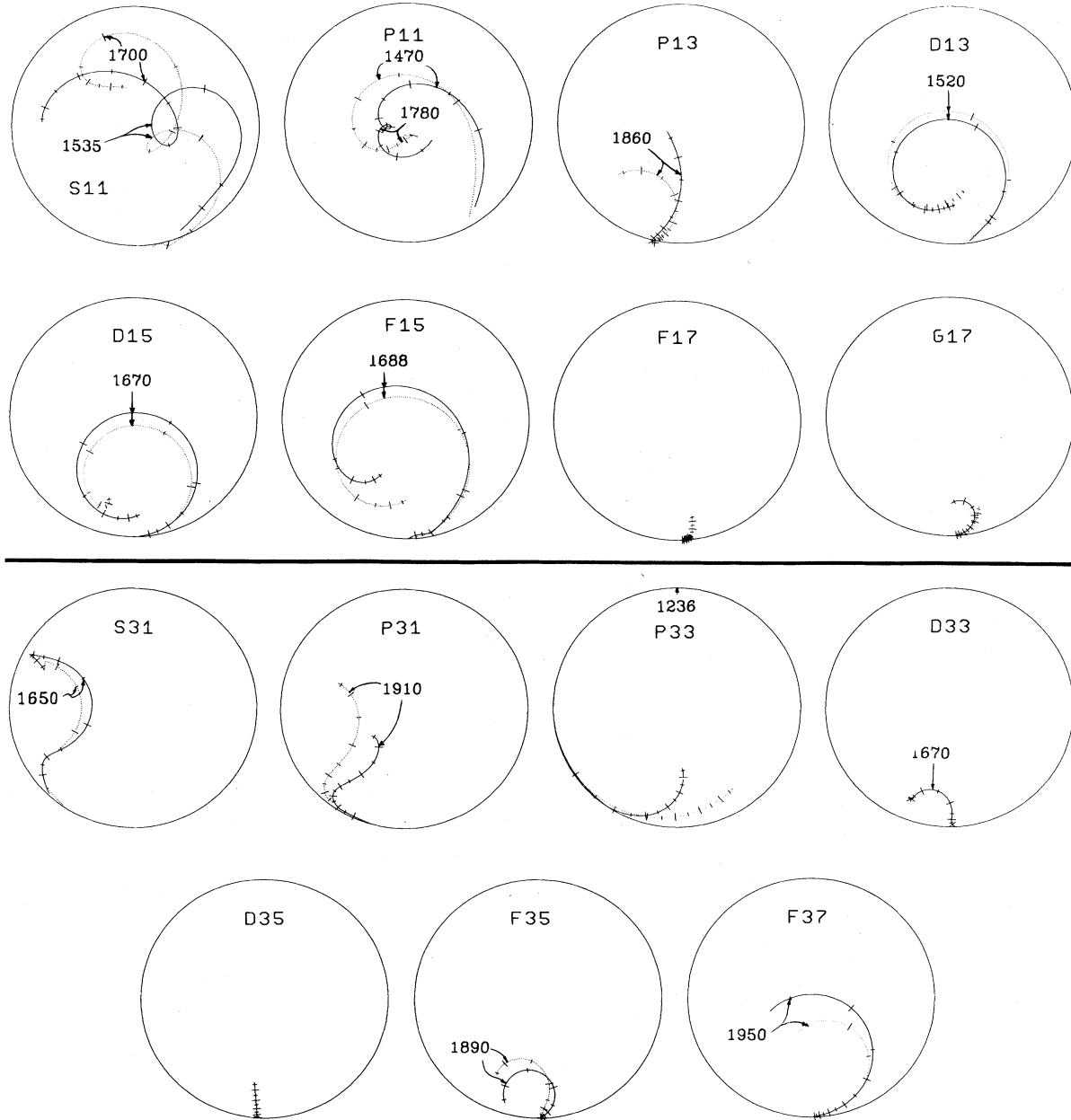


Fig. 2. Phase-shift analysis results of DAVIES 70. The curves plotted here are the results of the energy-dependent analysis. DAVIES A (solid curves) is obtained by starting from the set of phase shifts that is the best solution of the Glasgow group. DAVIES B (dashed curves) is obtained by starting from the CERN I phase shifts. DAVIES B is not shown when it is very close to DAVIES A. Scale marks are shown every 50 MeV. The first large mark is at $M = 1400$ MeV, the last large mark at $M = 1900$ MeV.

Baryons

PARTIAL-WAVE AMPLITUDES FOR SACLAY SOLUTIONS

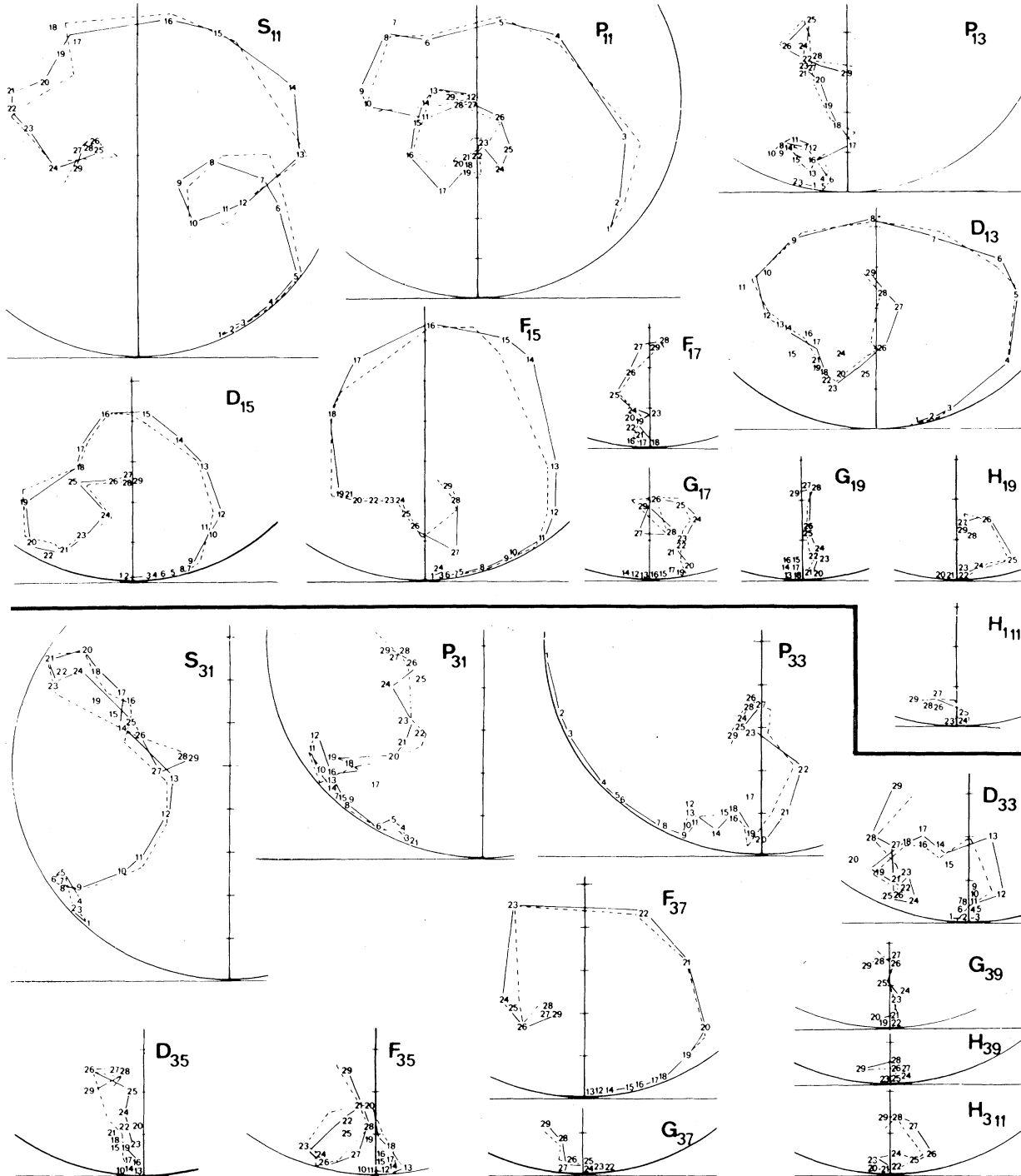


Fig. 3. Latest Saclay phase-shift analysis results [Phys. Letters 31B, 5981 (1970)]. The full and dashed lines connect points of "minimal path" and "minimal surface", respectively.

Baryons

Table I. Our evaluation of the status of all N and Δ resonances as seen in partial-wave analyses. D = definite, Pr = probable, Po = possible, A = ambiguous, No = not present. Notice that in the Glasgow fits the resonance hypothesis is built into the fit, so only the symbols D or No apply, except for one Pr at the upper end.

	Berkeley	CERN I	Saclay ^f	Glasgow	RBD ^a	CERN II	Our evaluation						
							ηn	KΛ	KΣ	πΔ	ρN	γN	
P ₁₁ ' (1470)	D	D	D	D			Good				D		
D ₁₃ ' (1520)	D	D	D	D			Good				D		D
S ₁₁ ' (1535)	D	D	D	D			Good	D				D	D
D ₁₅ (1670)	D	D	D	D			Good				D		
F ₁₅ (1688)	D	D	D	D			Good				D		D
S ₁₁ '' (1700)	D	D	D	D			Good	Po	D				D
D ₁₃ '' (1700)	Po	Po	A	No			Poor		Po				
P ₁₁ '' (1780)	Pr	Pr	Pr	D			Fair	Po	D				D
P ₁₃ ' (1860)	A ^c	A	Po	Pr	Pr		Fair		Po				
F ₁₇ (1990)	c	Pr	A	e	A		Poor						
D ₁₃ ''' (2040)	c	Pr	A	c	A		Poor						
G ₁₇ (2190)	c	D	D	c	Pr		Good						
H ₁₉ (2220)	c	c	D	c	c		Good						
P ₃₃ ' (1236)							Good						D
S ₃₁ '' (1650)	D	D	D	D		D	Good				Po ^b		
P ₃₃ '' (1690)	Po	A	Po ^g	No		A	Poor				Po ^b		
D ₃₃ (1670)	A	D	D	D		D	Fair				Po ^b		
F ₃₅ (1890)	Pr	Pr	Pr	D		Pr	Fair		Po				
P ₃₁ (1910)	Pr ^c	Pr	Po	D		D	Fair						
D ₃₅ (1960)	c	A	No	c	Po	A	Poor						
F ₃₇ (1950)	D	D	D	D		D	Good		D	D	D	D	
P ₃₃ ''' (2160)	Po ^c	Po	No	c		d	Poor						
H ₃₁₁ (2420)	c	c	D	c			Good						

a. RBD = LEA 69 (Lea et al., Ruth, Bristol, Daresbury) and APLIN 70.
 b. For these references see DONNACHIE-2, the latter part of the article.
 c. This state is very close to or beyond their highest energy.
 d. We can't say anything.
 e. Glasgow A has a G₁₇ state, Glasgow B may have an F₁₇. However, this region is very close to their highest energy.
 f. Evaluation based on new analysis (AYED 70).
 g. Around 1800 MeV, however.

Baryons

Table II. Mean Values of Resonance Parameters for N's and Δ's.

Method					M	Γ	x
● P ₁₁ ⁺ (1470)	1	Ayed	70	1464	164	0.56	
	2	Bareyre	68	1470	255	0.68	
	3	Bareyre	68	Speed	1505	205	0.68
	4	Berkeley	67	D			
	5	Donnachie-1	68	Abs.	1466	211	0.658
	6	Donnachie-2	68	Abs.	1470	211	0.66
	7	Kirsopp	68	Abs.	1466	211	0.66
	8	Glasgow	68	A	1462	391	0.49
	9	Glasgow	68	B	1436	224	0.46
Average				1467	234	0.61	
± Ind. ext. error				±18	±64	±.08	
● D ₁₃ ⁺ (1520)	1	Ayed	70	1523	131	0.59	
	2	Bareyre	68	σ	1510	125	0.54
	3	Bareyre	68	Speed	1515	110	0.54
	4	Berkeley	67	1526 ^b	114 ^b	0.57 ^b	
	5	Donnachie-1	68	Abs.	1541	149	0.509
	6	Donnachie-2	68	Abs.	1520	114	0.57
	7	Kirsopp	68	Abs.	1526	115	0.57
	8	Glasgow	68	A	1512	106	0.45
	9	Glasgow	68	B	1512	125	0.49
Average				1521	121	0.54	
± Ind. ext. error				±9	±12	±.04	
● S ₁₁ ⁺ (1535)	1	Ayed	70	1534	96	0.40	
	2	Bareyre	68	σ	1535	155	—
	3	Bareyre	68	Speed	1515	105	—
	4	Berkeley	67	1548 ^b	116 ^b	0.326 ^b	
	5	Donnachie-1	68	Abs.	1591	(268)	0.696
	6	Donnachie-2	68	Abs.	1550	116	0.33
	7	Kirsopp	68	Abs.	1540	160	0.3
	8	Glasgow	68	A	1502	(36)	0.36
	9	Glasgow	68	B	1499	53	0.35
Average				1535	114	0.40	
± Ind. ext. error				±27	±34	±.13	
● D ₁₅ ⁺ (1670)	1	Ayed	70	1675	143	0.39	
	2	Bareyre	68	σ	1680	135	0.41
	3	Bareyre	68	Speed	1655	105	0.41
	4	Berkeley	67	D			
	5	Donnachie-1	68	Abs.	1678	173	0.391
	6	Donnachie-2	68	Abs.	1680	173	0.391
	7	Kirsopp	68	Abs.	1678	175	0.391
	8	Glasgow	68	A	1669	115	0.50
	9	Glasgow	68	B	1667	115	0.43
Average				1673	142	0.41	
± Ind. ext. error				±8	±27	±.04	
● F ₁₅ ⁺ (1688)	1	Ayed	70	1682	109	0.59	
	2	Bareyre	68	σ	1690	110	0.64
	3	Bareyre	68	Speed	1680	105	0.64
	4	Berkeley	67	1692 ^b	132 ^b	0.68 ^b	
	5	Donnachie-1	68	Abs.	1687	177	0.56
	6	Donnachie-2	68	Abs.	1690	132	0.68
	7	Kirsopp	68	Abs.	1692	130	0.68
	8	Glasgow	68	A	1685	104	0.54
	9	Glasgow	68	B	1684	123	0.54
Average				1688	125	0.62	
± Ind. ext. error				±4	±21	±.06	
● S ₁₁ ⁺ (1700)	1	Ayed	70	1689	166	0.64	
	2	Bareyre	68	σ	1710	260	—
	3	Bareyre	68	Speed	1665	110	—
	4	Berkeley	67	1709 ^b	300 ^b	0.786 ^b	
	5	Donnachie-1	68	Abs.	—	—	—
	6	Donnachie-2	68	Abs.	1740	300	0.79
	7	Kirsopp	68	Abs.	1709	300	0.79
	8	Glasgow	68	A	1766	404	0.56
	9	Glasgow	68	B	1671	121	0.51
Average				1704	245	0.68	
± Ind. ext. error				±29	±96	±.12	
● D ₁₃ ⁺ (1700)	1	Ayed	70	A			
	2	Bareyre	68	σ			
	3	Bareyre	68	Speed	Po		
	4	Berkeley	67	Po			
	5	Donnachie-1	68	Abs.	—	—	—
	6	Donnachie-2	68	Abs.	1730		
	7	Kirsopp	68	Abs.	1680		
	8	Glasgow	68	A	No		
	9	Glasgow	68	B	No		
Average				1705			
± Ind. ext. error				±25			
Method					M	Γ	x
● P ₁₁ ⁺ (1780)	1	Ayed	70	1645	50	0.15	
	2	Bareyre	68	σ	Fr		
	3	Bareyre	68	Speed	Fr		
	4	Berkeley	67	Fr			
	5	Donnachie-1	68	Abs.	1751	327	0.32
	6	Donnachie-2	68	Abs.	1750	327	0.32
	7	Kirsopp	68	Abs.	1860	270	0.32
	8	Glasgow	68	A	1770	445	0.43
	9	Glasgow	68	B	(1867)	(525)	0.30
Average				1755	284	0.31	
± Ind. ext. error				±68	±130	±.08	
● P ₁₃ ⁺ (1860)	1	Ayed	70	1766	182	0.15	
	2	Bareyre	68	σ	A ^b		
	3	Bareyre	68	Speed	A ^b		
	4	Berkeley	67	A ^b			
	5	Donnachie-1	68	Abs.	1863	296	0.207
	6	Donnachie-2	68	Abs.	1860	296	0.21
	7	Kirsopp	68	Abs.	1900	325	0.25
	8	Glasgow	68	A	1844	449	0.40
	9	Glasgow	68	B	1854	307	0.26
	10	Lea	69	1860	—	—	
Average				1849	309	0.25	
± Ind. ext. error				±38	±78	±.08	
● F ₁₇ ⁺ (1990)	1	Ayed	70	No			
	2	Bareyre	68	σ	b		
	3	Bareyre	68	Speed	b		
	4	Berkeley	67	b			
	5	Donnachie-1	68	Abs.	1983	225	0.128
	6	Donnachie-2	68	Abs.	—	—	—
	7	Kirsopp	68	Abs.	1995	250	0.09
	8	Glasgow	68	A	c		
	9	Glasgow	68	B	c		
	10	Lea	69	~2000	—	—	
Average				1993	238	0.109	
± Ind. ext. error				±7	±13	±.019	
● D ₁₃ ⁺ (2040)	1	Ayed	70	No			
	2	Bareyre	68	σ	b		
	3	Bareyre	68	Speed	b		
	4	Berkeley	67	b			
	5	Donnachie-1	68	Abs.	2057	293	0.26
	6	Donnachie-2	68	Abs.	2030	290	0.11
	7	Kirsopp	68	Abs.	2040	240	0.15
	8	Glasgow	68	A	b		
	9	Glasgow	68	B	b		
	10	Lea	69	2030	—	—	
Average				2039	274	0.17	
± Ind. ext. error				±11	±24	±.06	
● G ₁₇ ⁺ (2190)	1	Ayed	70	2158	325	0.15	
	2	Bareyre	68	σ	b		
	3	Bareyre	68	Speed	b		
	4	Berkeley	67	b			
	5	Donnachie-1	68	Abs.	2265	298	0.349
	6	Donnachie-2	68	Abs.	2190	300	0.35
	7	Kirsopp	68	Abs.	2265	300	0.35
	8	Glasgow	68	A	(1906) ^c	(319) ^c	(0.14) ^c
	9	Glasgow	68	B	c		
	10	Lea	69	~2000			
Average				2176	306	0.30	
± Ind. ext. error				±97	±11	±.09	
● S ₃₁ ⁺ (1650)	1	Ayed	70	1614	142	0.32	
	2	Bareyre	68	σ	1695	250	—
	3	Bareyre	68	Speed	1650	130	—
	4	Berkeley	67	D			
	5	Donnachie-1	68	Abs.	1635	177	0.284
	6	Donnachie-2	68	Abs.	1640	177	0.28
	7	Kirsopp	68	Abs.	1635	180	0.28
	8	Glasgow	68	A	1617	141	0.28
	9	Glasgow	68	B	1623	140	0.25
Average				1639	167	0.28	
± Ind. ext. error				±24	±36	±.02	
● P ₃₃ ⁺ (1690)	1	Ayed	70	1801	598	0.14	
	2	Bareyre	68	σ	A		
	3	Bareyre	68	Speed	A		
	4	Berkeley	67	Po			
	5	Donnachie-1	68	Abs.	1688	281	0.098
	6	Donnachie-2	68	Abs.	1690	281	0.1
	7	Kirsopp	68	Abs.	1690	240	0.08
Method					M	Γ	x
8	Glasgow	68	A	No			
9	Glasgow	68	B	No			
Average				1717	350	0.10	
± Ind. ext. error				±48	±144	±.02	
● D ₃₃ ⁺ (1670)	1	Ayed	70	1722	258	0.22	
	2	Bareyre	68	σ	Po		
	3	Bareyre	68	Speed	Po		
	4	Berkeley	67	A			
	5	Donnachie-1	68	Abs.	1691	269	0.14
	6	Donnachie-2	68	Abs.	1690	269	0.14
	7	Kirsopp	68	Abs.	1690	300	0.13
	8	Glasgow	68	A	1649	188	0.12
	9	Glasgow	68	B	1650	174	0.13
Average				1682	243	0.15	
± Ind. ext. error				±26	±46	±.03	
● F ₃₅ ⁺ (1890)	1	Ayed	70	1837	198	0.15	
	2	Bareyre	68	σ	Po		
	3	Bareyre	68	Speed	Po		
	4	Berkeley	67	Pr			
	5	Donnachie-1	68	Abs.	1913	350	0.16
	6	Donnachie-2	68	Abs.	1910	350	0.16
	7	Kirsopp	68	Abs.	1910	380	0.15
	8	Glasgow	68	A	1841	136	0.2
	9	Glasgow	68	B	1852	150	0.19
Average				1877	261	0.17	
± Ind. ext. error				±34	±102	±.02	
● P ₃₁ ⁺ (1910)	1	Ayed	70	1783	308	0.13	
	2	Bareyre	68	σ	A ^b		
	3	Bareyre	68	Speed	A ^b		

Baryons

For notation, see illustrated key at beginning of data card listings.

CODE EVENTS QUANTITY ERROR+ ERROR- REFERENCE YR TECN STGN COMMENTS DATE ABOVE PUNCHED BACKGROUND

P 15 PROTON (938, J=1/2) I=1/2
SEE LISTINGS OF STABLE PARTICLES

n 17 NEUTRON (939, J=1/2) I=1/2
SEE LISTINGS OF STABLE PARTICLES

N(1470) 61 Σ^+ (1470, JP=1/2+) I=1/2 **P₁₁**
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE ABOVE

THE MASS AND WIDTH ARE BEST DETERMINED FROM PHASE-SHIFT ANALYSES. WE LIST PRODUCTION EXPERIMENTS SEPARATELY--SEE BELOW.

61 Σ^+ (1470) MASS (MEV)

M	(1340.0)	RUPFER	65	RVUE	PHASE-SHIFT ANAL	9/66
M	(1370.0)	BRANDSEN	65	RVUE	PHASE-SHIFT ANAL	9/66
M	(1470.0)	BARFVRE	68	RVUE	PHASE-SHIFT ANAL	11/67
M	(1505.0)	BARFVRE	68	RVUE	EYEBALL FIT	6/68
M	(1466.0)	DONNACH2	68	RVUE	PHASE-SHIFT ANAL	10/69
M	(1470.0)	DONNACH2	68	RVUE	PHASE-SHIFT-CERN1	10/69
M	(1465.0)	KIRSOPP	68	RVUE	PHASE SHIFT ANAL	10/69
M	(1462.0)	DONNACH1, 2	68	RVUE	KIRSOPP EYEBALL FIT CERN 1	10/69
M	(1436.0)	DAVIES	70	RVUE	P-S ANAL SOL A	8/69
M	(1224.0)	DAVIES	70	RVUE	P-S ANAL SOL A	8/69
M	(1461.0)	AYED	70	IPWA		1/71*
M	(1464.0)	AYED	70	IPWA		1/71*

61 Σ^+ (1470) WIDTH (MEV)

W	1	(285.0)	BARFVRE	68	RVUE	11/67	
W	2	(205.0)	BARFVRE	68	RVUE	11/67	
W	3	(211.0)	DONNACH1	68	RVUE	6/68	
W	3	(211.0)	DONNACH2	68	RVUE	PHAS.SHIFT-CERN1	10/69
W	3	(210.0)	KIRSOPP	68	RVUE	PHASE-SHIFT ANAL	10/69
W	4	(391.0)	DAVIES	70	RVUE	P-S ANAL SOL A	8/69
W	5	(224.0)	DAVIES	70	RVUE	P-S ANAL SOL B	8/69
W	6	(164.0)	AYED	70	IPWA		1/71*

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

61 Σ^+ (1470) PARTIAL DECAY MODES

PARTIAL DECAY MODES		DECAY MASSES	
P1	Σ^+ (1470) INTO Σ^+ P1	139+ 938	
P2	Σ^+ (1470) INTO Σ^+ (SIGMA MESON)	938+ 410	
P3	Σ^+ (1470) INTO Σ^+ (1236) P1	1236+ 130	
P4	Σ^+ (1470) INTO GAMMA N	0+ 938	
P6	Σ^+ (1470) INTO N RHO (P1P1, I=1)	938+ 765	

61 Σ^+ (1470) BRANCHING RATIOS

BRANCHING RATIOS		(P1)/TOTAL				
R1	Σ^+ (1470) INTO Σ^+ P1/TOTAL	(0.68)	11/67			
R1	(0.68)	BARFVRE	68	RVUE	6/68	
R1	(0.65R)	DONNACH1	68	RVUE	10/69	
R1	(.64)	DONNACH2	68	RVUE	PHAS.SHIFT-CERN1	10/69
R1	(.64)	KIRSOPP	68	RVUE	PHASE-SHIFT ANAL	10/69
R1	(0.49)	DAVIES	70	RVUE	P-S ANAL SOL A	8/69
R1	(0.46)	DAVIES	70	RVUE	P-S ANAL SOL B	8/69
R1	(0.56)	AYED	70	IPWA		1/71*
R1	(0.67) (0.18)	SAXON	70	HRC	AT 1400 MEV	6/70
R1	(0.78) (0.09)	SAXON	70	HRC		6/70

A AND B CORRESPOND TO THE 2 BEST SOLUTIONS. ANALYSIS IS DONE ON THREE BODY DECAYS. ASSUMING ONLY P1, P2 AND P3 DECAYS PRESENT. SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

61 Σ^+ (1470) BRANCHING RATIOS CONT.

BRANCHING RATIOS CONT.		(P2)/TOTAL				
R2	Σ^+ (1470) INTO Σ^+ P2/TOTAL	(0.32)	11/67			
R2	(0.32)	THURNAUER	65	RVUE	-	11/67
R2	(0.32)	NAMYSLOW	66	RVUE	-	11/67
R2	(0.30)	ROSENFIELD	67	RVUE	-	11/67
R2	(0.30)	MORGAN	68	RVUE	ISOBAR MODEL	6/68
R2	(0.20)	SAXON	70	HRC		6/70
R2	(0.20)	SAXON	70	HRC		6/70

A AND B CORRESPOND TO THE 2 BEST SOLUTIONS. SEE NOTE IN R1. ANALYSIS IS DONE ON THREE BODY ANALYSIS. ASSUMING R1=0.61

61 Σ^+ (1470) BRANCHING RATIOS CONT.

BRANCHING RATIOS CONT.		(P3)/TOTAL				
R3	Σ^+ (1470) INTO Σ^+ P3/TOTAL	(0.21)	6/70			
R3	(0.21) (0.12)	SAXON	70	HRC		6/70
R3	(0.21) (0.12)	SAXON	70	HRC		6/70

A AND B CORRESPOND TO THE 2 BEST SOLUTIONS. SEE NOTE IN R1. ANALYSIS IS DONE ON THREE BODY ANALYSIS. ASSUMING R1=0.61

61 Σ^+ (1470) BRANCHING RATIOS CONT.

BRANCHING RATIOS CONT.		(P5)/(P1)				
R4	Σ^+ (1470) INTO Σ^+ P5/P1	(0.07)	7/70			
R4	(0.07)	LODI-RIZZ	70	DBC	GAM-N TO PI-P	7/70

61 Σ^+ (1470) BRANCHING RATIOS CONT.

BRANCHING RATIOS CONT.		(P6)/TOTAL				
R5	Σ^+ (1470) INTO Σ^+ P6/TOTAL	(0.07)	1/71*			
R5	(0.07)	DIFM	70	IPWA	3 BODY ANALYSIS	1/71*

ASSUMING R1=0.61

REFERENCES -- Σ^+ (1470)

RUPFER	65	PR	138	R190	LD RUPFER, RM WRIGHT, BT FELD (LRL-LVMR-MIT) IJP
BRANDSEN	65	PR	139	R1566	+DONNELL, MODRHOUSE (DURHAM-RTHF) IJP
THURNAUER	65	PR	14	985	P. G. THURNAUER (RICH)
NAMYSLOW	66	PR	157	1328	NAMYSLOWSKI, RAZMI, ROBERTS (STAN, EDIN, IC)
ROSENFIELD	67	PR	17	CONF	A. H. ROSENFIELD, P. SODING (LRL)

BARFVRE	68	PR	165	1731	P. BARFVRE, C. BRICMAN, G. VILLET (SACLAY) IJP
DONNACH1	68	PL	248	141	A. DONNACHIE, R. G. KIRSOPP, C. LOVELACE (CERN) IJP
DONNACH2	68	VIENNA	139		DONNACHIE, RAPPOORTEUR'S TALK (GLAS)
KIRSOPP	68	THESES			R. G. KIRSOPP (EDIN)
MORGAN	68	PR	166	1731	D. MORGAN (RTHF)

AYED	70	KTEV	CONF		R. AYED, P. BARFVRE, G. VILLET (SACL) IJP
DAVIES	70	MP	R21	359	A. SHADJA, CHAVANON, DELFER, DOLBEAU+ (SACL)
DALITZ	65	PL	14	159	+FLORE, PIAZZA+... (PAVIA+ROMA+CEN+NAPOLI)
LODI-RIZ	70	LNC	3	697	SAXON, MULVEY, CHENOWSKY (OXF-LRL)
SAXON	70	PR	D2	1790	

PAPERS NOT REFERRED TO IN DATA CARDS.

BARFVRE	64	PL	R	137	+BRICMAN, VALLADAS, VILLET, + (SACLAY, CERN) IJP
BARFVRE	65	PL	R	342	+BRICMAN, STIRLING, VILLET (SACLAY) IJP
DALITZ	65	PL	14	159	R. H. DALITZ, R. G. MODRHOUSE (OXF-RTHF)
JOHNSON	67	UCAL	L7683	THESIS	C. H. JOHNSON (LRL)
DONNACH1	69	NP	108	433	A. DONNACHIE, R. KIRSOPP (GLAS+EDIN)
AYED	70	PL	318	598	+BARFVRE, VILLET (SACLAY)
BERARDO	70	PL	24	419	+HADDOCK, WFKENS+... PARSONS+... (UCLA-LRL)

FOR PHOTONIC COUPLINGS SEE RANKIN (GLAS), WALKER (CALT).

THE FOLLOWING ARE THEORETICAL PAPERS CONCERNING THE Σ^+ (1470) --
RESNICK 66 PR 150 1292 L. RESNICK (NIELS RHU) (LRL)
SCHWARZ 66 PR 152 1325 J. H. SCHWARZ (GLAS)
GOLDBERG 67 PR 155 1725 H. S. BALL, GL. SHAW, DY WONG (UCLA, UCL, UCSD) (CORNELL)
SCHWARZ 67 PR 154 1558 M. GOLDBERG (CORNELL)

1470 MEV REGION -- PRODUCTION EXPERIMENTS

61 Σ^+ (1470) PROD. EXP.

PROD. EXP.									
M	(1425.0)	APPROX	ADELMAN	65	HRC	+ K-P	1.45	GEV/C	7/66
M	(1400.0)	APPROX	COCCONI	64	CNTR	+ PP	3.6-12	GEV/C	7/66
M	(1420.0)	APPROX	ANKENBRAN	66	CNTR	+ PP	7.1	GEV/C	7/66
M	(1400.0)	APPROX	RELETTIN	65	SPRK	+ PP, D	10-26	GEV/C	7/66
M	(1405.0)	(15.0)	ANDERSON	66	SPRK	+ PP, A	6-30	GEV/C	7/66
M	(1410.0)	(15.0)	BLAIR	66	CNTR	+ PP	2.0-7.0	GEV/C	7/66
M	(1400.0)	(30.0)	FOLLEY	67	CNTR	P1+	P AND PP		11/67
M	(1450.0)	(17.0)	ALMEIDA	68	HRC	+ PP-P2P1	1.0	GEV/C	10/69
M	(1420.0)	APPROX	RELL	68	HRC	P1+	P, A	GEV/C	6/68
M	(1400.0)	APPROX	LAMSA	68	HRC	P1+	P, A	GEV/C	6/68
M	(1400.0)	(30.0)	SHAPIRA	68	DBC	INTO P1, P N	7.0	10/69	10/69
M	S	TAN 68, SHAPIRA 68	ARF ONLY	PRODUCTION EXPERIMENT TO SEE P1 DECAY					
M	S	(1390.0) (15.0)	TAN	68	HRC	PP T3 PIP, 6-1		10/69	10/69
M	S	(120(1403.0) (15.0)	RHOFF	69	HRC	PP 22	GEV/C		
M	M	1410.0 15.0	ANDERSON	70	MMS	- P1- P TO P1- MMS		2/71*	2/71*

61 Σ^+ (1470) WIDTH (MEV) PROD. EXP.

W	S	(100.0)	RELL	68	HRC	P1+	P AND PP	6/68	
W	S	(175 (100.0) (140.0)	SHAPIRA	68	DBC			10/69	
W	S	(150.0) (60.0)	TAN	68	HRC	+		10/69	
W	S	(120 (100.0) (15.0)	RHOFF	69	HRC	- PP	22	GEV/C	10/69
W	W	210.0 15.0	ANDERSON	70	MMS	- P1- P TO P1- MMS		2/71*	2/71*

61 Σ^+ (1470) BRANCHING RATIOS PROD. EXP.

BRANCHING RATIOS PROD. EXP.								
R1	Σ^+ (1470) INTO Σ^+ P1/TOTAL	(.66)	10/69					
R1	(.66)	TAN	68	HRC	PP TO PIP, 6-1	10/69		
R2	Σ^+ (1470) INTO Σ^+ P2/TOTAL		11/68					
R2	(.30)	PROBABLY SEEN	LAMSA	68	HRC	PP 22	GEV/C	11/68
R2	(.30)	PROBABLY SEEN	JESPERSEN	68	HRC	PP 22	GEV/C	11/68

61 Σ^+ (1470) BRANCHING RATIOS PROD. EXP. CONT.

BRANCHING RATIOS PROD. EXP. CONT.						
R3	Σ^+ (1470) INTO Σ^+ P3/TOTAL		1/71*			
R3	(.07)	DIFM	70	IPWA	3 BODY ANALYSIS	1/71*

REFERENCES -- Σ^+ (1470) -- PROD. EXP.

COCCONI	64	PL	R	134	+LILLETHUN, SCANLON, STAHLBRANDT, + (CERN)
ADELMAN	65	PL	14	1043	S. L. ADELMAN (CAMBRIDGE (CERN))
ANKENBRAN	65	NC	35	1052	ANKENBRANDT, CLYDE, CORK, KEEFE, KERTH+ (LRL)
BELLETTIN	65	PR	18	167	BELLETTIN, COCCONI, DIDDENS + (CERN)
ANDERSON	66	PL	16	855	+BLESER, COLLINS, FUJII + (BNL, CERN)
BLAIR	66	PL	17	789	+TAYLOR, CHAPMAN, + (HARWELL, QUEENMARY, UCL)
GELLERT	66	PL	17	884	+SMITH, VOJTECKI, COLTON, SHELFIN + (LRL, UCLA)
FOLLEY	67	PL	19	397	+JONES, LINDENBAUM, L'OVE, OZART + (BNL)
ALBERT	68	PR	176	1631	+APPEL, RUDNITZ, CHEN, DUNNING, GOTTEN+ (HARV)
ALMEIDA	68	PR	174	1638	+RUSHBRONKH, SCHARFENBERG + (CAVE, DESY)
BELL	68	PR	20	1644	+KRENEHL, HROBEK, KARSHOV, LAI + (BNL, CERN)
CLEGG	68	PREPRINT			A. B. CLEGG (LANC)
JESPERSEN	68	PL	21	1368	JESPERSEN, KANG, KERMAN+ (IOWA STATE)
LAMSA	68	PR	166	1395	+CASON, BISWAS, DERADO, GROVES, + (NOTRE DAME)
SHAPIRA	68	PL	21	1835	+HENARY, EISENBERG, RONAT, YAFFE + (REHO)
TAN	68	PL	288	195	TAN, PERL, MARTIN, CHINOWSKY + (SLAC+LRL+UCI)
WALKER	68	PL	20	133	+THOMPSON, ROBERTSON, HOWLIE, HARTUNG, + (WISC)
RHOFF	69	PR	187	1844	RHOFF, LEADOCK, KERMAN, JESPERSEN+ (LANC)
ANDERSON	70	PL	25	699	+BLESER, RLIFDIN, COLLINS+ (BNL, CERN)

END PRODUCTION EXPERIMENTS

Baryons

For notation, see illustrated key at beginning of data card listings.

N(1520) 62 N*1/2(1520, JP=3/2-) I=1/2 D13
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

Table with columns for mass (MEV), author, and resonance parameters for N(1520). Includes entries for ROPER, BRANDSFN, BAREYRE, etc.

Table with columns for width (MEV), author, and resonance parameters for N(1520). Includes entries for BAREYRE, DONNACHI, etc.

Table with columns for partial decay modes, author, and resonance parameters for N(1520). Includes entries for INTO PI N, INTO N PI, etc.

Table with columns for branching ratios, author, and resonance parameters for N(1520). Includes entries for INTO (PI N) TOTAL, INTO (PI) TOTAL, etc.

ALMOST THE ENTIRE INELASTICITY IS IN N PI PI (ONLY N ETA COULD COMPETE, AND IT DOESN'T). THE N PI PI SEEMS TO BE MAINLY N*3/2(1236) PI, IN BOTH S AND D WAVES.

Table with columns for partial decay modes, author, and resonance parameters for N(1520). Includes entries for INEL DECAY, INTO N PI, INTO N SIGMA, etc.

Table with columns for branching ratios, author, and resonance parameters for N(1520). Includes entries for INTO (PI N) TOTAL, INTO (PI) TOTAL, etc.

REFERENCES -- N*1/2(1520)

Table of references for N(1520) listing authors and publication details.

PAPERS NOT REFERRED TO IN DATA CARDS.
KIRZ 63 PR 130 24R1 J KIRZ, J SCHWARTZ, R D TRIPP (LRL)
BAREYRE 65 PL 18 342 + BRICMAN, STIRLING, VILLET (SACLAY) IJP
CROUCH 65 DESY CONF 11 21 + (BRUNNEN, CE, HARVARD, MIT, PADOVA, WPI, MANN)

N(1535) 63 N*1/2(1535, JP=1/2-) I=1/2 S11
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

Table with columns for mass (MEV), author, and resonance parameters for N(1535). Includes entries for HFNDRY, MICHAEL, UCHIYAMA, etc.

Table with columns for width (MEV), author, and resonance parameters for N(1535). Includes entries for HENDRY, MICHAEL, UCHIYAMA, etc.

Table with columns for partial decay modes, author, and resonance parameters for N(1535). Includes entries for INTO PI N, INTO N ETA, INTO N PI, etc.

Table with columns for branching ratios, author, and resonance parameters for N(1535). Includes entries for INTO (PI N) TOTAL, INTO (PI) TOTAL, etc.

Table with columns for partial decay modes, author, and resonance parameters for N(1535). Includes entries for INEL DECAY, INTO N PI, INTO N SIGMA, etc.

Table with columns for branching ratios, author, and resonance parameters for N(1535). Includes entries for INTO (PI N) TOTAL, INTO (PI) TOTAL, etc.

REFERENCES -- N*1/2(1535)

Table of references for N(1535) listing authors and publication details.

Baryons

For notation, see illustrated key at beginning of data card listings.

BARREYRE 68 PR 165 1731 P BARREYRE, C BRICMAN, G VILLET (SACLAY)IJP
DORNACHI 68 PL 268 161 A DONNACHIE, P G KIRSOPP, C LOVELACE (CERN)IJP
DONNACH2 68 VIENNA 139 DONNACHIE, RAPPOORTEUR,S TALK (GLAS)
KIRSOPP 68 THESIS R G KIRSOPP (EDIN)
DELCOURT 69 PL 298 75 DFLCOURT, LFFRANCOSI, PEREZ-Y-JORBA, + (ORSA)
DEANS 69 PR 185 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)

PAPERS NOT REFERRED TO IN DATA CARDS.

BRANDSEN 55 PR 139 81566 +DUNNELL, MOORHOUSE (DURHAM, RTHFD)IJP
-- BASIS OF NUMBERS WE QUOTE FROM HENDRY 65.
BARREYRE 65 PL 19 342 + BRICMAN, STIRLING, VILLET (SACLAY)IJP
LOVFLACE 67 HEIDELBERG C, 79 C LOVFLACE (CERN)IJP
JOHNSON 67 UCL-17683 THESIS C H JOHNSON (LRL)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYFD 70 PL 318 598 +BARREYRE+VILLET (SACLAY)

FOR PHOTONIC COUPLINGS SEE MOORHOUSE 70 NP 823,181 +RANKIN (GLAS)
WALKER 69 PR 182,1729 R.L.WALKER (CALT)

THE FOLLOWING ARTICLES DEAL WITH THE REACTIONS PI- P TO ETA N AND GAMMA P TO ETA P NEAR THRESHOLD. THE DATA AND THE THEORETICAL ARTICLES ARE USEFUL IN UNDERSTANDING THE BEHAVIOR OF THE S11 AMPLITUDE AS DETERMINED IN PI P PHASE-SHIFT ANALYSES. FURTHER REFERENCES MAY BE FOUND IN THEM.

MAINLY EXPERIMENTAL --
RULDS 64 PRL 13 486 + (BROWN, BRANDEIS, HARVARD, MIT, PADOVA) I
RICHARDS 65 PRL 16 1221 +CHIU, FANDI, HILMHOLTZ, KENNY, + (LRL, HAWAII) IJ
JONES 66 PL 23 597 +RININI, DUANE, HORSEY, MASON, + (IMPCCOL, RTHFD) IJP
BACCI 66 NC 454 983 +PENSO, SALVINI, MENGUCCINI, + (ROME, FRASCATI) IJP
PREPOST 67 PRL 18 32 R PREPOST, D LUNDOQUIST, D QUINN (STANFORD)
ALDYM 68 PRL 21 1130 +HUSG, PRESCOTT, ROCHESTER (CALTECH)
RULDS 69 PR 187 1827 +LANDI, BORONER, BASTIEN, + (U+HAD+MIT+PA)
HEUSCH 70 PRL 25,1381 +PRESCOTT, ROCHESTER, WINSTEIN (CALT)
MAINLY THEORETICAL --
MASON 66 PR 146 1022 P N DODSON (HAWAII)
MINAMI 66 PR 167 1123 S MINAMI (OSAKA)
BALL 66 PR 149 1191 J S BALL (UCLA)
LOGAN 67 PR 153 1634 R K LOGAN, F UCHIYAMA-CAMPRELL (ILL)
MENGUCCI 67 NC 484 579 C MENGUCCINI, A REALE (FRASCATI)
DEANS 67 PR 161 1465 S R DEANS, W G HULLADAY (VANDERBILT)
MINAMI 67 PR 162 1619 S MINAMI (OSAKA)
MOSS 67 PR 163 1785 T A MOSS (LSU)
DEANS 68 PR 165 1986 S R DEANS, W G HULLADAY (VANDERBILT)
PAL 68 PR 167 1950 R K PAL (NPL NEW DELHI)
BALL 69 PR 177 2257 +GARG, SHAW (UCLA+UCI)
LEFEBVRE 70 NC 66A 349 +LERUSTE (CDF)

N(1520) BUMPS

1520 MEV REGION - PRODUCTION EXPERIMENTS

Table with columns for N*(1520) PRODUCTION EXPERIMENTS, MASS (MEV), WIDTH (MEV), BRANCHING RATIOS, and INTRODUCTION EXPERIMENTS. Includes rows for R1, R2, R3, R4, R5, R6, R7, R8.

REFERENCES --N*(1520)-- PROD. EXP.
ALLES-BORRELLI, FRENCH, FRISK, MICHEJDA (CERN)
ALEXANDER, BERNARD, CZAPEK, + (WEIZMANN(CERN))
BASSOMPIERRE, + (CERN, BRUXELLES)
LEF 67 PR 159 1156 +MEERS, ROE, STINCLAIR, VANDER VELDE (MICH)
ANDERSON 70 PRL 25,699 +BLESEER, BLIEDEN, COLLINS, + (BNL, CERN)

FNO PRODUCTION EXPERIMENTS

N(1670) D15

64 N*1/2(1670, JP=5/2-) I=1/2
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1670).

Table for N*1/2(1670) MASS (MEV) with columns for mass, width, and branching ratios. Includes rows for M 1, M 2, M 3, M 4, M 5, M 6.

Table for N*1/2(1670) WIDTH (MEV) with columns for width, branching ratios, and decay masses. Includes rows for W 1, W 2, W 3, W 4, W 5, W 6.

Table for N*1/2(1670) PARTIAL DECAY MODES with columns for decay mode and branching ratios. Includes rows for P1, P2, P3, P4, P5.

Table for N*1/2(1670) BRANCHING RATIOS with columns for branching ratios and decay masses. Includes rows for R1, R2, R3, R4, R5, R6.

Table for N*1/2(1670) INTRODUCTION EXPERIMENTS with columns for introduction experiments and branching ratios. Includes rows for R1, R2, R3, R4, R5, R6, R7, R8.

SEE NOTE PRECEDING THE N*1/2(1670) INELASTIC DECAY MODE MEASUREMENTS.

REFERENCES -- N*1/2(1670)

BRANDSEN 65 PL 19 420 +DUNNELL, MOORHOUSE (DURHAM, RTHFD)IJP
TRIPP 67 NP 83 10 + LEITH, + (LRL, SLAC, CERN, HEIDEL, SACLAY)
BARREYRE 68 PR 165 1731 P BARREYRE, C BRICMAN, G VILLET (SACLAY)IJP
DONNACHI 68 PL 268 161 A DONNACHIE, P G KIRSOPP, C LOVELACE (CERN)IJP
DONNACH2 68 VIENNA 139 DONNACHIE, RAPPOORTEUR,S TALK (GLAS)
DUKE 68 PR 166 1444 +JONES, KEMP, MURPHY, THRESHER, + (RTHFD, OXF)IJP
KIRSOPP 68 THESIS R G KIRSOPP (EDIN)
RUSH 68 PR 173 1774 J E RUSH (UNIV ALABAMA)
BOTKE 69 PR 180 1417 J C BOTKE (UCSB)
DEANS 69 PR 185 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)
CARRERAS 70 NP 16A 35 B CARRERAS, A DONNACHIE (DARF, MCHS)
DAVIES 70 NP 821 359 A DAVIES (GLAS)
AYFD 70 KIEV CONF R AYFD, P BARREYRE, G VILLET (SACLAY)IJP
WAGNER 70 PREPRINT TH 1227 F WAGNER, C LOVFLACE (CFRN)
RRODY 71 SUBM. TO PRL +CASHMORE, +, HERNDON, +. (SLAC+LRL)

PAPERS NOT REFERRED TO IN DATA CARDS.
DUKE 65 PRL 15 468 +JONES, KEMP, MURPHY, PRENTICE, + (RTHFD, OXF)IJP
BARREYRE 65 PL 19 342 + BRICMAN, STIRLING, VILLET (SACLAY)IJP
JOHNSON 67 UCL-17683 THESIS C H JOHNSON (LRL)
DEANS 69 PRL 177 2623 S R DEANS (FLORIDA)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYFD 70 PL 318 598 +BARREYRE+VILLET (SACLAY)

FOR PHOTONIC COUPLINGS SEE WALKER 69 PR 182,1729 R.L.WALKER (CALT)

Baryons

For notation, see illustrated key at beginning of data card listings.

N(1688)

65 N*1/2(1688, JP=5/2+) I=1/2
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

F15

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for BRANSDEN, DUKE, BARREYRE, etc.

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for W 1, W 2, W 3, etc.

Table with columns for partial decay modes (P1, P2, P3, P4, P5) and decay masses.

Table with columns for branching ratios (R1, R2, R3, R4, R5, R6) and (P1)/TOTAL.

MORE INFORMATIONS ON THE INELASTIC DECAY MODES OF THE 1690 MEV RUMP, AS SEEN IN PRODUCTION EXPERIMENTS, MAY BE FOUND BELOW

Table with columns for branching ratios (R2, R3, R4, R5) and (P2)/TOTAL, (P3)/TOTAL, (P4)/TOTAL.

REFERENCES -- N*1/2(1688)

List of references for N*1/2(1688) including authors like BRANSDEN, HEUSCH, TRIPP, etc.

N(1700)

66 N*1/2(1700, JP=1/2-) I=1/2
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

S11

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for M 1, M 2, M 3, etc.

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for W 1, W 2, W 3, etc.

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for W 1, W 2, W 3, etc.

Table with columns for partial decay modes (P1, P2, P3, P4) and decay masses.

Table with columns for branching ratios (R1, R2, R3, R4, R5, R6) and (P1)/TOTAL.

Table with columns for branching ratios (R2, R3, R4) and (P3)/TOTAL.

Table with columns for branching ratios (R3, R4, R5, R6) and (P3)/TOTAL, (P2)/TOTAL, (P4)/TOTAL.

REFERENCES -- N*1/2(1700)

List of references for N*1/2(1700) including authors like BRANSDEN, MICHAEL, BARREYRE, etc.

N(1700)

66 N*1/2(1700, JP=3/2-) I=1/2
FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

D13

Table with columns for mass (MEV), author, and resonance parameters. Includes entries for M 3, M 4, M 5.

Baryons

For notation, see illustrated key at beginning of data card listings.

18 N*(1700) WIDTH (MEV)

19 N*(1700) PARTIAL DECAY MODES

PI N*(1700) INTO PI N

DECAY MASSES
137* 93*

REFERENCES -- N*(1700)

DONNACH 69 VIENNA 139 DONNACHIE RAPPOURTEUR,S TALK (GLAS)
KIRSOPP 68 THESIS R G KIRSOPP (EDIN)
WAGNER 70 PREPRINT TH 1227 F WAGNER, C LOVELACE (CERN)

N(1700) BUMPS

1700 MEV REGION - PRODUCTION EXPERIMENTS

20 N*(1700) PRODUCTION EXPERIMENTS

PARTIAL WAVE ANALYSIS REQUIRES AT LEAST FOUR l=1/2 STATES IN THE 1670 T L780 REGION (D15, F15, S11, P11) AND AT LEAST ONE l=3/2 STATE (D33), OBVIOUSLY DIFFERENT EXPERIMENTS ARE OFFERING DIFFERENT STATES AND OFTEN IT IS NOT CLEAR WHAT SPIN STATE IS BEING OBSERVED. NO EFFORT WAS MADE TO SEPARATE THESE EXPERIMENTS ACCORDING TO JP, SINCE NONE OF THE REPORTED JP IS FIRMLY ESTABLISHED. WE LIST ALL THE INFORMATION HERE, BUT WE HAVE NOT USED IT IN THE BARYON TABLE.

20 N*(1700) MASS (MEV)

M	(1695,0)	(7,0)	A-BORELLI 67 HRC	+ PRAR P 5,7 BEV/C	8/67
M	(1734,0)	(21,0)	ALMEIDA 68 HRC	+ PP 10 BEV/C	9/69
M	(1730,0)	(18,0)	GALLOWAY 68 HRC	PI-P A 5EV/C	8/69
M	100(1693,1)	(15,1)	RHODE 69 HBC	PP 22 GEV/C	10/69
M	JP IS PROBABLY 5/2+				
M	1 (1712,0)	(6,0)	BARNES 69 HRC	K-P TO K-P 2PI	7/70
M	1 [JP CONSISTENT WITH S11(1700) OR P11(1780) IN FORMATION]				
M	A (1667,0)	(5,0)	RENVENUTI 69 DBC	0 PI-D 2,26 GEV	5/770
M	A (1719,0)	(6,0)	WILLMANN 70 HBC	+ PI+P 13, GEV	5/770
M	A 1691,1	6,6	ANDERSON 70 HRC	- PI- P TO PI- HMS	2/71*
M	177 1710,	10,10	CIRBA 70 HRC	+ PI+ P TO P+3PI	2/71*
M	40 1763,	25,25	COOPER 70 HBC	+ LAMB, K PRD,	2/71*
M	505 1730,0	15,0	GREENELL 70 HBC	+ PI-P/PI+P 6 GEV	1/71*
M	60(1710,0)		KUZNETSOV 70 HRC	- LAMB, K PRD,	2/71*
M	AVG	1697,1	7,5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.1)	

20 N*(1700) WIDTH (MEV)

W	(70,0)	(20,0)	A-BORELLI 67 HRC		9/69
W	(140,0)	(57,0)	ALMEIDA 68 HRC		9/69
W	(155,0)	(15,0)	GALLOWAY 68 HRC		8/69
W	B 190 (235,0)	(50,0)	RHODE 69 HRC	PP 22 GEV/C	10/69
W	1 (70,0)	(15,0)	BARNES 69 HRC	K-P TO K-P 2PI	7/70
W	A (105,0)	(14,0)	RENVENUTI 69 DBC	0	5/770
W	A 130,	10,	ANDERSON 70 HRC	- PI- P TO PI- HMS	2/71*
W	177 66,	26,	CIRBA 70 HRC	+ PI+ P AT 5 GEV/C	2/71*
W	102,	40,	COOPER 70 HRC	+ PI+P, 5,5 GEV/C	2/71*
W	505 130,0	30,0	GREENELL 70 HBC	+ PI-P/PI+P 6 GEV	1/71*
W	40 (229,1)		KUZNETSOV 70 HRC	- PI-P, 4 GEV/C	2/71*
W	A (63,0)	(12,0)	WILLMANN 70 HRC		5/70
W	AVG	121,5	11,0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

20 N*(1700) BRANCHING RATIOS

R1	N*	(1700) INTO (PI N1)/(PI N1+PI N2)	PRD. EXP.			
R1 A	(0,77) OR LESS	LEE 67 HRC	+ PP 5,5 BEV/C	11/67		
	(9,0) OR MORE	RENVENUTI 69 DBC		5/770		
R2	N*	(1700) INTO (N ETA)/(N PI + N PI PI)	PRD. EXP.			
R2	(0,025) OR LESS	KRAEMER 64 DBC	+ PI+D 1,23 BEV/C	9/66		
R2	(0,042)OR LESS (95% CL)	A-BORELLI 67 HRC	+ PRAR P 5,7 BEV/C	9/69		
R3	N*	(1700) INTO (LAMBDA K1)/(P1+ PI-)	PRD. EXP.			
R3	(0,034) OR LESS	ALFAXANDER 67 HRC	+ PP 5,5 BEV/C	11/67		
R3	(0,07) OR LESS CL+.95	CIRBA 70 HRC	PI+P AT 5 GEV/C	2/71*		
R4	N*	(1700) INTO (LAMBDA K1)/(N PI + N PI PI)	PRD. EXP.			
R4	(0,013)OR LESS (95% CL)	A-BORELLI 67 HRC		8/67		
R4	SFN	CHINOWSKY 68 HRC	+ PP TO K-Y N	6/68		
R4 1	LIMITS 0,025 TO 0,11	BARNES 69 HRC	K-P TO K-P 2PI	7/70		
R4 A	LESS THAN 0,025	WILLMANN 70 HRC	PI+P TO 3PI P	6/70		
R4	25	0,025	0,005	GREENELL 70 HRC	+ PI+P TO 3PI P	1/71*
R5	N*	(1700) INTO (N PI)/(N PI PI)	PRD. EXP.			
R5	(1,26)OR LESS (95% CL)	A-BORELLI 67 HRC		8/67		
R5	0,025	0,13	GREENELL 70 HRC		1/71*	
R6	N*	(1700) INTO (N*3/2(1236)) PI1/(N PI PI)	PRD. EXP.			
R6	NO EVIDENCE	A-BORELLI 67 HRC		8/67		
R6	SEE MERLO 66 FOR A REVIEW.					
R7	N*	(1700) INTO (NEUTRON PI+1)/(PI+ PI-)	PRD. EXP.			
R7	0,67	0,40	ALEXANDER 67 HRC	+ PP 5,5 BEV/C	11/67	
R7	0,47	0,25	A-BORELLI HRC	PRAR P 5,5 GEV/C	7/70	
R7	AVG	0,53	0,21	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1,0)		
R8	N*	(1700) INTO (N*1236)++ PI-1)/(P1+ PI-)	PRD. EXP.			
R8	0,74	0,14	ALEXANDER 67 HRC	+ PP 5,5 BEV/C	11/67	
R8	(1,1,0)	0,3	ALMEIDA 68 HRC	+ PP 10 BEV/C	9/66	
RR	(0,83)		KAYAS 68 HRC	PP 8,1 REV/C	11/68	
RR 1	LESS THAN 0,15		BARNES 69 HRC	K-P TO K-P 2PI	7/70	
RR	(0,50) OR LESS	CL+.95	CIRBA 70 HRC	PI+P AT 5 GEV/C	2/71*	
RR	NO EVIDENCE		GREENELL 70 HRC		1/71*	
RR A	(2,3)	(1,6)	WILLMANN 70 HRC	PI+P TO 3PI P	6/70	
R9	N*(1700) INTO (SIG K1)/(LAMB K1)		PRD. EXP.			
R9	LESS THAN .20		COOPER 70 HRC	+ PI+P, 5,5 GEV/C	2/71*	

20 N*(1700) BRANCHING RATIOS (continued)

R10	N*	(1700) INTO (SIG K1)/(LAMB K1)	PRD. EXP.		
R10	LESS THAN .20		COOPER 70 HRC	+ PI+P, 5,5 GEV/C	2/71*

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

REFERENCES -- N* IN PRODUCT EXPERIMENTS

KRAEMER 64 PR 136 8496	+MADANSKY, + [J HOPKINS, W. WESTERN, WOODSTOCK] I
ALEXANDER 67 PR 154 1284	ALEXANDER, BENARY, CZAPEK, + (WEIZMANN(CERN))
A-BORELLI 67 NC 47 232	ALLES-BORELLI, FRENCH, FRISK, MICHEJDA (CERN)
LEE 67 PR 159 1156	+MOEBS, ROE, SINCLAIR, VANDER VELDE (MICH)
ALMEIDA 68 PR 174 1634	+RUSHROOKE, + [CAVNDISH, DESY(CERN)]
CHINOWSKY 68 PR 165 1466	CHINOWSKY, KINSEY, KLEIN, + (LRL, SLAC)
KAYAS 68 NP 85 169	+GIYADER, SENE, YIDOU, ALITTI, + (DARSAY, SAQIAY)
GALLOWAY 68 PL 270 250	GALLOWAY, ALYEA, CRITTENDEN, PRICKETT, + (IND)
BARNES 69 PRL 23 1516	+BASSANO+CHUNG+EISNER+FLAMINTO+KINSON (RNLIIJ)
RHODE 69 PR 187 1844	RHODE, LEACOCK, KERNAN, JESPERSEN, + (AMES)
RENVENUTI 69 PR 187 1852	RENVENUTI, MARQUIT, OPPENHEIMER (MINN, COLO)
ANDERSON 70 PRL 25,699	+BLESER, BLIEDEN, COLLINS + (BNL, CERN)
CIRBA 70 NP 823,533	+VANDERHAGEN + (EPOL,DURH,NIJH,TORI,PIUB)
COOPER 70 NP 823,605	+MANNER, MUGSGRAVE, POLLARO, VOYVODIC (IANL)
GREENELL 70 PRL 25 187	+LAI, LOUVEY, SCARRI, SINS (BNL)
KUZNETSOV 70 SNP 10,332	+MELNIKOV, RYLSOVA, CHADRAA, BALINTP (JINR)
WILLMANN 70 PRL 24 1260	+LAMS, GAIDOS, EZELL (PURDUIJ)

PAPERS NOT REFERRED TO IN DATA CARDS.

MERLO 66 P RYS SOC 289 489 J P MERLO, G VALLAOS (SACLAY)

N(1780)

14 N*(1780, JP=1/2+) l=1/2

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*(1740).

14 N*(1780) MASS (MEV)

M	3 (1751,0)	DONNACH 68 RVUE	PHASE-SHIFT ANAL	8/69
M	3 (1750,1)	DONNACH 68 RVUE	PHAS. SHIFT-CERNI	10/69
M	3 (190,1)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69
M	3 WHERE MAX. ABSORPTION IS -DONNACH, 2	KIRSOPP EYEBALL FIT CERN 1		10/69
M	(1640,0) (70,0)	ORITO 69 RVUE	K LAMBDA PS ANAL	9/69
M	4 (1770,0)	DAVIES 70 RVUE	P-5 ANAL SOL A	8/69
M	5 (1867,0)	DAVIES 70 RVUE	P-5 ANAL SOL B	8/69
M	6 SOL B IS E-D FIT TO SAME DATA START FROM CERN I EXPER. (DONNACH 68)			
M	6 (1645,0)	AYED 70 IPWA		1/71*
M	6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM			
M	A THERE ARE 3 SIMILAR SOLUTIONS	WAGNER 70 IPWA	PI-P TO K LAMB	1/71*

14 N*(1780) WIDTH (MEV)

W	3 (327,0)	DONNACH 68 RVUE		8/69
W	3 (327,1)	DONNACH 68 RVUE	PHAS. SHIFT-CERNI	10/69
W	3 (270,1)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69
W	4 (310,0) (50,0)	ORITO 69 RVUE	PHASE SHIFT ANAL	9/69
W	5 (525,0)	DAVIES 70 RVUE	SOL A	8/69
W	6 (50,0)	DAVIES 70 RVUE	SOL B	8/69
W	6 (160,0)OR(220,0)	AYED 70 IPWA		1/71*
W	A	WAGNER 70 IPWA	PI-P TO K LAMB	1/71*

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED. P EY

14 N*(1780) PARTIAL DECAY MODES

PI N*(1780) INTO PI N

139* 93*

P4 N*(1780) INTO N GAMMA

938* 0

14 N*(1780) BRANCHING RATIOS

R1	N*	(1780) INTO (PI N1)/TOTAL	(PI1)/TOTAL		
R1 3	(0,32)	DONNACH 68 RVUE		8/69	
R1 3	(1,32)	DONNACH 68 RVUE	PHAS. SHIFT-CERNI	10/69	
R1 3	(1,32)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69	
R1 4	(0,43)	DAVIES 70 RVUE	SOL A	8/69	
R1 5	(0,30)	DAVIES 70 RVUE	SOL B	8/69	
R1 6	(0,14-9)	AYED 70 IPWA		1/71*	
R2	N*(1780) INTO (LAMBDA K1)/(PI N1)/TOTAL**2		(P2*PI1)/TOTAL**2		
R2	0,004	0,003	ORITO 69 RVUE		8/69
R2 A	(0,025)OR 0,043	WAGNER 70 IPWA	PI-P TO K LAMB		1/71*
R3	N*(1780) INTO (LAMBDA K1)/TOTAL		(P21)/TOTAL		
R3	N*(1780) INTO (LAMBDA K1)/TOTAL		(P21)/TOTAL		
R3 B	(0,0031) TO 0,065	RUSH 68 MPWA	T POLE + RESON.		8/69
R3 B	PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING				
R4	N*(1780) INTO (N ETA)/TOTAL		(P31)/TOTAL		
R4 B	(0,19)	BOTKE 69 MPWA	T POLE + RESON.		10/69
R4 B	(0,015)OR 0,035	CARRERAS, 70 MPWA	T POLE + RESON.		5/70
R4 B	PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING				
R5	N*(1780) INTO (N GAMMA1)/TOTAL (PERCENT)		(P41)/TOTAL		
R5	(0,0096)	ORITO 69 CNTR	K-LAMB. PHOTOPRO.		

REFERENCES -- N*(1780)

DONNACH 68 PL 248 161	A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
DONNACH 68 VIENNA 139	DONNACHIE RAPPOURTEUR,S TALK (GLAS)
KIRSOPP 68 THESIS	R G KIRSOPP (EDIN)
RUSH 68 PR 173 1776	J F RUSH (UNIV ALABAMA)
HOTKE 69 PR 180 1417	J C HOTKE (UCS)
DEANS 69 PR 195 1787	S DEANS, J WOOTEN (UNIV S FLORIDA)
ORITO 69 LNC I 936	S ORITO,S SASAKI (TOKYO-OSAKA)
ORITO 69 INS J 113	S ORITO (THEIST)
CARRERAS 70 NP 168 35	B CARRERAS, A DONNACHIE (DAPE, MCHS)
DAVIES 70 NP 821 359	A DAVIES (GLAS)
AYED 70 KIEV CONF	R AYED, P BAREVRE, G VILLET (SACLAYIJP)
WAGNER 70 PREPRINT TH 1227	F WAGNER, C LOVELACE (CERN)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623	S R DEANS (FLORIDA)
DONNACH 68 NP 104 433	A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598	+BAREVRE+VILLET (SACLAY)

For notation, see illustrated key at beginning of data card listings.

N(1860)

15 N*1/2(1860, JP=3/2+) I=1/2 **P13**

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

15 N*1/2(1860) MASS (MEV)

M	3	(1860.0)	DDNNACH1	6R RVUE	PHASE-SHIFT ANAL	6/69
M	3	(1860.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
M	3	(1860.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
M	3	WHERE MAX. ABSORPTION IS -DDNNACH1, 2	KIRSOPP	EYEBALL FIT CERN 1		10/69
M	4	(1864.0)	DAVIES	70 RVUE	P-S ANAL SOL B	8/69
M	5	(1854.0)	DAVIES	70 RVUE	P-S ANAL SOL B	8/69
M	X	SOL R IS EAD FIT TO SAME DATA START FROM CERN 1 EXPR. (DDNNACH1 6R)				
M	X	(1860.0)	APPROX	LEA	69 CNTR	PI-P ELASTIC 8/69
M	X	SEE ALSO APLIN 70				
M	6	(1786.0)	AYED	70 IPWA		1/71*
M	6	FROM ENER. DEP. FIT OF ARGAND DIAGRAM	WAGNER	70 IPWA	PI-P TO K LAMB	1/71*
M	A	(1800.0)				
M	A	P13 RESONATES ONLY IN ONE OUT OF 3 POSSIBLE SOLUTIONS				

15 N*1/2(1860) WIDTH (MEV)

W	3	(296.00)	DDNNACH1	6R RVUE		8/69
W	3	(296.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
W	3	(296.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
W	4	(449.0)	DAVIES	70 RVUE	SOL A	8/69
W	5	(307.0)	DAVIES	70 RVUE	SOL B	8/69
W	6	(182.0)	AYED	70 IPWA		1/71*
W	A	(220.0)	WAGNER	70 IPWA	PI-P TO K LAMB	1/71*

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

15 N*1/2(1860) PARTIAL DECAY MODES

DECAY MASSES

P1 N*1/2(1860) INTO PI N 139+ 938

P2 N*1/2(1860) INTO LAMBDA K 1115+ 497

P3 N*1/2(1860) INTO N ETA 939+ 548

P4 N*1/2(1860) INTO N PI PI 939+ 139+ 139

15 N*1/2(1860) BRANCHING RATIOS

R1 N*1/2(1860) INTO (PI N)/TOTAL (P1)/TOTAL 8/69

R1 3 (0.21) DDNNACH1 6R RVUE T POLE + RESON. 10/69

R1 3 (2.21) DDNNACH2 6R RVUE PHAS-SHIFT-CERN1 10/69

R1 3 (.25) KIRSOPP 6R RVUE PHASE SHIFT ANAL 10/69

R1 4 (0.40) DAVIES 70 RVUE SOL A 8/69

R1 5 (0.26) DAVIES 70 RVUE SOL B 8/69

R1 6 (0.149) AYED 70 IPWA 1/71*

R2 N*1/2(1860) INTO (LAMBDA K)/TOTAL (P2)/TOTAL 8/69

R2 8 (0.01417) 0.16 T POLE + RESON.

R2 8 PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

R3 N*1/2(1860) INTO (N ETA)/TOTAL (P3)/TOTAL 10/69

R3 8 (0.0364) ROTKE 69 MPWA T POLE + RESON. 5/70

R3 8 (0.0031) (0.003) DEANS 69 MPWA T POLE + RESON. 5/70

R3 8 (0.0301) 0.094 CARRERAS 70 MPWA T POLE + RESON. 5/70

R3 8 PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

R4 N*1/2(1860) INTO (LAMBDA K)+PI N)/TOTAL*2 (P2*P1)/TOTAL*2 1/71*

R4 A (0.015) WAGNER 70 IPWA PI-P TO K LAMB

REFERENCES -- N*1/2(1860)

DDNNACH1 6R PL 269 161 A DDNNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP

DDNNACH2 6R VIENNA 139 DDNNACHIF RAPPORTEUR-S TALK (GLAS)

KIRSOPP 6R THESIS R G KIRSOPP (EDIN)

RUSH 6R PR 173 1776 J E RUSH (UNIV ALABAMA)

ROTKE 69 PR 180 1417 J C ROTKE (UCSB)

DEANS 49 PR 195 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)

LEA 69 PL 298 584 LEA, DADES, WARD, COWAN, + (RHEL, BRISTOL, DAREF)

APLIN 70 RPPH/67 *COWAN, GIBSON, GILMORE+ (RHEL, BRISTOL, DAREF)

AYED 70 KIEV CONF R AYEY, P BAREYRE, G VILLET (SACLAY)IJP

CARRERAS 70 NP 16R 35 B CARRERAS, A DDNNACHIE (DARE, MCHS)

DAVIES 70 NP B21 35 A DAVIES (GLAS)

WAGNER 70 PREPRINT TH 1227 F WAGNER, C LOVELACE (CERN)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623 S R DEANS (FLORIDA)

DDNNACH1 69 NP 10B 433 A DDNNACHIE, R KIRSOPP (GLAS+EDIN)

AYED 70 PL 31B 598 *BAREYRE, VILLET (SACLAY)

N(1990)

17 N*1/2(1990, JP=7/2+) I=1/2 **F17**

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

17 N*1/2(1990) MASS (MEV)

M	3	(1983.0)	DDNNACH1	6R RVUE	PHASE-SHIFT ANAL	10/69
M	3	(1995.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
M	3	WHERE MAX. ABSORPTION IS -DDNNACH1, 2	KIRSOPP	EYEBALL FIT CERN 1		10/69
M	X	(2000.0)	APPROX	LEA	69 CNTR	PI-P ELASTIC 8/69
M	X	SEE ALSO APLIN 70				

17 N*1/2(1990) WIDTH (MEV)

W	3	(225.0)	DDNNACH1	6R RVUE		8/69
W	3	(250.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69

17 N*1/2(1990) PARTIAL DECAY MODES

DECAY MASSES

P1 N*1/2(1990) INTO PI N 139+ 938

P2 N*1/2(1990) INTO N PI PI 939+ 139+ 139

P3 N*1/2(1990) INTO N ETA 939+ 548

REFERENCES -- N*1/2(1990)

DDNNACH1 6R PL 269 161 A DDNNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP

DDNNACH2 6R VIENNA 139 DDNNACHIF RAPPORTEUR-S TALK (GLAS)

KIRSOPP 6R THESIS R G KIRSOPP (EDIN)

RUSH 6R PR 173 1776 J E RUSH (UNIV ALABAMA)

ROTKE 69 PR 180 1417 J C ROTKE (UCSB)

DEANS 49 PR 195 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)

LEA 69 PL 298 584 LEA, DADES, WARD, COWAN, + (RHEL, BRISTOL, DAREF)

APLIN 70 RPPH/67 *COWAN, GIBSON, GILMORE+ (RHEL, BRISTOL, DAREF)

AYED 70 KIEV CONF R AYEY, P BAREYRE, G VILLET (SACLAY)IJP

CARRERAS 70 NP 16R 35 B CARRERAS, A DDNNACHIE (DARE, MCHS)

DAVIES 70 NP B21 35 A DAVIES (GLAS)

WAGNER 70 PREPRINT TH 1227 F WAGNER, C LOVELACE (CERN)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623 S R DEANS (FLORIDA)

DDNNACH1 69 NP 10B 433 A DDNNACHIE, R KIRSOPP (GLAS+EDIN)

AYED 70 PL 31B 598 *BAREYRE, VILLET (SACLAY)

17 N*1/2(1990) BRANCHING RATIOS

R1 N*1/2(1990) INTO (PI N)/TOTAL (P1)/TOTAL 10/69

R1 3 (0.09) KIRSOPP 6R RVUE PHASE SHIFT ANAL 10/69

R2 N*1/2(1990) INTO (N ETA)/TOTAL (P3)/TOTAL 5/70

R2 8 (0.02) (0.02) DEANS 69 MPWA T POLE + RESON.

R2 8 PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

REFERENCES -- N*1/2(1990)

DDNNACH1 6R PL 269 161 A DDNNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP

KIRSOPP 6R THESIS R G KIRSOPP (EDIN)

LEA 69 PL 298 584 LEA, DADES, WARD, COWAN, + (RHEL, BRISTOL, DAREF)

DEANS 69 PR 195 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)

APLIN 70 RPPH/67 *COWAN, GIBSON, GILMORE+ (RHEL, BRISTOL)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623 S R DEANS (FLORIDA)

AYED 70 PL 31B 598 *BAREYRE, VILLET (SACLAY)

N(2040)

16 N*1/2(2040, JP=3/2-) I=1/2 **D13**

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

16 N*1/2(2040) MASS (MEV)

M	3	(2057.0)	DDNNACH1	6R RVUE	PHASE-SHIFT ANAL	6/69
M	3	(2030.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
M	3	(2040.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
M	3	WHERE MAX. ABSORPTION IS -DDNNACH1, 2	KIRSOPP	EYEBALL FIT CERN 1		10/69
M	X	(2030.0)	APPROX	LEA	69 CNTR	PI-P ELASTIC 8/69
M	X	SEE ALSO APLIN 70				

16 N*1/2(2040) WIDTH (MEV)

W	3	(293.0)	DDNNACH1	6R RVUE		8/69
W	3	(290.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
W	3	(240.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

16 N*1/2(2040) PARTIAL DECAY MODES

DECAY MASSES

P1 N*1/2(2040) INTO PI N 139+ 938

P2 N*1/2(2040) INTO N PI PI 938+ 139+ 139

P3 N*1/2(2040) INTO N ETA 939+ 548

16 N*1/2(2040) BRANCHING RATIOS

R1 N*1/2(2040) INTO (PI N)/TOTAL (P1)/TOTAL 10/69

R1 3 (.26) DDNNACH1 6R RVUE T POLE + RESON. 10/69

R1 3 (1.15) DDNNACH2 6R RVUE PHAS-SHIFT-CERN1 10/69

R1 3 KIRSOPP 6R RVUE PHASE SHIFT ANAL 10/69

R2 N*1/2(2040) INTO (N ETA)/TOTAL (P3)/TOTAL 5/70

R2 8 (0.009) 0.009 CARRERAS 70 MPWA T POLE + RESON.

R2 8 PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

REFERENCES -- N*1/2(2040)

DDNNACH1 6R PL 269 161 A DDNNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP

DDNNACH2 6R VIENNA 139 DDNNACHIF RAPPORTEUR-S TALK (GLAS)

KIRSOPP 6R THESIS R G KIRSOPP (EDIN)

LEA 69 PL 298 584 LEA, DADES, WARD, COWAN, + (RHEL, BRISTOL, DAREF)

APLIN 70 RPPH/67 *COWAN, GIBSON, GILMORE+ (RHEL, BRISTOL, DAREF)

CARRERAS 70 NP 16R 35 B CARRERAS, A DDNNACHIE (DARE, MCHS)

PAPERS NOT REFERRED TO IN DATA CARDS.

DDNNACH1 69 NP 10B 433 A DDNNACHIE, R KIRSOPP (GLAS+EDIN)

AYED 70 PL 31B 598 *BAREYRE, VILLET (SACLAY)

N(2190)

71 N*1/2(2190, JP=7/2-) I=1/2 **G17**

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

71 N*1/2(2190) MASS (MEV)

M		(2190.0)	DIDDENS	63 CNTR	PI+ P TOTAL	
M		(2210.0)	HUEBLER	64 RVUE	DATA + DISP REL	
M		(2210.0)	APPROX	YOKOSAWA	64 CNTR	PI-P DISC + PDL 7/66
M	3	(2265.0)	DDNNACH1	6R RVUE	PHASE-SHIFT ANAL	6/69
M	3	(2190.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
M	3	(2265.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
M	3	WHERE MAX. ABSORPTION IS -DDNNACH1, 2	KIRSOPP	EYEBALL FIT CERN 1		10/69
M	6	(2000.0)	APPROX	LEA	69 CNTR	PI-P ELASTIC 8/69
M	6	(2158.0)	AYED	70 IPWA		1/71*
M	6	FROM ENER. DEP. FIT OF ARGAND DIAGRAM	ANDERSON	70 MMS	PI-P TO PI- MMS	2/71*
M	6	25.	HULL	70 MPWA	SMALL ANGLE PI-P	1/71*

71 N*1/2(2190) WIDTH (MEV)

W		(200.0)	DIDDENS	63 CNTR		
W		(200.0)	HUEBLER	64 RVUE		7/66
W		(220.0)	APPROX	YOKOSAWA	64 CNTR	7/66
W	3	(299.0)	DDNNACH1	6R RVUE		6/69
W	3	(300.)	DDNNACH2	6R RVUE	PHAS-SHIFT-CERN1	10/69
W	3	(300.)	KIRSOPP	6R RVUE	PHASE SHIFT ANAL	10/69
W	3	275.	ANDERSON	70 MMS	PI-P TO PI- MMS	2/71*
W	6	(325.0)	AYED	70 IPWA		1/71*
W	6	(239.0)	HULL	70 MPWA	SMALL ANGLE PI-P	1/71*

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

REFERENCES -- N*1/2(2190)

DIDDENS 63 PR 10B 433 DIDDENS (GLAS)

HUEBLER 64 PR 10B 433 HUEBLER (GLAS)

YOKOSAWA 64 PR 10B 433 YOKOSAWA (GLAS)

DDNNACH1 6R PL 269 161 A DDNNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP

DDNNACH2 6R VIENNA 139 DDNNACHIF RAPPORTEUR-S TALK (GLAS)

KIRSOPP 6R THESIS R G KIRSOPP (EDIN)

LEA 69 PL 298 584 LEA, DADES, WARD, COWAN, + (RHEL, BRISTOL, DAREF)

APLIN 70 RPPH/67 *COWAN, GIBSON, GILMORE+ (RHEL, BRISTOL, DAREF)

ANDERSON 70 NP 16R 35 B ANDERSON (DARE, MCHS)

HULL 70 PR 10B 433 HULL (GLAS)

PAPERS NOT REFERRED TO IN DATA CARDS.

Baryons

For notation, see illustrated key at beginning of data card listings.

----- 71 N*1/2(2190) PARTIAL DECAY MODES -----
 P1 N*1/2(2190) INTO PI N 139+ 938
 P2 N*1/2(2190) INTO LAMBDA K 1115+ 497
 P3 N*1/2(2190) INTO N PI PI 938+ 139+ 139

----- 71 N*1/2(2190) BRANCHING RATIOS -----
 R1 N*1/2(2190) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 (0.3) APPROX DIDDS 63 CNTR 7/66
 R1 (0.3) APPROX YOKOSAWA 66 CNTR 7/66
 R1 3 (0.349) DONNACHI 6R RVUE 6/68
 R1 3 (.35) DONNACH2 6R RVUE PHAS. SHIFT-CERN1 10/69
 R1 3 (.35) KIRSOPP 6R RVUE PHASE SHIFT ANAL 10/69
 R1 6 (0.150) AYEY 70 IPWA 1/71*
 R1 (0.09) HULL 70 MPWA SMALL ANGLE PI-P 1/71*

 REFERENCES -- N*1/2(2190)
 +JENKINS, KYCIA, RILFY (BNL) I
 G HOHLER, J GIESECKE (KARLSRUHE) I
 +SUWA, HILL, ESTERLING, BOOTH (ARG,CHI) JP
 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN) IJP
 DONNACHIE, RAPPORTEUR, S TALK (GLAS) (EDIN)
 R G KIRSOPP (EDIN)
 LEA, DODES, WARD, GOWAN, + (RHUL, ARISTOL, OARF)
 +RLESFR, BLIEDEN, COLLINS+ (BNL, CERN)
 R AYEY, P BAREYRE, G VILLET (SACLAY) IJP
 HULL, R LEACOCK (ISU)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN DATA CARDS.
 CARROLL 66 PRL 16 289 +CORBETT, DAMERELL, MIDDLEMAS, + (RTHFD, OXF) J-L
 CARROLL 66 PRL 17 1274 +CORBETT, DAMERELL, MIDDLEMAS, + (RTHFD, OXF) J-L
 --- ERBATION CHANGING THE RATHER WEAK DETERMINATION OF J-L TO +1/2.
 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 KIRBY 58 PRL 16 709 W BARBER, D CLINE (WISC) P
 RUSZA 66 PRL 16 913 +DAVIS, DUFF, HEYMANN, + (UNICOL, WESTFELD)
 RUSZA 67 NC 52A 331

PAPERS NOT REFERRED TO IN DATA CARDS.
 AYEY 70 PL 31A 598 +BAREYRE, VILLET (SACLAY)

 REFERENCES -- N*1/2(2190)
 +CORBETT, DAMERELL, MIDDLEMAS, + (RTHFD, OXF) J-L
 +CORBETT, DAMERELL, MIDDLEMAS, + (RTHFD, OXF) J-L
 --- ERBATION CHANGING THE RATHER WEAK DETERMINATION OF J-L TO +1/2.
 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 W BARBER, D CLINE (WISC) P
 +DAVIS, DUFF, HEYMANN, + (UNICOL, WESTFELD)

 REFERENCES -- N*1/2(2190)
 +RLESFR, BLIEDEN, COLLINS+ (BNL, CERN)
 R AYEY, P BAREYRE, G VILLET (SACLAY) IJP
 HULL, R LEACOCK (ISU)

----- 90 N*1/2(2220) MASS (MEV) -----
 M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM (2200.0)
 M 6 (2221.0) APPROX. BUSZA 67 ODPK LEG. POLYN. ANAL. 2/71*
 M 6 (2245.0) HULL 70 IPWA SMALL ANGLE PI-P 1/71*
 M 6 (2245.0) HULL 70 MPWA SMALL ANGLE PI-P 1/71*

----- 90 N*1/2(2220) WIDTH (MEV) -----
 W 6 (258.0) AYEY 70 IPWA 1/71*
 W 6 (329.0) HULL 70 MPWA SMALL ANGLE PI-P 1/71*

----- 90 N*1/2(2220) PARTIAL DECAY MODES -----
 P1 N*1/2(2220) INTO PI N 139+ 938
 P2 N*1/2(2220) INTO N PI PI 938+ 139+ 139

----- 90 N*1/2(2220) BRANCHING RATIOS -----
 R1 N*1/2(2220) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 6 (0.140) AYEY 70 IPWA 1/71*
 R1 (0.15) HULL 70 MPWA SMALL ANGLE PI-P 1/71*

 REFERENCES -- N*1/2(2220)
 R AYEY, P BAREYRE, G VILLET (SACLAY) IJP
 +DAVIS, DUFF, HEYMANN, NIMMON, + (UCL+MCL)
 J HULL, R LEACOCK (ISU)

PAPERS NOT REFERRED TO IN DATA CARDS.
 AYEY 70 PL 31B 598 +BAREYRE, VILLET (SACLAY)

 REFERENCES -- N*1/2(2220)
 R AYEY, P BAREYRE, G VILLET (SACLAY) IJP
 +DAVIS, DUFF, HEYMANN, NIMMON, + (UCL+MCL)
 J HULL, R LEACOCK (ISU)

PAPERS NOT REFERRED TO IN DATA CARDS.
 AYEY 70 PL 31B 598 +BAREYRE, VILLET (SACLAY)

M > 2200 MEV - PRODUCTION AND σ_{TOTAL} EXPERIMENTS

BUMPS

----- 72 N*1/2(2650, JP= -) I=1/2 -----
 72 N*1/2(2650) MASS (MEV)
 M (2700.0)
 M (2600.0)
 M (2660.0)
 M 2649.0 10.0 ALVAREZ 64 CNTR PI PHOTOPRD 7/66
 M (2633.0) APPROX WAHLIG 64 ODPK 0 PI-P CH EX 11/67
 HOHLER 64 RVUE DATA + DISP REL
 CITRON 66 CNTR PI-P TOTAL + CH EX
 BARGER 66 FIT

----- 72 N*1/2(2650) WIDTH (MEV) -----
 W (100.0) ALVAREZ 64 CNTR 7/66
 W (200.0) HOHLER 64 RVUE 7/66
 W 360.0 20.0 CITRON 66 CNTR 7/66
 W (425.0) BARGER 66 FIT TOTAL + CH EX 11/67

----- 72 N*1/2(2650) PARTIAL DECAY MODES -----
 P1 N*1/2(2650) INTO PI N 139+ 938
 P2 N*1/2(2650) INTO LAMBDA K 1115+ 497
 P3 N*1/2(2650) INTO N PI PI 938+ 139+ 139

----- 72 N*1/2(2650) BRANCHING RATIOS -----
 R1 N*1/2(2650) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 ONLY (J+1/2)* (PI N)/TOTAL MEASURED FOR THIS STATE
 R1 (0.436) 0.028 CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
 R1 9 (0.456) (0.019) BARGER 66 RVUE TOTAL + CH EX. 11/67
 R1 B (0.30) BARGER 67 CNTR USES KORMANYOS66 11/67
 B USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
 B FOR CRITICISM OF THIS METHOD, SEE DDLEN 68.
 R1 D (0.24) DIKMEN 67 RVUE USES KORMANYOS66 11/67
 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
 KORMANYOS 67 CNTR PI-P AT 180 DEGS. 11/67

 REFERENCES -- N*1/2(2650)
 +BAR-YAM, KERN, LUCKEY, DORRNE, + (MIT, CEA)
 +MANNFELL, SODICKSON, FAGLER, WARD, + (MIT)
 G HOHLER, J GIESECKE (KARLSRUHE) I
 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 V BARGER, M DLSSON (WISC)
 V BARGER, D CLINE (WISC) P
 F N DIKMEN (MICH)
 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 R DDLEN, D HORN, C SCHMID (CAL TECH)

PAPERS NOT REFERRED TO IN DATA CARDS.
 BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE, ORSAY) J-L
 WAHLIG 68 PR 168 1915 M A WAHLIG, I MANNELLI (MIT, PISA)
 --- FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH
 CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
 COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

----- 72 N*1/2(2650) PARTIAL DECAY MODES -----
 P1 N*1/2(2650) INTO PI N 139+ 938
 P2 N*1/2(2650) INTO LAMBDA K 1115+ 497
 P3 N*1/2(2650) INTO N PI PI 938+ 139+ 139

----- 72 N*1/2(2650) BRANCHING RATIOS -----
 R1 N*1/2(2650) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 ONLY (J+1/2)* (PI N)/TOTAL MEASURED FOR THIS STATE
 R1 (0.436) 0.028 CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
 R1 9 (0.456) (0.019) BARGER 66 RVUE TOTAL + CH EX. 11/67
 R1 B (0.30) BARGER 67 CNTR USES KORMANYOS66 11/67
 B USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
 B FOR CRITICISM OF THIS METHOD, SEE DDLEN 68.
 R1 D (0.24) DIKMEN 67 RVUE USES KORMANYOS66 11/67
 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
 KORMANYOS 67 CNTR PI-P AT 180 DEGS. 11/67

 REFERENCES -- N*1/2(2650)
 +BAR-YAM, KERN, LUCKEY, DORRNE, + (MIT, CEA)
 +MANNFELL, SODICKSON, FAGLER, WARD, + (MIT)
 G HOHLER, J GIESECKE (KARLSRUHE) I
 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 V BARGER, M DLSSON (WISC)
 V BARGER, D CLINE (WISC) P
 F N DIKMEN (MICH)
 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 R DDLEN, D HORN, C SCHMID (CAL TECH)

PAPERS NOT REFERRED TO IN DATA CARDS.
 BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE, ORSAY) J-L
 WAHLIG 68 PR 168 1915 M A WAHLIG, I MANNELLI (MIT, PISA)
 --- FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH
 CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
 COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

 REFERENCES -- N*1/2(2650)
 +BAR-YAM, KERN, LUCKEY, DORRNE, + (MIT, CEA)
 +MANNFELL, SODICKSON, FAGLER, WARD, + (MIT)
 G HOHLER, J GIESECKE (KARLSRUHE) I
 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 V BARGER, M DLSSON (WISC)
 V BARGER, D CLINE (WISC) P
 F N DIKMEN (MICH)
 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 R DDLEN, D HORN, C SCHMID (CAL TECH)

PAPERS NOT REFERRED TO IN DATA CARDS.
 BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE, ORSAY) J-L
 WAHLIG 68 PR 168 1915 M A WAHLIG, I MANNELLI (MIT, PISA)
 --- FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH
 CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
 COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

----- N(3030) 73 N*1/2(3030, JP=) I=1/2 -----
 73 N*1/2(3030) MASS (MEV)
 M (3080.0) HOHLER 64 RVUE 64 RVUE DATA + DISP REL 7/66
 M (3030.0) CITRON 66 CNTR PI-P TOTAL 7/66

----- 73 N*1/2(3030) WIDTH (MEV) -----
 W (400.0) CITRON 66 CNTR 7/66

----- 73 N*1/2(3030) PARTIAL DECAY MODES -----
 P1 N*1/2(3030) INTO PI N 139+ 938
 P2 N*1/2(3030) INTO N PI PI 938+ 139+ 139

----- 73 N*1/2(3030) BRANCHING RATIOS -----
 R1 N*1/2(3030) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 ONLY (J+1/2)* (PI N)/TOTAL MEASURED FOR THIS STATE
 R1 (0.66) CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
 R1 B (0.098) (0.016) BARGER 66 RVUE TOTAL + CH EX. 11/67
 R1 B (0.12) BARGER 67 CNTR USES KORMANYOS66 11/67
 B USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
 B FOR CRITICISM OF THIS METHOD, SEE DDLEN 68.
 R1 D (0.015) DIKMEN 67 RVUE USES KORMANYOS67 11/67
 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

 REFERENCES -- N*1/2(3030)
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 BARGER 66 PR 151 1123 V BARGER, M DLSSON (WISC)
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 DDLEN 68 PR 166 1768 R DDLEN, D HORN, C SCHMID (CAL TECH)

 REFERENCES -- N*1/2(3030)
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 BARGER 66 PR 151 1123 V BARGER, M DLSSON (WISC)
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 DDLEN 68 PR 166 1768 R DDLEN, D HORN, C SCHMID (CAL TECH)

PAPERS NOT REFERRED TO IN DATA CARDS.
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 BARGER 66 PR 151 1123 V BARGER, M DLSSON (WISC)
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 DDLEN 68 PR 166 1768 R DDLEN, D HORN, C SCHMID (CAL TECH)

 REFERENCES -- N*1/2(3030)
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONITIC, PHILLIPS, + (BNL) I
 BARGER 66 PR 151 1123 V BARGER, M DLSSON (WISC)
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P
 DDLEN 68 PR 166 1768 R DDLEN, D HORN, C SCHMID (CAL TECH)

----- N(3245) 74 N(3245, JP= +) -----
 74 N(3245) MASS (MEV)
 M 3245.0 10.0 KORMANYOS 67 CNTR PI-P 180 DEG FL 6/68
 W 86 (35.0) OR LESS KORMANYOS 67 CNTR

----- 74 N(3245) WIDTH (MEV) -----
 W 86 (35.0) OR LESS KORMANYOS 67 CNTR

----- 74 N(3245) PARTIAL DECAY MODES -----
 P1 N(3245) INTO PI N 139+ 938
 P2 N(3245) INTO LAMBDA K 1115+ 497
 P3 N(3245) INTO N PI PI 938+ 139+ 139

----- 74 N(3245) BRANCHING RATIOS -----
 R1 N(3245) INTO (PI N)/TOTAL (PI)/TOTAL
 R1 J IS NOT KNOWN. FOLLOWING IS (J+1/2)* (PI N)/TOTAL
 KORMANYOS 67 CNTR 6/68

 REFERENCES -- N(3245)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

 REFERENCES -- N(3245)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

 REFERENCES -- N(3245)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

 REFERENCES -- N(3245)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

 REFERENCES -- N(3245)
 KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

Baryons

For notation, see illustrated key at beginning of data card listings.

N(3690) 75 N*1/2(3690, JP=) I=1/2
 A BUMP SEEN IN THE INVARIANT MASS OF A VERY COMPLICATED STATE IN SEVEN PIS, SO AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. NOT INCLUDED IN TABLE.

 75 N*1/2(3690) MASS (MEV) -----
 M 3690.0 10.0 BARTKE 67 HRC + PI+P R PRONGS 8/67

 75 N*1/2(3690) WIDTH (MEV) -----
 W 50.0 30.0 BARTKE 67 HRC + 8/67

 75 N*1/2(3690) PARTIAL DECAY MODES -----
 PL N*1/2(3690) INTO N + 7 PIS 0 DECAY MASSES

 REFERENCES -- N*1/2(3690)
 BARTKE 67 PL 248 118 +CZYZEWSKI, DANYSZ, + (CRACOW, ORSAY) I

N₇(3755) 76 N* /2(3755, JP=) I
 A SMALL PEAK IN THE (P P PBAR) INVARIANT MASS FROM 8.4 MEV/C PI+ P TO PI+ P P PBAR EVENTS. AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. OMITTED FROM TABLE.

 76 N* /2(3755) MASS (MEV) -----
 M 3755.0 8.0 EHRLICH 68 HBC + PI+ P P PBAR 6/68

 76 N* /2(3755) WIDTH (MEV) -----
 W 40.0 20.0 EHRLICH 68 HBC + 6/68

 76 N* /2(3755) PARTIAL DECAY MODES -----
 PL N* /2(3755) INTO PI+ P P PBAR 0 DECAY MASSES

 REFERENCES -- N* /2(3755)
 EHRLICH 68 PRL 20 696 R EHRLICH, R J PLANO, J B WHITTAKER (RUTGERS)

 END PRODUCTION EXPERIMENTS

Δ(1236) 33 N*3/2(1236, JP=3/2+) I=3/2 **P₃₃**
 CARTER 70 REPORT NEW PRECISE CROSS SECTION MEASUREMENTS FOR PI+P, PI-P AND CHARGE EXCHANGE. THEIR ANALYSIS COMBINES TOTAL CROSS SECTION DATA WITH THE PHASE SHIFTS OF DONNACHIE 68 (USED FOR THE BACKGROUND UNDER THE P33) THE CHARGE EXCHANGE DATA WERE NOT USED.
 DONNACHIE 68 HAS DONE A SIMILAR ANALYSIS ON OLDER DATA, USING ROPER 65 PHASE SHIFTS WITH A FREE OVERALL NORMALIZATION.

Comments on the Mass and Width of Δ(1236)

In a paper submitted to the Kiev Conference, CARTER 70 present new data on the π[±]p total cross sections and the charge exchange cross sections in the region of Δ(1236). The same group also has data on the differential cross sections, but they have not yet been fully analyzed. The data of CARTER 70 seem both statistically and systematically much better than the older results (which date back to the early 50's). When compared with the previous measurements, the most interesting feature of the new data is an apparent small shift of the peak values of the cross sections to lower momenta.

CARTER 70 then calculate the δ₃₃ phase shift from the total π[±]p cross-section data. This they do by using the phase-shift solutions of DONNACHIE 68 for the S₃₁ and P₃₁ partial waves, which they subtract to obtain the δ₃₃ phases. The errors assigned to the phases do not include uncertainties in the background subtraction.

The δ₃₃ phase shifts were next fit by CARTER 70 with the so-called Layson form of the resonance line shape:

$$\tan \delta_R = \frac{\Gamma(E_\pi)}{2(E_{\pi 0} - E_\pi)}, \quad \Gamma(E_\pi) = \frac{4m_p(qR)^3 \gamma^2}{(E_\pi + E_{\pi 0})[1 + (qR)^2]}$$

where E_π is the pion c.m. energy and an additional subscript 0 denotes the resonant value. The mass M at resonance, reduced width γ², and range R are determined by a fit to the data. The width at resonance is given by Γ(E_{π0})M/(M-E_{π0}). The results are shown in the accompanying table (the values quoted in the

Comparison of analyses of the δ₃₃ phase shifts which CARTER 70 obtain from their π[±]p total cross section data. a₃₃ is the background scattering length; a dash indicates no background was permitted.

	"Breit-Wigner"	a ₃₃ (1/m _π ³)	M (MeV)	Γ ₀ (MeV)	R (fermi)	χ ² (14 points)	N _{param.}
CARTER 70	Layson	-	1231±0.6	111±3	?	?	?
Our studies	Standard	-	1231	109	1.09	48	3
	Layson	-	1231	110	0.92	35	3
	Standard	0.0100	1234	120	0.89	9.7	4
	Layson	0.0099	1234	120	0.75	10	4
	Standard with two radii	-	1231.4	112	0.83, 0.62	12.8	4
	3-parameter polynomial	-	1230.8	110	-	35	3
	4-parameter polynomial	-	1231.4	112	-	12.3	4

Baryons

For notation, see illustrated key at beginning of data card listings.

listing are slightly different because the additional Coulomb correction has been applied to obtain the final values).

The results of this fit are considerably lower than the values we had quoted earlier. These were the results of OLSON 65 who did the same type of analysis on the old data, using the ROPER 65 phase shift for S_{31} and P_{31} with an additional overall constant scaling factor for the background. How much of the disagreement between CARTER 70 and OLSON 65 is due to the difference in the data and how much to the difference in the background is not clear to us at this time.

There is, however, an additional problem, which we now discuss. We have repeated the fit of CARTER 70 to their δ_{33} 's with both the Layson and what we shall call the Standard Breit-Wigner form:

$$\tan \delta_R = \frac{M\Gamma(E_{CM})}{M^2 - E_{CM}^2}, \quad \Gamma(E_{CM}) = \Gamma_0 \left(\frac{q}{q_0}\right)^3 \frac{1 + (q_0 R)^2}{1 + (qR)^2}$$

As seen in the table, our results for the Standard and Layson forms agree with those of CARTER 70. We conclude from this that there is little difference between three-parameter resonant forms so far as the values of masses and widths are concerned.

While we obtain the same low values as do CARTER 70, the disturbing feature of these fits is the poor χ^2 . This shortcoming can be removed by permitting a small background in the P_{33} partial wave itself:

$$\tan \delta_B = a_{33} \left(\frac{q}{m_\pi}\right)^3$$

We have then added δ_R and δ_B as required by unitarity, assuming the π^+p channel is perfectly elastic. Thus the overall amplitude becomes

$$T = T_B + e^{2i\delta_B} T_R$$

where $T = 1/(\cot \delta - i)$, with similar expressions for T_B and T_R . As shown in the table, a small 33 background scattering length yields an excellent χ^2 . The contribution of T_B to the total cross section ($8\pi\lambda^2 \text{Im}T_B$) is less than ~ 1 mb from threshold to ~ 1300 MeV. The primary effect of the background is to rotate the resonance in a counterclockwise sense in the Argand plot. What is surprising is that the mass and width of the $\Delta(1236)$ are much closer to the older values we gave in our tables. (Olson's scaling factor in the background tends to move M, Γ in the same direction as our background.)

It is not clear to us at the moment if such a background is needed or if simply the Breit-Wigner form used is not adequate to fit the data, now that very precise experiments are available. We illustrate this

point by replacing the width in the Standard form by a parametrization in terms of two radii:

$$\Gamma(E_{CM}) = \Gamma_0 \left(\frac{q}{q_0}\right)^3 \frac{1 + (q_0 R_1)^2 + (q_0 R_2)^4}{1 + (q R_1)^2 + (q R_2)^4}$$

Both statistically and in view of the manner in which the experimental phases were obtained, the χ^2 of 12.8 obtained with this form should be considered quite good (see the table). Finally we have used the following "polynomial" expression to fit the δ_{33} phases:

$$q^3 \cot \delta = \sum_{n=1}^N a_n q^{2n-2}$$

In this case the resonant mass corresponds to the energy ($s = E_{CM}^2 = M^2$) where $q^3 \cot \delta = 0$. The width is given by

$$\Gamma_0 = \frac{q_0^3}{M} \left[\frac{d}{ds} (q^3 \cot \delta) \right]_{s=M^2}^{-1}$$

As seen from the table, the three-parameter fit ($N=3$) yields a poor χ^2 , whereas the four-parameter fit has a reasonable χ^2 and gives essentially the same mass and width as the "two radii" Standard form.

The central problem we have raised is how seriously we should take the usual energy dependence of the width (other than q^{2l+1}). In the fits that combine background and resonance, it should be noted that $q^3 \cot(\delta_B + \delta_R)$ passes through zero at precisely the same mass as the four-parameter "polynomial" fit. Thus the inclusion of background assumes, in the context of these simple fits, that the parametrization of the widths in either the Layson or Standard forms is correct. At the very least, an analysis of the forthcoming angular distributions of Carter et al. will determine the precise energy at which $\cot \delta_{33} = 0$. More optimistically, it may put to rest the problem of parameters for the senior citizen of resonances.

We conclude that mass and width of $\Delta(1236)$ are in a state of flux; therefore we do not quote any errors in the table.

----- 33 N*3/2(1236) MASS (MEV) -----					
M	(1234.)		ROPER	65 DPWA 0++ PHASE SHIFT AN.	
M++	1230.0	0.6	CARTER	70 MPWA ++ PI+P SIG. TOTAL	1/71*
M++	1236.0	0.55	OLSSON	65 RVUE ++ TOTAL-SIGMA DATA	
M++ AVG	1233.3	3.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 7.4)		
MO	1232.9	0.6	CARTER	70 MPWA 0 PI-P SIG. TOTAL	1/71*
MO	1236.45	0.65	OLSSON	65 RVUE 0	
MO AVG	1234.5	1.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 4.0)		
----- 33 N*3/2(1236) WIDTH (MEV) -----					
W	(120.)		ROPER	65 DPWA 0++ PHASE SHIFT AN.	
W++	112.8	3.0	CARTER	70 MPWA ++ PI+P SIG TOT.	1/71*
W++	129.0	2.0	OLSSON	65 RVUE ++	
W++ AVG	117.8	3.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)		
WO	114.7	3.0	CARTER	70 MPWA 0 PI-P SIG TOT.	1/71*
WO	119.6	2.4	OLSSON	65 RVUE 0	
WO AVG	117.7	2.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		

Baryons

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle type (D, R), mass (e.g., 10.451), and other parameters (OLSSON, 65 RVUE, etc.).

Table for 33 N*3/2(1236) PARTIAL DECAY MODES, listing decay masses and branching ratios.

Table for 33 N*3/2(1236) BRANCHING RATIOS, listing ratios for various decay channels.

Table for 33 N*3/2(1236) BRANCHING RATIOS (continued), listing further ratios and references.

FOR EXTENSIVE REFERENCES TO DATA AND PHASE-SHIFT ANALYSES TILL 1965, SEE RDP#P 65, ESPECIALLY APPENDIX II.

Table for 81 N*3/2(1236, JP=3/2+) I=3/2 PROD. EXP., listing bump parameters.

Table for 81 N*3/2(1236) MASS (MEV), listing mass values and uncertainties.

Table for 81 N*3/2(1236) MASS DIFFERENCE (MEV), listing mass differences.

Table for 81 N*3/2(1236) WIDTH (MEV), listing width values.

REFERENCES -- N*3/2(1236) PROD. EXPERIMENTS, listing experimental references.

Table for 82 N*3/2(1650, JP=1/2-) I=3/2 S31, listing bump parameters.

Table for 82 N*3/2(1650) MASS (MEV), listing mass values.

Table for 82 N*3/2(1650) WIDTH (MEV), listing width values.

Table for 82 N*3/2(1650) PARTIAL DECAY MODES, listing decay masses.

Table for 82 N*3/2(1650) BRANCHING RATIOS, listing ratios for various channels.

REFERENCES -- N*3/2(1650), listing references for this resonance.

Table for 82 N*3/2(1650) BRANCHING RATIOS (continued), listing further ratios.

PAPERS NOT REFERRED TO IN DATA CARDS, listing references not in data cards.

Table for 10 N*3/2(1670, JP=3/2-) I=3/2 D33, listing bump parameters.

FOR DISCUSSION CONCERNING RESONANT PARAMETERS, SEE NOTE PRECEDING N*1/2(1470).

Table for 10 N*3/2(1670) MASS (MEV), listing mass values.

Table for 10 N*3/2(1670) WIDTH (MEV), listing width values.

Table for 10 N*3/2(1670) PARTIAL DECAY MODES, listing decay masses.

Table for 10 N*3/2(1670) BRANCHING RATIOS, listing ratios for various channels.

Table for 82 N*3/2(1670) INTO (K SIGMA)/TOTAL, listing ratios.

REFERENCES -- N*3/2(1670), listing references for this resonance.

Table for 19 N*3/2(1690, JP=3/2+) I=3/2 P33, listing bump parameters.

Table for 19 N*3/2(1690) MASS (MEV), listing mass values.

Table for 19 N*3/2(1690) BRANCHING RATIOS, listing ratios for various channels.

Baryons

For notation, see illustrated key at beginning of data card listings.

Table for particle Δ(1910) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1690)
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
AYED 70 KIEV CONF

Table for particle Δ(1890) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1890)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1950) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1890)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1910) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1890)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1910) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1910)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1890) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1910)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1950) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

REFERENCES -- N*3/2(1890)
DONNACH1 6R PL 268 161
DONNACH2 6R VIENNA 139
KIRSOPP 6R THESES
FEUERBACH 70 NP 168 85
DAVIES 70 NP B21 359
AYED 70 KIEV CONF
MEHTANI 70 KIEV CONF.

Table for particle Δ(1910) with columns for mass (MEV), width (MEV), and branching ratios. Includes data for various decay modes and theoretical considerations.

Baryons

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle ID (P1-P7), decay mode (e.g., INTO PI N, INTO SIGMA K), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (R1-R6), branching ratios (e.g., (PI N)/TOTAL), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (R2-R3), branching ratios (e.g., (SIGMA K)/(PI N)/TOTAL**2), and production experiments (e.g., SEEN, BORREANI).

MORE INFORMATIONS ON ELASTIC DECAY MODES OF BUMPS, SEEN IN PRODUCTION EXPERIMENTS AROUND 1950 MEV, MAY BE FOUND IN THE NEXT ENTRY

Table with columns for particle ID (DUKE, YOKOSAWA, etc.), references (e.g., JONES, KEMP, MURPHY), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (LAYSON, HOHLER, etc.), references (e.g., W M LAYSON, G HOHLER), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (M 3, M X), mass (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (M 3, M X), width (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (P1-P2), decay mode (e.g., INTO PI N, INTO SIGMA K), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (R1-R3), branching ratios (e.g., (PI N)/TOTAL), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (R2-R3), branching ratios (e.g., (K SIGMA)/TOTAL), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (R3), branching ratios (e.g., (D(1236) PI)/(PI N)/TOTAL**2), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (DONNACHI, KIRSOPP, etc.), decay mode (e.g., PL 268 161), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (DONNACHI, KIRSOPP, etc.), decay mode (e.g., NP 108 433), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (M, N), mass (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (W), width (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (R1-R8), branching ratios (e.g., (PI N)/TOTAL), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (COOL, RISSON, etc.), references (e.g., COOL, PICCIONI), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (M 3), mass (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (M 3), width (MEV), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (M 3), decay mode (e.g., KIRSOPP 6R RVUE), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (P1), decay mode (e.g., INTO PI N), and decay masses (e.g., 139+, 938).

Table with columns for particle ID (R1-R3), branching ratios (e.g., (PI N)/TOTAL), and production experiments (e.g., APPROX, YOKOSAWA).

Table with columns for particle ID (KIRSOPP), references (e.g., R G KIRSOPP), and production experiments (e.g., APPROX, YOKOSAWA).

Δ(1960)

D₃₅

Δ(1950) BUMPS

Δ(2160)

D₃₃

Baryons

For notation, see illustrated key at beginning of data card listings.

Δ(2420) **H₃ 11**

84 N*3/2(2420, JP=11/2+) I=3/2

84 N*3/2(2420) MASS (MEV) -----

M 6 (2312.0) AYED 70 IPWA 1/71*

6 FROM ENR. DEP. FIT OF ARGAND DIAGRAM

84 N*3/2(2420) WIDTH (MEV) -----

M 6 (347.0) AYED 70 IPWA 1/71*

84 N*3/2(2420) PARTIAL DECAY MODES -----

P1 N*3/2(2420) INTO PI N 139+ 938

P2 N*3/2(2420) INTO SIGMA K 1197+ 493

84 N*3/2(2420) BRANCHING RATIOS -----

R1 N*3/2(2420) INTO (PI N)/TOTAL (P11)/TOTAL 1/71*

R1 6 (0.113) AYED 70 IPWA

REFERENCES -- N*3/2(2420)

AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACLAY) JJP

PAPERS NOT REFERRED TO IN DATA CARDS.

BFLMAYO 67 PRL 19 476 +BUCKLEY, DOBSON, + (WESTFIELD, UNICOL) JP

AYED 70 PL 318 598 +BAREYRE-VILLET (SACLAY)

M > 2200 MEV - PRODUCTION AND σ_{TOTAL} EXPERIMENTS

BUMPS

Δ(2420) 84 N*3/2(2420) I=3/2

84 N*3/2(2420) MASS (MEV) -----

M (2360.0) DIDDENS 63 CNTR PI+ P TOTAL

M (2520.0) (40.0) ALVAREZ 64 CNTR PI PHOTOPROD 7/66

M (2400.0) APPROX WAHLIG 64 OSPK 0 PI-P CH EX

M (2440.0) HOHLER 64 RVUE DATA + DISP REL 7/66

M 7423.0 10.0 CITRON 66 CNTR PI+ P TOTAL 11/67

M R (2452.0) BARGER 66 RVUE TOTAL + CH EX

R USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE

R FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.

84 N*3/2(2420) WIDTH (MEV) -----

W (200.0) DIDDENS 63 CNTR 7/66

W (245.0) HOHLER 64 RVUE 7/66

W 310.0 20.0 CITRON 66 CNTR 7/66

W B (275.0) BARGER 66 RVUE TOTAL + CH EX 11/67

84 N*3/2(2420) PARTIAL DECAY MODES -----

P1 N*3/2(2420) INTO PI N 139+ 938

P2 N*3/2(2420) INTO SIGMA K 1197+ 493

P3 N*3/2(2420) INTO NEUTRON PI+ PI+ 1236+ 139

P4 N*3/2(2420) INTO NEUTRON PI+ PI+ 939+ 139+ 139

84 N*3/2(2420) BRANCHING RATIOS -----

R1 N*3/2(2420) INTO (PI N)/TOTAL (P11)/TOTAL 7/66

R1 (0.067) APPROX DIDDENS 63 CNTR ASSUMING J=11/2

R1 0.113 0.0036 CITRON 66 CNTR ASSUMING J=11/2 7/66

R1 B (0.12) BARGER 67 FIT ASSUMING J=11/2 11/67

R1 D (0.163) DIKMEN 67 FIT ASSUMING J=11/2 11/67

R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

R1 (0.06) KORMANYOS 67 CNTR ASSUMING J=11/2 11/67

R2 N*3/2(2420) INTO (PI N)*(NEUTRON PI+ PI+)/(TOTAL**2) (P1*P4)/TOTAL**2 6/68

R2 0.0195 0.0048 GALLOWAY 68 RVUE

REFERENCES -- N*3/2(2420)

DIDDENS 63 PRL 10 262 +JENKINS, KYCIA, RILEY (BNL) I

ALVAREZ 64 PRL 12 710 +BAR-YAM, KERN, LUCKEY, OSBORNE, + (MIT, GEA)

WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)

HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I

CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I

BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC) P

BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P

DIKMEN 67 PR 18 798 F N DIKMEN (MICH) P

KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CAL TECH)

GALLOWAY 68 PL 268 334 P F GALLOWAY (INDIANA) I

PAPERS NOT REFERRED TO IN DATA CARDS.

DOBROWOL 67 PL 248 203 DOBROWOLSKI, GUSKOV, LIKHACHEV, + (DUBNA) P

BAECKE 67 NC 51A 761 J BAECKE, M YVERT (KARLSRUHE, ORSAY) J-L

WAHLIG 68 PR 168 1515 M WAHLIG, I MANNELLI (MIT, PISA)

--- FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

Δ(2850) 85 N*3/2(2850, JP= +) I=3/2

85 N*3/2(2850) MASS (MEV) -----

M (2700.0) APPROX WAHLIG 64 OSPK 0 PI-P CH EX

M (2870.0) HOHLER 64 RVUE DATA + DISP REL 7/66

M 2850.0 12.0 CITRON 66 CNTR PI+ P TOTAL 7/66

M (2850.0) BARDADIN 66 HRC ++ N4 TO P + 3 PIS

85 N*3/2(2850) WIDTH (MEV) -----

W 400.0 40.0 CITRON 66 CNTR 7/66

W (150.0) BARDADIN 66 HRC ++ 7/66

85 N*3/2(2850) PARTIAL DECAY MODES -----

P1 N*3/2(2850) INTO PI N 139+ 938

P2 N*3/2(2850) INTO P PI PI 938+ 139+ 139

P3 N*3/2(2850) INTO N PI PI 938+ 139+ 139

85 N*3/2(2850) BRANCHING RATIOS -----

R1 N*3/2(2850) INTO (PI N)/TOTAL (P11)/TOTAL 11/67

R1 ONLY (J+1/2)*(PI N)/TOTAL MEASURED FOR THIS STATE

R1 (0.251) 0.048 CITRON 66 CNTR TOTAL CROSS-SEC. 11/67

R1 B (0.224) (0.016) BARGER 66 RVUE TOTAL + CH EX. 11/67

R1 B (0.40) BARGER 67 RVUE USES KORMANYOS 66 11/67

R1 D USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE

R1 D FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.

R1 D (0.49) DIKMEN 67 RVUE USES KORMANYOS 67 11/67

R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

R1 (0.10) KORMANYOS 67 CNTR PI+ P AT 180 DEG. 11/67

R1 (0.39) DOBROWOLSKI 67 CNTR PI+ P AT 180 DEG.

REFERENCES -- N*3/2(2850)

WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)

HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I

CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I

BARDADIN 66 PL 21 357 BARDADIN-DIMONOWSKA, DANYSZ, + (WARSAW)

BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)

BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P

DIKMEN 67 PR 18 798 F N DIKMEN (MICH) P

DOBROWOLSKI 67 PL 248 203 DOBROWOLSKI, GUSKOV, LIKHACHEV, + (DUBNA) P

KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CAL TECH)

PAPERS NOT REFERRED TO IN DATA CARDS.

BAECKE 67 NC 51A 761 J BAECKE, M YVERT (KARLSRUHE, ORSAY) J-L

WAHLIG 68 PR 168 1515 M WAHLIG, I MANNELLI (MIT, PISA)

--- FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

Δ(3230) 86 N*3/2(3230, JP=) I=3/2

86 N*3/2(3230) MASS (MEV) -----

M (3230.0) CITRON 66 CNTR PI+ P TOTAL 7/66

86 N*3/2(3230) WIDTH (MEV) -----

W (440.0) CITRON 66 CNTR 7/66

86 N*3/2(3230) PARTIAL DECAY MODES -----

P1 N*3/2(3230) INTO PI N 139+ 938

P2 N*3/2(3230) INTO N PI PI 934+ 139+ 139

86 N*3/2(3230) BRANCHING RATIOS -----

R1 ONLY (J+1/2)*(PI N)/TOTAL MEASURED FOR THIS STATE

R1 (0.061) CITRON 66 CNTR TOTAL CROSS-SEC. 11/67

R1 B (0.03) (0.01) BARGER 66 RVUE TOTAL + CH EX. 11/67

R1 B (0.03) TO 0.1 BARGER 67 CNTR USES KORMANYOS 66 11/67

R1 B USES REGGE AMP. + RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE

R2 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.

R1 D (0.25) DIKMEN 67 RVUE USES KORMANYOS 67 11/67

R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

REFERENCES -- N*3/2(3230)

CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I

BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC) P

BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P

DIKMEN 67 PR 18 798 F N DIKMEN (MICH) P

KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ARG) P

DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CAL TECH)

END PRODUCTION EXPERIMENTS

EXOTIC NUCLEONS CROSS SECTIONS LIMITS FOR EXOTIC NUCLEONS

THIS IS NOT A COMPLETE LIST. WE WILL TABULATE EXOTICS FROM NOW ON

CS LT. 40 MB I=5/2, M=1540-1750 BANNER 70 NSPK +++ PI+P, 1.9 GEV/C 7/70

REFERENCES --- EXOTIC NUCLEONS

BANNER 70 NP B15 205 +CHEZE, HAMEL, TEIGER, ZACCONE + (SACLAY)

For notation, see illustrated key at beginning of data card listings.

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EX(1640)  92 EX(1640),JP= 11=
-----
M  A  2911627.1  112.1  92 EX(1640) MASS (MEV)  70 DRC  -- K-D AT 4.91GEV/C  3/71*
7  A  FOUR S. D. EFFECT  PRICE
-----
M  R  29  30.  OR LESS .00 CL  92 EX(1640) WIDTH (MEV)  70 DRC  -- PI-PI-N BUMP  3/71*
W  R  CROSS SECTION 13.0+-3.9 MICROBARN  PRICE
-----
*****
REFERENCES -- EX(1640)
*****
PRICE  70 PL 33R,533  +REFG,SALANT,WATERS,WEBSTER,WEINBERG (VAND)
*****

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Note on Possible Z^* 's

Although much work has been done on the strange-ness +1 reactions during the past year, it is not yet clear whether the peaks seen in total KN cross sections near 1 GeV/c are resonances; see Fig. 1. Since positive-strangeness baryons cannot be made from three quarks, it is very important to find out if these peaks are resonances.

(a) I = 0 System. New total-cross-section data have been reported by the Arizona group in the 0.36 to 0.72 GeV/c region (BOWEN 70). The old K^+p and K^+d total cross sections have been reanalyzed by ABRAMS 69. They made a new unfolding of the isospin-0 total cross section and found that the well-known peak in this cross section just above 1.0 GeV/c ($M \sim 1865$ MeV) has a companion at about 0.8 GeV/c ($M \sim 1780$ MeV). COOL 70 adds the new data of BOWEN 70 to the analysis and obtains the two possible results shown in Fig. 2a for the I = 0 cross section. DOWELL 70 also reports similar results. The choice of two possibilities depends mostly on the disagreement between the data and very little on the method of unfolding, as discussed by LYNCH 70.

Until the K^+p and K^+d cross sections are re-measured in this region, the lower peak, though probable, cannot be considered definitely established. There is, however, no doubt about there being a large broad peak in the isospin-0 elastic cross section. The inelastic cross section is very small until the K^*N threshold at 1.08 GeV/c is approached, then it rises rapidly (HIRATA 68 and HIRATA 70), as shown in Fig. 2c. Subtracting this from total cross sections gives the two possible results of Fig. 2b. (The elastic cross section has not been independently measured.)

The $Z_0^*(1780)$ in the listings below is the result of fitting the elastic peak with a Breit-Wigner resonance form. The width is anomalously large, nearly 600 MeV. The resonance (if it exists) must be nearly entirely elastic because the inelastic cross section is small at the peak. If J were greater than 1/2, the resonant peak would exceed the observed height of $4\pi\lambda^2$. This fixes the spin as 1/2, and means that there is little cross section left over for other partial waves. Of course it is quite possible that in fact the peak is not caused by a resonance at all.

Phase shift analyses of the elastic scattering data in the I = 0 state have been done only up to 810 MeV/c;¹ therefore it is not clear whether the structure (either double-peaked or not) is associated with a definite resonance-like structure in one or more partial waves.

(b) I = 1 System. New K^+p data have been reported by REBKA 70 from 1.54 to 1.74 GeV/c and by BARBER 70 from 1.4 to 2.3 GeV/c. As for new analyses at the Duke Conference on Hyperons, a total of nine partial-wave analyses were presented. Each of these analyses gets more than one solution to choose from, which indicates that the data are not good enough to eliminate some of the possibilities. Certainly new data and the simultaneous analysis of elastic and inelastic channels (copious inelastic data is desirable at the moment) could improve the understanding of this system.

The P_{13} amplitude still remains the best candidate for a resonance in the K^+p system. The preferred P_{13} Argand plots obtained by the various groups are shown in Fig. 3. Each analysis in one way or another gets at least one solution with a counterclockwise P_{13} amplitude. Resonant P_{13} is claimed by REBKA 70, KATO 70, ERNE 70, and GIACOMELLI 70; the results of the other five analyses are not so clear cut. AYED 70 claim that the speed of variation of the amplitude as a function of the energy is nearly as well fit by a smoothly varying background as by a broad (500-MeV) Breit-Wigner amplitude superimposed on a background. For this reason they claim one is not obliged to invoke a resonance. BENNETT 70 got two solutions: the preferred one shows no resonant behavior; the other one might be resonant, but the errors on the amplitudes are too large to believe any circle drawn among them.

CARRERAS-1 70 get a P_{13} amplitude similar to that found in other analyses, which can be interpreted

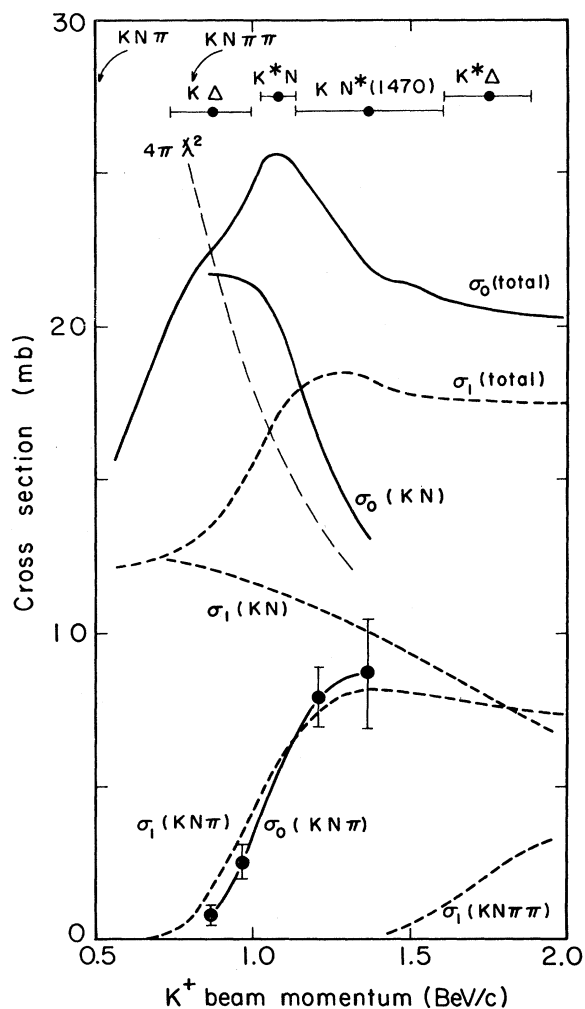
Baryons


Fig. 1.

KN total and partial cross sections. Subscripts indicate isospin. Total cross sections are from CARTER 68, which uses data from COOL 66 and BUGG 68. Isospin-1 partial cross sections are adapted from a compilation made by BLAND 68. Isospin-0 partial cross sections are from HIRATA 68. Thresholds for various processes are indicated at the top.

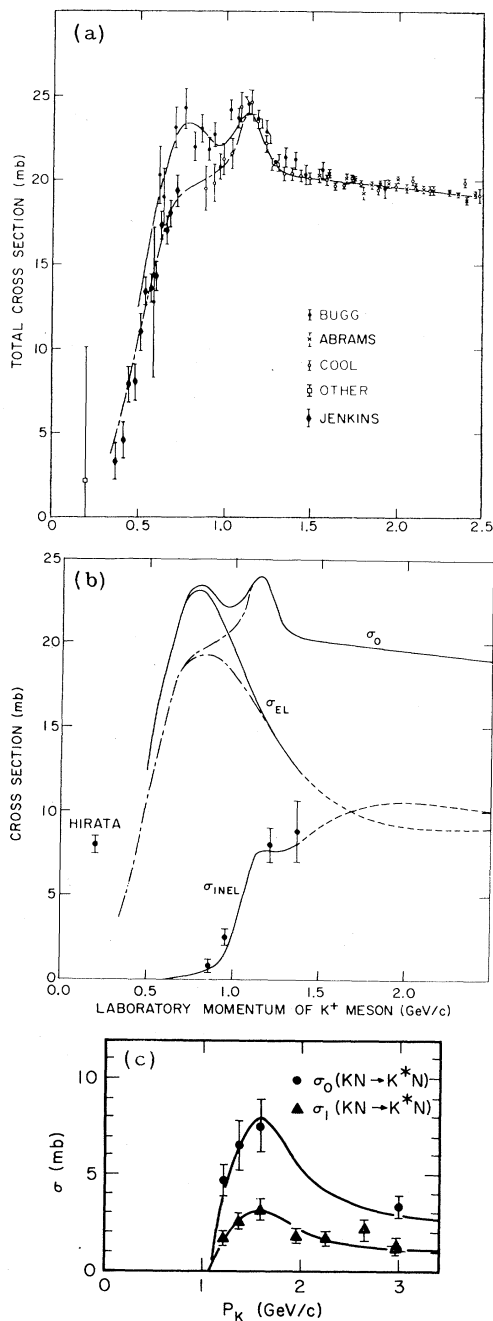


Fig. 2.

(a) The $I = 0$ KN cross section evaluated by COOL 70. The new data of Jenkins are reported in BOWEN 70.

(b) The $I = 0$ cross sections as obtained by the two possible interpretations of the data. This shows the uncertainties in the 0.8 to 1.2 GeV/c region.

(c) Energy dependence of the I-spin 0 and I-spin 1 cross sections for the reaction $KN \rightarrow K^* N$

Baryons

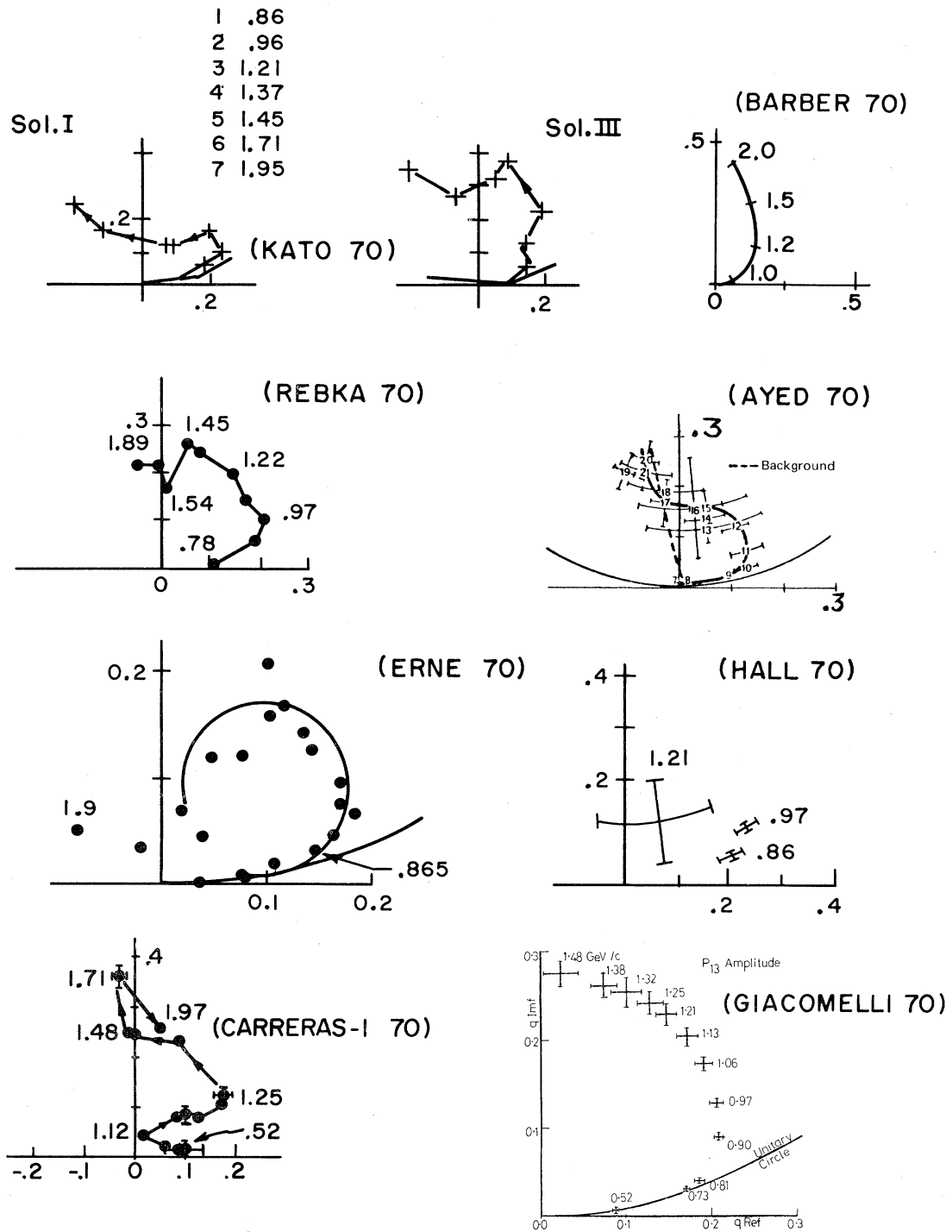


Fig. 3. Argand plots for the P_{13} partial wave as obtained in partial-wave analyses performed by the authors indicated. (CARRERAS-1 70 plotted by us from η , δ .)

Baryons

For notation, see illustrated key at beginning of data card listings.

as having resonant behavior, but do not feel that the resonance interpretation is compelling.

BARBER 70 include the new data up to 2.3 GeV/c. Their preferred solution is shown in Fig. 3 and shows no structure in the speed. They also report a second solution which resembles the Argonne solution I (KATO 70 in Fig. 3) below 1.8 GeV/c.

HALL 70 have presented an analysis that includes the elastic and inelastic data. The analysis was done at three momenta, and the P_{13} amplitude of their best solution is shown in Fig. 3. Here again the errors are really too large to detect any possible Breit-Wigner behavior.

The S_{11} amplitude. It should be mentioned that AYED 70 find a solution with the P_{13} negative and moving clockwise, and the S_{11} positive and moving counterclockwise. In this solution the S wave is essentially zero below 520 MeV/c, and a negative P_{11} wave rather than S wave would fit the angular distributions at low energy. The trouble with this solution is that the P_{11} phase shifts do not show the q^3 dependence at low energy as required from a P wave. CARRERAS 70 report an attractive S wave at low energy. This is in contradiction with the constructive interference between the elastic and Coulomb scattering suggested by the data. CARRERAS-1 70 mention the possibility that S and P_{11} waves contribute equally at low energy (below 300 MeV/c). This is also contradicted by the data, which do not show the required SP interference term. See GOLDHABER 70 for more discussions on the sign of the S_{11} .

Threshold effects. Another possible way to describe the P_{13} amplitude would be in terms of a coupled-channel threshold effect: the KN amplitude becomes rapidly absorptive as it feeds the rapidly increasing $K\Delta$ channel. The main question still remains: Is it also a resonance? If it is, its elasticity is small (≈ 0.2) and it decays mainly into $K\Delta$. But a definite conclusion has yet to be made and awaits much more data.

One more comment on exotic resonances is that, as pointed out by ERNE 70, the present upper limits for the cross sections for production of broad exotic resonances are not very small; that is, they are of the same order as cross sections for Y^* or N^* production.

Conclusions. The existence of exotic resonances is still an open question. It looks to us that the following are needed:

- 1) More high-precision data in the critical region

where the peak is seen.

- 2) The fit to the data should extend over a very wide region, down to threshold and up to 2 GeV/c as already done by ERNE 70, BARBER 70, and AYED 70.

- 3) Simultaneous study of elastic and inelastic channels should certainly help.

- 4) More emphasis on the speed plot.²

- 5) For the $I = 0$ states it would be very useful, however difficult it may be, to measure the elastic cross section.

References

1. V. J. Stenger, W. E. Slater, D. H. Stork, H. K. Ticho, G. Goldhaber, and S. Goldhaber, Phys. Rev. **134**, B1111 (1964).
2. See the January 1970 edition of the Review of Particle Properties for some earlier speed plots and definitions.

Z₀(1780)		95 Z*(1780, JP=1/2 ⁺ I=0 THIS EFFECT, IF A RESONANCE, MUST HAVE SPIN=1/2, BECAUSE THE INELASTIC CROSS SECTION IS VERY SMALL AND THE TOTAL CROSS SECTION IS ABOUT 4*PI/K**2.	

95 Z*(1780) MASS (MEV)		-----	
M	1780.0 10.0	COOL	70 CNTR + K*P, D TOTAL 1/71*
M	SEP*	DOWELL	70 RVUE 7/70
M	SEE ALSO DISCUSSION OF LYNCH 70		7/70

95 Z*(1780) WIDTH (MEV)		-----	
W	(565.0)	COOL	70 CNTR + K*P, D TOTAL 1/71*

95 Z*(1780) PARTIAL DECAY MODES		-----	
P1	Z*(1780) INTO K N		DECAY MASSES 493+ 939

95 Z*(1780) BRANCHING RATIOS		-----	
R1	Z*(1780) INTO (K N)/TOTAL		(P1)/TOTAL
R1	(0.95)	COOL	70 CNTR + K*P, D TOTAL 1/71*

REFERENCES -- Z*(1780)			
SEE REFERENCES FOR THE Z*(1900)			

Z₀(1865)		96 Z*(1865, JP=) I=0 THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K* N THRESHOLD. SEE HIRATA 69 AND 70.	

96 Z*(1865) MASS (MEV)		-----	
M	(1868.0) (10.0)	KYCIA	67 CNTR K*P, D TOTAL 8/67
M	(1860.0) (15.0)	CARTER	67 THEO DISPERSION REL. 8/67

96 Z*(1865) WIDTH (MEV)		-----	
W	(160.0) (30.0)	KYCIA	67 CNTR 8/67
W	(200.0) (50.0)	CARTER	67 THEO 8/67

96 Z*(1865) PARTIAL DECAY MODES		-----	
P1	Z*(1865) INTO K N		DECAY MASSES 493+ 939
P2	Z*(1865) INTO N K*(890)		938+ 897

96 Z*(1865) BRANCHING RATIOS		-----	
R1	Z*(1865) INTO (K N)/TOTAL		(P1)/TOTAL
R1	(0.40) (0.05)	KYCIA	67 CNTR IF J=1/2 8/67
R1	(0.31) (0.05)	CARTER	67 THEO IF J=1/2 8/67
R2	Z*(1865) INTO N K*(890)		(P2)
R2	MAIN INELASTIC DECAY	HIRATA	68 HBC 11/68

REFERENCES -- Z*(1865)			
SEE REFERENCES FOR THE Z*(1900)			

Baryons

For notation, see illustrated key at beginning of data card listings.

Z ₁ (1900)		97 Z ⁺ 1(1900, JP=) I=1	
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K ⁺ N ⁺ THRESHOLD. IF A RESONANCE, THE SPIN-PARITY IS ALMOST CERTAINLY 3/2 ⁺ .			
----- 97 Z ⁺ 1(1900) MASS (MEV) -----			
M	1900.0	COOL	70 CNTR ++ K+P TOTAL 1/71*
M	1 (1922.0)	AYED	70 IPWA P13,SOL.I 6/70
M	1 (1999.0)	AYED	70 IPWA P13,SOL.II 6/70
M	1 (2030.0)	AYED	70 IPWA S11,SOL.III 6/70
THREE SOLNS IN ORDER OF DECREASING SIGNIFICANCE, THROUGH AYED TO 1 GIVE PARAMETERS, THEY CONCLUDE RESONANT INTERPRETATION DOUBTFUL.			
M	2 (1940.0)	BARNETT	70 IPWA P13,SOLN III 7/70
RESONANCE SIGNAL RARELY ABOVE BACKGROUND DUE TO THE LARGE ERRORS IN THE AMPLITUDES RESULTING FROM THE ANALYSIS			
M	2 (1959.0)	ERNE	70 IPWA P13 6/70
M	1 (1880.0) (90.0)	KATO	70 IPWA P13 IN SOLN I 6/70
M	1 (1950.0) (100.)	KATO	70 IPWA SOL. III 7/70
----- 97 Z ⁺ 1(1900) WIDTH (MEV) -----			
W	1 (240.0)	COOL	70 CNTR ++ K+P TOTAL 1/71*
W	1 (220.0)	AYED	70 IPWA K+P 6/70
W	1 (200.0)	AYED	70 IPWA K+P 6/70
W	1 (157.0)	AYED	70 IPWA K+P 6/70
W	2 (80.0)	BARNETT	70 K+P EIPWA 7/70
W	2 (251.0)	ERNE	70 IPWA K+P 6/70
W	ABOUT 130.0	KATO	70 IPWA P13 IN SOLN I 6/70
W	(180.) APPROX.	KATO	70 IPWA SOL. III 7/70
SEE THE NOTES ACCOMPANYING MASSES QUOTED.			
----- 97 Z ⁺ 1(1900) PARTIAL DECAY MODES -----			
P1	Z ⁺ 1(1900) INTO K N	DECAY MASSES 493+ 938	
P2	Z ⁺ 1(1900) INTO N*3/2(1236) K	1236+ 493	
----- 97 Z ⁺ 1(1900) BRANCHING RATIOS -----			
R1	Z ⁺ 1(1900) INTO (K N)/TOTAL	(P1)/TOTAL	
R1	(0.10) OR LESS	CARTER	67 THFO DISPERSION REL. 8/67
R1	(0.12) (ASSUMING J=3/2)	COOL	70 CNTR ++ K+P TOTAL 1/71*
R1	(0.16)	AYED	70 IPWA 6/70
R1	(0.20)	AYED	70 IPWA 6/70
R1	(0.17)	AYED	70 IPWA 6/70
R1	(0.15)	ERNE	70 IPWA 6/70
R1	ABOUT 0.1	KATO	70 IPWA P13 IN SOLN I 6/70
R1	(0.28) APPROX.	KATO	70 IPWA SOL. III 7/70
R1	(0.09)	BARNETT	70 K+P EIPWA 7/70
SEE NOTES ACCOMPANYING THE MASSES QUOTED.			
R2	Z ⁺ 1(1900) INTO K N*3/2(1236)	(P2)	
R2	MAIN INELASTIC DECAY	BLAND	67 HBC ++ 8/67
----- Z ⁺ 1 CROSS SECTION LIMITS (MICROBARN) -----			
SEE MINIREVIEW PRECEDING Z*0			
CS	LESS THAN 50.	BASSOMPI 68 HBC	K+P TO Z++ P1+ 10/69
CS	A LESS THAN 2 +.3	-1 ANDERSSON 69 ASPK +	PI-P TO K-Z** 10/69
CS	A ABOVE LIMIT FOR M=1.2 TO 1.4 GEV -	CL= 99 P.C.	
CS	A LESS THAN 1.4 +1.9	-5 ANDERSSON 69 ASPK +	PI-P TO K-Z** 10/69
CS	B ABOVE LIMIT FOR M=1.5 TO 2.5 GEV		

REFERENCES -- Z ⁺ 1(1900)			
TOTAL-CROSS-SECTION EXPERIMENTS ---			
COOL	66 PRL 17 102	+GIACOMELLI, KYCIA, LEONTIC, LI, LUNDRY, +	(BNL) I
	SLIGHTLY REVISED RESULTS FROM KYCIA 67 REPLACE COOL 66 ---		(BNL) I
KYCIA	67 PRIVATE COMM.	T F KYCIA	(BNL) I
ARRAMS	67 PRL 19 259	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, +	(BNL) I
BUGG	68 PR 169 1466	+GILMORE, KNIGHT, + (RTHFD, BRMGH, CVNDSH) I	
	COOL 66, ARRAMS 67, AND ARRAMS 69 ARE SUPERCEDED BY COOL 70 AND ARRAMS 70		
ARRAMS	69 PL 304 564	+COOL, GIACOMELLI, KYCIA, LI, MICHAEL (BNL) I	
ARRAMS	70 PR 01 1917	+COOL, GIACOMELLI, KYCIA, LEONTIC, +	(BNL) I
BOWEN	70 PP 02 2599	+CALDWELL, DIKMEN, JENKINS, KALBACH, + (ARIZ) I	
COOL	70 DUKE CONF 47	R F COOL	(BNL) I
COOL I	70 PR 01 1887	+GIACOMELLI, KYCIA, LEONTIC, LI, +	(BNL) I
OWELL	70 DUKE 53	J.D. OWELL	(RIRN)
LYNCH	70 DUKE 9	G LYNCH (REVIEWER OF CR. SEC. DATA)	(LRL)
DISPERSION-RELATION CALCULATION USING TOTAL-CROSS-SECTION DATA ---			
CARTER	67 PRL 18 801	A A CARTER	(CAVENDISH)
CARTER	68 PREPRINT	A A CARTER	(CAVENDISH)
A K-MATRIX ANALYSIS OF SOME OF THE EARLY K+P DATA --- (ILLINOIS)			
HITE	67 THESIS	G E HITE	(ILLINOIS)
A RFGG-POLE ANALYSIS EXTENDED TO LOW ENERGIES -- (DARESURY, MANCH)			
CARRERAS	70 NP 819 349	B CARRERAS, A DONNACHIE	(DARESURY, MANCH)
EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS ---			
BLAND	67 PRL 18 1077	+BOWLER, BROWN, G+S GOLDHABER, SEEGER, +	(LRL)
BLAND	68 PRL 18 1315	THESES R W BLAND	(LRL)
HIRATA	68 PRL 21 1485	HIRATA, WOHL, GOLDHABER, TRILLING	(LRL)
BLAND	69 NP P13 595	+BOWLER, BROWN, KADYK, GOLDHABER, +	(LRL)
BLAND	70 NP 818 537	+BOWLER, BROWN, GOLDHABER, +	(LRL)
	BLAND 69 AND BLAND 70 REPLACE BLAND 67 AND BLAND 68.		
HIRATA	70 DUKE 429	+GOLDHABER, SEEGER, TRILLING, WOHL	(LRL)
RGR	70 KIEV CONF	(BOLOGNA, GLASGOW, POME, TRIESTE COLLAB) I	
THE MAIN ELASTIC-SCATTERING AND POLARIZATION EXPERIMENTS AND PHASE-SHIFT ANALYSES --			
CARROLL	68 PRL 21 1282	+FISCHER, LUNDRY, PHILLIPS, +	(BNL, ROCH)
ANDERS-1	69 PL 288 611	ANDERSSON, DAUM, ERNE, LAGNAUX, +	(CERN)
ASBURY	69 PRL 23 194	+DONNELLY, KATO, LUNDQUIST, NOVY, +	(ARG-MO)
BLAND	69 PL 298 618	R W BLAND, G GOLDHABER, G H TRILLING	(LRL)
ANDERS-2	69 PL 308 56	ANDERSSON, DAUM, ERNE, LAGNAUX, +	(CERN)
AYED	70 PL 328 404	+BAREVRE, FELTESSE, VILLET	(SACLAY) IJP
BARNER	70 PL 328 214	+BONNE, DUFF, HEYMANN, IMRIE, +	(LOUG, RHEI) IJP
BARNETT	70 DUKE 443	+GOLDMAN, LAASANEN, STEINBERG	(MARYLAND) IJP
CARRERAS	70 NP 823 525	B CARRERAS, A DONNACHIE	(DARE) IJP
ALSD	70 DUKE 447	+DONNACHIE, WIRSOPP	(DARE+MCHS+U-COIN)
ERNE	70 DUKE 375	+SENS, WAGNER	(CERN) IJP
GIACOMEL	70 NP P20 301	GIACOMELLI, GRIFFITHS, IRGNA, GLAS, ROMA, TRSTI IJP	
HALL	70 DUKE 435	+BLAND, GOLDHABER, TRILLING	(LRL)
KATO	70 PRL 24 615	+KOEHLER, NOVY, YOKOSAWA, +	(ARG, WJES) IJP
REKA	70 PRL 24 160	+ROTHBERG, ETKINS, GLODIS, +	(YALE) IJP

EARLIER ANALYSES THAT DO NOT INCLUDE RECENT POLARIZATION DATA --

LEA	68 PR 165 1770	LEA, MARTIN, OADES	(PTHFO, BNL, CERN)
MARTIN	68 PRL 21 1286	B R MARTIN	(BNL)
HALL	69 UCRL-19231	HALL, BLAND, GOLDHABER, TRILLING	(LRL)
LEA	69 LUND PAPER 362	LEA, MARTIN, OADES	(RHEL+UCL)
CUTKOSKY	70 PR D1 2547	R E CUTKOSKY, R B DEO	(CARNEGIE-MELLON) I

PRODUCTION EXPERIMENTS THAT LOOK FOR A Z* ---

TYSON	67 PRL 19 255	+GREENBERG, HUGHES, LU, MINEHART, MORI, (YALE)
BASSOMPI	68 PL 278 468	BASSOMPIERRE, + (CERN, BRUXELLES)
MORI	68 PL 288 152	+GREENBERG, HUGHES, LU, ROTHBERG, + (YALE)
ANDERSON	69 PL 298 136	+RLESER, BLIEDEN, COLLINS, + (BNL, CARNEGIE)
MORI	69 PR 185 1687	+GREENBERG, HUGHES, LU, MINEHART, + (YALE)
	-- MORI 69 REPLACES TYSON 67 AND MORI 68.	

LATEST REVIEW TALKS

LEVISETT	69 LUND CONF 341	R LEVI SETTI (RAPPORTFURI)	(CHICAGO)
GOLDHABER	70 DUKE 407	G. GOLDHABER (REVIEWER)	(LRL)

Note on Y*¹s

The number of known or suspected Y*¹ states has increased considerably in the last few years, following closely a similar increase in the number of N* states. Just as the recently discovered N*¹s are only weakly coupled in the πN → πN reaction, so also are the recently discovered Y*¹s only weakly coupled in the $\bar{K}N \rightarrow \bar{K}N$, $\bar{K}N \rightarrow \Lambda\pi$, and $\bar{K}N \rightarrow \Sigma\pi$ reactions. The older, well-established resonances are usually clearly visible as peaks in cross sections, as characteristic variations of angular distributions of 2-body final states, and (or) as peaks in invariant-mass distributions of subsets of particles in 3-or-more-body final states. Although some of the newer and less-well-established resonances are seen as small peaks in invariant-mass distributions, many of them make no direct appearance at all, often because there are many states at the same mass and it is not clear which ones (or how many) are being observed. Rather than the 2-body reactions are partial-wave analyzed, some of the amplitudes are found to traverse resonance-like counterclockwise circles. Clearly the results of partial-wave analysis give the J^P information, whereas a peak seen in an invariant mass distribution or a total cross section usually cannot be analyzed for its quantum numbers. We will keep information coming from formation experiments and from production experiments separate, whenever necessary.

Production experiments. These types of experiments are often difficult to analyze. Information on I = 0 states is possible only when there is no I = 1 state at similar mass. The main controversies at the present time concern resonances in the 1600 to 1700 MeV region. See the mini-reviews on Σ(1620) and on Σ(1670) in the listings. A good review is given by MILLER 70.¹ Also, the branching ratios of Σ(1915) F₁₅ as measured in formation and production experiments do not agree. This is probably

Baryons

due to two facts: 1) the elasticity is small, 2) the nearby $D_{13}(1940)$ may contribute to production experiments.

Formation experiments. Partial-wave analyses have been performed on $\bar{K}N$, $\Lambda\pi$, $\Sigma\pi$ and ΞK channels. Given the present accuracy of the data it is not possible to perform a completely energy-independent analysis, that is, solve for the partial-wave amplitudes at each energy in a model-independent way. Usually many solutions are found and even when it is required that solutions at neighboring energies join smoothly, it is not possible to select a unique overall solution. To overcome this, one specifies the form of the energy dependence of some or all of the partial-wave amplitudes. Analyses in which the energy dependence of all the amplitudes is specified are called energy dependent. When referring to results of this type of analysis, the technique listed is DPWA. Thus an amplitude known to resonate will be given a Breit-Wigner form, whereas an amplitude not a priori known to resonate may be tried alternately with a resonance form and with some simple nonresonant form, the choice between these then being made by comparing the goodness-of-fit for the two fits. See the Note on Resonance Parameters preceding the Baryon listings for additional discussion of DPWA and of the K-Matrix parametrization of KIM 74.

Analyses in which most of the amplitudes are left unspecified are called (not quite correctly) energy independent. The technique used for these analyses is listed as IPWA. Figure 1 shows results of such an analysis of the reaction $K^-p \rightarrow \Lambda\pi$ by ARMENTEROS 70. The D_{15} amplitude was fixed as the $\Sigma(1765)$, with resonance parameters obtained from an earlier energy-dependent analysis. This amplitude acts as an analyzer for the other amplitudes, which were allowed to vary freely. The S_{11} and D_{13} amplitudes appear to resonate. Figure 2 shows results of a similar analysis, also by ARMENTEROS 70, of the reaction $K^-p \rightarrow \Sigma\pi$. Here the $D_{03}\Lambda(1520)$, $D_{13}\Sigma(1670)$, $D_{03}\Lambda(1690)$, $D_{15}\Sigma(1765)$, $F_{05}\Lambda(1820)$, and $F_{15}\Sigma(1910)$ were fixed. It appears that several of the other amplitudes may resonate too. It should be clear from the figures that it is not always possible to decide whether or not an amplitude resonates. Neither is it possible to determine very accurately the parameters of the amplitudes that do resonate, nor to assign meaningful errors to the parameters. The state of knowledge of the newer

Y^* 's is rather more qualitative than quantitative.

Partial-wave analyses above 1.1 GeV/c.

One of us² has reviewed this energy region at the Duke Conference on Hyperon Resonances. The $\Lambda\pi$ channel has been analyzed by five groups, using different sets of data points that are in very good agreement with each other. LITCHFIELD 70 has used all the data, which plotted together show absolutely no discrepancies. In spite of this fact the five analyses agree only on three of the eight partial waves used to fit the data! And all five groups claim to obtain a good fit. Table I, taken from Ref. 2, shows the resonances included in the fits; the other partial waves were parametrized as background amplitudes with the indicated number of parameters. The three partial waves where agreement was found were D_{15} , F_{15} , F_{17} .

The problem here, more than insufficient data, is that the $\Lambda\pi$ channel alone is not constrained very much and a large number of solutions can be found. For the $\Sigma\pi$ channel, where two I-spin states are present, the interference terms certainly help to reduce the number of possible solutions. Table II, again taken from Ref. 2, is a summary of the states claimed above 1.1 GeV/c where, in view of the above difficulties, most of the new states have been classified "wait". The signs of the amplitudes at resonance follow the convention advocated by Levi Setti.³ Clearly, resonances have to show up very strongly in at least one channel before they are accepted without reservation.

Errors on masses and widths. Often the quoted errors are only statistical, but, as discussed in the Note on Resonance Parameters, at the beginning of the Baryon listings, the values of masses and widths can change well above these errors when a new parametrization is used. For this reason we report the values of M , Γ , and x_1 obtained by different authors even if they analyze the same data. The spread of these masses and widths is certainly a better estimate of the uncertainties than the statistical errors.

Recently it has become the custom to quote errors as obtained by inspection of various fits done with different hypotheses [see for example BERTHON 70 and GALTIERI 70 under $\Sigma(1915)$]. These errors are probably more realistic. In conclusion, we chose not to give errors on masses and widths determined in partial-wave analyses, but, whenever necessary, we give a range of values.

Baryons

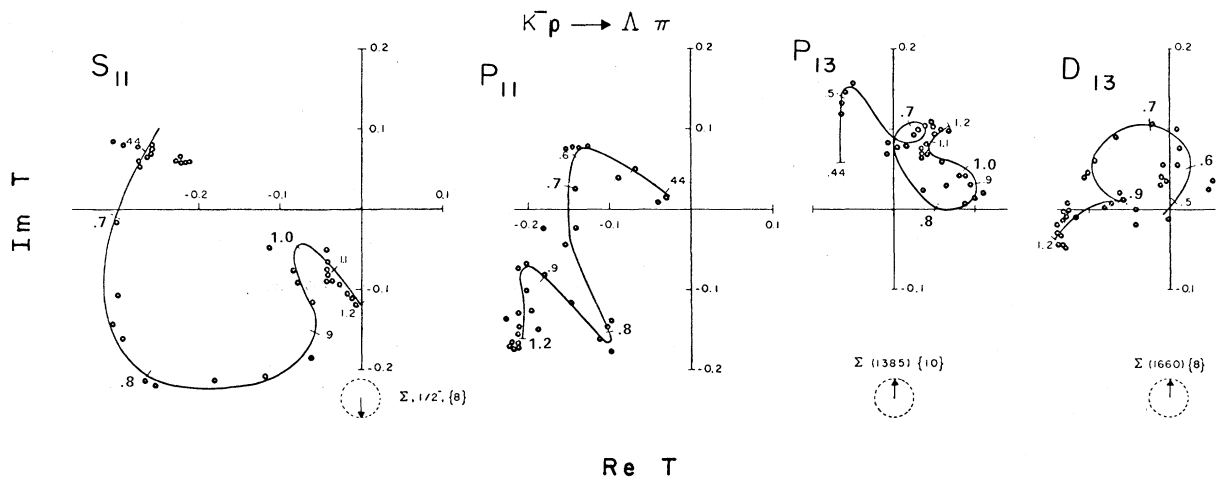


Fig. 1. Partial-wave amplitudes for the reaction $K^-p \rightarrow \Lambda \pi$, as determined in the energy-independent analysis of ARMENTEROS 69. The K^- laboratory momenta are indicated. The arrows in a circle, drawn in the lower part of the imaginary axes, fix the sign convention used. See LEVI SETTI 69. Notice that the sign convention used here is different from the one of the Argand plots of our 1969 edition [Rev. Mod. Phys. 41, 109 (1969)].

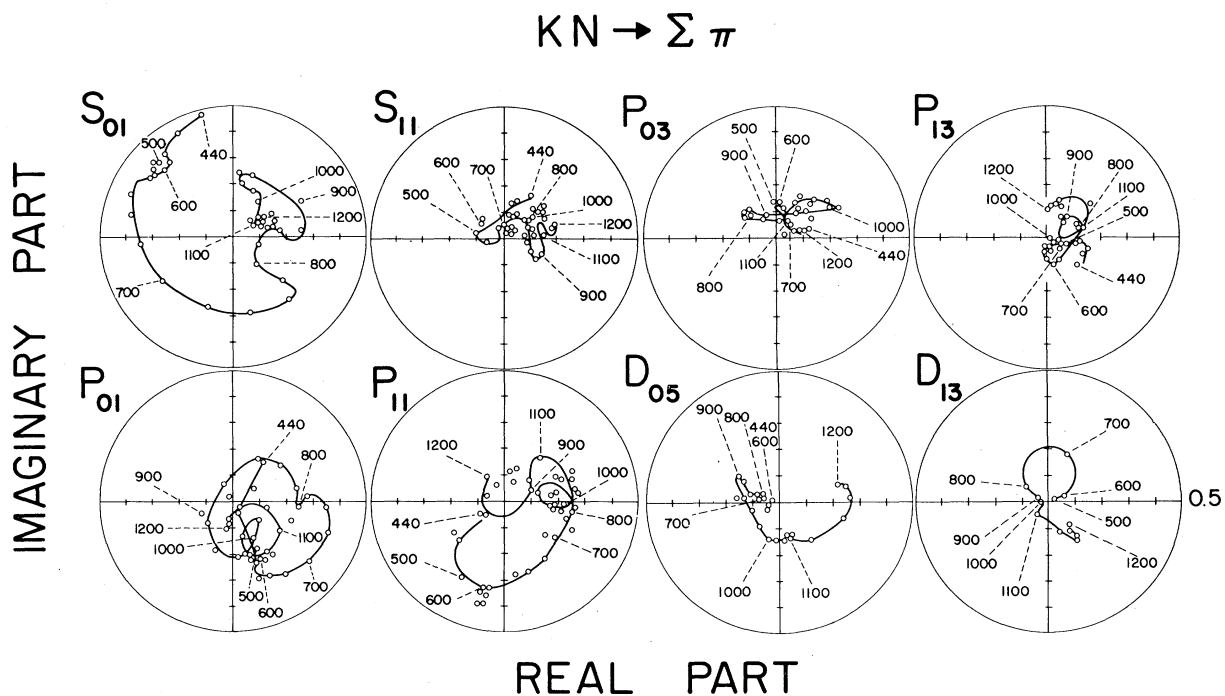


Fig. 2. Partial-wave amplitudes for the reaction $K^-p \rightarrow \Sigma \pi$, as determined in the energy-independent analysis of ARMENTEROS 70. The K^- laboratory momenta are indicated. Here again the sign convention follows LEVI SETTI 69.

Baryons

Table I. Results of partial-wave analyses in the $\Lambda\pi$ channel. In the second column $N = 50$ means that 50 parameters were used in the overall fit. In the first row to the right, N_{par} indicates how many parameters were used for that partial wave. R(4) means that the resonant amplitude included four parameters (M, Γ , $|t|, \Phi$).

		N_{par}	P_{11}	P_{13}	D_{13}	F_{15}	F_{17}	G_{17}	$^d G_{19}$
^a SMART 66	$\Lambda\pi^-, \Lambda\pi^0$	N_{par}	4+R(4)	6	6	4+R(4)	R(4)	3	3
	.75-1.9 GeV/c	M	1882±40			1902±11	2032±6		
	N=50	Γ	222±150			52±25	160±16		
	$\chi^2/\text{DF} = 418/390$	t	-.11±.03			-.08±.02	.21±.01		
BERTHON 70	$\Lambda\pi^0$	N_{par}	4	8	4	8+R(4)	R(3)	4	---
	1.134-1.84 GeV/c	M				1940±20	2030±10		
	N=39	Γ				60±20	165±30		
	$\chi^2/\text{DF} = 319/295$	t				-.10±.02	.20±.02		
^b COX 70	$\Lambda\pi^-$	N_{par}	6	3+R(3)	3	3+R(3)	R(3)	3	3
	1.18 -1.71 GeV/c	M		^b 2400±20		1903±10	2027±6		
	N=40	Γ		87±20		77±27	158±16		
	$\chi^2/\text{DF} = 258/199$	t		-.16±.03		-.09±.02	.19±.01		
^b LITCHFIELD 70	$\Lambda\pi^-, \Lambda\pi^0$	N_{par}	5+R(3)	4+R(3)	4+R(3)	5+R(3)	R	7	---
	1.0 -1.85 GeV/c	M	1920±30	2070±30	1940±30	1895±10	2022±4		
	N=54	Γ	170±40	250±40	280±40	70±15	170±15		
	$\chi^2/\text{DF} = 705/636$	t	-.14±.03	-.09±.03	-.14±.03	-.070±.015	.200±.008		
^c GALTIERI 70	$\Lambda\pi^0$	N_{par}	4+R(4)	4	R(4)	4+R(4)	R(3)	R(3)	R(1)
	1.0 -1.85 GeV/c	M	1950±50		1940±50	1905±30	2010±15	2060±20	[2250]
	N=38	Γ	200±50		200±50	70±20	115±15	70±30	[140]
	$\chi^2/\text{DF} = 299/245$	t	-.09±.04		-.12±.04	-.11±.03	.16±.03	-.07±.02	-.18

- a. The D_{15} state $\Sigma(1765)$ was also included in the analysis; the following parameters were found: $M = 1775 \pm 7$, $\Gamma = 146 \pm 9$, $t = -.266 \pm .017$.
- b. Parameters reported here are the ones of fit 13, except for the P_{13} where mass and errors have been estimated by us, examining the various fits reported by the authors. The other errors are only statistical.
- c. The results quoted are the ones from fit A.
- d. $\Sigma(2250)$, G_{19} is really outside the range of most analysis, but its lower tail can be included instead of a G_{19} background.
- e. Quantities in square brackets have been kept fixed. This paper also reports a fit from 0.61 to 1.85 GeV/c.

Table II. Summary of the results obtained for eleven states above 1800 MeV. For the $\bar{K}N$ channel, the elasticity x is listed. For the other channels, the quantity listed is t , the amplitude at resonance ($t = \sqrt{x_0 x}$). The last column is our rating for the resonance.

		M (MeV)	Γ (MeV)	$\bar{K}N$ x	$\Lambda\pi$ t	$\Sigma\pi$ t	ΞK t	Status
$\Lambda(1870)$	$F_{07}^?$	1870	40	.10				Poor
$\Lambda(2015)$	F_{07}	2020	160			-.15		Wait
$\Lambda(2040)$	D_{03}	2010	130			-.20 (?) ^a		Wait
$\Lambda(2100)$	G_{07}	2100	60-145	.29		.06 seen ^b		Good
$\Sigma(1900)$	P_{11}	~1900	200		-.11			Poor
$\Sigma(1915)$	F_{15}	1905	60	.10	-.09	-.008		Fair
$\Sigma(1940)$	D_{13}	1940	200		-.13	-.12 (?) ^a		Wait
$\Sigma(2020)$	F_{17}	2020	100-170	.1-.25	.19	-.05	.023	Good
$\Sigma(2070)$	P_{13}	~2100	87-250		-.12			Poor
$\Sigma(2120)$	G_{17}	~2100	~100		-.07	.13		Wait
$^c \Sigma(2250)$	$G_{19}^?$	2040	~160	.05 (-.18) ^c		(.07) ^c		Fair

- a. The analysis of BURGUN 68 suggests a $J = 3/2$ state with either l-spin in this mass region.
- b. The two analysis of the ΞK channel do not agree on this value (see BURGUN 68 and MULLER 69).
- c. This state is really outside the region where the analyses have been done. Here it has been assumed to be a G_{19} state, as suggested by the analysis of DAUBER 66.

Conclusions. Table III is an attempt to evaluate the status of the various Y^{*1} s. The evaluations are of course partly subjective. A blank indicates that there is no corresponding evidence at all. This may mean either that the relevant couplings are small or that the resonance does not really exist. The Baryon Table includes only the well-established resonances. It seems clear, however, that whereas any particular one of the questionable resonances may disappear with the next analysis, there definitely are many new resonances underlying those we are more familiar with.

Changes in format. When determining branching ratios from partial wave analyses, the quantity most accessible is the so-called "amplitude at resonance". This is essentially the diameter of the resonant circle in the Argand plot. For a resonance going from channel 1 to channel 2, this is simply $i\sqrt{x_1 x_2}$

Baryons

For notation, see illustrated key at beginning of data card listings.

Table with 2 columns: W (width) and Y*0(1405) WIDTH (MEV). Rows include ALSTON 41 HBC, ALEXANDER 62 HBC, ALSTON 62 HBC, MUSGRAVE 65 HBC, ENGLER 65 HBC, BIRMINGHAM 66 HBC, GALTIERI 68 HBC.

Table with 2 columns: P1 Y*0(1405) INTO SIGMA PI and DECAY MASSES. Values: 1197+ 139.

REFERENCES -- Y*0(1405). Lists various experiments and authors like ALVAREZ, EBERHARD, GOOD, GRAZIANO, etc.

Λ(1405) EXTRAPOLATION BELOW THRESHOLD

SEE NOTE IN Y*0(1405) PRODUCTION EXPERIMENTS - THE DIFFICULTIES IN EXTRAPOLATING FROM THE PHYSICAL REGION TO THE RESONANCE LOCATION ARE DISCUSSED BY DALITZ ET AL.

Table with 2 columns: M (mass) and Y*0(1405) MASS (MEV). Rows include KIM 65 HRC, SAKITT 65 HBC, KITTEL 66 HBC, KIM 67 HBC, MARTIN 69 HRC, MARTIN 70 RVUE.

Table with 2 columns: W (width) and Y*0(1405) WIDTH (MEV). Rows include KIM 65 HRC, SAKITT 65 HRC, KITTEL 66 HRC, KIM 67 HRC, MARTIN 69 HRC, MARTIN 70 RVUE.

REFERENCES -- Y*0(1405) FROM EXTRAPOLATIONS. Lists authors like J K KIM, DAY, GLASSER, SEFMAN, FRIEDMAN, etc.

Table with 2 columns: W (width) and Y*0(1405) WIDTH (MEV). Rows include ABRAMS 65 PR 139 R454, DALITZ 67 PR 153 1617, DONALD 66 PR 22 711, KADYK 65 PR 17 1592.

END -EXTRAPOLATION BELOW THRESHOLD-

Λ(1520)

D03

PRODUCTION AND FORMATION EXPERIMENTS AGREE QUITE WELL WITH EACH OTHER, THEREFORE THEY HAVE NOT BEEN SEPARATED FOR THIS PARTICLE

Table with 2 columns: M (mass) and Y*0(1520) MASS (MEV). Rows include WATSON 63 HBC, GALTIERI 63 HBC, ALMEIDA 64 HBC, MUSGRAVE 65 HBC, BURKHARDT 69 HBC, KIM 71 DPWA.

Table with 2 columns: W (width) and Y*0(1520) WIDTH (MEV). Rows include WATSON 63 HBC, MUSGRAVE 65 HBC, BIRMINGHAM 66 HBC, BURKHARDT 69 HBC, KIM 71 DPWA.

Table with 2 columns: P1 Y*0(1520) PARTIAL DECAY MODES and DECAY MASSES. Lists various decay channels like INTO KBAR N, INTO SIGMA PI, INTO LAMBDA PI PI, etc.

Table with 2 columns: W (width) and Y*0(1520) PARTIAL WIDTHS (MEV). Rows include WATSON 63 HBC, DAHLER 67 HBC, DAHL 67 HBC, SCHEUER 68 HRC, BURKHARDT 69 HBC.

Table with 2 columns: R1 Y*0(1520) INTO (SIGMA PI)/(KBAR N) and (P2)/(P1). Rows include MUSGRAVE 65 HBC, DAHLER 67 HBC, DAHL 67 HBC, SCHEUER 68 HRC, BURKHARDT 69 HBC.

Table with 2 columns: R2 Y*0(1520) INTO (LAMBDA PI PI)/(KBAR N) and (P3)/(P1). Rows include DAHLER 67 HBC, DAHL 67 HBC, SCHEUER 68 HRC, BURKHARDT 69 HBC, KIM 71 DPWA.

Table with 2 columns: R3 Y*0(1520) INTO (SIGMA PI)/(LAMBDA PI PI) and (P2)/(P3). Rows include ARMENTEROS 65 HBC, BIRMINGHAM 66 HBC, UHLIG 67 HBC.

Table with 2 columns: R4 Y*0(1520) INTO (LAMBDA GAMMA)/TOTAL (PERCENT) and (P4)/TOTAL. Rows include MAST 68 HBC.

Table with 2 columns: R5 Y*0(1520) INTO (SIGMA GAMMA)/TOTAL (PERCENT) and (P5)/TOTAL. Rows include MAST 68 HBC.

Table with 2 columns: R6 Y*0(1520) INTO (KBAR N)/TOTAL and (P1)/TOTAL. Rows include GALTIERI 69 HBC, KIM 71 DPWA.

Table with 2 columns: R7 Y*0(1520) INTO (SIGMA PI)/TOTAL and (P2)/TOTAL. Rows include GALTIERI 69 HBC, KIM 71 DPWA.

Table with 2 columns: R8 Y*0(1520) INTO (SIGMA PI PI)/TOTAL and (P6)/TOTAL. Rows include GALTIERI 69 HBC.

Table with 2 columns: R9 Y*0(1520) INTO (Y*(1385)+PI)/(LAMBDA PI PI) and (P7)/(P3). Rows include BURKHARDT 70 HBC, CLINE 70 HBC.

Table with 2 columns: R10 Y*0(1520) INTO (Y*(1385)+PI)/TOTAL and (P7)/TOTAL. Rows include CHAN 70 HRC.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS. The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_i δP_i), while the off-diagonal elements are the normalized correlation coefficients (δP_i δP_j) / (δP_i δP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of branching fractions P1 to P6. Values include 0.458±0.010, -0.794, -0.289, -0.063, -0.158, -0.067, etc.

REFERENCES -- Y*0(1520). Lists various experiments and authors like WATSON, GALTIERI, ALMEIDA, MUSGRAVE, BURKHARDT, ARMENTEROS, BIRMINGHAM, DAHLER, DAHL, UHLIG, SCHEUER, BURKHARDT, GALTIERI, CLINE, KIM, CHAN, BURKHARDT.

Table with 2 columns: W (width) and Y*0(1520) WIDTH (MEV). Rows include BIRMINGHAM 66 PR 152 1148, DAHLER 67 PR 163 1377, DAHLER 67 PR 248 925, UHLIG 67 PR 155 1448, MAST 68 PR 21 1715, SCHEUER 68 PR 88 503.

Table with 2 columns: W (width) and Y*0(1520) WIDTH (MEV). Rows include BURKHARDT 69 NP B14 106, GALTIERI 69 LUND PAPER 91, CLINE 70 LNC 2 407, KIM 71 HARVARD PREPRINT, CHAN 70 PREPRINT, BURKHARDT TO CERN 70/58 SUR.NP.

Baryons

For notation, see illustrated key at beginning of data card listings.

$\Lambda(1670)$

40 Y*0(1670, JP=1/2-1) I=0
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THIS RESONANCE IS WELL ESTABLISHED.
 (SEE THE NOTE FOR THE Y*0(1330)).

S_{01}

40 Y*0(1670) MASS (MEV)
 M (1666.0)OR(1675.0) BERLEY 65 HBC 0 K-P TO LAM ETA 7/66
 M THE FIRST VALUE ASSUMES THE BRANCHING RATIO INTO LAMBDA ETA IS SMALL, THE SECOND THAT IT IS LARGE, BECAUSE THE RESONANCE IS NEAR THE LAMBDA ETA THRESHOLD, THE BRANCHING RATIO AFFECTS THE MOMENTUM DEPENDENCE OF THE TOTAL WIDTH, AND THUS ALSO THE RESONANCE PARAMETERS OBTAINED BY FITTING TO THE DATA.
 M N (1673.0) (3.0) ARMENT-1 69 HBC 0 ELASTIC, CH EXCH 11/68
 M N (1676.0) (2.0) ARMENT-2 68 HBC 0 K-P TO SIGMA PI 11/68
 M N (1662.0) (3.0) ARMENT-4 69 HBC 0 ELAST, CH EXC. ED 9/69
 M N (1680.0) (1.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 9/69
 M A (1674.0) (5.0) ARMENT-3 69 HBC 0 MULTICHANNEL, 9/69
 M N (1679.0) (5.0) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 M (1683.0) (5.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 M 1670. KIM 71 DPWA K-MATRIX ANAL. 3/71*
 M A THE MULTICHANNEL ANALYSIS INCLUDES ELASTIC AND SIGMA PI. 10/69
 M N THE APPARENT DISCREPANCY BETWEEN THESE RESULTS IS PROBABLY NOT SERIOUS. THE ERRORS GIVEN ARE JUST STATISTICAL. THE SYSTEMATIC ERRORS THAT RESULT FROM THE RESTRICTIVE PARAMETERIZATION FORCED ON THE PARTIAL-WAVE AMPLITUDES ARE NOT INCLUDED, AND CAN BE LARGE.

40 Y*0(1670) WIDTH (MEV)
 W M (22.0)OR(15.0) BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
 W N (24.0) (1.0) ARMENT-1 69 HBC 0 SEE NOTE M ABOVE 11/68
 W A (26.0) (7.0) ARMENT-2 68 HBC 0 SEE NOTE M ABOVE 11/68
 W A (24.0) (3.0) ARMENT-3 69 HBC 0 9/69
 W N (34.0) (15.0) ARMENT-4 69 HBC 0 ELAST, CH EXC. ED 9/69
 W N (31.0) (5.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 9/69
 W (31.0) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 W (25.0) (5.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 W 35. KIM 71 DPWA K-MATRIX ANAL. 3/71*
 SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

40 Y*0(1670) PARTIAL DECAY MODES
 P1 Y*0(1670) INTO KBAR N 497+ 939
 P2 Y*0(1670) INTO LAMBDA ETA 115+ 548
 P3 Y*0(1670) INTO SIGMA PI 1189+ 139

40 Y*0(1670) BRANCHING RATIOS
 BELOW X1 = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1 Y*0(1670) INTO (KBAR N)/TOTAL X1
 R1 P (0.14) (0.04) ARMENT-1 69 HBC 0 OLD DATA 11/68
 R1 P (0.14) (0.04) ARMENT-4 69 HBC 0 NEW DATA 9/69
 R1 A (0.39) (0.05) ARMENT-3 69 HBC 0 9/69
 R1 A (0.39) (0.05) CONFORTO 70 HBC 0 K-P, ELAST, CEX 6/70
 R1 A (0.28) KIM 71 DPWA K-MATRIX ANAL. 3/71*
 R1 A EFFECT BELOW REGION ANALYSED. VALUE OF .18 DOES NOT AFFECT FIT OR VALUES OF OTHER PARAMETERS
 R1 P THIS IS THE DIAMETER OF THE CIRCLE IN THE ARGAND PLOT. IT IS SUPERIMPOSED ON A LARGE BACKGROUND.

R2 Y*0(1670) INTO (KBAR N)(LAMBDA ETA)/TOTAL**2 (P1*P2)/TOTAL**2
 R2 M (0.039)OR(0.053) BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
 R2 (0.064) ARMENT-3 69 HBC 0 9/69
 R2 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71*
 SEE THE NOTES ACCOMPANYING MASSES QUOTED

R3 Y*0(1670) FROM KBAR N TO SIGMA PI SQRT(X1*X3)
 R3 (-0.25) (0.06) ARMENT-2 68 HBC 0 OLD DATA 9/69
 R3 (-0.27) ARMENT-3 69 HBC 0 9/69
 R3 (-0.30) ARMENT-4 69 HBC 0 NEW DATA 6/70
 R3 (-0.27) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 R3 (-0.29) (0.03) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 R3 (-0.39) KIM 71 DPWA K-MATRIX ANAL. 3/71*

REFERENCES -- Y*0(1670)
 BERLEY 65 PRL 15 641 (CONNOLLY, HART, RAUM, STONEHILL, + (BNL) IJP
 ARMENT-1 68 NP 18 195 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-2 68 NP 18 223 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-3 69 LUND PAPER 229 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-4 69 NP 114 91 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 BERLEY 69 PL 308 430 + HART, RAHM, WILLIS, YAMAMOTO (BNL) IJP
 CONFORTO 70 EPL 70-2R(SUB NP) +HARSEN+LASINSKI++ (EFI+HEID) IJP
 GALTIERI 70 DIKE 173 A. BARBARO GALTIERI (ILL) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DIKE 161 J. K. KIM (HARV) IJP

PAPERS NOT REFERRED TO IN DATA CARDS
 BIRMING - 66 PR 152 1148 (BIRMINGHAM, GLASGOW, IMPCOL, OXFORD, RUTHERF) I
 LEVISETT 69 LUND 339 R LEVI SETTI (RAPPORTEUR) (CHICAGO)

$\Lambda(1690)$

45 Y*0(1690, JP=3/2-1) I=0
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THIS RESONANCE IS WELL ESTABLISHED.

D_{03}

45 Y*0(1690) MASS (MEV)
 M (1695.0) (4.0) BUGG 68 CNTR 0 K-P, D TOTAL 7/68
 M (1696.0) (3.0) ARMENT-1 68 HBC 0 ELASTIC, CH EXCH 11/68
 M (1681.0) (2.0) ARMENT-3 68 HBC 0 K-P TO SIGMA PI 11/68
 M (1681.0) (8.0) BARTLEY 68 DBC 0 K-P AND K-D DATA 11/68
 M N (1697.0) (2.0) CONFORTO 68 HBC 0 ELASTIC, CH EXCH 11/68
 M (1701.0) (4.0) BERTANZA 69 HBC 0 ELASTIC, CH EXCH 9/69
 M A (1691.0) (2.0) ARMENT-4 69 HBC 0 ELAS, CH EXC. ED 9/69
 M A (1688.0) (2.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 9/69
 M (1689.0) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 M (1688.0) (3.0) CONFORTO 70 HBC 0 K-P, ELAST, CEX 6/70
 M (1680.0) (5.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 M 1690. KIM 71 DPWA K-MATRIX ANAL. 3/71*

M M THE Y*0(1690) IS AT THE EDGE OF THE ENERGY REGION ANALYZED BY CONFORTO. THE SAME DATA AS WELL AS OTHERS EXTENDING TO LOWER ENERGIES ARE INCLUDED IN ARMENTEROS I.
 M A ANALYSIS INCLUDES OLD AND NEW DATA OF CHS COLLAB. 43-8 GEV/C 10/69
 M A THE APPARENT DISCREPANCY BETWEEN THE SIGMA PI AND OTHER RESULTS IS PROBABLY NOT SERIOUS. THE ERRORS GIVEN ARE JUST STATISTICAL. THE A SYSTEMATIC ERRORS THAT RESULT FROM THE RESTRICTIVE PARAMETERIZATION OF THE PARTIAL-WAVE AMPLITUDES ARE NOT INCLUDED, AND CAN BE LARGE.

45 Y*0(1690) WIDTH (MEV)
 W (35.0) (7.0) ARMENT-1 68 HBC 0 OLD DATA 11/68
 W (85.0) (7.0) ARMENT-3 68 HBC 0 OLD DATA 11/68
 W 40.0 (7.0) RUGG 68 CNTR 0 7/68
 W (15.0) (15.0) BARTLEY 68 DBC 0 K-P AND K-D DATA 11/68
 W M (27.0) (5.0) CONFORTO 68 HBC 0 SEE NOTE M ABOVE 11/68
 W A (31.0) (7.0) ARMENT-4 69 HBC 0 ELAS, CH EXC. ED 9/69
 W A (72.0) (4.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 9/69
 W 29.0 (6.0) BERTANZA 69 HBC 0 9/69
 W 57.0 BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 W 85.0 (10.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 W 64.0 (5.0) CONFORTO 70 HBC 0 K-P, ELAST, CEX 6/70
 W K. KIM 71 DPWA K-MATRIX ANAL. 3/71*
 SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

45 Y*0(1690) PARTIAL DECAY MODES
 P1 Y*0(1690) INTO KBAR N 497+ 939
 P2 Y*0(1690) INTO LAMBDA PI 1189+ 139
 P3 Y*0(1690) INTO SIGMA PI 1189+ 139+ 139

45 Y*0(1690) BRANCHING RATIOS
 THE SUM OF ALL THE QUOTED BRANCHING RATIOS IS MORE THAN 1.0. THE TWO-BODY RATIOS ARE FROM PARTIAL WAVE ANALYSES, AND THUS PROBABLY ARE MORE RELIABLE THAN THE THREE-BODY RATIOS, WHICH ARE DETERMINED FROM BUMPS IN CROSS SECTIONS. OF THE LATTER, THE SIGMA PI PI BUMP LOOKS MORE SIGNIFICANT (THE ERROR GIVEN FOR THE LAMBDA PI PI RATIO LOOKS UNREASONABLY SMALL). HARDLY ANY OF THE SIGMA PI PI DECAY CAN BE VIA Y*(1385), FOR THEN NINE TIMES AS MUCH LAMBDA PI PI DECAY WOULD BE REQUIRED.

BELOW X1 = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1 Y*0(1690) INTO (KBAR N)/TOTAL X1
 R1 (0.23) BUGG 68 CNTR 0 ASSUMING J=3/2 7/68
 R1 (0.18) (0.03) ARMENT-1 68 HBC 0 11/68
 R1 M (0.22) (0.03) CONFORTO 68 HBC 0 SEE NOTE M ABOVE 11/68
 R1 (0.28) (0.04) BERTANZA 69 HBC 0 9/69
 R1 (0.18) (0.02) ARMENT-4 69 HBC 0 NEW DATA 9/69
 R1 N (0.34) (0.02) CONFORTO 70 HBC 0 K-P, ELAST, CEX 6/70
 R1 N (0.22) KIM 71 DPWA K-MATRIX ANAL. 3/71*
 R1 N EFFECT IS AT END OF REGION ANALYZED. THIS COULD AFFECT VALUE OF X1.

R2 Y*0(1690) FROM KBAR N TO SIGMA PI SORT(X1*X2)
 R2 (-0.33) (0.02) ARMENT-3 68 HBC 0 OLD DATA 11/68
 R2 (-0.36) (0.02) ARMENT-4 69 HBC 0 NEW DATA 9/69
 R2 (-0.27) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
 R2 (-0.31) (0.03) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
 R2 (-0.40) KIM 71 DPWA K-MATRIX ANAL. 3/71*

R3 Y*0(1690) FROM KBAR N TO LAMBDA PI PI SORT(X1*X3)
 R3 0.25 (0.02) BARTLEY 68 HBC 0 K-N TO LAM PI PI 11/68

R4 Y*0(1690) FROM KBAR N TO SIGMA PI PI SORT(X1*X4)
 R4 (0.21) ARMENT-2 68 HBC 0 K-N TO SIG PI PI 11/68

REFERENCES -- Y*0(1690)
 DAVIES 67 PRL 18 62 +DOWELL, + (BRNGMH, CVNSH, RTHFRD) I
 -- REPLACED BY BUGG 68.
 ARMENT-1 68 NP 18 195 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-2 68 NP 18 216 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-3 68 NP 18 223 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 BARTLEY 68 PRL 21 1111 +CHU, DDMO, GREENE, + (TUFTS, FLOR, ST. BRANDEIS) I
 RUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (KTHFRD, BRNGMH, CVNSH) I
 CONFORTO 68 NP 18 265 +HARSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 ARMENT-4 69 NP 114 91 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 BERTANZA 69 PR 177 2036 +BIGI, CARRARA, CASALI, + (PTSA, BNL, YALE) IJP
 BERLEY 69 PL 308 430 + HART, RAHM, WILLIS, YAMAMOTO (BNL) IJP
 CONFORTO 70 EPL 70-2R(SUB NP) +HARSEN+LASINSKI++ (EFI+HEID) IJP
 GALTIERI 70 DIKE 173 A. BARBARO GALTIERI (ILL) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DIKE 161 J. K. KIM (HARV) IJP

$\Lambda(1750)$

77 Y*0(1750, JP=1/2+1) I=0
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

P_{01}

THE EVIDENCE FOR THIS STATE IS SOMEWHAT CONFUSED. IT WAS FIRST SUGGESTED IN A PARTIAL WAVE ANALYSIS OF KBAR N DATA BY THE BEHAVIOR OF THE POI AMPLITUDE WHEN IT WAS PARAMETERIZED AS A TWO-STRAIGHT-LINE BACKGROUND. WHEN IT WAS REPARAMETERIZED AS A RESONANCE SUPERIMPOSED ON A ONE-STRAIGHT-LINE BACKGROUND, A BROAD RESONANCE RESULTED (ARMENTEROS 69). A REANALYSIS OF ESSENTIALLY THE SAME DATA, BUT THIS TIME WITH THE POI AMPLITUDE UNCONSTRAINED, SUGGESTED A MUCH NARROWER RESONANCE AT HIGHER ENERGY (ARMENTEROS 70).

A WIDER AND MORE ELASTIC POI RESONANCE AT ABOUT THE SAME MASS IS SUGGESTED BY THE ANALYSIS OF BAILEY 69. THIS USES CONSIDERABLY LESS DATA THAN THE ARMENTEROS ANALYSES. FOR THIS REASON WE DO NOT QUOTE ANY PARAMETERS FOR THE OTHER PARTIAL WAVES OBTAINED IN THIS ANALYSIS.
 ARMENTEROS 70, GALTIERI 70, AND KIM 71 PRESENT EVIDENCE FOR A POI STATE IN THE SIGMA PI CHANNEL. IN ADDITION THE ANALYSIS OF KIM 71 INDICATES A SECOND POSSIBLE POI STATE AT ABOUT 1570. WE TENTATIVELY LIST THESE EFFECTS TOGETHER.

Baryons

For notation, see illustrated key at beginning of data card listings.

Table with columns for particle name, mass (MEV), and various decay channels. Includes entries for Armentero, Bailey, Galtieri, and Kim.

Table with columns for particle name, width (MEV), and various decay channels. Includes entries for Armentero, Bailey, Galtieri, and Kim.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero and Galtieri.

Table with columns for particle name, branching ratios, and various decay channels. Includes entries for Armentero, Bailey, Galtieri, and Kim.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero, Galtieri, and Kim.

Lambda(1815)

F05

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS. THIS STATE IS WELL ESTABLISHED. MOST OF THE QUOTED ERRORS ARE STATISTICAL ONLY. THE SYSTEMATIC ERRORS DUE TO THE PARTICULAR PARAMETERIZATION USED IN THE P.W.A. ARE NOT INCLUDED. FOR THIS REASON WE DO NOT CALCULATE WEIGHTED AVERAGES FOR MASS AND WIDTH.

Table with columns for particle name, mass (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, width (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, branching ratios, and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, mass (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i + delta P_i, where delta P_i = sqrt(delta P_i^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of branching fractions with columns P1, P2, P3, P4 and rows of numerical values.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero, Galtieri, and Kim.

THE FOLLOWING PAPERS ARE NOW OF ONLY HISTORICAL INTEREST -- CHAMBERLAIN, CROWE, KEEFE, KEITH, + (LRL) I

Table with columns for particle name, mass (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THE BEST EVIDENCE FOR THIS RESONANCE COMES FROM THE SIGMA PI CHANNEL. IT APPEARS TO BE WELL ESTABLISHED.

Table with columns for particle name, mass (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, width (MEV), and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Table with columns for particle name, branching ratios, and various decay channels. Includes entries for Armentero, Bugg, Brichman, and Galtieri.

Lambda(1830)

D05

Baryons

For notation, see illustrated key at beginning of data card listings.

56 $\Lambda(1830)$ BRANCHING RATIOS

BELOW, $X_1 = (\text{PARTIAL-WIDTH } 1/\text{TOTAL WIDTH, ETC.})$ A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Ψ MINI-REVIEW.

R1	$\Psi(1830)$ INTO (KBAR N)/TOTAL	X1	
R1	0.09 (0.01)	ARMENFERO 68 HBC 0	ELASTIC, CH EXCH 11/68
R1	0.03 (0.02)	BRICMAN 70 DPWA 0	SIGTOT, ELAS, CHEX 1/71*
R1	0.05 (0.02)	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH 6/70
R1	(0.24)	KIM 71 DPWA	K-MATRIX ANAL. 3/71*

R2 $\Psi(1830)$ FROM KBAR N INTO SIGMA PI SORT(X1*X2)

R2	0.15 (0.02)	ARMENFERO 67 DPWA 0	K-P TO SIGMA PI 8/67
R2	0.19 (0.01)	BELL 67 DPWA 0	K-P TO SIGMA PI 11/67
R2	-0.16 (0.03)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI 7/70
R2	0.15	KIM 71 DPWA	K-MATRIX ANAL. 3/71*

REFERENCES -- $\Psi(1830)$

ARMENFERO 67 PL 248 199 ARMENFERO, F-LUZZI, + (CERN, HEIDEL, SACLAY) IJP
 RELL 67 PRL 19 936 R R RELL (LRL) IJP
 ARMENFERO 68 NP BR 195 ARMENFERO, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 CONFORTO 68 NP BR 245 HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 --- CONFORTO 68 IS SUPERSEDED BY CONFORTO 70.
 BRICMAN 70 PL 338 511 FERRO-LUZZI, LAGNAUX (CERN) IJP
 GALTIERI 70 DPKE CONF 173 A BARBARO-GALTIERI (LRL) IJP
 CONFORTO 70 (UNPUBLISHED) HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV) IJP

60 $\Lambda(1860)$, JP=3/2+ I=0 **P₀₃**

SEE THE MINI-REVIEW AT THE START OF THE Ψ LISTINGS.

THE QUANTUM NUMBERS OF THIS STATE ARE PROBABLY 3/2+.
 AN F07 WAS FIRST SUGGESTED IN THE PHASE SHIFT ANALYSIS OF THE KBAR-DATA BY ARMENFERO 67. THE ISOSPIN-0 TOT. CROSS SECTION HAS A SHOULDER ON THE HIGH SIDE OF $\Psi(1815)$ THAT IS COMPATIBLE WITH A STATE AT THIS MASS (RUGG 68). THE ARMENFERO 68 AND CONFORTO 68 ANALYSES OF IMPROVED KBAR-N DATA INCLUDED THE F07 STATE. HOWEVER IN THE NEW ANALYSES OF CONFORTO 70 AND LATER, BRICMAN 70 TO THE F07 IS NOT REQUIRED BY THE NEW DATA AND A P03 RESONANCE IS ACCEPTABLE. IN ADDITION BOTH THESE NEW ANALYSES INCLUDED RECENT POLARIZATION DATA.

60 $\Lambda(1860)$ MASS (MEV)

M	N	1870.0	5.0	RUGG 68 CNTR 0	K-P TOTAL	7/68
M	A	F07 1854.0	2.0	ARMENFERO 68 DPWA 0	ELASTIC, CH EXCH	11/68
M	A	F07 1877.0	6.0	BRICMAN 70 CNTR 0	TOTAL AND CH EX	6/70
M	N	P03 1870.0	6.0	BRICMAN 70 DPWA 0	SIGTOT, ELAS, CHEX	1/71*
M	N	P03 1881.0	10.0	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH	6/70
M	I	P03 1710.		KIM 71 DPWA	K-MATRIX ANAL.	3/71*

M A THESE TWO ANALYSES GAVE THE F07 ASSIGNMENT. THEY HAVE TO BE DISCARDED IN VIEW OF CONFORTO 70 AND BRICMAN 70
 M N DUE TO PARTICULAR PARAMETERIZATION USED, ERROR CAN BE LARGE
 M I POSSIBLE EFFECT MAINLY IN SIGMA PI. WE TENTATIVELY LIST IT HERE.

60 $\Lambda(1860)$ WIDTH (MEV)

W	N	40.0	10.0	RUGG 68 CNTR 0	K-P TOTAL	7/68
W	A	F07 23.0	7.0	ARMENFERO 68 DPWA 0	ELASTIC, CH EXCH	11/68
W	A	F07 24.0	15.0	BRICMAN 70 CNTR 0	TOTAL AND CH EX	6/70
W	N	P03 37.0	10.0	BRICMAN 70 DPWA 0	SIGTOT, ELAS, CHEX	1/71*
W	N	P03 80.0	20.0	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH	6/70
W	I	P03 20.		KIM 71 DPWA	K-MATRIX ANAL.	3/71*

SEE THE NOTES ACCOMPANYING MASSES QUOTED

60 $\Lambda(1860)$ PARTIAL DECAY MODES

P1	$\Psi(1860)$ INTO KBAR N	DECAY MASSES
P2	$\Psi(1860)$ INTO SIGMA PI	497+ 939 1199+ 139

60 $\Lambda(1860)$ BRANCHING RATIOS

BELOW, $X_1 = (\text{PARTIAL-WIDTH } 1/\text{TOTAL WIDTH, ETC.})$ A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Ψ MINI-REVIEW.

R1	$\Psi(1860)$ INTO (KBAR N)/TOTAL	X1	
R1	(J+1/2)X = 0.40	RUGG 68 CNTR 0	7/68
R1	A F07 0.12 0.02	ARMENFERO 68 HBC 0	ELASTIC, CH EXCH 11/68
R1	A F07 0.07 0.02	BRICMAN 70 CNTR 0	TOTAL AND CH EX 6/70
R1	P03 0.14 0.02	BRICMAN 70 DPWA 0	SIGTOT, ELAS, CHEX 1/71*
R1	N P03 0.25 0.03	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH 6/70

SEE THE NOTES ACCOMPANYING MASSES QUOTED

R2 $\Psi(1860)$ INTO SIGMA PI (P2)

R2	P	PROBABLY SEEN	GALTIERI 68 DRC 0	K-N TO SIG PI PI 11/68
R2	P	POSSIBLY THIS BUMP SPEN AT 1840-10 MEV WITH A WIDTH OF 35+10 MEV		
R2		IS THE $\Psi(1830)$, WHICH DECAYS STRONGLY TO SIGMA PI. HOWEVER THE NARROW WIDTH HERE ARGUES FOR ITS BEING THE $\Psi(1860)$.		

REFERENCES -- $\Psi(1860)$

ARMENFERO 68 NP BR 195 ARMENFERO, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 RUGG 68 NP BR 148 1446 GILMORE, KNIGHT, + (ETHRO, BRANNGH, CVOSHI) I
 GALTIERI 68 PRL 21 573 BARBARO-GALTIERI, MATISON, + (LRL, SLAC)
 BRICMAN 70 PL 318 152 FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
 BRICMAN 70 PL 338 511 FERRO-LUZZI, LAGNAUX (CERN)
 CONFORTO 70 EPI 70-43 HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV) IJP

PAPERS NOT REFERRED TO IN DATA CARDS

ARMENFERO 67 NP BR 592 ARMENFERO, F-LUZZI, + (CERN, HEIDEL, SACLAY) IJP
 --- ARMENFERO 67 IS REPLACED BY ARMENFERO 68 AND CONFORTO 68
 CONFORTO 68 NP BR 245 HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 --- CONFORTO 68 IS SUPERSEDED BY CONFORTO 70.
 LEVISTFT 69 LUND 339 R. LEVI SETTI (RAPPORTEUR) (EPFINS)

$\Lambda(1870)$ **S₀₁**

36 $\Psi(1870)$, JP=1/2- I=0

THE S01 AMPLITUDE SHOWS A SECOND RESONANCE BEHAVIOR AT ABOUT 1800 MEV IN BOTH THE BRICMAN 70 AND KIM 71 ANALYSES. THE ELASTICITIES OF THESE TWO ANALYSES ARE IN RAD DISAGREEMENT.

36 $\Psi(1870)$ MASS (MEV)

M	N	(1872.0)	(10.0)	BRICMAN 70 DPWA	TOT. ELAS. CHEX	1/71*
M	N	(1780.)		KIM 71 DPWA	K-MATRIX ANAL.	3/71*

36 $\Psi(1870)$ WIDTH (MEV)

W	N	(100.0)	(20.0)	BRICMAN 70 DPWA	TOT. ELAS. CHEX	1/71*
W	N	(40.)		KIM 71 DPWA	K-MATRIX ANAL.	3/71*

36 $\Psi(1870)$ PARTIAL DECAY MODES

P1	$\Psi(1870)$ INTO KBAR N	DECAY MASSES
P2	$\Psi(1870)$ INTO SIGMA PI	497+ 939 1197+ 139

36 $\Psi(1870)$ BRANCHING RATIOS

R1	$\Psi(1870)$ INTO (KBAR N)/TOTAL	X1	
R1	(0.19) (0.02)	BRICMAN 70 DPWA	TOT. ELAS. CHEX 1/71*
R1		KIM 71 DPWA	K-MATRIX ANAL. 3/71*
R2	$\Psi(1870)$ FROM KBAR N TO SIGMA PI	SORT(X1*X2)	
R2	(0.24)	KIM 71 DPWA	K-MATRIX ANAL. 3/71*

REFERENCES -- $\Psi(1870)$

BRICMAN 70 PL 338 511 C BRICMAN, M FERRO-LUZZI, J P LAGNAUX(CERN) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV) IJP

89 $\Psi(2010)$, JP=3/2- I=0 **D₀₃**

SEE THE MINI-REVIEW AT THE START OF THE Ψ LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY ONE BUT NOT OTHER PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

89 $\Psi(2010)$ MASS (MEV)

M	N	(2010.0)	(30.0)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI	7/70
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89 $\Psi(2010)$ WIDTH (MEV)

W	N	(130.0)	(50.0)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI	7/70
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89 $\Psi(2010)$ PARTIAL DECAY MODES

P1	$\Psi(2010)$ INTO KBAR N	DECAY MASSES
P2	$\Psi(2010)$ INTO SIGMA PI	497+ 939 1197+ 139

89 $\Psi(2010)$ BRANCHING RATIOS

BELOW, $X_1 = (\text{PARTIAL-WIDTH } 1/\text{TOTAL WIDTH, ETC.})$ A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Ψ MINI-REVIEW.

R1	$\Psi(2010)$ FROM KBAR N TO SIGMA PI	SORT(X1*X2)	
R1	(-0.20) (0.04)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI 7/70

REFERENCES -- $\Psi(2010)$

GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJP

$\Lambda(2020)$ **F₀₇**

27 $\Psi(2020)$, JP=7/2+ I=0

SEE THE MINI-REVIEW AT THE START OF THE Ψ LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY ONE BUT NOT OTHER PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

27 $\Psi(2020)$ MASS (MEV)

M	N	(2020.0)	(20.0)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI	7/70
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27 $\Psi(2020)$ WIDTH (MEV)

W	N	(160.0)	(30.0)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI	7/70
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27 $\Psi(2020)$ PARTIAL DECAY MODES

P1	$\Psi(2020)$ INTO KBAR N	DECAY MASSES
P2	$\Psi(2020)$ INTO SIGMA PI	497+ 939 1197+ 139

27 $\Psi(2020)$ BRANCHING RATIOS

BELOW, $X_1 = (\text{PARTIAL-WIDTH } 1/\text{TOTAL WIDTH, ETC.})$ A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Ψ MINI-REVIEW.

R1	$\Psi(2020)$ FROM KBAR N TO SIGMA PI	SORT(X1*X2)	
R1	(-0.15) (0.02)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI 7/70

REFERENCES -- $\Psi(2020)$

GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJP

Baryons

For notation, see illustrated key at beginning of data card listings.

$\Lambda(2100)$

41 $\gamma^0(2100, JP=7/2-1) I=0$ **G07**
SEE THE MINI-REVIEW AT THE START OF THE γ^* LISTINGS.

THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES. PARAMETERS OF PEAKS SEEN IN CROSS-SECTIONS AND INVARIANT-MASS DISTRIBUTIONS AROUND 2100 MEV ARE GIVEN IN THE NEXT ENTRY. EVENTUALLY THE PARTIAL-WAVE ANALYSES SHOULD GIVE THE BEST RESULTS, AS THEY ISOLATE THE G07 WAVE. THIS SUPERIORITY IS, HOWEVER, PROBABLY NOT YET ATTAINED, AND WE RELY ON BOTH ENTRIES FOR PARAMETERS GIVEN IN THE MAIN BARYON TABLE.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes entries for M, W, and P1-P4.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes entries for W, W, W, W.

Table with columns for partial decay modes and decay masses. Includes entries for P1, P2, P3, P4.

41 $\gamma^0(2100)$ BRANCHING RATIOS
BELOW, $X_1 = (\text{PARTIAL-WIDTH})/(\text{TOTAL WIDTH, ETC.})$. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE γ^* MINI-REVIEW.

Table with columns for branching ratios and decay masses. Includes entries for R1, R2, R3, R4.

SEE THE NOTES ACCOMPANYING MASSES QUOTED

REFERENCES -- $\gamma^0(2100)$
WOHL 66 PRL 17 107 C G WOHL, F T SOLMITZ, M L STEVENSON (LRL)IJP
TRIPP 67 NP B3 10 LEITH, + (LRL, SLAC, CERN, HEIDEL, SACLAY)
DAUM 68 NP B7 19 ERNE, LAGNAUX, SEFS, STEUER, UDD (CERN)JP
BURGUN 69 NP BR 457 MEYER, PAULI, + (SACLAY, COLFRANCE, RTHFD)
MULLER 69 THESIS, UCLR 1972 R A MULLER (LRL)
GALTIERI 70 DIKE CNF 173 A BARRARD-GALTIERI (LRL)IJP
BERTHONI 70 NP B24 417 VRANA, RUTTERWORTH, + (CDEF, RHEL, SACLAY)IJP

$\Lambda(2110)$

35 $\gamma^0(2110, JP=5/2+1) I=0$ **F05**

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes entries for M, W, and P1-P4.

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REFERENCES -- $\gamma^0(2110)$
BERTHONI 70 NP B24 417 VRANA, RUTTERWORTH, + (CDEF, RHEL, SACLAY)IJP

M > 2100 MEV - PRODUCTION AND TOTAL EXPERIMENTS

BUMPS

$\Lambda(2100)$

25 $\gamma^0(2100)$ CROSS-SECTION AND INVARIANT-MASS PEAKS
SEE THE MINI-REVIEW AT THE START OF THE γ^* LISTINGS.

SEE THE NOTE TO THE G07 $\gamma^0(2100)$, WHICH PRECEDES THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE $\gamma^0(2100)$, BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes entries for M, W, and P1-P4.

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Table with columns for mass (MEV), width (MEV), and partial decay modes. Includes entries for M, W, and P1-P4.

Baryons

For notation, see illustrated key at beginning of data card listings.

$\Lambda(2585)$

7 Y*(0)2585, JP= 1 I=0
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

7 Y*(0)2585 MASS (MEV)

M	2585.0 (2530.0)	45.0	ABRAMS	70 CNTR	K-P, D TOTAL	10/70*
			LU	70 CNTR 0	GAMMA P TO K+ Y*	1/71*

7 Y*(0)2585 WIDTH (MEV)

W	(300.0)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*
	(150.0)		LU	70 CNTR 0	GAMMA P TO K+ Y*	1/71*

7 Y*(0)2585 PARTIAL DECAY MODES

PI Y*(0)2585 INTO KBAR N
DECAY MASSES
497+ 999.

7 Y*(0)2585 BRANCHING RATIOS

RI Y*(0)2585 INTO KBAR N/TOTAL
J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*X1

RI	(1.0)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*
RI C	(0.12)	(0.12)	BRICHMAN	70 CNTR	TOTAL AND CH EX	10/70*

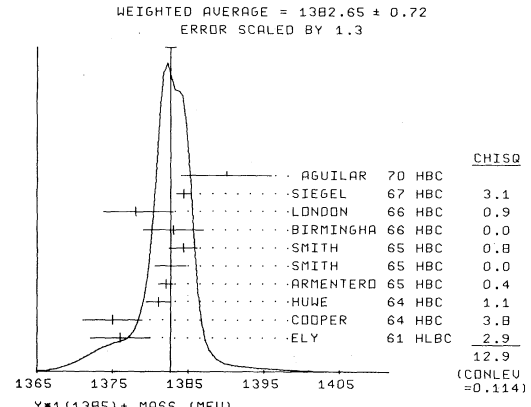
RI C FIT OF TOTAL CROSS SECTION GIVEN BY BRICHMAN IN THIS REGION
RI C IS PDDR.

REFERENCES -- Y*(0)2585)

ARRAMS 70 PR 10 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
BRICHMAN 70 PL 31B 152 +FERRO LUZZI, PERROFINI, + (CERN/CNRS/SACLAY)
LU 70 PR 02 1946 +GREENBERG, HUGHES, MINEHART, MORI, + (YALE)

PAPERS NOT REFERRED TO IN DATA CARDS

COOL 66 PRL 16 122R +GIACOMELLI, KYCIA, LEONTIC, LUNDBY + (BNL) I



Σ^+

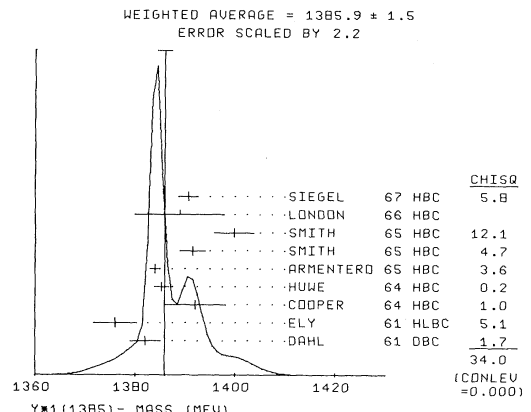
19 SIGMA + (1199, JP=1/2+) I=1
SEE LISTINGS OF STABLE PARTICLES

Σ^-

20 SIGMA - (1198, JP=1/2+) I=1
SEE LISTINGS OF STABLE PARTICLES

Σ^0

21 SIGMA 0 (1193, JP=1/2+) I=1
SEE LISTINGS OF STABLE PARTICLES



$\Sigma(1385)$

43 Y*(1)1385, JP=3/2+ I=1 **P13**
FOR DISCUSSION OF INCONSISTENCY OF ERRORS AND OUR MODIFICATIONS, SEE NOTE ON K*(890)

FOR THE TABLES WE USE ONLY THE UNSTARRED DATA, WHICH ATTEMPTS TO OBTAIN THE SEPARATE CHARGE-STATE MASSES AND WIDTHS. SEE HOWEVER THE IDEOGRAMS INSERTED IN LISTING THESE INDICATE SERIOUS SYSTEMATICS, PERHAPS ARISING FROM INTERFERENCE EFFECTS THAT CHANGE WITH PRODUCTION MECHANISM AND BEAM MOMENTUM.

43 Y*(1)1385 MASS (MEV)

M	141(1384.0)		ALSTON	60 HBC	+ K-P 1.15 BEV/C	
	331(1384.0)		MADOTIN	61 HBC	0+ K20 P 1.0R BEV/C	
	(1395.0)		BERGE	61 HBC	+ K-P 1.4-1.85 BEV/C	
	(1392.0)	(7.0)	COLLEY	62 HLBC	0-PI- PRP 2. BEV/C	
MD	106(1391.0)	(4.0)	CURTIS	63 DSPK	0 PI-P 1.5 BEV/C	
	(1392.0)	(10.0)	MUSGRAVE	65 HBC	+ OPBAR P 3.4 BEV/C	7/66
	(1385.0)	(3.0)	BALTAY	65 HBC	+ PBAR P 3.7 BEV/C	7/66

M+ F 154 1376.0 3.0 ELY 61 HLBC + K-P 1.11 BEV/C
ERROR OF 3.0 ENLARGED TO 3.0 BY US, BECAUSE LT STATIST. ERR. 10/69

M+ 170 1374.0 3.0 COOPER 64 HBC + K-P 1.45 BEV/C
M+ 859 1381.0 1.6 HUWE 64 HBC + K-P 1.22 BEV/C
M+ 750 1382.0 1.0 ARMENTERO 65 HBC + K-P 1.9-1.2 BEV/C
M+ S 250 1382.6 2.1 SMITH 65 HBC + K-P 1.95 BEV/C 9/66
M+ S 250 1384.3 1.9 SMITH 65 HBC + K-P 1.8 BEV/C 9/66
M+ S ERROR OF 1.4 ENLARGED TO 2.1 BY US, BECAUSE LT STATIST. ERR. 10/69
M+ S ERROR OF 1.1 ENLARGED TO 1.9 BY US, BECAUSE LT STATIST. ERR. 10/69
M+ B 40 1383.0 4.0 BIRMINGHAM 66 HBC + 3.5 K-P 9/67
M+ B ERROR OF 2.0 ENLARGED TO 4.0 BY US, BECAUSE LT STATIST. ERR. 10/69
M+ 1378.0 5.0 LONDON 66 HBC + K-P 2.24 BEV/C 7/66
M+ 1260 1384.4 1.0 SIEGEL 67 HBC + K-P AT 2.1 GEV/C 10/69
M+ 1390.0 6.0 AGUILAR 70 HBC + K-P 4 GEV/SIG.PI 5/70

M+ AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
(SEE IDEOGRAM BELOW)

M- 93 1382.0 3.0 DAHL 61 DBC - K-D 0.45 BEV/C
M- E 224 1376.0 6.4 ELY 61 HLBC 10/69

M- F ERROR OF 3.0 ENLARGED TO 4.4 BY US, BECAUSE LT STATIST. ERR. 10/69

M- 200 1392.0 6.2 COOPER 64 HBC -
M- 1086 1385.3 1.5 HUWE 64 HBC -
M- 1390 1384.0 1.0 ARMENTERO 65 HBC -
M- S 120 1391.5 2.5 SMITH 65 HBC - K-P 1.8 BEV/C 9/66
M- S 58 1399.8 4.0 SMITH 65 HBC - K-P 1.95 BEV/C 9/66
M- S ERROR OF 1.8 ENLARGED TO 2.6 BY US, BECAUSE LT STATIST. ERR. 10/69
M- S ERROR OF 1.4 ENLARGED TO 4.0 BY US, BECAUSE LT STATIST. ERR. 10/69
M- 1389.0 9.0 LONDON 66 HBC - 7/66
M- 370 1390.7 2.0 SIEGEL 67 HBC - K-P AT 2.1 GEV/C 10/69

M- AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.2)
(SEE IDEOGRAM BELOW)

43 Y*(1)1385 - Y*(1)1385 MASS DIFFERENCE (MEV)

D	R	(10.0)	(4.2)	ELY	61 HLBC	+ K-P 1.11 BEV/C	8/66
D	R	(17.1)	(7.1)	COOPER	64 HBC		10/69
D	R	(4.3)	(2.2)	HUWE	64 HBC	+ K-P 1.22 BEV/C	8/66
D	R	(2.0)	(1.5)	ARMENTERO	65 HBC	+ K-P 1.9-1.2 BEV/C	8/66
D	R	(17.2)	(2.1)	SMITH	65 HBC	+ K-P 1.8 BEV/C	9/66
D	R	(17.2)	(2.0)	SMITH	65 HBC	+ K-P 1.95 BEV/C	9/66
D	R	(11.0)	(9.0)	LONDON	66 HBC	+ K-P 2.24 BEV/C	8/66
D	R	9.0	6.0	LONDON	66 HBC	+ LAMDA 3 PI EVTS	7/66
D	R	(6.3)	(2.0)	SIEGEL	67 HBC	+ K-P AT 2.1 GEV/C	10/69

R REDUNDANT WITH DATA IN MASS LISTING.

43 Y*(1)1385 WIDTH (MEV)

W	[64.0]		ALSTON	60 HBC	+
	(20.0)		MARTIN	61 HBC	0+
W	(40.0)		BERGE	61 HBC	+
W	(80.0)	(10.0)	COLLEY	62 HLBC	0-
W	(30.0)	(9.0)	CURTIS	63 DSPK	0
W	(38.0)	(9.0)	MUSGRAVE	65 HBC	+0
W	(26.0)	(5.0)	BALTAY	65 HBC	+ -

OR LESS

W+	48.0	8.0	ELY	61 HLBC	+	
W+	51.0	10.0	COOPER	64 HBC	+	
W+	46.5	3.0	HUWE	64 HBC	+	
W+	30.3	3.1	ARMENTERO	65 HBC	+	
W+	33.1	3.8	SMITH	65 HBC	+ K-P 1.8 BEV/C	
W+	40	25.0	6.0	BIRMINGHAM	66 HBC	+ 3.5 K-P
W+	1260	36.0	3.0	SIEGEL	67 HBC	+ K-P AT 2.1 GEV/C
W+				AGUILAR	70 HBC	+ K-P 4 GEV/SIG.PI

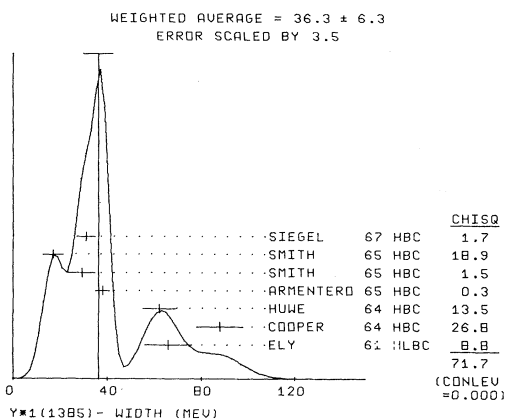
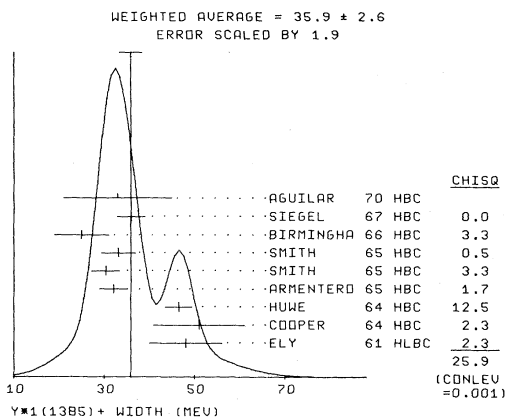
M+ AVG 35.9 2.6 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)
(SEE IDEOGRAM BELOW)

W-	(40.0)		DAHL	61 DBC	-	
W-	66.0	10.0	ELY	61 HLBC	-	
W-	88.0	10.0	COOPER	64 HBC	-	
W-	62.0	7.0	HUWE	64 HBC	-	
W-	39.0	3.0	ARMENTERO	65 HBC	-	
W-	29.2	5.7	SMITH	65 HBC	- K-P 1.80 BEV/C	
W-	17.1	4.4	SMITH	65 HBC	- K-P 1.95 BEV/C	
W-	370	31.0	4.0	SIEGEL	67 HBC	- K-P AT 2.1 GEV/C

W- AVG 36.3 6.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.5)
(SEE IDEOGRAM BELOW)

Baryons

For notation, see illustrated key at beginning of data card listings.



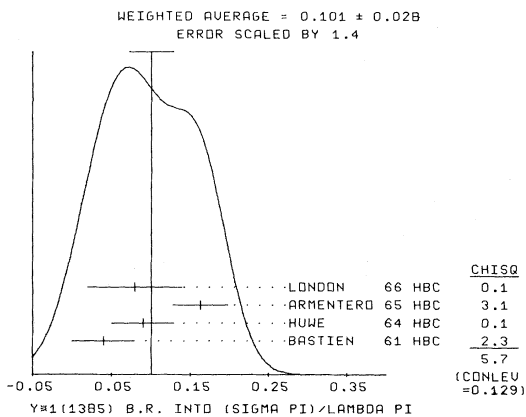
43 Y*1(1385) PARTIAL DECAY MODES

P1	Y*1(1385) INTO	DECAY MASSES
P1	INTO LAMBDA PI	1115+ 139
P2	INTO SIGMA PI	1197+ 139

43 Y*1(1385) BRANCHING RATIOS

R1	Y*1(1385) INTO (SIGMA PI)/(LAMBDA PI)	(P2)/(P1)
R1	0.04	OR LESS
R1	0.09	0.04
R1	0.163	0.035
R1	0.08	0.06
R1	0.101	0.028

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
(SEE IDEOGRAM BELOW)



REFERENCES -- Y*1(1385)

ALSTON 60 PRL 5 520 +ALVAREZ, EBERHARD, GOOD, GRAZIANO, + (LRL) I
 DAHL 61 PRL 6 142 +HORWITZ, MILLER, MURRAY, WHITE (LRL)
 MARTIN 61 PRL 6 283 +LEIPUNFA, CHINDENSKY, SHIVELY, + (BNL, YALE)
 BERGE 61 PRL 6 557 +BASTIEN, DAHL, FERRO-LUZZI, KIRZ, + (LRL)
 RASTIEN 61 PRL 6 702 P BASTIEN, M FERRO-LUZZI, A H ROSENFELD (LRL)
 ELY 61 PRL 7 461 +FUNG, GDAL, PAN, POWELL, WHITE (LRL) J

ALSTON 62 CERN CONF 311 +ALVAREZ, FERRO-LUZZI, ROSENFELD, + (LRL)
 COLLEY 62 PR 128 1930 +GELFAND, NAUENBERG, + (COLUMBIA, RUTGERS) JP
 CURTIS 63 PR 132 1771 +COFFIN, MEYER, TERWILLIGER (MICH) J
 COOPER 64 PL 8 365 +FILTHUTH, FRIDMAN, MALAMUD, + (CERN, AMSTR)
 HUWE 64 UCRL-11291 THESIS D O HUWE (LRL) JP
 ALSO 69 PR 180 1824 D O HUWE (LRL)

MUSGRAVE 65 NC 35 735 +PETMEZAS, + (BIRMINGHAM, CERN, EP, IMP, COL, SACLAY)
 ARMENTEROS 65 PL 19 75 (CERN, HEIDEL, SACLAY)
 BALTAY 65 PR 140 81027 +SANDWEISS, TAFT, CULWICK, KOPP, + (YALE, BNL)
 SMITH 65 THESIS (UCLA) L T SMITH (UCLA)

BIRMINGHAM 66 PR 152 1148 BIRMINGHAM, GLASGOW, I.C., OXFORD, RUTHERFORD
 LONDON 66 PR 143 1034 +RAU, SAMIOS, YAMAMOTO, GOLDBERG, + (BNL, SYR) J
 SIEGEL 67 UCRL 18041 THESIS D M SIEGEL (LRL)
 AGUILAR 70 PRL 25 58 +HARNES, BASSANO, CHUNG, EISNER, + (BNL, SYR)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN DATA CARDS.

SHAFFER 66 PR 134 11372 J B SHAFFER, D O HUWE (LRL) JP
 MALAMUD 64 PL 10 145 E MALAMUD, P E SCHLEIN (CERN, UCLA) JP

M < 1500 MEV - PRODUCTION EXPERIMENTS

80 Y*1(1440, JP=) I=1

Σ(1440) BUMPS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

CLINE 68 FIND A NARROW PEAK AT 1440 MEV (JUST ABOVE THE KBAR N THRESHOLD) IN THE LAMBDA PI INVARIANT MASS FOR K- D TO LAMBDA PI- P EVENTS. THEY DISCUSS ALTERNATE INTERPRETATIONS -- THAT IT IS A RESONANCE OR A KINEMATIC EFFECT. IN CLINE 68 THE K- BEAM MOMENTUM IS 0.4 GEV/C. IN A STUDY OF THE SAME REACTION WITH A MOMENTUM OF 1.1 GEV/C, ALEXANDER 69 FIND NO PEAK. IN ADDITION, THEY ARE ABLE TO EXPLAIN THE RESULTS OF BOTH EXPERIMENTS WITHOUT INVOKING A NEW RESONANCE.

REFERENCES -- Y*1(1440)

CLINE 68 PRL 21 1372 D CLINE, R LAUMANN, J MAPP (WISCONSIN) I
 ALEXANDER 69 PRL 22 483 ALEXANDER, HALL, JEM, + (LRL, RIVERSIDE)

23 Y*1(1480, JP=) I=1

Σ(1480) BUMPS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

PEAKS ARE SEEN IN LAMBDA PI AND SIGMA PI SPECTRA IN THE REACTION PI+P TO K+ PI Y AT 1.7 GEV/C. ALSO THE POLARIZATION OSCILLATES IN THE SAME REGION.

SEE MILLER 70 FOR A DISCUSSION OF THIS STATE. HE SUGGESTS A POSSIBLE ALTERNATE EXPLANATION IN TERMS OF A REFLECTION OF N*1/2(1670) DECAY TO LAMBDA K.

23 Y*1(1480) MASS (MEV)

M	O	(1480.0)	(15.0)	YU-LI PAN 69 HBC + P1+P TO K PI LAM	9/69
M	O	(1465.0)	(20.0)	YU-LI PAN 69 HBC + P1+P TO K PI SIG	9/69
M	O	1479.	10.	PAN 70 HBC + P1+P TO K PI LAM	3/71*
M	O	1465.	15.	PAN 70 HBC + P1+P TO K PI SIG	3/71*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

23 Y*1(1480) WIDTH (MEV)

W	O	(35.0)	YU-LI PAN 69 HBC + P1+P TO K PI LAM	9/69	
W <th>O</th> <th>(25.0)</th> <th>YU-LI PAN 69 HBC + P1+P TO K PI SIG</th> <th>9/69</th>	O	(25.0)	YU-LI PAN 69 HBC + P1+P TO K PI SIG	9/69	
W <th>O</th> <th>31.</th> <th>15.</th> <th>PAN 70 HBC + P1+P TO K PI LAM</th> <th>3/71*</th>	O	31.	15.	PAN 70 HBC + P1+P TO K PI LAM	3/71*
W <th>O</th> <th>30.</th> <th>20.</th> <th>PAN 70 HBC + P1+P TO K PI SIG</th> <th>3/71*</th>	O	30.	20.	PAN 70 HBC + P1+P TO K PI SIG	3/71*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

23 Y*1(1480) PARTIAL DECAY MODES

P1	Y*1(1480) INTO	DECAY MASSES
P1	INTO KBAR N	497+ 939
P2	INTO LAMBDA PI	1115+ 139
P3	INTO SIGMA PI	1189+ 139

23 Y*1(1480) BRANCHING RATIOS

R1	Y*1(1480) INTO (SIGMA PI)/(LAMBDA PI)	(P3)/(P2)
R1	(0.72)	(0.49)
R1	0.82	0.51
R2	Y*1(1480) INTO (PROTON KOBARI)/(LAMBDA PI)	(P1)/(P2)
R2	(0.36)	(0.25)
R2	0.36	0.25

REFERENCES -- Y*1(1480)

YU-LI PA 69 PRL 23 806 YU-LI PAN, F L FORMAN (PENN) I
 YU-LI PA 69 PRL 23 808 YU-LI PAN, F L FORMAN (PENN) I
 MILLER 70 DUKE 229 D H MILLER (REVIEW TALK) (PURDUE)
 PAN 70 PR D2,449 +FORMAN, KO, HAGOPIAN, SELOVE (PENN)

END PRODUCTION EXPERIMENTS

Baryons

For notation, see illustrated key at beginning of data card listings.

Σ(1620) 32 Y*(1620, JP=1/2-1) I=1 **S'11**
 THE K-MATRIX ANALYSIS OF KIM 71 SUGGESTS THIS STATE.
 THE ELASTICITY IS REALLY SMALL.

 32 Y*(1620) MASS (MEV) -----
 M (1620.) KIM 71 DPWA K-MATRIX ANAL. 3/71*

 32 Y*(1620) WIDTH (MEV) -----
 W (40.) KIM 71 DPWA K-MATRIX ANAL. 3/71*

 32 Y*(1620) PARTIAL DECAY MODES -----
 P1 Y*(1620) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(1620) INTO SIGMA PI 1197+ 139
 P3 Y*(1620) INTO LAMBDA PI 1115+ 134

 32 Y*(1620) BRANCHING RATIOS -----
 BFLOW, XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1 Y*(1620) INTO KBAR N XI
 R1 (0.05) KIM 71 DPWA K-MATRIX ANAL. 3/71*

R2 Y*(1620) FROM KBAR N TO SIGMA PI
 R2 (0.08) KIM 71 DPWA SORT(XI*X2) K-MATRIX ANAL. 3/71*

R3 Y*(1620) FROM KBAR N TO LAMBDA PI
 R3 (0.15) KIM 71 DPWA SORT(XI*X3) K-MATRIX ANAL. 3/71*

 REFERENCES -- Y*(1620)
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV) IJP

Σ(1620) 79 Y*(1620, JP=1/2+) I=1 **P'11**
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THE PARTIAL-WAVE ANALYSIS OF K- N TO SIGMA PI BY
 ARMENTEROS 70 SUGGESTS SUCH A RESONANCE. KIM 71 FINDS
 A SIGNAL IN BOTH KBAR-N AND SIGMA PI.

 79 Y*(1620) MASS (MEV) -----
 M 1500. -- 1600. ARMENTEROS 70 HD8C 0- K-N TO SIGMA PI 6/70
 M (1670.) KIM 71 DPWA K-MATRIX ANAL. 3/71*

 79 Y*(1620) WIDTH (MEV) -----
 W (50.0) ARMENTEROS 70 HD8C 0- K-N TO SIGMA PI 6/70
 W (50.) KIM 71 DPWA K-MATRIX ANAL. 3/71*

 79 Y*(1620) PARTIAL DECAY MODES -----
 P1 Y*(1620) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(1620) INTO SIGMA PI 1197+ 139

 79 Y*(1620) BRANCHING RATIOS -----
 BFLOW XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1 Y*(1620) FROM KBAR N TO SIGMA PI
 R1 (40.2) ARMENTEROS 70 HD8C 0- K-N TO SIGMA PI 6/70
 R1 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71*

R2 Y*(1620) INTO KBAR N XI
 R2 (0.14) KIM 71 DPWA K-MATRIX ANAL. 3/71*

 REFERENCES -- Y*(1620)
 ARMENTEROS 70 DUKE 123 ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV) IJP

Note on Σ(1620)

This state was first suggested by the BNL-CCNY collaboration (CRENNELL 68) who presented evidence for it in the reaction $K^-n \rightarrow \Sigma(1620)^+ \pi^- \pi^-$ with $\Sigma(1620)^+$ decaying into $\Lambda \pi^+$. Since then there have been conflicting reports about this state. A good review of the production experiments has been recently given by MILLER 70. We summarize here the situation.

Formation experiments. The CHS energy independent (see note on Y^* 's) partial-wave analysis (ARMENTEROS 70) shows evidence for a P_{11} state at 1620 MeV in the $\Sigma \pi$ channel. They had previously

reported evidence for this state in the $\Lambda \pi$ channel (ARMENTEROS 69), which they now disclaim. This is because of the fact that new data have been added at the low-energy end of their previous analysis (i.e., around the 1600-MeV region) and now the preferred solution shows structure at a higher mass, in the $\Lambda \pi$ channel.

Production experiments. Here the evidence is only in the $\Lambda \pi$ channel. The BNL-CCNY collaboration, with increased data, CRENNELL 69, still claim the effect in the $\Lambda \pi$ channel (no evidence seen in $\bar{K}N$ or $\bar{K}N\pi$). SABRE 70 studied the same reaction at 3.0 GeV/c with comparable statistics and do not see any evidence for it in the $\Lambda \pi$ channel; on the contrary, they believe it to be a spurious peak resulting from misidentified Σ^0 from the production of $\Sigma(1670)$ decaying into $\Sigma^0 \pi^+$. CRENNELL 69 give counter arguments to show that this is not the case in their data and the controversy goes on. AMMANN 70 studied the same reaction at 4.5 GeV/c and report a state at 1640 MeV, again decaying only into $\Lambda \pi$ (no evidence seen in $\Sigma \pi$ or $\bar{K}N$ channels). The closeness of this mass to 1670 MeV is suggestive that this effect may be related to what goes on in that region (see discussion below).

In conclusion, for $\Sigma(1620)$ we have to wait for more data and for a complete understanding of the entire mass region 1600 to 1700 MeV. The hope is that the determination of quantum numbers for each of these effects for each decay mode may eventually clarify the situation.

Σ(1620) BUMPS 78 Y*(1620, JP= 1 I=1 PROD.EXP.
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THIS RESONANCE NEEDS CONFIRMATION. THE RESULTS OF
 CRENNELL 69 AT 3.9 GEV/C ARE NOT CONFIRMED BY THE SABRE
 COLLABORATION AT 3.0 GEV/C (SABRE 70). HOWEVER IN AN EXPERIMENT AT
 4.5 GEV/C, AMMANN 70 SEE A PEAK AT 1642 MEV WHICH ON THE BASIS OF
 BRANCHING RATIOS THEY DO NOT ASSOCIATE WITH THE Y*(1670). SEE MILLER
 70 FOR A REVIEW OF THESE CONFLICTS.

THERE WAS AN INDICATION OF A Y*(1610) IN AN EARLY PHASE-SHIFT
 ANALYSIS OF K-P TO LAMBDA PI. HOWEVER MORE DETAILED ANALYSIS OF
 MORE EXTENSIVE DATA BY THE SAME (CERN, HEIDELBERG, SACLAY) GROUP
 FAILED TO CONFIRM THIS RESULT. HOWEVER THEY NOW SEE IT IN THE SIGMA PI
 CHANNEL (SEE PREVIOUS ENTRY). (OLD LAMBDA PI ANALYSIS LISTED AS
 ARMENTEROS 69, NEW ANALYSIS AS ARMENTEROS 70.)

 78 Y*(1620) MASS (MEV) -----
 M N (1616.0) (8.0) CRENNELL 68 DBC +- K-D 3.9 GEV/C 11/68
 N EVENTS OF CRENNELL 68 ARE IN THE LARGER SAMPLE OF CRENNELL 69.
 M 20 1618.0 3.0 BLUMENFEL 69 HRC + KN LONG + PROTON 9/69
 M 1619.0 8.0 CRENNELL 69 DBC +- K-N TO LAM 3 PI 9/69
 M 1642.0 12.0 AMMANN 70 DBC K-P 4.5 GEV/C 6/70
 M
 M AVG 1619.4 3.8 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)

 78 Y*(1620) WIDTH (MEV) -----
 W N (66.0) (16.0) CRENNELL 68 DBC +- SEE NOTE N ABOVE 11/68
 W 20 30.0 10.0 BLUMENFEL 69 HRC + 9/69
 W 72.0 22.0 CRENNELL 69 DBC +- 9/69
 W 55.0 24.0 AMMANN 70 DBC K-P 4.5 GEV/C 6/70
 W
 W AVG 41.3 12.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
 SEE THE NOTES ACCOMPANYING THE MASSES QUOTE)

 78 Y*(1620) PARTIAL DECAY MODES -----
 P1 Y*(1620) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(1620) INTO LAMBDA PI 1115+ 139
 P3 Y*(1620) INTO Y*(1385) PI 1395+ 139
 P4 Y*(1620) INTO LAMBDA PI PI 1115+ 139+ 139
 P5 Y*(1620) INTO SIGMA PI 1197+ 139
 P6 Y*(1620) INTO Y*(1405) PI 1405+ 139

Baryons

For notation, see illustrated key at beginning of data card listings.

78 Y*(1620) BRANCHING RATIOS

R1	Y*(1620) INTO (LAMBDA PI P1)/(LAMBDA PI)	(P4)/(P3)	
R1	14 (2.5) APPROX	RLUMENFEL 69 HRC	+
R2	Y*(1620) INTO (KBAR N)/(LAMBDA PI)	(P1)/(P2)	
R2	10.0 (0.1)	CRFNELL 68 DRC	+
R2	0.4	AMMANN 70 DRC	K-P 4.5 GEV/C 6/70
R3	Y*(1620) INTO LAMBDA PI	(P2)	
R3	LARGE	CRFNELL 68 DRC	-- 11/68
R4	Y*(1620) INTO (Y*(1385) P1)/(LAMBDA PI)	(P3)/(P2)	
R4	10.2 (0.1)	CRFNELL 68 DRC	+
R4	10.33 OR LESS CL=.95	AMMANN 70 DRC	K-P 4.5 GEV/C 6/70
R5	Y*(1620) INTO (SIGMA PI)/(LAMBDA PI)	(P5)/(P2)	
R5	1.1 (0.9) PC UPPER LIMIT	AMMANN 70 DRC	K-P 4.5 GEV/C 6/70
R6	Y*(1620) INTO (Y*(1405) P1)/(LAMBDA PI)	(P6)/(P2)	
R6	0.7 0.4	AMMANN 70 DRC	K-P 4.5 GEV/C 6/70

REFERENCES -- Y*(1620)

CRFNELL 69 PRL 21 448 *DELANEY, FLAMINIO, KARSHON, + (BNL,CCNY) I
 RLUMENFEL 69 PL 298 58 *RLUMENFEL, KALDFELDSCH (BNL) I
 CRFNELL 69 LUND CONF *KARSHON, LAI, ONEIL, SCARR, + (BNL,CCNY) I
 AMMANN 70 PRL 24 327 ARE QUOTED IN LEVI SETTI 69, + GARFINKEL, CARMONY, GUTAY, + (PURDUE,IND) I
 PAPERS NOT REFERRED TO IN DATA CARDS

ARMENTEROS 68 NP 88 183 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY) I
 LEVISETT 69 LUND CONF R LEVI SETTI (RAPPORTFUR) EFINS (LRL) I
 TRIPP 69 UCPL 19361 R D TRIPP (LRL) I
 SARRÉ 70 NP 815 201 SARRÉ COLLAB. (SACL,AMST,BGNA,REHO,EPDL) I
 ARMENTEROS 70 DUKE 123 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY) I
 MILLER 70 DUKE 229 D H MILLER (REVIEW TALK) (PURDUE) I

 END PRODUCTION EXPERIMENTS

Note on $\Sigma(1670)$

Formation experiments show the presence of only one $I = 1$ state in this energy region with major decay modes into $\bar{K}N$ (9%), $\Delta\pi$ (10%), $\Sigma\pi$ (40%) and a very small branching fraction into $\Sigma\pi\pi$. Its quantum numbers are $J^P = 3/2^-$.

Production experiments are more confused.

When determined, the most likely quantum numbers are also $3/2^-$ [for $\Sigma\pi$ and $\Lambda(1405)\pi$]. The measured branching ratio $R = \Sigma\pi/\Sigma\pi\pi$ changes with the momentum transfer to the proton. This was first observed by EBERHARD 69 who suggested the existence of $2 Y_1^*$ with the same mass and quantum numbers; one object with a large $\Sigma\pi\pi$ [mainly $\Lambda(1405)\pi$] decay mode produced peripherally, and another one with a large $\Sigma\pi$ decay mode produced at larger angles. This observation has now been confirmed by AGUILAR-BENITEZ 70.

The other difficulty comes from the different $\Delta\pi/\Sigma\pi$ branching ratios reported by the various experiments. The possibility of a third Y_1^* state with a large $\Delta\pi/\Sigma\pi$ branching ratio still exists. This large branching ratio is the main justification for this hypothesis and needs confirmation. It relies on the separation between $K^-p \rightarrow \Lambda\pi^+\pi^-$ and $K^+p \rightarrow \Sigma^0\pi^+\pi^-$ that is experimentally difficult at high energy. These problems are reviewed by MILLER 70.

Two resonances of the same spin and parity have been hypothesized by EBERHARD 69 as the origin of much of the complexity observed in the 1600 to 1700 MeV region in production experiments. See also note on $\Sigma(1620)$.

$\Sigma(1670)$ 44 Y*(1670,JP=3/2-) I=1 D_{13}^+

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 SEE NOTE ABOVE
 WELL ESTABLISHED RESONANCE. IT HAS BEEN SEEN IN BOTH FORMATION AND PRODUCTION EXPERIMENTS. HOWEVER THE BRANCHING RATIOS OBTAINED BY THESE TWO METHODS SHOW LARGE INCONSISTENCIES.

SEE LISTING OF PRODUCTION EXPERIMENTS BELOW

AS FOR THE QUANTUM NUMBERS, THE ANALYSES OF LAMBDA PI CHANNEL (IN FORMATION EXP.) AS WELL AS THE SIGMA PI CHANNEL AGREE ON JP=3/2-.

44 Y*(1670) MASS (MEV)

M	1660.0	BERLEY 64 HRC	0	K-P TO LAM P10	7/66
M	1669.	ARMENTEROS 68 HRC	0	K-P ELAS.+CH,EX	11/68
M	(1661.0) (2.0)	ARMENTE2 68 HRC	0	K-P TO SIGMA PI	11/68
M	1680.	ARMENTE4 69 DRC	K-N TO SIG- P10	12/68	
M	1663.0	ARMENT-5 69 HRC	0	K-P TO SIGMPI ED	9/69
M	1672.0	BERLEY 69 HRC	K-P TO SIG PI	5/70	
M	1660.	ARMENTER 70 HRC	OK-P TO LAM, PI EI	5/70	
M	1662.0	GALTIERI 70 HRC	0	SIG PI,EDPWA	7/70
M	1665.	GALTIERI 70 HRC	0	LAM, PI, EDPWA	7/70
M	1670.	KIM 71 DPWA	K-MATRIX ANAL.	3/71*	

44 Y*(1670) WIDTH (MEV)

W	60.0	BERLEY 64 HRC	0	K-P ELAS.+CH,EX	7/66
W	56.	ARMENTEROS 68 HRC	0	K-P ELAS.+CH,EX	11/68
W	(44.0) (4.0)	ARMENTE2 68 HRC	0	K-P TO SIGMA PI	11/68
W	47.0	ARMENTE4 69 DRC	K-N TO SIG- P10	12/68	
W	49.0	ARMENT-5 69 HRC	0	K-P TO SIGMPI ED	9/69
W	34.0	BERLEY 69 HRC	K-P TO SIG PI	5/70	
W	50.	ARMENTER 70 HRC	0	K-P TO LAM, PI	5/70
W	(48.0) (5.0)	GALTIERI 70 HRC	0	SIG PI,EDPWA	7/70
W	50.	GALTIERI 70 HRC	0	LAM, PI, EDPWA	7/70
W	40.	KIM 71 DPWA	K-MATRIX ANAL.	3/71*	

44 Y*(1670) PARTIAL DECAY MODES

P1	Y*(1670) INTO KBAR N	DECAY MASSES
P2	Y*(1670) INTO LAMBDA PI	497+ 939
P3	Y*(1670) INTO SIGMA PI	1115+ 139
P4	Y*(1670) INTO LAMBDA PI PI	1197+ 139
P5	Y*(1670) INTO SIGMA PI PI	1115+ 139+ 139
P6	Y*(1670) INTO Y*(1385) PI	1197+ 139+ 139
P7	Y*(1670) INTO Y*(1405) PI	1385+ 139
		1405+ 139

44 Y*(1670) BRANCHING RATIOS

BELOW X1 = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1	Y*(1670) INTO (KBAR N)/TOTAL	X1	
R1	(0.14) (0.02)	ARMENTEROS 68 HRC	0
R1	0.08 (0.02)	ARMENT-5 69 HRC	0
R1	0.07	KIM 71 DPWA	K-MATRIX ANAL. 9/69
R2	Y*(1670) INTO (LAMBDA PI P1)/TOTAL	(P4)/TOTAL	
R2	(0.11) OR LESS	ARMENTE3 68 HRC	K-P (P1=.09) 9/69
R3	Y*(1670) INTO (SIGMA PI P1)/TOTAL	(P5)/TOTAL	
R3 A	(0.14) OR LESS	ARMENTE3 68 HRC	K-P AND 0-P1+.09 11/68
R3 A	RATIO ONLY FOR (SIGMPI) SYSTEM IN I=1, WHICH CANNOT BE Y*(1385)		11/68
R4	Y*(1670) INTO (Y*(1405) P1)/TOTAL	(P7)/TOTAL	
R4	(0.06) OR LESS	ARMENTE3 68 HRC	K-P AND 0-P1+.09 11/68
R5	Y*(1670) FROM KBAR N TO LAMBDA PI	SORT(X1)*X2	
R5	+0.1	ARMENTER 70HRC	K-P TO LAMB PI 5/70
R5	+0.02 (0.02)	GALTIERI 70 HRC	0 LAM, PI, EDPWA 7/70
R5	0.08	KIM 71 DPWA	K-MATRIX ANAL. 3/71*
R6	Y*(1670) FROM KBAR N TO SIGMA PI	SORT(X1)*X3	
R6	(+0.21) (0.01)	ARMENTE2 68 HRC	0 OLD DATA 11/68
R6	+0.19	ARMENTE4 69 DRC	0 NEW DATA 9/69
R6	+0.20 (0.01)	ARMENT-5 69 HRC	0 NEW DATA 9/69
R6	+0.18	BERLEY 69 HRC	0 SIG PI,EDPWA 5/70
R6	+0.18 (0.06)	GALTIERI 70 HRC	0 SIG PI,EDPWA 7/70
R6	0.15	KIM 71 DPWA	K-MATRIX ANAL. 3/71*
R7	Y*(1670) INTO (Y*(1385)P1)*(KBAR N)/TOTAL**2	(P6*P1)/TOTAL**2	
R7	(0.031) (0.004)	SIMS 68 DRC	- 11/68
R8	Y*(1670) INTO (Y*(1405)P1)*(KBAR N)/TOTAL**2	(P7*P1)/TOTAL**2	
R8	(0.03) OR LESS	BERLEY 69 HRC	0 K-P .6-.92 REV/C 5/70

 REFERENCES -- Y*(1670)

BERLEY 64 OJUNA CONF I 565 *CONNOLLY,HART,RAHM,STONEHILL, + (BNL)IJP
 ARMENTEROS 68 NP 88 195 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY)IJP
 ARMENTE1 68 NP 88 183 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY)IJP
 ARMENTE2 68 NP 88 223 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY)IJP
 ARMENTE3 68 PL 288 521 *ARMENTEROS,BAILLON + (CERN,HEID+SALAY)IJP
 SIMS 68 PRL 21 1413 SIMS,ALBRIGHT,BARTLEY,MEER+ (FLO+TAFTS+BRK)
 ARMENTE4 69 NP 810 450 *ARMENTEROS,BAILLON,MINTEN + (CERN+SACLAY) J
 ARMENT-5 69 NP 814 91 *ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
 BERLEY 69 PL 308 430 BERLEY,HART,RAHM,WILLIS,YAMAMOTO (BNL) I
 ARMENTER 70 DUKE 123 *ARMENTEROS, BAILLON, + (CERN,HEID) I
 GALTIERI 70 DUKE 173 A. BARBARO GALTIERI (LRL)IJP
 KIM 71 HARVARD PREPRINT J.K.KIM (HARV) IJP
 SEE ALSO DUKE 161 J. K. KIM (HARV)IJP

PAPERS NOT REFERRED TO IN DATA CARDS

BASTIEN1 63 PRL 10 188 P L BASTIEN, J P BERGE (LRL) IJ
 --- REPLACED BY BASTIEN 2, BUT SIMILAR AND MORE READILY AVAILABLE.
 BASTIEN2 63 UCRL-10779 THESIS P L BASTIEN (LRL) IJ
 T-ZADEH 63 PRL 11 470 TAHER-ZADEH,PROWSE,SCHLEIN,SLAFER,+ (UCLA) IJP
 --- SEE NOTE FOLLOWING SCHLEIN 66.
 SCHLEIN 66 UCRL-1016 P.E. SCHLEIN, T.G. TRIPPE (UCLA) JP
 --- REANALYSES DATA OF TAHER-ZADEH 63. BASTIEN 63 AND ALL PUBLISHED LAMBDA PI CROSS SECTION DATA IN THE LIGHT OF THE NOW KNOWN Y*(1765). REVERSES THE MODEL-DEPENDENT CONCLUSION OF TAHER-ZADEH ON THE PREFERRED JP ASSIGNMENT (FROM 3/2+ TO 3/2-).
 SMART 66 PRL 17 556 W Y SMART,A KERNAN,G E KALMUS,R P FLY (LRL)IJP
 ARMENTEROS 67 NP 83 592 *ARMENTEROS,FERRI-LUZZI+ (CERN,HEID,SACLAY) I

Baryons

For notation, see illustrated key at beginning of data card listings.

57 Y*(1765) WIDTH (MEV)

W	ABOUT 50.0	MEYER	67 RVUE	9/69
W	ABOUT 80.0	ARMENTERO	70 DPWA 0	6/70
W	(55.0) (10.0)	CONFORTO	70 DPWA 0	6/70
W	(50.1)	KIM	71 DPWA	3/71*

57 Y*(1765) PARTIAL DECAY MODES

P1	Y*(1765) INTO KBAR N	497+ 930
P2	Y*(1765) INTO SIGMA ETA	1197+ 548
P3	Y*(1765) INTO LAMBDA PI	1115+ 134
P4	Y*(1765) INTO SIGMA PI	1197+ 139

57 Y*(1765) BRANCHING RATIOS

BELOW, X1 = (PARTIAL-WIDTH 1)/TOTAL WIDTH, ETC. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1	Y*(1765) INTO (KBAR N)/TOTAL	X1	9/69
R1	(0.12) (0.05)	CONFORTO 70 DPWA 0	6/70
R1	(0.8)	KIM 71 DPWA	3/71*
R2	Y*(1765) FROM KBAR N INTO SIGMA ETA	ELASTIC, CH EXCH	6/70
R2	SEFN	CLINE 69 DRC - THRESHOLD BUMP.	9/69
R3	Y*(1765) FROM KBAR N INTO LAMBDA PI	ELASTIC, CH EXCH	6/70
R3	(-0.25) (0.09)	ARMENTERO 70 DPWA 0	6/70
R3	(0.09)	KIM 71 DPWA	3/71*
R4	Y*(1765) FROM KBAR N TO SIGMA PI	ELASTIC, CH EXCH	6/70
R4	(0.16)	KIM 71 DPWA	3/71*

REFERENCES -- Y*(1765)

CLINE 67 PL 250 41
 MEYER 67 HEIDELBERG C 117
 ARMENTERO 70 DUKE 123
 CONFORTO 70 NP (SUBMITTED)
 KIM 71 HARVARD PREPRINT
 SEE ALSO DUKE 161

WISCONSIN IJJP (SACLAY) IJJP
 (CERN, HEIDEL) IJJP
 (CHICAGO, HEIDEL) IJJP
 (HARV) IJJP (HARV) IJJP

PAPERS NOT REFERRED TO IN DATA CARDS

FERRI-LU 56 BROOKLYN CONF 183 M
 ARMENTERO 68 NP BR 183
 ARMENTERO 69 LUND CONF PAPER

Σ(1765)

45 Y*(1765, JP=5/2-) I=1

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

45 Y*(1765) MASS (MEV)

W	1765.0	10.0	GALTIERI 63 DRC 0	K-D 1.51 BEV/C	7/66
W	1765.0	10.0	ARMENTERO 65 HBC 0	K-P TO Y*1520 PI	7/66
W	1760.0	10.0	RELL 2 66 DRC -	K-N TO Y*1520 PI	7/66
W	1765.0	8.0	FENSTER 66 HBC 0	K-P TO Y*1520 PI	9/66
W	1775.0	7.0	BUGG 68 CNTR	K-P, 0 TOTAL	11/68
W	1768.0	2.0	SMART 68 RVUE 0	K-N TO LAMBDA PI	7/68
W	1770.0	10.0	ARMENT-1 68 DPWA 0	ELASTIC, CH EXCH	11/68
W	1765.0	10.0	COOL 70 CNTR	K-P, 0 TOTAL	10/70*
W	1770.0	3.0	GALTIERI 70 DPWA 0	K-P TO LAMBDA PI	7/70
W	(1765.)		CONFORTO 70 DPWA 0	ELASTIC, CH EXCH	6/70
W			KIM 71 DPWA	K-MATRIX ANAL.	3/71*
W					1/71*

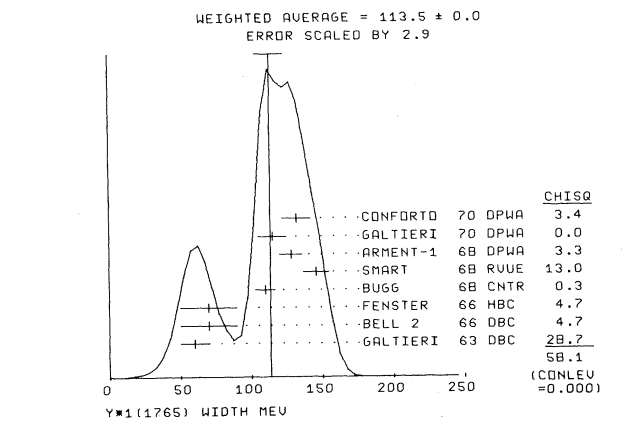
N ERROR STATIST. ONLY- NO ERROR DUE TO PARTICULAR P.W. ANAL. INCLUDED

AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)

45 Y*(1765) WIDTH (MEV)

W	60.0	10.0	GALTIERI 63 DRC 0	7/66
W	70.0	20.0	RELL 2 66 DRC -	7/66
W	70.0	20.0	FENSTER 66 HBC 0	9/66
W	110.0	7.0	BUGG 68 CNTR	K-P, 0 TOTAL
W	148.0	9.0	SMART 68 RVUE 0	7/68
W	128.0	8.0	ARMENT-1 68 DPWA 0	ELASTIC, CH EXCH
W	(100.0)		COOL 70 CNTR	K-P, 0 TOTAL
W	115.0	10.0	GALTIERI 70 DPWA 0	K-P TO LAMBDA PI
W	132.0	10.0	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH
W	(100.)		KIM 71 DPWA	K-MATRIX ANAL.
W				
W	113.5	10.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.9)	

(SEE IDEOGRAM BELOW)



45 Y*(1765) PARTIAL DECAY MODES

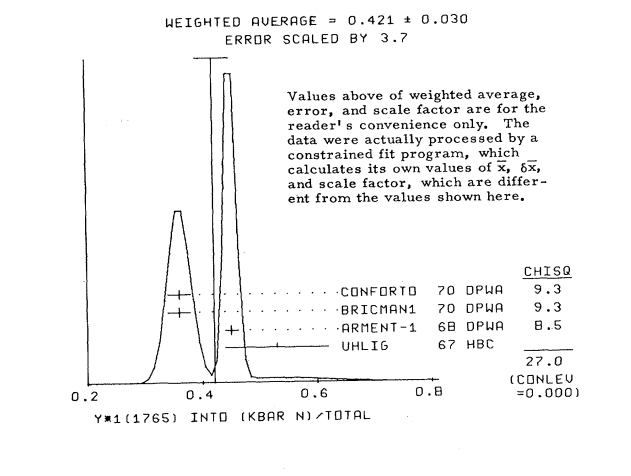
P1	Y*(1765) INTO KBAR N	497+ 930
P2	Y*(1765) INTO LAMBDA PI	1115+ 134
P3	Y*(1765) INTO Y*(01520) PI	1520+ 139
P4	Y*(1765) INTO Y*(11385) PI	1395+ 139
P5	Y*(1765) INTO SIGMA PI	1197+ 139
P6	Y*(1765) INTO SIGMA ETA	1197+ 548
P7	Y*(1765) INTO SIGMA PI PI	1197+ 139+ 139

45 Y*(1765) BRANCHING RATIOS

BELOW, X1 = (PARTIAL-WIDTH 1)/TOTAL WIDTH, ETC. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

ERRORS QUOTED BY EXPERIMENTERS DO NOT INCLUDE UNCERTAINTY DUE TO PARAMETRIZATION USED IN THE P.W.A. THEY SHOULD BE INCREASED.

R1	Y*(1765) INTO (KBAR N)/TOTAL	X1	9/66
R1	(0.6)	GALTIERI 63 HBC 0	K-P RVUE
R1	0.93	UHLIG 67 HBC 0	11/66
R1	(0.37)	BUGG 68 CNTR	ELASTIC, CH EXCH
R1	0.45	ARMENT-1 68 DPWA 0	SIGTOT, ELAS, CH EXCH
R1	0.36	BRICHMANI 70 DPWA	K-P, 0 TOTAL
R1	(0.41)	COOL 70 CNTR	K-P, 0 TOTAL
R1	0.36	CONFORTO 70 DPWA 0	ELASTIC, CH EXCH
R1	(0.42)	KIM 71 DPWA	K-MATRIX ANAL.
R1			3/71*
R1	0.421	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.7)	
R1	0.424	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 3.3)	
R1		(SEE IDEOGRAM BELOW)	



45 Y*(1765) FROM KBAR N INTO LAMBDA PI

R2	Y*(1765) FROM KBAR N INTO LAMBDA PI	ELASTIC, CH EXCH	6/70
R2	(-0.266) 0.017	SMART 68 DPWA 0	K-N TO LAMBDA PI
R2	(-0.22) 0.03	GALTIERI 70 DPWA 0	K-P TO LAMBDA PI
R2	(0.30)	KIM 71 DPWA	K-MATRIX ANAL.
R2			3/71*
R2	0.285	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)	
R2	0.249	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

45 Y*(1765) FROM KBAR N INTO Y*(01520) PI

R3	Y*(1765) FROM KBAR N INTO Y*(01520) PI	ELASTIC, CH EXCH	6/70
R3	0.27 0.03	ARMENTERO 65 HBC 0	K-P TO Y*1520 PI
R3	0.35 0.05	FENSTER 66 HBC 0	K-P TO Y*1520 PI
R3			9/66
R3	0.291	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
R3	0.243	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)	

45 Y*(1765) FROM KBAR N INTO Y*(11385) PI

R4	Y*(1765) FROM KBAR N INTO Y*(11385) PI	ELASTIC, CH EXCH	6/70
R4	0.24 0.03	ARMENT-2 67 HBC 0	K-P TO LAM PI PI
R4	0.32 0.06	SIMS 68 DRC -	K-N TO LAM PI PI
R4			11/68
R4	0.256	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)	
R4	0.236	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	

45 Y*(1765) FROM KBAR N INTO SIGMA PI

R5	Y*(1765) FROM KBAR N INTO SIGMA PI	ELASTIC, CH EXCH	6/70
R5	+0.06 0.03	ARMENTERO 67 DPWA 0	K-P TO SIGMA PI
R5	(0.06)	GALTIERI 70 DPWA 0	K-P TO SIGMA PI
R5		KIM 71 DPWA	K-MATRIX ANAL.
R5			3/71*
R5	0.067	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R5	0.066	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	

45 Y*(1765) INTO (LAMBDA PI)/(KBAR N)

R6	Y*(1765) INTO (LAMBDA PI)/(KBAR N)	UHLIG 67 HBC 0	K-P, 9 GEV/C
R6	0.33		9/66
R6			
R6	0.344	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)	

45 Y*(1765) INTO (Y*(01520))/(KBAR N)

R7	Y*(1765) INTO (Y*(01520))/(KBAR N)	UHLIG 67 HBC 0	K-P, 9 GEV/C
R7	0.28		9/66
R7			
R7	0.327	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

45 Y*(1765) INTO (Y*(11385))/(KBAR N)

R8	Y*(1765) INTO (Y*(11385))/(KBAR N)	UHLIG 67 HBC 0	K-P, 9 GEV/C
R8	0.25		9/66
R8			
R8	0.310	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

45 Y*(1765) INTO (SIGMA PI PI)/TOTAL

R9	Y*(1765) INTO (SIGMA PI PI)/TOTAL	ARMENT-2 68 HBC 0	K-N TO SIG PI PI
R9	(0.12)		11/68
R9			
R9		FOR ABOUT 3/4 OF THIS, THE SIGMA PI SYSTEM HAS I=0 AND IS ALMOST ENTIRELY Y*(01520). FOR THE OTHER 1/4, THE SIGMA PI HAS I=1. THIS P IS ABOUT WHAT IS EXPECTED FROM THE KNOWN RATE Y*(1765) TO Y*(11385) P PI, AS SEEN IN LAMBDA PI PI.	

Baryons

For notation, see illustrated key at beginning of data card listings.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i^2)}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 5	P 6
P 1	.424+-0.026					
P 2	-.124	.146+-0.014				
P 3	-.249	-.031	.139+-0.020			
P 4	-.032	.004	-.009	.131+-0.024		
P 5	-.095	.012	-.023	.003	.010+-0.005	
P 6	-.555	-.252	-.594	-.559	-.068	.149+-0.042

REFERENCES -- *Y*(1765)

GALTIERI 63 PL 6 294 A BARBARO-GALTIERI, A MUSSAIN, RO TRIPP (LRL) IJ
 ARMENTER 65 PL 19 338 ARMENTEROS, + (CERN, HEIDELBERG, SACLAY) IJP
 RELL 1 66 PRL 16 203 R B BELL, R W BIRGE, Y-L PAN, R T PU (LRL) IJP
 RELL 2 66 UCL-16936 THESIS R B BELL (LRL) IJP
 FENSTER 65 PRL 17 841 *GELFAND, HARMSEN, L-SETTI, + (CHI, ARG, CERN) IJP
 ARMENTER 67 PL 249 198 ARMENTEROS, FERRO-LUZZI + (CERN, HEID, SACLAY) IJP
 ARMENT-2 67 ZEIT, PHYS. 202, 486 ARMENTEROS, FERRO-LUZZI + (CERN, HEID, SACLAY) IJP
 UNLIG 67 PR 155 1448 *CHARLTON, CONDON, GLASSER, YODH, + (MD, USNL)

ARMENT-1 6R NP 89 195 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
 ARMENT-2 6R NP 89 216 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) I
 BUGG 6R PR 169 1466 *GILMORE, KNIGHT, DAVIES + (BIRM, CAMB, RUTH) I
 SIMS 6R PL 21 1413 SIMS, ALBRIGHT, BARTLEY, MFER + (FLO+TAFS, ORR) I
 SMART 6R PR 169 1330 W M SMART (LRL) IJP
 BRICMANI 70 PL 338 511 *FERRO-LUZZI, LAGNAUX (CERN)
 COOL 70 PR 10 1887 *GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJP
 CONFORTO 70 NP (SUBMITTED) *HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 KIM 71 HARVARD PREPRINT J. K. KIM (HARV) IJP
 SEE ALSO DUKE 161 (HARV) IJP

PAPERS NOT REFERRED TO IN DATA CARDS

CONFORTO 6R NP 89 265 *HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 -- CONFORTO 6R IS SUPERSEDED BY CONFORTO 70.

$\Sigma(1880)$

67 *Y*(1880, JP=1/2+) I=1 **P₁₁**
 SEE THE MINI-REVIEW AT THE START OF THE *Y* LISTINGS.
 SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

----- 67 *Y*(1880) MASS (MEV) -----

M	1882.0	40.0	SMART	6R DPWA 0- K-N TO LAM PI	7/6R
W	(1880.0)		BAILEY	69 DPWA 0 ELASTIC, CH EXCH	10/70*
W	1920.0	30.0	LITCHFIELD	70 DPWA 0- K-N TO LAM PI	6/70
W	AROUT 1850.0		ARMENTERO	70 IPWA 0- ELASTIC, CH EXCH	6/70
M	1950.0	50.0	GALTIERI	70 DPWA 0- K-N TO LAM PI	7/70
M	AVG	194.5	21.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

----- 67 *Y*(1880) WIDTH (MEV) -----

W	222.0	150.0	SMART	6R DPWA 0- K-N TO LAM PI	7/6R
W <td>(200.0)</td> <td></td> <td>BAILEY</td> <td>69 DPWA 0 ELASTIC, CH EXCH</td> <td>10/70*</td>	(200.0)		BAILEY	69 DPWA 0 ELASTIC, CH EXCH	10/70*
W <td>170.0</td> <td>40.0</td> <td>LITCHFIELD</td> <td>70 DPWA 0- K-N TO LAM PI</td> <td>6/70</td>	170.0	40.0	LITCHFIELD	70 DPWA 0- K-N TO LAM PI	6/70
W <td>AROUT 30.0</td> <td></td> <td>ARMENTERO</td> <td>70 IPWA 0- ELASTIC, CH EXCH</td> <td>6/70</td>	AROUT 30.0		ARMENTERO	70 IPWA 0- ELASTIC, CH EXCH	6/70
W <td>200.0</td> <td>50.0</td> <td>GALTIERI</td> <td>70 DPWA 0- K-N TO LAM PI</td> <td>7/70</td>	200.0	50.0	GALTIERI	70 DPWA 0- K-N TO LAM PI	7/70
W <td>AVG</td> <td>193.4</td> <td>30.5</td> <td>AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)</td> <td></td>	AVG	193.4	30.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	

----- 67 *Y*(1880) PARTIAL DECAY MODES -----

P1	*Y*(1880) INTO KBAR N	497+ 939	DECAY MASSES
P2	*Y*(1880) INTO LAMBDA PI	1115+ 134	

67 *Y*(1880) BRANCHING RATIOS

BELOW, X1 = (PARTIAL-WIDTH 1)/TOTAL WIDTH, ETC. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE *Y* MINI-REVIEW.

R1	*Y*(1880) INTO (KBAR N)/TOTAL	X1		
R1	(0.20)	BAILEY	69 DPWA 0 ELASTIC, CH EXCH 10/70*	
R1		ARMENTERO	70 IPWA 0- ELASTIC, CH EXCH 6/70	
R2	*Y*(1880) FROM KBAR N INTO LAMBDA PI	SORT(X1*X2)		
R2	-0.14	0.03	SMART	6R DPWA 0- K-N TO LAM PI 7/6R
R2	-0.09	0.04	LITCHFIELD	70 DPWA 0- K-N TO LAM PI 6/70
R2	-0.09	0.04	GALTIERI	70 DPWA 0- K-N TO LAM PI 7/70
R2	AVG MOD	0.117	0.019	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)

REFERENCES -- *Y*(1880)

SMART 6R PR 169 1330 W M SMART (LRL) IJP
 BAILEY 69 THESIS UCL-50617 DAVID SAAL BAILEY (LRL LIVERMORE) IJP
 LITCHIE 70 NP 822 269 P J LITCHFIELD (RUTHERFORD) IJP
 ARMENTER 70 DUKE CONF 123 ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJP
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJP

$\Sigma(1915)$

46 *Y*(1915, JP=5/2+) I=1 **F₁₅**
 SEE THE MINI-REVIEW AT THE START OF THE *Y* LISTINGS.
 THIS RESONANCE WAS FIRST SEEN IN THE TOTAL-CROSS-SECTION MEASUREMENTS OF CODL 66. IN THIS ENTRY, HOWEVER, WE LIST ONLY THE RESULTS FROM PARTIAL-WAVE ANALYSES. SEE THE NEXT ENTRY FOR THE PARAMETERS OF PEAKS SEEN AROUND 1900-1950 MEV IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. WE MAKE THIS SEPARATION BECAUSE ONLY THE PARTIAL-WAVE ANALYSES ISOLATE THE F15 WAVE (OR AT LEAST ATTEMPT TO -- THE SIGNAL IS WEAK). THIS MASS REGION IS COMPLICATED AND POORLY UNDERSTOOD AND THE PEAKS MAY CONTAIN MORE THAN JUST THE *Y*(1915). SEE ALSO THE NOTE TO THE NEXT ENTRY.

----- 46 *Y*(1915) MASS (MEV) -----

M	1902.0	11.0	SMART	6R DPWA 0- K-N TO LAMBDA PI	7/6R
W	1903.0	10.0	COX	70 DPWA 0- K-N TO LAMBDA PI	6/70
W	1910.0	20.0	BERTHON	70 DPWA 0 K-P TO LAMBDA PI	7/70
W	1900.0	15.0	BERTHON	70 DPWA 0 K-P TO SIGMA PI	10/70*
M	N 1936.0	(3.0)	BRICMANI	70 DPWA 0 SIGTOT, ELAS, CHEX	1/71*
M	1905.0	30.0	GALTIERI	70 DPWA 0 K-P TO LAMBDA PI	7/70
M	1895.0	10.0	LITCHFIELD	70 DPWA 0- K-N TO LAMBDA PI	6/70
M	N	ERROR STATIST. ONLY- NO ERROR DUE TO PARTICULAR P.W. ANAL.	INCLUDED	1/71*	
M	AVG	MEANINGLESS (SCALE FACTOR = 1.0)			

----- 46 *Y*(1915) WIDTH (MEV) -----

W	A (50.0)	(20.0)	ARMENTERI <th>67 DPWA 0 ELASTIC, CH EXCH</th> <th>11/67</th>	67 DPWA 0 ELASTIC, CH EXCH	11/67
W	52.0	25.0	SMART	6R DPWA 0- K-N TO LAMBDA PI	7/6R
W	60.0	20.0	BERTHON	70 DPWA 0 K-P TO LAMBDA PI	7/70
W	75.0	20.0	BERTHON	70 DPWA 0 K-P TO SIGMA PI	10/70*
W	135.0	12.0	BRICMANI	70 DPWA 0 SIGTOT, ELAS, CHEX	1/71*
W	77.0	27.0	COX	70 DPWA 0- K-N TO LAMBDA PI	6/70
W	70.0	20.0	GALTIERI	70 DPWA 0 K-P TO LAMBDA PI	7/70
W	70.0	15.0	LITCHFIELD	70 DPWA 0- K-N TO LAMBDA PI	6/70
W	A	LACK OF DATA PREVENTS FROM DETERMINING UNAMB. THIS AMPLITUDE			11/67
W	AVG	MEANINGLESS (SCALE FACTOR = 1.9)			

46 *Y*(1915) PARTIAL DECAY MODES

P1	*Y*(1915) INTO KBAR N	497+ 939	DECAY MASSES
P2	*Y*(1915) INTO LAMBDA PI	1115+ 139	
P3	*Y*(1915) INTO SIGMA PI	1174+ 139	

46 *Y*(1915) BRANCHING RATIOS

BELOW, X1 = (PARTIAL-WIDTH 1)/TOTAL WIDTH, ETC. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE *Y* MINI-REVIEW.

R1	*Y*(1915) INTO (KBAR N)/TOTAL	X1		
R1	A (0.12)	(.01)	ARMENTERI	67 DPWA 0 ELASTIC, CH EXCH 11/67
R1	0.18	(0.02)	BRICMANI	70 DPWA 0 SIGTOT, ELAS, CHEX 1/71*
R1	0.11	(0.03)	CONFORTO	70 DPWA 0 ELASTIC, CH EXCH 6/70
R2	*Y*(1915) FROM KBAR N INTO LAMBDA PI	SORT(X1*X2)		
R2	-0.08	(0.02)	SMART	6R DPWA 0- K-N TO LAMBDA PI 7/6R
R2	-0.1	(0.02)	BERTHON	70 DPWA 0 K-P TO LAMBDA PI 7/70
R2	-0.09	(0.02)	COX	70 DPWA 0- K-N TO LAMBDA PI 6/70
R2	-0.11	(0.03)	GALTIERI	70 DPWA 0 K-P TO LAMBDA PI 7/70
R2	-0.07	(0.015)	LITCHFIELD	70 DPWA 0- K-N TO LAMBDA PI 6/70
R3	*Y*(1915) FROM KBAR N INTO SIGMA PI	SORT(X1*X3)		
R3	A (0.00)	(0.01)	ARMENTERO	67 DPWA 0 K-P TO SIGMA PI 11/67
R3	-0.13	(0.03)	BERTHON	70 DPWA 0 K-P TO SIGMA PI 10/70*
R3	-0.06	(0.03)	GALTIERI	70 DPWA 0 K-P TO SIGMA PI 7/70

REFERENCES -- *Y*(1915)

ARMENTER 67 PL 249 198 ARMENTEROS, FERRO-LUZZI + (CERN, HEID, SACLAY)
 ARMENTER 67 NP 83 592 ARMENTEROS, FERRO-LUZZI + (CPN, HEID, SACLAY)
 SMART 6R PR 169 1330 W M SMART (LRL) IJP
 BERTHON 70 NP 824 417 *RANGAN, VRANA, + (COL FRANCE, RHEL, SACLAY) IJP
 BRICMANI 70 PL 338 511 *FERRO-LUZZI, LAGNAUX (CERN)
 CONFORTO 70 NP (SUBMITTED) *HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
 COX 70 NP 819 61 *ISLAM, GOLLEY, + (BIRM, EDIN, GLAS, LOIC) IJP
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJP
 LITCHFIELD 70 NP 822 269 P J LITCHFIELD (RUTHERFORD) IJP

PAPERS NOT REFERRED TO IN DATA CARDS

SMART 66 PRL 17 556 W M SMART, A KERNAN, G E KALMUS, R P FLY (LRL) IJP
 -- SMART 66 IS SUPERSEDED BY SMART 68.
 CONFORTO 6R NP 89 265 *HARMSEN, LASINSKI, + (CHICAGO, HEIDEL)
 -- CONFORTO 6R IS SUPERSEDED BY CONFORTO 70.

$\Sigma(1900)$ BUMPS

29 *Y*(1900-1950) CROSS-SECTION AND PRODUCTION PEAKS
 SEE THE MINI-REVIEW AT THE START OF THE *Y* LISTINGS.
 SEE THE NOTES TO THE *Y*(1915) AND *Y*(1940), WHICH IMMEDIATELY PRECEED AND FOLLOW THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS SEEN IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE MOST CERTAINLY ASSOCIATED WITH THE F15 *Y*(1915) SEEN IN PARTIAL-WAVE ANALYSES. THE INVARIANT-MASS PEAKS SEEM MORE LIKELY TO BE ASSOCIATED WITH THE NOT-COMpletely-ESTABLISHED D13 *Y*(1940).

----- 29 *Y*(1900-1950) MASS (MEV) -----

M	CROSS-SECTION PEAKS --			
M	1905.0	5.0	RUGG	6R CNTR K-P, D TOTAL 11/66
M	1906.0	6.0	BRICMAN	70 CNTR 0 TOTAL AND CH EX 6/70
M	1912.0	10.0	COOL	70 CNTR K-P, D TOTAL 10/70*
M	INVARIANT-MASS-DISTRIBUTION	PEAKS --		
M	(1942.0)	(9.0)	ROCK	65 HBC PBAR P 5.7 BEV/C
M	1940.0	11.0	AGUILAR	70 HBC + 3.9-6.6 GEV/C K-
M	AVG	MEANINGLESS (SCALE FACTOR = 1.7)		

Baryons

For notation, see illustrated key at beginning of data card listings.

29 Y*(1900-1950) WIDTH (MEV)

CROSS-SECTION PEAKS --		RUGG	68 CNTR	11/66
W	40.0	10.0	BRICMAN	70 CNTR 0 TOTAL AND CH EX 6/70
W	50.0	12.0	COHL	70 CNTR K-P, D TOTAL 10/70*
W	130.0	10.0	COHL	70 CNTR K-P, D TOTAL 10/70*
INVARIANT-MASS-DISTRIBUTION PEAKS --		ROCK	65 HBC	5/70
W	(36.0)	(20.0)	(36.0)	AGUILAR
W	90.0	20.0	AGUILAR	70 HRC + 3.9-4.6 GEV/C K-
W	90.0	20.0	AGUILAR	70 HRC + 3.9-4.6 GEV/C K-
W	90.0	20.0	AGUILAR	70 HRC + 3.9-4.6 GEV/C K-
AVERAGE MEANINGLESS (SCALE FACTOR = 1.2)				

29 Y*(1900-1950) PARTIAL DECAY MODES

DECAY MASSES		497+ 939
P1	Y*(1900-1950) INTO KBAR N	1115+ 134
P2	Y*(1900-1950) INTO LAMBDA PI	1197+ 139
P3	Y*(1900-1950) INTO SIGMA PI	

29 Y*(1900-1950) BRANCHING RATIOS

R1	Y*(1900-1950) INTO (KBAR N)/TOTAL	(P1)/TOTAL
R1	THESE VALUES OF ELASTICITIES ASSUME J=5/2 --	
R1	COHL 0.04	RUGG 48 CNTR ASSUMING J=5/2 6/68
R1	0.07	BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
R1	0.07	COHL 70 CNTR K-P, D TOTAL 10/70*
R2	Y*(1900-1950) INTO (KBAR N)/(SIGMA PI)	(P1)/(P3)
R2	(.37) OR LESS	BARNES 69 HBC + 1 STAN. DEV. 10/69
R3	Y*(1900-1950) INTO (LAMBDA PI)/(SIGMA PI)	(P2)/(P3)
R3	(.29) OR LESS	BARNES 69 HBC + 1 STAN. DEV. 10/69

REFERENCES -- Y*(1900-1950)

ROCK 65 PL 17 166 *COOPER, FRENCH, KINSON, + (CERN, SACLAY) I
 COHL 66 PRL 16 1228 *GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, + (BNL) I
 -- COHL 66 IS SUPERSEDED BY COHL 70.
 RUGG 68 PR 16A 1466 *GILMORE, KNIGHT, DAVIES + (BIRM, CAMB, RUTHI) I
 BARNES 69 PRL 22 479 *FLAMINI, MONTANFI, SAMIOS + (BNL-SYR) I
 BRICMAN 70 PL 31B 152 *FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY) I
 COHL 70 PR 19 1887 *GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 AGUILAR 70 PRL 25 58 *AGUILAR-BENITEZ, BARNES, + (BNL, SYR) I

PAPERS NOT REFERRED TO IN DATA CARDS

PRIMER 68 PRL 20 610 *GOLDBERG, JAEGER, BARNES, DORNAN + (SYR, BNL) I
 -- PRIMER 68 IS SUPERSEDED BY BARNES 69 AND AGUILAR-BENITEZ 70.

98 Y*(1940, JP=3/2-1 I=1) **D₁₃**

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE. THIS EFFECT IS PERHAPS ASSOCIATED WITH THE BUMPS SEEN IN PRODUCTION EXPERIMENTS NEAR THIS MASS. SEE THE PRECEDING ENTRY.

98 Y*(1940) MASS (MEV)

M	1940.0	50.0	GALTIERI	70 DPWA	K-N TO LAM PI	7/70
M	1940.0	40.0	GALTIERI	70 DPWA	K-P TO SIGMA PI	7/70
M	1940.0	30.0	LITCHFIEL	70 DPWA	K-N TO LAM PI	7/70
M	1940.0	21.6	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			

98 Y*(1940) WIDTH (MEV)

W	200.0	50.0	GALTIERI	70 DPWA	K-N TO LAM PI	7/70
W	200.0	50.0	GALTIERI	70 DPWA	K-P TO SIGMA PI	7/70
W	200.0	40.0	LITCHFIEL	70 DPWA	K-N TO LAM PI	7/70
W	235.1	28.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)			

98 Y*(1940) PARTIAL DECAY MODES

DECAY MASSES		497+ 939
P1	Y*(1940) INTO KBAR N	1115+ 139
P2	Y*(1940) INTO LAMBDA PI	1197+ 139
P3	Y*(1940) INTO SIGMA PI	

98 Y*(1940) BRANCHING RATIOS

BELOW, X1 = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1	Y*(1940) FROM KBAR N INTO LAMBDA PI	SQRT(X1*X2)
R1	-0.12	0.04
R1	-0.14	0.03
R1	AVG MOD	0.133
R1	0.133	0.024
AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R2	Y*(1940) FROM KBAR N INTO SIGMA PI	SQRT(X1*X3)
R2	-0.12	0.03

REFERENCES -- Y*(1940)

GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRLI)JJP
 LITCHFIE 70 NP R22 269 P J LITCHFIELD (RUTHERFORD)JJP

98 Y*(1940, JP=7/2+ I=1) **F₁₇**

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES. PARAMETERS OF PEAKS SEEN IN CROSS-SECTIONS AND INVARIANT-MASS DISTRIBUTIONS AROUND 2030 MEV ARE GIVEN IN THE NEXT ENTRY. EVENTUALLY THE PARTIAL-WAVE ANALYSES SHOULD GIVE THE BEST RESULTS, AS THEY ISOLATE THE F17 WAVE. THIS SUPERIORITY IS, HOWEVER, PROBABLY NOT YET ATTAINED, AND WE RELY ON BOTH ENTRIES FOR PARAMETERS GIVEN IN THE MAIN BARYON TABLE.

47 Y*(2030) MASS (MEV)

M	(2030.0)	(20.0)	WOHL	66 HBC 0	K-P TO LAM PTO	7/66
M	2032.0	6.0	SMART	68 DPWA	K-N TO LAMBDA PI	6/68
M	2030.0	10.0	BERTHON	70 DPWA 0	K-P TO LAMBDA PI	7/70
M	2035.0	10.0	BERTHON	70 DPWA 0	K-P TO SIGMA PI	10/70*
M	2027.0	6.0	COX	70 DPWA	K-N TO LAMBDA PI	6/70
M	2010.0	15.0	GALTIERI	70 DPWA 0	K-P TO LAMBDA PI	7/70
M	2000.0	20.0	GALTIERI	70 DPWA 0	K-P TO SIGMA PI	7/70
M	2022.0	4.0	LITCHFIE	70 DPWA 0	K-N TO LAMBDA PI	6/70
AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)						

47 Y*(2030) WIDTH (MEV)

W	(170.0)	16.0	WOHL	66 HBC 0	7/66
W	160.0	16.0	SMART	68 DPWA	K-N TO LAMBDA PI 6/68
W	165.0	30.0	15.0	BERTHON	70 DPWA 0 K-P TO LAMBDA PI 7/70
W	150.0	20.0	BERTHON	70 DPWA 0	K-P TO SIGMA PI 10/70*
W	158.0	16.0	COX	70 DPWA	K-N TO LAMBDA PI 6/70
W	115.0	15.0	GALTIERI	70 DPWA 0	K-P TO LAMBDA PI 7/70
W	100.0	40.0	GALTIERI	70 DPWA 0	K-P TO SIGMA PI 7/70
W	170.0	15.0	LITCHFIE	70 DPWA 0	K-N TO LAMBDA PI 6/70
AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)					

47 Y*(2030) PARTIAL DECAY MODES

DECAY MASSES		497+ 939
P1	Y*(2030) INTO KBAR N	1115+ 134
P2	Y*(2030) INTO LAMBDA PI	1197+ 139
P3	Y*(2030) INTO SIGMA PI	1321+ 497
P4	Y*(2030) INTO XI K	

47 Y*(2030) BRANCHING RATIOS

R1	Y*(2030) INTO (KBAR N)/TOTAL	X1
R1	(0.25)	WOHL 66 HBC 0 K-P CH EX 7/66
R1	0.17	0.04 CAMPBELL 71 DRC - K- NEUTRON ELAST 1/71*
R1	0.11	DAUM 68 CNTR - K-P ELA, POL, SIGT 7/70
R1	DAUM 68 ASSUMES (J1/2)*X VALUE SEEN IN TOTAL CROSS SECTION.	
R2	Y*(2030) FROM KBAR N INTO LAMBDA PI	SQRT(X1*X2)
R2	+0.21	0.01 WOHL 66 HBC 0 K-P TO LAMBDA PI 7/66
R2	+0.19	0.01 SMART 68 DPWA - K-N TO LAMBDA PI 6/68
R2	+0.2	0.02 BERTHON 70 DPWA 0 K-N TO LAMBDA PI 7/70
R2	+0.20	0.00R LITCHFIE 70 DPWA 0 K-N TO LAMBDA PI 6/70
R2	+0.16	0.03 GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)		
R3	Y*(2030) FROM KBAR N INTO SIGMA PI	SQRT(X1*X3)
R3	-0.052	0.010 GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
R3	-0.09	0.02 BERTHON 70 DPWA 0 K-P TO SIGMA PI 10/70*
AVERAGE MEANINGLESS (SCALE FACTOR = 1.7)		
R4	Y*(2030) FROM KBAR N INTO XI K	SQRT(X1*X4)
R4	(0.05) OR LESS	TRIPP 67 RVUE 0 K-P TO XI K 8/67
R4	(0.05) OR LESS	BURGIN 68 DPWA 0 K-P TO XI K 10/69
R4	(0.023)	MULLER 69 DPWA 0 7/70

REFERENCES -- Y*(2030)

WOHL 65 PRL 17 107 C G WOHL, F Y SOUMITZ, M L STEVENSON (LRL)JJP
 TRIPP 67 NP R3 10 + LEITH, (LRL, SLAC, CERN, HEIDEL, SACLAY) I
 BURGIN 68 NP BR 447 *MEYER, PAULI, TALLINI + (SACL+CF+RHEL) I
 -- DAUM 68 CONFIRMS THE SPIN-PARITY ASSIGNMENT.
 DAUM 68 NP R7 19 *ERNE, LAGNAUX, SENS, STEUER, UDD (CERN)JJP
 SMART 68 PR 169 1336 W M SMART (LRL)JJP
 MULLER 69 THESIS, UCRL 19372 R A MULLER (LRL)JJP
 COX 70 NP R19 61 *ISLAM, COLLEY, + (BIRM, EDIN, GLAS, LOIC)JJP
 BERTHON 70 NP R20 476 *RANGAN, VRANA, + (COL FRANCE, RHEL, SACLAY)JJP
 LITCHFIE 70 NP R22 269 P J LITCHFIELD (RUTHERFORD)JJP
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL)JJP
 BERTHON 70 NP R24 417 *VRANA, RUTHERWORTH, + (COEF, RHEL, SACLAY)JJP
 CAMPBELL 71 NP R25 75 *MORTON, NEGUS, GOYAL, MILLER (GLAS, LOIC)JJP

47 Y*(2030) CROSS-SECTION AND INVARIANT-MASS PEAKS

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE THE NOTE TO THE F17 Y*(2030), WHICH PRECEEDS THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE Y*(2030), BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

28 Y*(2030) MASS (MEV) -- AS SEEN IN PEAKS

M	(2022.0)	(20.0)	BLANPIED	68 CNTR 0	GAMMA P TO K+ Y*	6/68
M	2026.0	7.0	RUGG	68 CNTR	K-P, D TOTAL	6/70
M	2024.0	4.0	BRICMAN	70 CNTR 0	TOTAL AND CH EX	6/70
M	2025.0	10.0	COHL	70 CNTR	K-P, D TOTAL	10/70*
M	(2025.0)	(20.0)	LU	70 CNTR 0	GAMMA P TO K+ Y*	1/71*
AVERAGE MEANINGLESS (SCALE FACTOR = 2.8)						

28 Y*(2030) WIDTH (MEV) -- AS SEEN IN PEAKS

W	(120.0)	(20.0)	BLANPIED	65 CNTR 0	6/68
W	130.0	10.0	RUGG	68 CNTR	6/70
W	126.0	11.0	BRICMAN	70 CNTR 0	TOTAL AND CH EX 6/70
W	165.0	10.0	COHL	70 CNTR	K-P, D TOTAL 10/70*
W	(80.0)		LU	70 CNTR 0	GAMMA P TO K+ Y* 1/71*
AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)					

28 Y*(2030) PARTIAL DECAY MODES

DECAY MASSES		497+ 939
P1	Y*(2030) INTO KBAR N	497+ 939+ 139
P2	Y*(2030) INTO KBAR N PI	

Σ(2030)

Σ(2030) BUMPS

Baryons

For notation, see illustrated key at beginning of data card listings.

28 Y*(2030) BRANCHING RATIOS -- AS SEEN IN PEAKS

R1 Y*(2030) INTO (KBAR N)/TOTAL XI
 THESE VALUES OF ELASTICITIES ASSUME J=7/2 --
 R1 0.131 RUGG 68 CNTR 6/68
 R1 0.27 (0.02) BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
 R1 0.12 COOL 70 CNTR K-P, D TOTAL 10/70*

R2 Y*(2030) INTO KBAR N PI P2
 R2 SPFN HOCK HRC

 REFERENCES -- Y*(2030) AS SEEN IN PEAKS
 BLANPIED 65 PRL 14 741 +GREENBERG, HUGHES, KITCHING, LU, + (YALF(CFA))
 COOL 66 PRL 16 1229 +GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, + (BNL) I
 -- COOL 66 IS SUPERSEDED BY COOL 70.
 BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RTHFD, BRGHM, CVNDSH) I
 BRICMAN 70 PR 318 152 +FERRO LUZZI, PERRÉAU, + (CERN, CAEN, SACLAY) I
 COOL 70 PR 10 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 LU 70 PR 02 1946 +GREENBERG, HUGHES, MINEHART, MORI, + (YALF)

 REFERENCES -- Y*(2070)

Σ(2070) 34 Y*(2070, JP=5/2+) I=1 **F₁₅**
 THIS STATE HAS BEEN SUGGESTED BY ONLY ONE PARTIAL
 WAVE ANALYSIS ACROSS THIS REGION. IT NEEDS CONFIRMATION

34 Y*(2070) MASS (MEV) 1/71*
 M (2070.) (10.) BERTHONI 70 DPWA - K- P TO SIG PI
 34 Y*(2070) WIDTH (MEV) 1/71*
 W (140.) (20.) BERTHONI 70 DPWA - K- P TO SIG PI
 34 Y*(2070) PARTIAL DECAY MODES

P1 Y*(2070) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(2070) INTO SIGMA PI 1197+ 139
 BELOW, XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

34 Y*(2070) BRANCHING RATIOS 1/71*
 R2 Y*(2070) FROM KBAR N TO SIGMA SORT(X1*X2)
 R2 (+.12) (.02) BERTHONI 70 DPWA - K- P TO SIG PI

 REFERENCES -- Y*(2070)
 BERTHONI 70 NP B24 417 +VRANA, BUTTERNORTH, + (COEF, RHEL, SACLAY) JIP

Σ(2080) 88 Y*(2080, JP=3/2+) I=1 **P₁₃**
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL
 PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE
 IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN
 BARYON TABLE.

88 Y*(2080) MASS (MEV) 6/70
 M (2082.0) (4.0) COX 70 DPWA - K- N TO LAM PI
 M (2070.0) (30.0) LITCHFIELD 70 DPWA 0- K- N TO LAM PI 6/70
 88 Y*(2080) WIDTH (MEV) 6/70
 W (87.0) (20.0) COX 70 DPWA - K- N TO LAM PI
 W (250.0) (40.0) LITCHFIELD 70 DPWA 0- K- N TO LAM PI 6/70
 88 Y*(2080) PARTIAL DECAY MODES

P1 Y*(2080) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(2080) INTO LAMBDA PI 1115+ 139
 88 Y*(2080) BRANCHING RATIOS
 BELOW, XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.

R1 Y*(2080) FROM KBAR N TO LAMBDA PI SORT(X1*X2)
 R1 (-0.16) (0.03) COX 70 DPWA - K- N TO LAM PI 6/70
 R1 (-0.09) (0.03) LITCHFIELD 70 DPWA 0- K- N TO LAM PI 6/70

 REFERENCES -- Y*(2080)
 COX 70 NP R19 61 +ISLAH, COLLEY, + (BIRM, EDIN, GLAS, LOIC) JIP
 LITCHFIELD 70 NP R22 269 P J LITCHFIELD (RUTHERFORD) JIP

Σ(2100) 26 Y*(2100, JP=7/2-) I=1 **G₁₇**
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL
 PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE
 IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN
 BARYON TABLE.

26 Y*(2100) MASS (MEV) 7/70
 M (2060.0) (20.0) GALTIERI 70 DPWA 0 K-P TO LAMBDA PI
 M (2120.0) (30.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI

26 Y*(2100) WIDTH (MEV)
 W (70.0) (30.0) GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
 W (135.0) (30.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
 26 Y*(2100) PARTIAL DECAY MODES
 P1 Y*(2100) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(2100) INTO LAMBDA PI 1115+ 139
 P3 Y*(2100) INTO SIGMA PI 1197+ 139

26 Y*(2100) BRANCHING RATIOS
 BELOW, XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.
 R1 Y*(2100) FROM KBAR N TO LAMBDA PI SORT(X1*X2)
 R1 (-0.07) (0.02) GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
 R2 Y*(2100) FROM KBAR N TO SIGMA PI SORT(X1*X3)
 R2 (+0.13) (0.02) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70

 REFERENCES -- Y*(2100)
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) JIP

M > 2200 MEV - PRODUCTION AND σ_{TOTAL} EXPERIMENTS

BUMPS

48 Y*(2250, JP=) I=1
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THE PHASE-SHIFT-ANALYSIS RESULTS ARE TOO WEAK TO
 WARRANT SEPARATING THEM FROM THE PRODUCTION AND CROSS-
 SECTION EXPERIMENTS. IN AN ANALYSIS OF ELASTIC AND
 POLARIZATION DATA, OUM 68 COULD NOT EXCLUDE ANY
 POSSIBILITY FOR JP= 5/2+ TO JP= 11/2+ FOR THIS STATE. BRICMAN 70
 SUGGESTS 7/2-.

48 Y*(2250) MASS (MEV)
 M (2245.0) BLANPIED 65 CNTR GAMMA P TO K+ Y*
 M (2299.0) (6.0) BOCK 65 HRC PBAR P 5.7 REV/C
 M 2250.0 7.0 RUGG 68 CNTR K-P, D TOTAL 6/68
 M 2237.0 11.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
 M 2255.0 10.0 COOL 70 CNTR K-P, D TOTAL 10/70*
 M 2280. 14.0 AGUILAR 70 HRC + K- 3.9-4.6 GEV/C 5/70
 M (2250.0) (20.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71*
 M AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)

48 Y*(2250) WIDTH (MEV)
 W (150.0) BLANPIED 65 CNTR GAMMA P TO K+ Y*
 W (21.0) (17.0) (21.0) BOCK 65 HRC PBAR P 5.7 GEV/C
 W 230.0 20.0 RUGG 68 CNTR K-P, D TOTAL 6/68
 W 194.0 50.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
 W (170.0) COOL 70 CNTR K-P, D TOTAL 10/70*
 W 100.0 20.0 AGUILAR 70 HRC + K- 3.9-4.6 GEV/C 5/70
 W (125.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71*
 W AVERAGE MEANINGLESS (SCALE FACTOR = 3.3)

48 Y*(2250) PARTIAL DECAY MODES
 P1 Y*(2250) INTO KBAR N DECAY MASSES 497+ 939
 P2 Y*(2250) INTO LAMBDA PI 1115+ 139
 P3 Y*(2250) INTO SIGMA PI 1197+ 139
 P4 Y*(2250) INTO KBAR N PI 497+ 939+ 139

48 Y*(2250) BRANCHING RATIOS
 BELOW, XI = (PARTIAL-WIDTH I)/TOTAL WIDTH, ETC. A SIGN, WHERE
 EXPLICIT, IS IN THE CONVENTION DISCUSSED IN THE Y* MINI-REVIEW.
 R1 Y*(2250) INTO (KBAR N)/TOTAL XI
 R1 J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*XI.
 R1 (0.47) BUGG 68 CNTR 6/68
 R1 (0.16) (0.12) BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
 R1 (0.42) COOL 70 CNTR K-P, D TOTAL 10/70*

R2 Y*(2250) FROM KBAR N TO LAMBDA PI SORT(X1*X2)
 R2 THE FOLLOWING ASSUMES JP=9/2-. DATA INSUF. FOR DETERM. THIS AMP.
 R2 (-0.18) GALTIERI 70 DPWA K-P TO LAMBDA PI 10/70*

R3 Y*(2250) FROM KBAR N TO SIGMA PI SORT(X1*X3)
 R3 THE FOLLOWING ASSUMES JP=9/2-. DATA INSUF. FOR DETERM. THIS AMP.
 R3 (+0.07) GALTIERI 70 DPWA K-P TO SIGMA PI 10/70*

R4 Y*(2250) INTO (KBAR N)/(SIGMA PI) (P11)/(P3)
 R4 (0.18) OR LESS BARNES 69 HRC + 1 STAN DEV LIMIT 10/69
 R5 Y*(2250) INTO (LAMBDA PI)/(SIGMA PI) (P21)/(P3)
 R5 (0.18) OR LESS BARNES 69 HRC + 1 STAN DEV LIMIT 10/69

 REFERENCES -- Y*(2250)
 BLANPIED 65 PRL 14 741 +GREENBERG, HUGHES, KITCHING, + (YALF(CFA))
 BOCK 65 PR 17 166 +CODDER, FRENCH, KINSON, + (CERN, SACLAY)
 BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RTHFD, BRGHM, CVNDSH) I
 BARNES 69 PR 22 479 +FLAMINI, MONTANET, SAMIOS, + (BNL+SYR)
 BRICMAN 70 PR 318 152 +FERRO LUZZI, PERRÉAU, + (CERN, CAEN, SACLAY)
 COOL 70 PR 10 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) JIP
 AGUILAR 70 PR 25 58 +AGUILAR-BENITEZ, BARNES, + (BNL, SYR)
 LU 70 PR 02 1946 +GREENBERG, HUGHES, MINEHART, MORI, + (YALF)

Baryons

For notation, see illustrated key at beginning of data card listings.

PAPERS NOT REFERRED TO IN DATA CARDS
 CODL 66 PRL 16 122B +GIACOMELLI, KYCIA, LEONTIC, LI, LUNDY, + (BNL) I
 -- CODL 66 IS SUPERSEDED BY CODL 70,
 DAUBER 66 PL 23 154 +SCHLEIN, SLATER, SYORR, TICHU (UCLA(LRL)) J
 -- SUGGESTS J=9/2 RESONANT BEHAVIOR IN SIGMA- P+, BUT APPEARS
 INCONSISTENT WITH PARAMETERS OF CODL 66.
 DAUM 69 NP 87 19 +FERRE, LAGNAUX, SENS, STEUER, UDD (CERN)JP

Σ(2455)

53 Y*(2455, JP=) I=1
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THERE IS ALSO SOME SLIGHT EVIDENCE FOR Y* STATES IN THIS MASS REGION FROM THE REACTION GAMMA + P TO K+ + MISSING MASS -- SEE GREENBERG 68.

----- 53 Y*(2455) MASS (MEV) -----						
M	2455.0	7.0	BUGG	68 CNTR	K-P, D TOTAL	6/68
M	2455.0	10.0	ABRAMS	70 CNTR	K-P, D TOTAL	10/70*
M
M	AVG	2455.0	5.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		

----- 53 Y*(2455) WIDTH (MEV) -----						
W	100.0	20.0	BUGG	68 CNTR	K-P, D TOTAL	6/68
W	140.0		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*

----- 53 Y*(2455) PARTIAL DECAY MODES -----						
P1	Y*(2455) INTO KBAR N			DECAY MASSES		
				497+ 939		

----- 53 Y*(2455) BRANCHING RATIOS -----						
P1	Y*(2455) INTO (KBAR N)/TOTAL			DECAY MASSES		
	J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*X1			497+ 939		
R1	0.31		BUGG	68 CNTR	K-P, D TOTAL	6/68
R1	0.39		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*
R1	(0.05)	(0.05)	BRICMAN	70 CNTR	D TOTAL AND CH EX	6/70*
R1	FIT OF TOTAL CROSS SECTION GIVEN BY BRICMAN 70 IS POOR IN THIS REGION.					

 REFERENCES -- Y*(2455)
 ABRAMS 70 PR 1D 1917 +CODL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
 BUGG 68 PR 16A 1466 +GILMORE, KNIGHT, + (RTHFD, BRGMH, CVNDNH) I
 BRICMAN 70 PL 31B 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
 PAPERS NOT REFERRED TO IN DATA CARDS
 ABRAMS 67 PRL 19 67R +CODL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 -- ABRAMS 67 IS SUPERSEDED BY ABRAMS 70.
 GREENBERG 68 PRL 20 221 GREENBERG, HUGHES, LU, MINEHART, + (YALE)

Σ(2620)

54 Y*(2620, JP=) I=1
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

----- 54 Y*(2620) MASS (MEV) -----						
M	2620.0	15.0	ABRAMS	70 CNTR	K-P, D TOTAL	10/70*

----- 54 Y*(2620) WIDTH (MEV) -----						
W	(175.0)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*

----- 54 Y*(2620) PARTIAL DECAY MODES -----						
P1	Y*(2620) INTO KBAR N			DECAY MASSES		
				497+ 939		

----- 54 Y*(2620) BRANCHING RATIOS -----						
P1	Y*(2620) INTO (KBAR N)/TOTAL			DECAY MASSES		
	J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*X1			497+ 939		
R1	0.36	0.12	BRICMAN	70 CNTR	D TOTAL AND CH EX	6/70
R1	(0.32)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70*

 REFERENCES -- Y*(2620)
 ABRAMS 67 PRL 19 67R +CODL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 -- ABRAMS 67 IS SUPERSEDED BY ABRAMS 70.
 ABRAMS 70 PR 1D 1917 +CODL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
 BRICMAN 70 PL 31B 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)

Σ(3000)

59 Y*(3000, JP=) I=1
 SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

ENHANCEMENT IN LAMBDA PI AND KBAR N INVARIANT MASS SPECTRA AND IN MISSING MASS OF NEUTRALS RECOILING AGAINST K0. EVIDENCE NOT CONCLUSIVE. OMITTED FROM TABLE.

----- 59 Y*(3000) MASS (MEV) -----						
M	(3000.0)		EHRlich	66 HBC	0 P1-P 7.91 BEV/C	9/66

----- 59 Y*(3000) PARTIAL DECAY MODES -----						
P1	Y*(3000) INTO KBAR N			DECAY MASSES		
P2	Y*(3000) INTO LAMBDA PI			497+ 939		
				1115+ 139		

REFERENCES -- Y*(3000)
 EHRlich 66 PR 152 1194 R EHRlich, W SELOVE, H YUTA (PENN(BNL)) I

EXOTIC HYPERONS CROSS SECTION LIMITS (MICROBARNS)

THIS IS NOT A COMPLETE LIST. WE WILL TABULATE EXOTICS FROM NOW ON
 CR G (20.) OR LESS GALTIERI 68 DRC -- K-N TO SG-PI-PI0 7/70
 G ABOVE LIMIT FOR MASS LT 2.15 GEV AND GAMMA LT 60 MEV- (2.1 GEV/C K-) 7/70
 CR A (40.) OR LESS GALTIERI 68 DRC -- K-N TO SG-PI-PI0 7/70
 CR A ABOVE LIMIT FOR MASS LT 2.3 GEV AND GAMMA LT 120 MEV- (2.7 GEV/C K-) 7/70

REFERENCES --- EXOTIC HYPERONS

GALTIERI 68 PRL 21 573 A. BARBARO-GALTIERI, CHADWICK + (LRL+SACLAY)

Ξ Resonances

The Ξ resonance situation has long been and will probably long remain unsettled. This is because 1) they can only be produced as part of a final state, K⁺p → Ξ* + others, and 2) they are so produced with very small cross sections (<50 μb). Thus the numbers of events available are small, and the analysis is more complicated than if direct formation were possible. Only the Ξ(1530) is really well established. There are at least two Ξ states in the 1800-2000 MeV region and there are indications of several more above 2000 MeV, but the situation is very unclear. We are forced to group together rather disparate observations and await new results. Figures in the listings point out disagreements among various experiments. The table following this note gives our evaluation of the status of the Ξ resonances, based on the meager data available at this time.

STATUS OF THE XI* RESONANCES. THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. SEE, HOWEVER, THE NOTES FOR THE INDIVIDUAL RESONANCES, PARTICULARLY FOR THE XI(1820) AND XI(1940), FOR ADDITIONAL WARNINGS.

PARTICLE	LJ	OVERALL STATUS	STATUS AS SEEN IN --			
			XI PI	LAM K	SIG K	XI* PI OTHER CHANNELS
XI(1320)	P11	****				LAMBDA PI
XI(1530)	P13	****				
XI(1630)		**	**			
XI(1820)		***	***	**	**	
XI(1940)		***	***	**	**	
XI(2030)		**	**	**	**	3-BODY DECAYS
XI(2250)		**	**	**	**	3-BODY DECAYS
XI(2500)		**	**	**	**	3-BODY DECAYS

**** GOOD, CLEAR, AND UNMISTAKABLE.
 *** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
 ** NEEDS CONFIRMATION.
 * WEAK.

Baryons

For notation, see illustrated key at beginning of data card listings.

Ξ^- 22 Ξ^- - (1321, JP=1/2⁻) I=1/2
SEE LISTINGS OF STARLE PARTICLES

Ξ^0 23 Ξ^0 (1314, JP=1/2⁻) I=1/2
SEE LISTINGS OF STARLE PARTICLES

$\Xi(1530)$ 49 $\Xi^*(1530)$, JP=3/2⁺ I=1/2 **P₁₃**
THIS IS THE ONLY REALLY WELL-ESTABLISHED Ξ^* . THE QUANTUM NUMBERS 3/2⁺ ARE FAVORED BY THE DATA.

WE DO NOT USE DETERMINATIONS OF THE MASS AND THE WIDTH OF THIS STATE UNLESS THEY ARE ACCOMPANIED BY SOME DISCUSSION OF SYSTEMATICS AND RESOLUTION.

49 $\Xi^*(1530)$ MASS (MEV)										
M	55(1529.0)	(5.0)	PJERRDU	62 HRC	0-	K-P	1.8	GEV/C		
M	(1532.0)	(2.0)	BADIER	64 HRC	0-	K-P	3	GEV/C		
M	38 1535.7	3.2	LONDON	66 HRC	-	K-P	2.24	GEV/C	7/66	
M	76 1529.7	1.1	LONDON	66 HRC	0	K-P	2.24	GEV/C	7/66	
49 $\Xi^*(1530)$ MASS DIFFERENCE (MEV)										
D	5.7	3.0	PJERRDU	65 HRC	0-	1.8-1.95	GEV/C	7/66		
D	R	(7.0)	(4.0)	LONDON	66 HRC	0-	2.24	GEV/C	7/66	
D	2.0	3.2	MERRILL	66 HRC	0-	1.7-2.7	GEV/C	7/66		
D	R	REFOUNDANT WITH DATA IN MASS LISTING.								
D	AVG	4.0	2.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
49 $\Xi^*(1530)$ WIDTH (MEV)										
W	7.0	2.0	SCHLEIN	63 HRC	0	1.8, 1.95	GEV/C			
W	8.5	3.5	LONDON	66 HRC	0	2.24	GEV/C	7/66		
W	7.0	7.0	BERGE	66 HRC	0	1.5-1.7	GEV/C	7/66		
W	AVG	7.3	1.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						

49 $\Xi^*(1530)$ PARTIAL DECAY MODES

PI $\Xi^*(1530)$ INTO Ξ PI 1321+ 139

OTHER STRONG DECAYS ARE FORBIDDEN BY ENERGY CONSERVATION.

REFERENCES -- $\Xi^*(1530)$

PJERRDU 52 PRL 9 114 +PROWSE, SCHLEIN, SLATER, STORK, TICHQ (UCLA) I
 SCHLEIN 63 PRL 11 167 +CARMONY, PJERRDU, SLATER, STORK, TICHQ (UCLA) JP
 BADIER 64 DUBNA I 593 +DEMOLIN, GOLDBERG, + (FP, SACLAY, AMSTR) I
 PJERRDU 65 PRL 14 275 +SCHLEIN, SLATER, SMITH, STORK, TICHQ (UCLA)
 BERGF 66 PR 147 945 +FRERHARD, HUBBARD, MERRILL, B-SHAFER, + (LRL) I
 LONDON 66 PR 143 1034 +RAU, SAMIINS, YAMAMOTO, GOLDBERG, + (BNL, SYCR) IJ
 MERRILL 66 UCRL-16455 THESIS D W MERRILL (LRL) JP

PAPERS NOT REFERRED TO IN DATA CARDS

SHAFER 66 PR 142 883 BUTTON-SHAFER, LINDSEY, MURRAY, SMITH (LRL) JP
 -- A SPIN-PARITY DETERMINATION.

$\Xi(1630)$ 21 $\Xi^*(1630)$, JP= 1 I=1/2
THIS IS A 3- OR 4-STANDARD-DEVIATION BUMP SEEN IN ONE CHANNEL IN ONE EXPERIMENT. BARTSCH 69 SEE A SMALL, BROAD ENHANCEMENT NEAR 1650 MEV - IT IS NOT CLEAR THAT IT IS THE SAME PHENOMENON AS $\Xi^*(1630)$. THIS EFFECT NEEDS CONFIRMATION.

21 $\Xi^*(1630)$ MASS (MEV)									
M	40 1635.	10.	BMST	70 HRC	0	INTO Ξ -PI+			7/70
21 $\Xi^*(1630)$ WIDTH (MEV)									
W	40 57.	18.	BMST	70 HRC	0				
21 $\Xi^*(1630)$ PARTIAL DECAY MODES									
PI	$\Xi^*(1630)$ INTO Ξ PI				1321+ 139				
SEEN IN K- P TO Ξ -PI+ K0.									

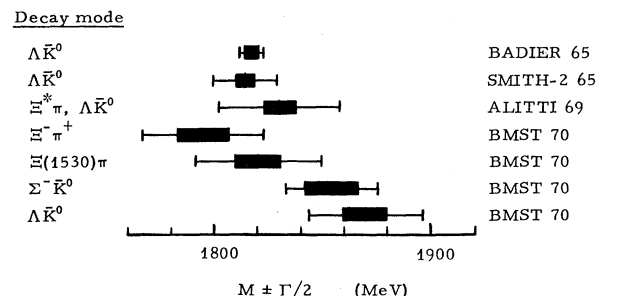
REFERENCES -- $\Xi^*(1630)$

BMST 70 DUKE 317 BRANDEIS+MARYLAND+SYRACUSE+TUFTS COLLAR.
 PAPERS NOT REFERRED TO IN DATA CARDS

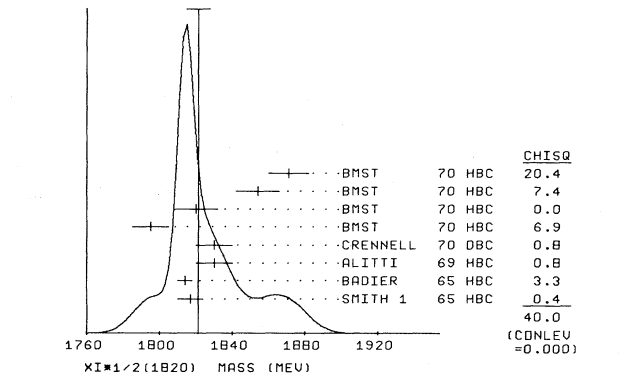
APSELL 69 PRL 23 884 + (BRANDEIS, MARYLAND, SYRACUSE, TUFTS)
 -- APSELL 69 IS SUPERSEDED BY BMST 70
 BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LONDON, VIENNA)
 KALBFLEI 70 DUKE CONF 331 G R KALBFLEISCH (BNL) I
 -- KALBFLEISCH 70 SUMMARIZES EVIDENCE FOR ISOSPIN 1/2.

$\Xi(1820)$ 50 $\Xi^*(1820)$, JP= 1 I=1/2
AS THE ACCOMPANYING IDEOGRAMS ILLUSTRATE, THE SITUATION IS CONFUSED. UNTIL SOME FUTURE CLARIFICATION, WE LIST UNDER $\Xi(1820)$ EVERYTHING REPORTED IN THE MASS RANGE 1750-1875 MEV. WHEN BRANCHING RATIOS ARE REPORTED, WE QUOTE THEM, BUT THERE ESPECIALLY ONLY THE MOST QUALITATIVE CONCLUSIONS (IF ANY AT ALL) ARE JUSTIFIED.

$\Xi(1820)$
Masses and widths of reported enhancements in the $\Xi(1820)$ region (solid rectangles indicate error on mass).



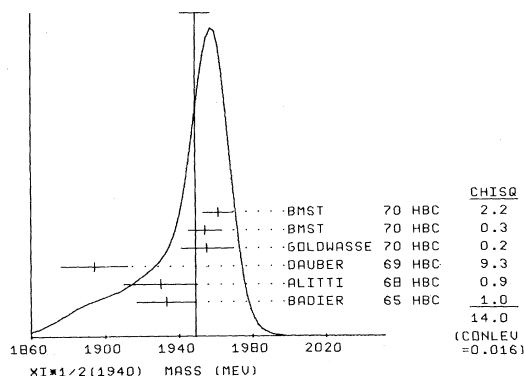
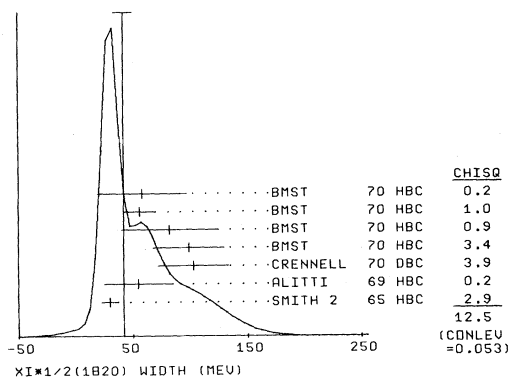
50 $\Xi^*(1820)$ MASS (MEV)									
M	(1770.0)		HALSTEINS	63 FRC	0-	K-FR	3.5	GEV/C	
M	29 1817.0	7.0	SMITH 1	65 HRC	0-	LAMBDA	KBAR		
M	30 1814.0	4.0	BADIER	65 HRC	0	LAMBDA	KBAR		
M	40 1830.0	10.0	ALITTI	69 HRC	-	LAM, SIG	KBAR	9/69	
M	25 1830.0	10.0	CRENNELL	70 DRC	0-	3.6, 3.9	GEV/C	10/70*	
M	65 1795.	10.	BMST	70 HRC	0	Ξ -PI+	(2.9 K-P)	7/70	
M	55 1820.	12.	BMST	70 HRC	Ξ (1530)	PI		7/70	
M	35 1854.	12.	BMST	70 HRC	-	SIGMA-	KOBAR	7/70	
M	65 1871.	11.	BMST	70 HRC	0	LAMBDA	KOBAR	7/70	
AVERAGE MEANINGLESS (SCALE FACTOR = 2.4) (SEE IDEOGRAM BELOW)									



50 $\Xi^*(1820)$ WIDTH (MEV)									
W	(80.0)	OR LESS	HALSTEINS	63 FRC	0-	K-FR	3.5	GEV/C	
W	(12.0)	(4.0)	BADIER	65 HRC	0	LAMBDA	KBAR		
W	30.0	7.0	SMITH 2	65 HRC	0-				
W	55.0	40.0	20.0	ALITTI	69 HRC	-	LAM, SIG	KBAR	9/69
W	103.0	38.0	24.0	CRENNELL	70 DRC	0-	3.6, 3.9	GEV/C	10/70*
W	65 99.	31.	BMST	70 HRC	Ξ -PI+	(2.9 K-P)		7/70	
W	55 82.	42.	BMST	70 HRC	Ξ (1530)	PI		7/70	
W	35 56.	14.	BMST	70 HRC	-	SIGMA-	KOBAR	7/70	
W	65 58.	39.	BMST	70 HRC	0	LAMBDA	KOBAR	7/70	
AVERAGE MEANINGLESS (SCALE FACTOR = 1.4) (SEE IDEOGRAM BELOW)									

Baryons

For notation, see illustrated key at beginning of data card listings.



50 XI*1/2(1820) PARTIAL DECAY MODES

Mode	Decay Masses
P1 XI*1/2(1820) INTO LAMBDA KBAR	1115+ 497
P2 XI*1/2(1820) INTO XI PI	1321+ 139
P3 XI*1/2(1820) INTO SIGMA KHAR	1197+ 497
P4 XI*1/2(1820) INTO XI*1/2(1530) PI	1530+ 139
P5 XI*1(1820) INTO XI PI PI (XI PI NOT XI*1(1530))	1321+ 139+ 139

50 XI*1/2(1820) BRANCHING RATIOS

Mode	Branching Ratio	Experiment
R1 XI*1/2(1820) INTO (LAMBDA KBAR)/TOTAL	0.3	ALITTI 69 HBC 9/69
R2 XI*1/2(1820) INTO (XI PI)/TOTAL	0.1	ALITTI 69 HBC 9/69
R3 XI*1/2(1820) INTO (SIGMA KHAR)/TOTAL	0.02	TRIPP 67 RVUE 8/67
R4 XI*1/2(1820) INTO (XI*1/2(1530) PI)/TOTAL	0.3	ALITTI 69 HBC 9/69
R5 XI*1/2(1820) INTO (XI PI)/(LAMBDA KBAR)	0.20	RADIER 65 HBC 7/66
R6 XI*1/2(1820) INTO (XI*(1530) PI)/(LAMBDA KBAR)	0.26	SMITH 1 65 HBC 7/66
R7 XI(1820) INTO (XI PI)/(XI*(1530) PI)	1.5	APSELL 70 HBC 6/70
R8 XI(1820) INTO (XI PI)/(XI*(1530) PI)	0.3	APSELL 70 HBC 6/70

REFERENCES -- XI*1/2(1820)

HALSTEIN 63 SIGMA CONF 173
SMITH 1 65 PRL 14 25
RADIER 65 PL 16 171
SMITH 2 65 ATHENS CONF 251
TRIPP 67 NP 83 10
-- USES DATA OF SMITH 1.

ALITTI 69 PRL 22 79
DAUBER 69 PR 179 1262
CRENNELL 70 PR 10 867
APSELL 70 PRL 24 777
BMST 70 DUKE 317

MERRILL 68 PP 167 1202
APSELL 69 PRL 23 884
-- APSELL 69 IS SUPERSEDED BY BMST 70.

52 XI*1/2(1940, JP= 1) I=1/2

WF LIST UNDER XI(1940) EVERYTHING REPORTED IN THE MASS RANGE 1875-2000 MEV. THE SITUATION IS PERHAPS NOT QUITE SO UNCLEAR AS IS THE CASE FOR THE XI(1820).

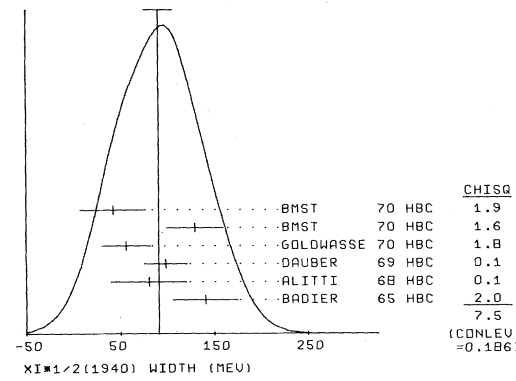
Mode	Decay Masses
M 35 1933.0	16.0
M 27 1930.0	20.0
M 46 1894.0	18.0
M 21 1955.0	14.0
M 110 1964.	9.
M 40 1961.	9.

AVERAGE MEANINGLESS (SCALE FACTOR = 1.7)
(SEE IDEOGRAM BELOW)

52 XI*1/2(1940) WIDTH (MEV)

Experiment	Width (MeV)
W 35 140.0	35.0
W 27 80.0	40.0
W 55 98.0	23.0
W 21 96.0	26.0
W 110 129.	30.
W 40 42.	35.

AVERAGE MEANINGLESS (SCALE FACTOR = 1.2)
(SEE IDEOGRAM BELOW)



52 XI*1/2(1940) PARTIAL DECAY MODES

Mode	Decay Masses
P1 XI(1940) INTO XI PI	1321+ 139
P2 XI(1940) INTO XI(1530) PI	1530+ 139
P3 XI(1940) INTO XI PI PI (XI PI NOT XI(1530))	1321+ 139+ 139

THE XI(1940) IS SEEN MAINLY IN XI PI AND SOME IN XI(1530) PI. IT HAS BEEN LOOKED FOR IN OTHER CHANNELS BUT NOT SEEN.

Mode	Branching Ratio	Experiment
R2 XI(1940) INTO (XI PI)/(XI(1530) PI)	2.8	APSELL 70 HBC 6/70
R3 XI(1940) INTO (XI PI)/(XI(1530) PI)	0.0	APSELL 70 HBC 6/70

REFERENCES -- XI*1/2(1940)

BADIER 65 PL 16 171
ALITTI 68 PRL 21 1110
DAUBER 69 PR 179 1262
APSELL 70 PRL 24 777
GOLDWASS 70 PR 10 1960
BMST 70 DUKE 317

APSELL 69 PRL 23 884
-- APSELL 69 IS SUPERSEDED BY BMST 70.

E(1940)

Baryons

For notation, see illustrated key at beginning of data card listings.

Xi(2030)

68 XI*1/2(2030, JP= 1 I=1/2

68 XI*1/2(2030) MASS (MEV)

M	42 2030.0	10.0	ALITTI	69 HRC	-	K-P 3.9-5 GEV/C	9/69
M	40 2059.0	17.0	BARTSCH	69 HRC	-	K-P 10GEV/C	9/69

AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)

69 XI*1/2(2030) WIDTH (MEV)

M	45.0	40.0	20.0	ALITTI	69 HRC	-	9/69
M	57.0	30.0		BARTSCH	69 HRC	-	9/69

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

69 XI*1/2(2030) PARTIAL DECAY MODES

P1	XI*1/2(2030)	INTO XI PI	1321+ 139
P2	XI*1/2(2030)	INTO LAMBDA KBAR	1115+ 497
P3	XI*1/2(2030)	INTO SIGMA KBAR	1197+ 497
P4	XI*1/2(2030)	INTO XI*1/2(1530) PI	1530+ 139
P5	XI*1/2(2030)	INTO LAMBDA (OR SIGMA) KBAR PI	1115+ 497+ 139

68 XI*1/2(2030) BRANCHING RATIOS

R1	XI*1/2(2030)	INTO (XI PI)/(MODES P1 THRU P4)	(P1)/(P1+P2+P3+P4)	9/69
R1	(0.30)	OR LESS	ALITTI 69 HRC	- 1 STD DEV LIMIT
R2	XI*1/2(2030)	INTO (LAM KBAR)/(MODES P1 THRU P4)	(P2)/(P1+P2+P3+P4)	9/69
R2	0.25	0.15	ALITTI 69 HRC	-
R3	XI*1/2(2030)	INTO (SIG KBAR)/(MODES P1 THRU P4)	(P3)/(P1+P2+P3+P4)	9/69
R3	0.75	0.20	ALITTI 69 HRC	-
R4	XI*1/2(2030)	INTO (XI* PI)/(MODES P1 THRU P4)	(P4)/(P1+P2+P3+P4)	9/69
R4	(0.15)	OR LESS	ALITTI 69 HRC	- 1 STD DEV LIMIT
R5	XI*1/2(2030)	INTO LAMBDA (OR SIGMA) KBAR PI	(P5)	9/69
R5	SEEN	BARTSCH 69 HRC		

REFERENCES -- XI*1/2(2030)

ALITTI 69 PRL 22 79 +BARNES,FLAMINIO,METZGER, + (BNL,SYRACUSE) I
 BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LONDON, VIENNA)

Xi(2250)

22 XI* (2250, JP=)

THE EVIDENCE FOR THIS STATE IS WEAK. BARTSCH 69 SEE A BUMP IN NOT MUCH STATISTICAL SIGNIFICANCE IN LAMBDA-KBAR-PI, SIGMA-KBAR-PI, AND XI-PI-PI MASS SPECTRA. GOLDWASSER TO SEE A NARROWER BUMP IN XI-PI-PI AT A HIGHER MASS. PERHAPS THEY ARE THE SAME STATE, PERHAPS THEY ARE NOT.

22 XI* (2250) MASS (MEV)

M	35 2244.0	52.0	BARTSCH	69 HRC	-	K-P 10 GEV/C	9/69
M	18 2295.0	15.0	GOLDWASSE	70 HRC	-	K-P 5.5 GEV/C	10/70

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

22 XI* (2250) WIDTH (MEV)

M	130.0	80.0	BARTSCH	69 HRC	-	9/69	
M	LESS THAN 30.0		GOLDWASSE	70 HRC	-	K-P 5.5 GEV/C	10/70

22 XI* (2250) PARTIAL DECAY MODES

P1	XI*1/2(2250)	INTO XI PI PI	1321+ 139+ 139
P2	XI*1/2(2250)	INTO LAMBDA KBAR PI	1115+ 497+ 139
P3	XI*1/2(2250)	INTO SIGMA KBAR PI	1197+ 497+ 139

REFERENCES -- XI* (2250)

BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LONDON, VIENNA)
 GOLDWASS 70 PR 10 1960 E L GOLDWASSER, P F SCHULTZ (ILLINOIS)

Xi(2500)

99 XI*1/2(2500, JP=) I=1/2

IT IS QUITE POSSIBLE THAT THE REASON THE EXPERIMENTS DISAGREE ABOUT THE MASS AND WIDTH IS THAT THEY ARE SEEING DIFFERENT XI'S. FOR NOW, HOWEVER, WE GROUP THEM TOGETHER.

99 XI*1/2(2500) MASS (MEV)

M	30 2430.0	20.0	ALITTI	69 HRC	-	K-P 4.5-5 GEV/C	9/69
M	45 2500.0	10.0	BARTSCH	69 HRC	0-	K-P 10 GEV/C	9/69

AVERAGE MEANINGLESS (SCALE FACTOR = 3.1)

99 XI*1/2(2500) WIDTH (MEV)

M	150.0	60.0	40.0	ALITTI	69 HRC	-	9/69
M	59.0	27.0		BARTSCH	69 HRC	0-	9/69

AVERAGE MEANINGLESS (SCALE FACTOR = 1.6)

99 XI*1/2(2500) PARTIAL DECAY MODES

P1	XI*1/2(2500)	INTO XI PI	1321+ 139
P2	XI*1/2(2500)	INTO LAMBDA KBAR	1115+ 497
P3	XI*1/2(2500)	INTO SIGMA KBAR	1197+ 497
P4	XI*1/2(2500)	INTO XI*1/2(1530) PI	1530+ 139
P5	XI*1/2(2500)	INTO LAMBDA (OR SIGMA) KBAR PI	1115+ 497+ 139
P6	XI*1/2(2500)	INTO XI PI PI	1321+ 139+ 139

99 XI*1/2(2500) BRANCHING RATIOS

R1	XI*1/2(2500)	INTO (XI PI)/(MODES P1 THRU P4)	(P1)/(P1+P2+P3+P4)	9/69
R1	(0.5) OR LESS	ALITTI 69 HRC	- 1 STD DEV LIMIT	
R2	XI*1/2(2500)	INTO (LAM KBAR)/(MODES P1 THRU P4)	(P2)/(P1+P2+P3+P4)	9/69
R2	0.5	0.2	ALITTI 69 HRC	-
R3	XI*1/2(2500)	INTO (SIG KBAR)/(MODES P1 THRU P4)	(P3)/(P1+P2+P3+P4)	9/69
R3	0.5	0.2	ALITTI 69 HRC	-
R4	XI*1/2(2500)	INTO (XI* PI)/(MODES P1 THRU P4)	(P4)/(P1+P2+P3+P4)	9/69
R4	(0.2) OR LESS	ALITTI 69 HRC	- 1 STD DEV LIMIT	
R5	XI*1/2(2500)	INTO LAMBDA (OR SIGMA) KBAR PI	(P5)	9/69
R5	SEEN	BARTSCH 69 HRC	0-	
R6	XI*1/2(2500)	INTO XI PI PI	(P6)	9/69
R6	SEEN	BARTSCH 69 HRC	0-	

REFERENCES -- XI*1/2(2500)

ALITTI 69 PRL 22 79 +BARNES,FLAMINIO,METZGER, + (BNL,SYRACUSE) I
 BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LONDON, VIENNA)

Omega

24 OMEGA - (1675, JP=3/2+) I=0

SEE LISTINGS OF STABLE PARTICLES

Appendix I

TEST OF $\Delta I=1/2$ RULE FOR K DECAYS

The quantities of interest for making tests of theoretical predictions regarding the $\Delta I=1/2$ rule for K decay are usually partial decay rates for single channels or special sums of channels. It is not possible to compute the errors on sums, differences, and ratios of partial decay rates from the information given in the Table of Stable Particles because of the presence of off-diagonal terms in the error matrix. For this reason we give some of these quantities below.

Table I. (000) or (+-0) refer to the sign of the pions into which the K decays.	
$\Gamma_{K_{\ell 3}^+} = \Gamma_{K_{e3}^+} + \Gamma_{K_{\mu 3}^+}$	$= (6.51 \pm .12) \times 10^6 \text{ sec}^{-1}$
$\Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}$	$= 0.659 \pm .022$
$\Gamma_{K_{\tau}^+} / \Gamma_{K_{\tau 1}^+}$	$= 3.329 \pm .086$
$\Gamma_{K_{\ell 3}^0} = \Gamma_{K_{e3}^0} + \Gamma_{K_{\mu 3}^0}$	$= (12.71 \pm .16) \times 10^6 \text{ sec}^{-1}$
$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0}$	$= 0.689 \pm .022$
$\Gamma_{K^0(000)} / \Gamma_{K^0(+0)}$	$= 1.704 \pm .072$

1. Leptonic decay rates

The $\Gamma_{K_{\ell 3}}$ rates are useful in testing the

leptonic $\Delta I = 1/2$ rule in the way suggested by Trilling.¹ The predictions are

$$\Gamma_{K_{\ell 3}^0} / 2\Gamma_{K_{\ell 3}^+} = 1.012, \text{ a phase-space}$$

factor,² and

$$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0} = \Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}.$$

From Table I,

$$\Gamma_{K_{\ell 3}^0} / 2\Gamma_{K_{\ell 3}^+} = 0.976 \pm .022$$

and

$$\frac{\Gamma_{K_{\mu 3}^0}}{\Gamma_{K_{e3}^0}} \left[\frac{\Gamma_{K_{\mu 3}^+}}{\Gamma_{K_{e3}^+}} \right]^{-1} = 1.046 \pm .048$$

These results seem to show a less than 2σ disagreement with the predictions, but the errors should be regarded with caution in view of the internal disagreements in the data. (Note the ideograms in the data listing for the charged K meson.)

2. Three-pion decays

We follow here the tests done by Mast et al.,³ based on the general analysis of K decays suggested by Zemach.⁴ Both decay rates and slopes (energy dependence of the Dalitz plot distributions) are used. The $\Delta I = 1/2$ rule gives the following predictions:

$$T_1 = \frac{2}{3} \frac{\Gamma_{K^0(000)}}{\phi_1} \left[\frac{\Gamma_{K^0(+0)}}{\phi_2} \right]^{-1} = 1,$$

$$T_2 = \frac{1}{4} \frac{\Gamma_{K_{\tau}^+}}{\phi_3} \left[\frac{\Gamma_{K_{\tau 1}^+}}{\phi_4} \right]^{-1} = 1,$$

$$T_3 = \frac{1}{2} \frac{\Gamma_{K_{\tau}^+}}{\phi_3} \left[\frac{\Gamma_{K^0(+0)}}{\phi_2} \right]^{-1} = 1,$$

$$T_4 = \frac{1}{2} g_{K_{\tau 1}^+} + g_{K_{\tau}^+} = 0,$$

$$T_5 = g_{K^0(+0)} + g_{K_{\tau}^+} - \frac{1}{2} g_{K_{\tau 1}^+} = 0,$$

where the ϕ_i are the phase-space factors. These factors have been calculated as described in Mast et al.³ by use of a relativistic formulation and the masses and slopes from this compilation. The factors labeled UDP are the relative areas of the Dalitz plots, assuming a uniform distribution. The NUDP include the observed slopes (see below). The CNUDP have been calculated by including the final-state Coulomb interaction. The values are:

	Method		
	UDP	NUDP	CNUDP
$\phi_1(000) =$	1.487	1.487	1.441
$\phi_2(+0) =$	1.249	1.293	1.278
$\phi_3(++-) =$	1.000	1.000	1.000
$\phi_4(+00) =$	1.247	1.182	1.146

The slopes for the various decays have not been tabulated in the Stable Particles Table. They are as follows:

$$\left. \begin{aligned} g_{K_{\tau}^+} &= -0.206 \pm 0.007 \\ g_{K_{\tau}^-} &= -0.194 \pm 0.007 \end{aligned} \right\} -0.200 \pm 0.005$$

$$g_{K_{\tau 1}^+} = 0.527 \pm 0.017,$$

$$g_{K_L^0(+0)} = 0.603 \pm 0.028 \text{ (Error scaled } \times 3.1)$$

A difference in the τ^+ and τ^- slopes would be an indication of CP violation in this decay. Since no

difference is present at this time, we average the two and use this value in T_4 .

Using the CNUDP and rates and slopes reported here, we get:

$$\begin{aligned} T_1 &= 1.006 \pm 0.043, \\ T_2 &= 0.954 \pm 0.025, \\ T_3 &= 1.184 \pm 0.028, \\ T_4 &= 0.064 \pm 0.010, \\ T_5 &= 0.139 \pm 0.030. \end{aligned}$$

The three-pion final state can be in isospin states $I = 1, 2, 3$. T_1 and T_2 test the existence of isospin $I=3$ in the final state and are consistent with no or very little $I=3$. T_4 is related to the $I=2$ amplitude in the final state and indicates the presence of $I=2$. T_3 and T_5 give information on the $\Delta I=3/2$ part of the $I=1$ amplitude relative to the $\Delta I=1/2$ part. Both tests indicate the presence of $\Delta I=3/2$.

The rate tests $T_1, T_2,$ and T_3 could differ from 1.0 by as much as 10% owing to the π mass differences and the occurrence of big slopes.⁵ The error on $g_{K_L^0(+0)}$ has been scaled by a factor of 3.1 to compensate for the disagreement among the determinations of this parameter. Since this error is the main contribution to the T_5 error, the reader is cautioned to refer to the $g_{K_L^0(+0)}$ ideogram in the data card listings before interpreting T_5 .

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2. N. Brene (CERN), private communication. In our Jan. 1968 edition we had erroneously used 1.04.
3. T. S. Mast, L. K. Gershwin, M. Alston-Garnjost, R. O. Bangerter, A. Barbaro-Galtieri, J. J. Murray, F. T. Solmitz, and R. D. Tripp, Phys. Rev. 183, 1200 (1969).
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Appendix II

A. SU(3) CLASSIFICATION OF BARYON RESONANCES

There are a few multiplets that have been studied and we report here the results. The relevant formulae are given below.

Mass Formulae

Decuplet $\Delta - \Sigma = \Sigma - \Xi^* = \Xi^* - \Omega$ GMO (1)

Octet $2(N + \Xi) = 3\Lambda + \Sigma$ GMO (2)

Nonet $\left\{ \begin{aligned} \sin^2 \theta &= \frac{\Lambda - M_8}{\Lambda - \Lambda'} && \text{Mixing angle}^\dagger \end{aligned} \right.$ (3)

$\left\{ \begin{aligned} M_8 &= \frac{2(N + \Xi) - \Sigma}{3} && \text{GMO} \end{aligned} \right.$ (4)

Here GMO stands for the Gell-Mann-Okubo formula; the particle symbol indicates its mass. The formulae would be the same if squared masses were used. For the nonet case, Λ is the "mostly-octet" particle, Λ' is the "mostly-singlet" particle.

Decay Rates

In terms of a relativistically invariant matrix element T , the decay rate for two-body decay of a resonance of mass M_R is

$$\Gamma \propto \frac{|T|^2 R_2}{M_R}, \tag{5}$$

where $R_2 = k/M_R$ is the two-body phase space factor. Since the numerator is an invariant, and since Γ must transform as $1/E$, we introduce the denominator $1/M_R$ (see FEYNMAN 62).

For meson decays (see below) the rates are calculated according to Eq. (5); for baryon resonance decays into $1/2^+$ baryons and 0^- mesons, one next takes into account the fact that spin sums in $|T|^2$ introduce another factor M_R , cancelling the $1/M_R$. We are then left with

$$\Gamma = \frac{|T|^2 k}{M_R} M_N, \text{ for baryons} \tag{5'}$$

$$= \frac{|T|^2 k}{M_R^2} M_N^2, \text{ for mesons} \tag{5''}$$

The powers of the nucleon mass M_N or M_N^2 have been introduced so that we can treat $|T|$ as dimensionless.

$|T|^2$ contains centrifugal barrier factors, which we call B_ℓ . We then have

Decuplet Singlet $\left. \right\} \Gamma = (cg)^2 B_\ell(k) \frac{M_N}{M_R} k$ (6)

Octet $\Gamma = (c_D g_D + c_F g_F)^2 B_\ell(k) \frac{M_N}{M_R} k$ (7)

Nonets $\left\{ \begin{aligned} G_8 &= \Lambda \cos \theta - \Lambda' \sin \theta \\ G_4 &= \Lambda \sin \theta + \Lambda' \cos \theta \end{aligned} \right.$ (8)

with $\begin{aligned} G_8 &= c_D g_D + c_F g_F \\ G_4 &= c_4 g_4. \end{aligned}$ (9)

Here c_i are the SU(3) coefficients with the sign convention adopted in this article [see note preceding the table of SU(3) isoscalar coefficients and Fig. 3, "Note on $Y^{*1} s''$ "]. M_N is the nucleon mass, M_R is the resonance mass for which Γ is calculated, k is the center-of-mass momentum for the channel being considered, g_i are the relevant couplings. For the case of singlet-octet mixing, formula (8) has to be used in conjunction with (6) and (7). G_8 and G_1 represent the couplings for the multiplet, and Λ and Λ' represent the couplings for the physical states. In a recent fit of all baryon states into SU(3) multiplets by the CHS collaboration (PLANE 70) two forms have been tried for B_ℓ .

(a) The form $B_\ell = k^2 D_\ell(kR)$, R being the radius of interaction and D_ℓ the polynomials in kR given by BLATT-WEISSKOPF 52. Usually R is taken to be 1 Fermi (TRIPP 68).

(b) The form $B_\ell = k^2 \theta$. It turns out that the values of θ and g_F/g_D obtained with the two forms for B_ℓ are the same, although the absolute values of g_F and g_D depend upon the form used.

The relation between g_D , g_F and the D (symmetric) and F (antisymmetric) couplings is as follows:

$$\frac{F}{D} = \sqrt{\frac{5}{3}} \frac{g_F}{g_D}. \quad (10)$$

Fits of baryon decay rates within SU(3) can be found, among others, in TRIPP 68 and 69 and LEVI SETTI 69.

The most recent published survey is by the CHS collaboration (PLANE 70); we report their results in Table I. A new fit by the BNL group has been reported by Samios at the Kiev Conference[†] (SAMIOS 70). The results of the two analyses are in good agreement, although the procedures used are somewhat different. We refer the reader to these papers for more details. Some comments are given below.

$\frac{1}{2}^-$ -Nonet (Baryon - Eta Resonances)

For this nonet relation, (7) was multiplied by the factor

$$\left[\frac{M_R - M_B}{\bar{M}_R - \bar{M}_B} \right]^2,$$

where M_B is the decay baryon and $\bar{M}_R - \bar{M}_B = 564$ MeV is the difference of the mean $1/2^-$ and $1/2^+$ baryon octet masses. This kinematic factor comes from PCAC arguments (i. e., the assumption that axial vector current remains an octet in presence of symmetry breaking) and it was advocated by Graham et al. (GRAHAM 67). For the $1/2^-$ nonet it has been used in this form first by Gell-Mann et al. (GELL-MANN 68).

The two values of F/D are dependent upon what value is used for the decay $N(1525) \rightarrow N\eta$. The first value is obtained by using 60% for the $N\eta$ branching ratio, the second value by using ~15%. SAMIOS 70 reports only the second value. However, there are now reservations about the second value (see footnote

Table I. SU(3) baryon multiplets with two or more known members.

The coupling constants are those for decay into baryon ($1/2^+$) octet \otimes pseudoscalar meson octet. Values of θ and F/D taken from PLANE 70. The analysis reported by SAMIOS 70 finds similar results.

J^P	Octet members ^a	Singlet	$\theta(\text{deg})^b$	F/D
$1/2^-$	$N'(1535)$ $\Lambda(1670)$ $\Sigma(1750)$ $\Xi(1825)$	$\Lambda(1405)$	18 ± 17	-2.9 -0.21
$3/2^-$	$N(1520)$ $\Lambda(1690)$ $\Sigma(1670)$ $\Xi(1800)$	$\Lambda(1520)$	-25 ± 6	2.3
$5/2^-$	$N(1670)$ $\Lambda(1830)$ $\Sigma(1765)$			-0.19
$5/2^+$	$N(1688)$ $\Lambda(1845)$ $\Sigma(1945)$			0.82 ^d
	Decuplet members	g_{10}^2 ^c		
$3/2^+$	$\Delta(1236)$ $\Sigma(1385)$ $\Xi(1530)$ Ω^-	0.94 to 2.38		
$7/2^+$	$\Delta(1950)$ $\Sigma(2030)$	0.25 to 0.97		

a. Masses in parentheses are the nominal masses used in the Baryon Table. For the Ξ members, the value of the mass is the one predicted by SU(3), using the mixing angle calculated by using the rates.

b. Values calculated from the decay rates (Eq. 8).

c. Values taken from TRIPP 68 who uses the convention of Eq. 6 with $B_\ell(k)$ being the BLATT-WEISSKOPF 52 form.

d. SAMIOS 70 finds F/D = 1.5 for this octet.

o to Baryon Table).

B. SU(3) CLASSIFICATION OF MESON RESONANCES

All of the discussion above applies, except that for Bosons the GMO formula is usually applied to the square of the masses, as opposed to the first power for fermions. Thus for example, Eq. (2) becomes

$$4\hat{K} = 3\hat{\eta} + \hat{\pi}. \quad (2')$$

The symbol \hat{K} was introduced by Glashow and Socolow[†] for the square of the K mass, etc.

Because of the difference between Eqs. (5') and (5''), there is also an extra factor of (M_N/M_R) in Eqs. (6) and (7). The three established nonets (0^- , 1^- , 2^+) and their mixing angles are listed in Table II.

Table II. Mixing angles from quadratic SU(3) mass formula. Of the two iso-singlets, the "mainly octet" one is written first, followed by a semicolon.

J^P	Octet members	θ
0^-	possible nonet [π , K, η ; η']	$10.4 \pm 0.2^\circ$
0^-	alternative nonet [π , K, η ; E]	$6.2 \pm 0.1^\circ$
1^-	[$\rho(765 \pm 10)$, K^* , ϕ ; ω]	$39.5 \pm 0.7^\circ$
2^+	[$A_2(1300 \pm 20)$, $K_N(1420)$, f' ; f]	$33.7 \pm 3.9^\circ$

Footnotes and References

[†]The formula has been calculated from analogy with the formula for mixing of meson states, first put in this form by S. L. Glashow and R. H. Socolow, Phys. Rev. Letters 15, 329 (1966). For the baryon formula see A. Barbaro-Galtieri, Phenomenology of Resonances and Particle Supermultiplets. UCRL-17054 (1966).

[‡]This is an updated version of the fits by Flaminio et al., BNL report 14572.

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