



CM-P00040025

PARTICLE PROPERTIES: January, 1969

Table S: STABLE PARTICLES. January, 1969.  
(Closing date for data: November 1, 1968)

From Review of Particle Properties, UCRL-8030.  
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Quantities in italics have changed by more than one (old) standard deviation since January, 1968.

Symbol	$J^P$	Mass (MeV)	Mass difference (MeV)	Mean life (sec)	Spin (cm)	Decays		Q (MeV)	P	P max (MeV/c)	General Atomic and Nuclear Constants <sup>a</sup>	
						Partial mode	Fraction				N	c
$\gamma$	$0, 1(1^-)$	0	$< 2 \times 10^{21}$ MeV	stable		0	stable				$= 6.02252 \times 10^{23} \text{ mole}^{-1}$ (based on $A_C 12=12$ )	
$\nu$	$J = \frac{1}{2}$	0	$< 0.2$ keV	stable		0	stable				$= 2.997925 \times 10^{10} \text{ cm sec}^{-1}$	
$e$	$J = \frac{1}{2}$	0.511006		stable		0.000	stable	$\mu_e = 1.001149575$			$= 4.80298 \times 10^{-10} \text{ esu} = 1.60210 \times 10^{-19} \text{ coulomb}$	
$\mu$	$J = \frac{1}{2}$	105.659		$2.1983 \times 10^{-6}$		0.011	$e\nu\bar{\nu}$	100	105	53	$= 1.60210 \times 10^{-6} \text{ erg}$	
$\pi^\pm$	$1(0^-)$	139.578		$2.604 \times 10^{-8}$		0.019	$e\nu$	$< 1.6$	105	53	$= 6.5849 \times 10^{-22} \text{ MeV sec}$	
$\pi^0$	$1^-(0^+)$	134.975		$0.89 \times 10^{-16}$		0.018	$e\nu$	$< 1.3$	104	53	$= 1.05449 \times 10^{-27} \text{ erg sec}$	
$K^\pm$	$\frac{1}{2}(0^-)$	493.82		$1.235 \times 10^{-8}$		0.244	$\mu\nu$	100	34	30	$= 1.9732 \times 10^{-14} \text{ MeV cm} = 197.32 \text{ MeV fermi}$	
$K_S^0$	$\frac{1}{2}(0^-)$	497.76		$0.862 \times 10^{-10}$		0.248	$\mu\nu$	$(63.65 \pm 0.29)\%$	236	236	$= 8.6171 \times 10^{-11} \text{ MeV}^2 = 1 \text{ eV}^2 / 11605^2 \text{ K}$	
$K_L^0$	$\frac{1}{2}(0^-)$	497.76		$5.38 \times 10^{-8}$		0.248	$\mu\nu$	$(21.03 \pm 0.30)\%$	205	205	$= e^2/hc = 1/137.038$	
$\eta$	$0^+(0^+)$	548.8		$5.38 \times 10^{-8}$		0.248	$\mu\nu$	$(1.74 \pm 0.04)\%$	135	67	$= h/m_e c = r_e \alpha^{-1} = 3.86144 \times 10^{-11} \text{ cm}$	
$p$	$\frac{1}{2}(\frac{1}{2}^+)$	938.256		stable		0.880	$\mu\nu$	$(1.24 \pm 0.03)10^{-4}$	139	70	$= \hbar^2/m_e c^2 = 2.81777 \text{ fermi} (1 \text{ fermi} = 10^{-13} \text{ cm})$	
$n$	$\frac{1}{2}(\frac{1}{2}^+)$	939.550		stable		0.882	$\mu\nu$	$(1.24 \pm 0.03)10^{-4}$	139	70	$= \hbar/m_e c^2 = r_e \alpha^{-2} = 0.529167 \text{ A} (1 \text{ A} = 10^{-8} \text{ cm})$	
$\Lambda$	$0(\frac{1}{2}^+)$	1115.60		$2.51 \times 10^{-10}$		1.245	$\mu\nu$	$(1.02 \pm 0.07)10^{-8}$	4	5	$= \frac{8}{3} \pi r_e^2 = 0.66516 \times 10^{-24} \text{ cm}^2 = 0.66516 \text{ barn}$	
$\Sigma^+$	$1(\frac{1}{2}^+)$	1189.40		$0.810 \times 10^{-10}$		1.412	$\mu\nu$	$(3.0 \pm 0.5)10^{-8}$	139	70	$= e\hbar/2m_p c = 0.578817 \times 10^{-14} \text{ Mev gauss}^{-1}$	
$\Sigma^0$	$1(\frac{1}{2}^+)$	1192.46		$2.51 \times 10^{-10}$		1.422	$\mu\nu$	$(98.83 \pm 0.04)\%$	135	67	$= \frac{1}{2} \omega_{\text{cyclotron}} (\text{rad sec}^{-1} \text{ gauss}^{-1}) = e/2m_p c = 8.79404 \times 10^6$	
$\Sigma^-$	$1(\frac{1}{2}^+)$	1197.32		$1.64 \times 10^{-10}$		1.434	$\mu\nu$	$(1.17 \pm 0.04)\%$	134	67	$= e/2m_p c = 4.7895 \times 10^3$	
$\Xi^0$	$\frac{1}{2}(\frac{1}{2}^+)$	1314.7		$3.03 \times 10^{-10}$		1.728	$\mu\nu$	$(3.4 \pm 0.3)10^{-5}$	214	203	$(\text{continued on other sheet})$	
$\Xi^-$	$\frac{1}{2}(\frac{1}{2}^+)$	1321.25		$3.03 \times 10^{-10}$		1.746	$\mu\nu$	$(3.4 \pm 0.3)10^{-5}$	214	203	$\text{Mod. Phys. 37, 537 (1965). Note that recent fine structure and ac Josephson effect measurements currently indicate that } \alpha \text{ is } \sim 20 \text{ ppm larger than the value given above. The corresponding readjustment of the fundamental constants will increase } e \text{ by } 60 \text{ ppm, } \hbar \text{ by } 100 \text{ ppm, } m_p \text{ by } 60 \text{ ppm, and decrease } N \text{ by } 60 \text{ ppm. For a preliminary analysis see W. H. Parker, B. N. Taylor, and D. N. Langenberg, Phys. Rev. Lett. 18, 287 (1967).}$	
$\Omega^-$	$0(\frac{3}{2}^+)$	1672.4		$1.3 \times 10^{-10}$		2.797	$\mu\nu$	$(3.4 \pm 0.3)10^{-5}$	214	203	$\text{CP violation parameters}$	

\* S = Scale factor =  $\sqrt{2/(N-1)}$  where N = number of experiments. S should be  $\approx 1$ . If  $S > 1$ , we have enlarged the error of the mean,  $\delta x$ , i.e.,  $\delta x \rightarrow S \delta x$ . This new convention is still inadequate, since if  $S > 1$ , the real uncertainty is probably even greater than  $S\delta x$ . See text of January 1967 edition.  
 † In decays with more than two bodies,  $P_{\text{max}}$  is the maximum momentum that any particle can have. ‡ Predicted from SU(3).  
 § Theoretical value, see also data card listings. d. Assumes rate for  $\Xi^- \rightarrow \Sigma^0 e \nu$  small compared with  $\Xi^- \rightarrow \Lambda e \nu$ .  
 ¶ See note in data card listings.  
 †† See data card listings for energy limits used in measuring this branching ratio.

## Footnotes for the MESON Table

- (f) Reported values range between 1% and 10%, and depend on assumptions on  $\rho$ - $\omega$  interference.
- (g) This  $\omega \rightarrow e^+e^-$  value is the average from a  $\pi^+p \rightarrow e^+e^-n$  experiment (giving  $0.0040 \pm 0.0015$  [ $\pm 0.0014$  from possible  $\rho\omega$  interference]) and an  $e^+e^- \rightarrow \pi^+\pi^-\pi^0$  experiment (giving  $0.0085 \pm 0.0016$ ).
- (h) This  $0^-$  meson was named  $\eta'$  on discovery, when it looked as if it completed the  $0^-$  nonet. With the recent evidence that the  $E(1420)$  is probably also  $0^-$ , it is no longer clear whether  $\eta'$  or  $E$  or both are mixed in with the  $\pi, \eta, K$  octet; so the name  $\eta'$  may be misleading.
- (i) Empirical limits on fractions for other decay modes of  $\eta'$  (958):  $\pi^+\pi^- < 7\%$ ,  $3\pi < 7\%$ ,  $4\pi < 1\%$ ,  $6\pi < 1\%$ ,  $\pi^+\pi^-e^+e^- < 0.6\%$ ,  $\pi^0e^+e^- < 1.3\%$ ,  $\eta e^+e^- < 1.1\%$ ,  $\pi^0\rho^0 < 4\%$ ,  $\pi^0\omega < 8\%$ .
- (j) Empirical limits on fractions for other decay modes of  $\phi(1019)$ :  $\pi^+\pi^- < 20\%$ ,  $\eta\gamma < 8\%$ ,  $\eta + \text{neutrals} < 13\%$ ,  $\pi^+\pi^-\gamma < 4\%$ ,  $\omega\gamma < 5\%$ ,  $\rho\gamma < 2\%$ .
- (k) Width of  $\eta_{0+}(1070) \rightarrow K_S K_S$ : Average value from two bubble chamber experiments is  $\Gamma = (72 \pm 13)$  MeV, whereas two spark chamber experiments give  $\Gamma > 100$  MeV. The latter also allow a scattering length fit. It is not clear whether the reported narrow ( $\Gamma \lesssim 25$  MeV)  $\pi^+\pi^-$  enhancements near 1070 MeV have anything to do with the  $\eta_{0+}(1070)$ .
- (l)  $\rho\pi$  fraction of  $3\pi$  mode difficult to distinguish because  $\rho$  bands cover most of the Dalitz plot.
- (m) Empirical limits on fractions for decay modes of  $B(1220)$ :  $\pi\pi < 30\%$ ,  $K\bar{K} < 2\%$ ,  $4\pi < 50\%$ ,  $\phi\pi < 1.5\%$ ,  $\eta\pi < 25\%$ ,  $(K\bar{K})^0 < 8\%$ ,  $K_S K_S \pi^\pm < 2\%$ ,  $K_S K_L \pi^\pm < 6\%$ .
- (n) Although the splitting of the  $A_2$  needs further confirmation, we give the results from the two published experiments that have observed a split  $A_2$ . Since most experiments have only seen one, rather wide,  $A_2$  enhancement, we here list its ("combined") properties:  $IG(J^P)C_n = 1^-(2^+)_{+}$ ;  $M = 1297 \pm 10$  MeV ( $S=1.8^*$ ) (§),  $\Gamma = 91 \pm 10$  MeV ( $S=1.1^*$ ) (§); partial decay modes:  $\rho\pi$   $86 \pm 2\%$ ,  $K\bar{K}$   $2.4 \pm 0.5\%$ ,  $\eta\pi$   $11 \pm 2\%$ ,  $\eta'\pi$   $0.5 \pm 0.4\%$  ( $S=2.1^*$ );  $\pi^+\pi^-\pi^0$  (excl.  $\rho\pi$ )  $< 17\%$ .
- (o) There is only a weak indication for a  $K^*\bar{K} + \bar{K}^*K$  mode of the  $f'(1514)$ . If this mode does not exist, the  $K\bar{K}$  branching fraction will have to be reported as  $(80 \pm 13)\%$  (rather than  $(72 \pm 12)\%$  as given in the table), and  $\eta\pi$  as  $(20 \pm 13)\%$ .
- (p) See the listings for many statistically weak  $Y=0$  bumps with  $M \geq 1700$  MeV, seen in bubble chambers. We tabulate here 9 statistically strong bumps seen with a missing mass spectrometer ( $\pi^+p \rightarrow p(MM)^-$ ) or in HBC or counter experiments on  $NN$  elastic scattering or total cross sections.

Name	I	M (MeV)	$\Gamma$ (MeV)	Decay Modes Observed
R1(1630)	$\geq 1$	$1630 \pm 15$	$\leq 21$	$1/3 / > 3$ charg. part. $\approx .37/.59/.04$
R2(1700)	$\geq 1$	$1700 \pm 15$	$\leq 30$	$1/3 / > 3$ charg. part. $\approx .43/.56/.01$
R3(1750)	$\geq 1$	$1748 \pm 15$	$< 38$	$1/3 / > 3$ charg. part. $> .14 / < .80 / .15$
? $\left\{ \begin{array}{l} NN(1925) \\ S(1930) \end{array} \right.$	$0, 1$	$\approx 1925$	$\approx 10$	structure in $p\bar{p}$ backw. el. scatt.
? $\left\{ \begin{array}{l} NN(1945) \\ NN(2190) \end{array} \right.$	$0, 1$	$\approx 1945$	$\approx 22$	$1/3 / > 3$ charg. part. $\approx 0/.92 / 0$
? $\left\{ \begin{array}{l} NN(2190) \\ T(2200) \end{array} \right.$	1	$2190 \pm 10$	$\approx 85$	structure in $p\bar{p}$ backw. el. scatt.
? $\left\{ \begin{array}{l} NN(2345) \\ U(2380) \end{array} \right.$	$\geq 1$	$2195 \pm 15$	$\approx 13$	structure in $NN$ total cross section
? $\left\{ \begin{array}{l} NN(2345) \\ U(2380) \end{array} \right.$	1	$2345 \pm 10$	$\approx 140$	$(MM)^- \rightarrow 3$ charged particl. $\approx 94\%$
? $\left\{ \begin{array}{l} NN(2345) \\ U(2380) \end{array} \right.$	$\geq 1$	$2382 \pm 24$	$\approx 30$	structure in $NN$ total cross section
$\left\{ \begin{array}{l} NN(2380) \\ NN(2380) \end{array} \right.$	0	$2380 \pm 10$	$\approx 140$	$(MM)^- \rightarrow 1/3 / > 3$ chrgd part. $\approx 30/45/25$
				structure in $NN$ total cross section

There is no evidence on the G, J, or P quantum numbers of these bumps (apart from the suggestion of  $\ell = \text{odd}$  for  $NN(1925)$ ,  $\ell = \text{even}$  for  $NN(1945)$ ), nor is there satisfactory agreement between them and the other bubble chamber claims. Further, the  $\sigma_{\text{tot}}(NN)$  bumps are broader than the  $(MM)^-$  bumps, and there is no evidence for or against their interpretation as resonances.

- (q) Taken from compilation by T. Ferbel, Proc. 1968 Philadelphia Conf. See the data listings for averages of the values given in the literature. Also see B. French's review of Mesons (Proc. 14th International Conf. High Energy Physics, Vienna (1968), p. 91) for possible differences between M and  $\Gamma$  of charged and neutral  $\rho_N(1650)$ .
- (r) See note in listings. Some investigators see a broad enhancement in mass ( $K\pi\pi$ ) from 1200 - 1350 MeV, and others see structure. A further bump at 1280 MeV,  $\Gamma = 80$  MeV, has been suggested. In light of this confusion, the masses, widths, quantum numbers, and branching ratios are at best tentative. For the mass region 1200 - 1350 MeV, the decay rate into  $K^*(890)\pi$  is large, and a  $K\rho$  decay is seen. The  $K\eta, K\omega$  and  $K\pi$  rates are less than a few percent.
- (s) This  $\eta' \rightarrow \gamma\gamma$  value is from a constrained fit under the assumption that  $\eta\pi\pi, \pi^+\pi^-\gamma$  (inclusive  $\rho^0\gamma$ ), and  $\gamma\gamma$  are the only existing decay modes. Note that direct measurement of the  $\eta' \rightarrow \gamma\gamma$  branching fraction gave the slightly different result of  $(5.5^{+3.6}_{-3.0})\%$ .

Mixing angles from Quadratic SU(3) Mass Formula:  $0^-$  nonet ( $\pi, K, \eta, \eta'$ )  $\theta = 10.4^\circ \pm 0.2^\circ$ ; alternative  $0^-$  nonet ( $\pi, K, \eta, E$ )  $\theta = 6.2^\circ \pm 0.1^\circ$ ;  $1^-$  nonet ( $\rho(m=765 \pm 15 \text{ MeV}), K^*, \phi, \omega$ )  $\theta = 39.9 \pm 1.1^\circ$ ;  $2^+$  nonet ( $A_2, K_N(1420), f', f$ )  $\theta = 29.9^\circ \pm 2.2^\circ$ .

## Footnotes for the BARYON Table

- \* Quoted error includes an S(scale) factor. See footnote to Table S.
- † For decay modes into  $\geq 3$  particles  $p_{\text{max}}$  is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated using the averaged central mass values, without taking into account the widths of the resonances.
- a. Square brackets indicate a sub-reaction of the previous unbracketed decay mode.
- b. J is not known; x is  $\Gamma_{e1}/\Gamma$ .
- c. This is only an educated guess; the error given is larger than the error of the average of the published values (see listings for the latter).
- † For the baryon states, the name [such as  $N(1470)$ ] contains the mass, which shifts by 5 or 10 MeV with each new analysis. We can't keep up with changing labels in the card-listing section, so we don't try. The name (col. 1) is the same as can be found in large print in the listings. The best current value of the mass (col. 4) is what we use to determine the beam parameters,  $M^2 \pm \Gamma M$ , c.m. decay momenta, etc., that are found in other columns.
- An arrow at the left of the Table indicates a candidate that has been omitted because the evidence for the existence of the effect and (or) for its interpretation as a resonance is open to considerable question. See listings for information on the following:  $\Delta(1690)P_{33}, N(1730)D_{13}, N(1860)F_{13}, N(1980)D_{13}, N(2080), N_2(3245), N(3690), N(3755), Z_0(1865), Z_1(1900), \Lambda(1327), \Lambda(1745)P_{01}, \Lambda(1750)S_{01}, \Lambda(1860)F_{07}, \Sigma(1440), \Sigma(1650)S_{11}, \Sigma(1780), \Sigma(1880),$  and  $\Xi(1705)$ .

Change in Notation. The subscript N stands for "normal spin-parity series" (J<sup>P</sup>=0<sup>+</sup>, 1<sup>-</sup>, 2<sup>+</sup>, ...), A for "abnormal" (J<sup>P</sup>=0<sup>-</sup>, 1<sup>+</sup>, 2<sup>-</sup>, ...). Quantities in italics have changed by more than one (old) standard deviation since January, 1968.

Main table with columns: Symbol (J<sup>P</sup>), I<sup>G</sup>(J<sup>P</sup>)C<sub>N</sub>, Mass (MeV), Width Γ (MeV), M<sup>2</sup> (GeV)<sup>2</sup>, Mode, Fraction %, p or P<sub>max</sub> (MeV/c), and Q (MeV). Rows include various meson states like π(140), η(549), ρ(765), ω(783), η'(958), φ(1019), η<sub>1</sub>(1070), A1(1070), B(1220), f(1260), D(1285), A2<sub>L</sub>(1270), A2<sub>H</sub>(1315), E(1420), f'(1515), etc.

The following bumps, excluded above, are listed among the data cards: σ(410); η<sub>1</sub>(720); H(990); A15(1170); ρ(1410); K<sub>S</sub>K<sub>S</sub>(1440); π/ρ(1550); φ(1650); R(1750); η or ρ(1830)→4π; φ or π(1830)→ωππ; S(1930); T(2200); NN<sub>0</sub>(2380); U(2380); κ(725); K<sub>N</sub>(1100-1200); K<sub>A</sub>(1-3/2)(1175); K<sub>A</sub>(1-3/2)(1265); K<sub>A</sub>(1-1/2)(1280); K<sub>N</sub>(1-1/2)(1660); K<sup>\*</sup>(2240)→πN. (See note (p).)

- \* Quoted error includes scale factor S = √(N-1). See footnote to Table S.
† Square brackets indicate a subreaction of the previous (unbracketed) decay mode.
‡ This is only an educated guess; the error given is larger than the error of the average of the published values (see listings for the latter).
(a) ΓM is the half-width of the resonance when plotted against M<sup>2</sup>.
(b) For decay modes into ≥ 3 particles P<sub>max</sub> is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated using the averaged central mass values, without taking into account the widths of the resonances.
(c) The values given for M(p) and Γ(p) and their errors are not average values from various experiments, but rather are intended to give the range where we believe the actual values are most likely to fall. Contrast the results tabulated in this note:
Energy-independent width is a narrow-resonance approximation which tends to give lower mass and width. Thus an analysis of all e<sup>+</sup>e<sup>-</sup>→π<sup>+</sup>π<sup>-</sup> data with Omnès' formula and energy-dependent width gives M = 770±4 MeV, Γ = 122±6 MeV (M. Roos, priv. comm.).
(d) The quoted value of the rate ρ<sup>0</sup>→e<sup>+</sup>e<sup>-</sup> is the average from two e<sup>+</sup>e<sup>-</sup>→π<sup>+</sup>π<sup>-</sup> experiments (which alone give an average of (0.0060±0.0006)%) and one photoproduction experiment of high mass resolution. Interference effects with ω decay are hoped to be small.
(e) Warning: The values given in the literature, and in our table, for the rate ρ<sup>0</sup>→μ<sup>+</sup>μ<sup>-</sup> may be somewhat too high, due to possible interference with ω decay.

(Footnotes continued on the back of Table S)

Summary table with columns: Meson, J<sup>P</sup>, I<sup>G</sup>(J<sup>P</sup>)C<sub>N</sub>, Mass (MeV), Width Γ (MeV), M<sup>2</sup> (GeV)<sup>2</sup>, Decay Modes, Fraction %, p or P<sub>max</sub> (MeV/c), Q (MeV). Rows include various meson states like K<sup>0</sup>, K<sup>+</sup>, K<sup>-</sup>, K<sub>S</sub><sup>0</sup>, K<sub>L</sub><sup>0</sup>, K<sub>N</sub><sup>0</sup>, K<sub>N</sub><sup>+</sup>, K<sub>N</sub><sup>-</sup>, K<sub>N</sub><sup>0</sup>(420), K<sub>N</sub><sup>+</sup>(420), K<sub>N</sub><sup>-</sup>(420), etc.

There is no well-established Y = 0 meson resonance with M > 1700 MeV. For a list of some possible states see footnote (p).

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Particle or resonance †	I(J <sup>P</sup> ) estab.	Beam π, K (BeV/c)	Mass (MeV)	Γ (MeV)	M <sup>2</sup> +ΓM (BeV <sup>2</sup> )	Partial decay modes			
						Mode	Fraction (%)	p or p <sup>†</sup> (MeV/c) <sup>max</sup>	4π <sup>2</sup> (mb)
p	1/2(1/2 <sup>+</sup> )		938.3 939.6		0.880 0.883	See Table S			
N <sup>+</sup> (1470)	1/2(1/2 <sup>+</sup> ) P <sub>11</sub>	T=0.52π p=0.64	1460	260	2.13 ±0.38	Nπ Nππ <sup>±</sup> [Nπ] <sup>±</sup> [Δπ] <sup>±</sup>	55 45 271 162	412 360 271 162	28.8
N(1518)	1/2(3/2 <sup>-</sup> ) D <sub>13</sub>	T=0.60 p=0.73	1515	115	2.30 ±0.17	Nπ Nππ [Δ(1236)π] <sup>±</sup> Nη	50 50 [domin] ~ 0.5	452 406 249 437	23.9
N(1550)	1/2(1/2 <sup>-</sup> ) S <sub>11</sub>	T=0.62 p=0.75	1525	80	2.33 ±0.12	Nπ Nππ Nη	35 65 small	459 161 414	23.1
N(1680)	1/2(5/2 <sup>-</sup> ) D <sub>15</sub>	T=0.88 p=1.01	1675	145	2.81 ±0.24	Nπ Nππ [Δ(1236)π] <sup>±</sup> AK Nη	45 55 [?] [?] [?]	564 530 209 374	15.4
N(1688)	1/2(5/2 <sup>+</sup> ) F <sub>15</sub>	T=0.90 p=1.03	1690	125	2.86 ±0.21	Nπ Nππ [Δ(1236)π] <sup>±</sup> AK Nη	60 40 [?] [?] [?]	574 540 374 234 390	14.9
N <sup>+</sup> (1710)	1/2(1/2 <sup>-</sup> ) S <sub>11</sub>	T=0.95 p=1.08	1715	280	2.94 ±0.48	Nπ	65	590	14.0
N <sup>+</sup> (1750)	1/2(1/2 <sup>+</sup> ) P <sub>11</sub>	T=1.08 p=1.21	1785	405	3.19 ±0.72	Nπ AK	34 seen	636	12.1
N(2190)	1/2(7/2 <sup>-</sup> ) G <sub>17</sub>	T=1.94 p=2.07	2190	300	4.80 ±0.66	Nπ	35	894	6.21
N(2650)	1/2( ? )	T=3.12 p=3.26	2650	360	7.02 ±0.95	Nπ	(J+1/2)κ=0.45 <sup>b</sup>	1154	3.67
N(3030)	1/2( ? )	T=4.26 p=4.40	3030	400	9.18 ±1.21	Nπ	(J+1/2)κ=0.05 <sup>b</sup>	1377	2.62
Δ(1236)	3/2(3/2 <sup>+</sup> ) P <sub>33</sub>	T=0.195 p=0.304 m <sub>0</sub> -m <sub>++</sub> = 0.45±0.85	(+) <sup>+</sup> 1236.0 +0.6 m <sub>-</sub> -m <sub>++</sub> = 7.9±6.8	120	1.53 ±0.15	Nπ Nπ <sup>±</sup> Nη	100 0 ~ 0.6	231 89 262	91.9
Δ(1640)	3/2(1/2 <sup>-</sup> ) S <sub>31</sub>	T=0.80 p=0.93	1630	160	2.69 ±0.26	Nπ Nππ	25 75	533	17.2
Δ(1690)	3/2(3/2 <sup>-</sup> ) D <sub>33</sub>	T=0.87 p=1.00	1670	225	2.79 ±0.38	Nπ	15	560	15.6
Δ(1910)	3/2(5/2 <sup>+</sup> ) F <sub>35</sub>	T=1.27 p=1.40	1880	250	3.53 ±0.47	Nπ	20	697	10.1
Δ(1930)	3/2(1/2 <sup>+</sup> ) P <sub>31</sub>	T=1.37 p=1.50	1905	300	3.63 ±0.57	Nπ	25	713	9.62
Δ(1950)	3/2(7/2 <sup>+</sup> ) F <sub>37</sub>	T=1.39 p=1.52	1940	210	3.76 ±0.41	Nπ ΣK Σ(1385)K Δ(1236)π Δ(1236)ρ	40 2.4 1.4 50 seen	735 450 215 564	8.90
Δ(2420)	3/2(1/2 <sup>+</sup> )	T=2.50 p=2.64	2420	310	5.86 ±0.75	Nπ Nππ	11 > 20	1024 1007	4.67
Δ(2850)	3/2( ? )	T=3.71 p=3.85	2850	400	8.12 ±1.14	Nπ	(J+1/2)κ=0.25 <sup>b</sup>	1266	3.05
Δ(3230)	3/2( ? )	T=4.94 p=5.08	3230	440	10.4 ±1.4	Nπ	(J+1/2)κ=0.05 <sup>b</sup>	1475	2.24
Λ	0(1/2 <sup>+</sup> )		1115.6		1.24	See Table S			
Λ(1405)	0(1/2 <sup>-</sup> ) S <sub>01</sub>	p < 0 K <sup>-</sup> p	1405 ±5 <sup>c</sup>	40 ±10 <sup>c</sup>	1.97 ±0.06	Σπ	100	140	
Λ(1520)	0(3/2 <sup>-</sup> ) D <sub>03</sub>	p=0.392	1518.8 ±1.5	16 ±2	2.31 ±0.02	NK Σπ Δππ Δη Δη	45±4 45±4 10±1 0.9±0.2	235 258 254 350	83.6
Λ <sup>+</sup> (1670)	0(1/2 <sup>-</sup> ) S <sub>01</sub>	p=0.74	1670	25	2.79 ±0.04	NK Δη Σπ	14 33 45	410 66 387	28.5
Λ <sup>+</sup> (1700)	0(3/2 <sup>-</sup> ) D <sub>03</sub>	p=0.78	1690	40	2.86 ±0.07	NK Σπ Δππ Σππ	25 35 20 20	429 403 419 350	26.1
Λ(1815)	0(5/2 <sup>+</sup> ) F <sub>05</sub>	p=1.05	1815 ±5 <sup>c</sup>	75 ±10 <sup>c</sup>	3.30 ±0.14	NK Σπ Σ(1385)π Δη	65 11 9 ~ 1	538 500 359 346	16.7
Λ(1830)	0(5/2 <sup>-</sup> ) D <sub>05</sub>	p=1.08	1830	80	3.35 ±0.15	NK Σπ	10 35	550 510	16.0
Λ(2100)	0(7/2 <sup>-</sup> ) G <sub>07</sub>	p=1.68	2100	140	4.41 ±0.29	NK Σπ Δη ΣK Δω	30 4 3 -1 <10	748 699 617 483 443	8.68
Λ(2350)	0( ? )	p=2.29 Seen in total c. s.	2350	210	5.52 ±0.49	NK	(J+1/2)κ=0.6 <sup>b</sup>	913	5.85
Σ	1(1/2 <sup>+</sup> )		(+)1189.4 (0)1192.5 (-)1197.3		1.41 1.42 1.43	See Table S			
Σ(1385)	1(3/2 <sup>+</sup> ) P <sub>13</sub>	p < 0 K <sup>-</sup> p S=4.8 <sup>±</sup> ↔ (-) 1388±3	(+)1382±1 S=2.1 <sup>±</sup>	(+)36±3 S=2.1 <sup>±</sup>	1.92 ±0.05	Δπ Σπ	90±3 10±3	208 117	
Σ(1610)	1( ? )	p=0.62	1615	65	2.61 ±0.10	NK Δπ Σ(1385)π	small domin. mode seen	355 406 171	37.9
Σ(1660)	1(3/2 <sup>-</sup> ) D <sub>13</sub>	p=0.72	1660	50	2.76 ±0.08	Λ(1405)π NK	large small for both	197 400	29.9
Σ(1700)	1( ? )	p=0.80	1700	110	2.89 ±0.19	Δπ Σπ	large not disentangled	470 411	25.1

(Footnotes are on the back of Table S)

# CLEBSCH-GORDAN COEFFICIENTS AND SPHERICAL HARMONICS

Note: A  $\sqrt{\quad}$  is to be understood over every coefficient; e.g., for  $-8/15$  read  $-\sqrt{8/15}$ .

Notation:

J	J	...
M	M	...
$m_1$	$m_2$	
$m_1$	$m_2$	Coefficients
.	.	.
.	.	.

$$1/2 \times 1/2$$

1	0	0
+1/2	+1/2	1
+1/2	-1/2	1/2
-1/2	+1/2	1/2
-1/2	-1/2	1

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left( \frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

$$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$$

$$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$$

$$2 \times 1/2$$

5/2	3/2
+2	1/2
+2	-1/2
+1	+1/2
+1	-1/2
0	+1/2
0	-1/2

$$3/2 \times 1/2$$

2	1
+3/2	+1/2
+3/2	-1/2
+1/2	+1/2
+1/2	-1/2
0	+1/2
0	-1/2

$$1 \times 1/2$$

3/2	1/2
+1	+1/2
+1	-1/2
0	+1/2
0	-1/2

$$2 \times 1$$

3	2	1
+2	+1	0
+2	0	1/3
+1	+1	2/3
+1	0	1/3
0	+1	2/3
0	0	1/3

$$3/2 \times 1$$

5/2	3/2
+3/2	+1
+3/2	0
+1/2	+1
+1/2	0
0	+1
0	0

$$1 \times 1$$

2	1
+1	+1
+1	0
0	+1
0	0

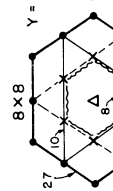
$$Y_l^{-m} = (-1)^m Y_l^{m*}$$

$$\langle j_1 j_2 m_1 m_2 | j_1 j_2 J M \rangle = (-1)^{J-j_1-j_2} \langle j_2 j_1 m_2 m_1 | j_2 j_1 J M \rangle$$

## SU(3) ISOSCALAR FACTORS

adapted from J.J. de Swart, Rev. Mod. Phys. 35, 916 (1963)

(8) ⊗ (8) = [27] ⊕ [10] ⊕ [10\*] ⊕ [8]₁ ⊕ [8]₂ ⊕ [1]₁  
 \* Five single-coefficient tables are omitted. The one involving a {10\*} has a negative coefficient, i.e. (NK|10\* = -1. The others, involving {27} and {10}, are all +1.



Multiplicity of 27:  $\bullet = 1, \times = 2, \Delta = 3$

$\xi_{1 \rightarrow 27}$	$\xi_{1 \rightarrow 10}$	$\xi_{1 \rightarrow 10^*}$	$\xi_{1 \rightarrow 8_a}$	$\xi_{1 \rightarrow 8_s}$	$\xi_{1 \rightarrow 10}$
$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$
$\begin{matrix} \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$

The phase factor  $\xi_1 = \pm 1$ , from de Swart's Table I, enters into his symmetry formula (14.3):  
 $\langle \mu_1 \mu_2 | \mu \rangle = \xi_1 (-1)^{I_1 + I_2 - I} \langle \mu_2 \mu_1 | \mu \rangle$   
 This factor is irrelevant if you are doing your own self-consistent calculations; it enters when you try to check someone else who chose  $\mu_2 \otimes \mu_1$  instead of  $\mu_1 \otimes \mu_2$ .

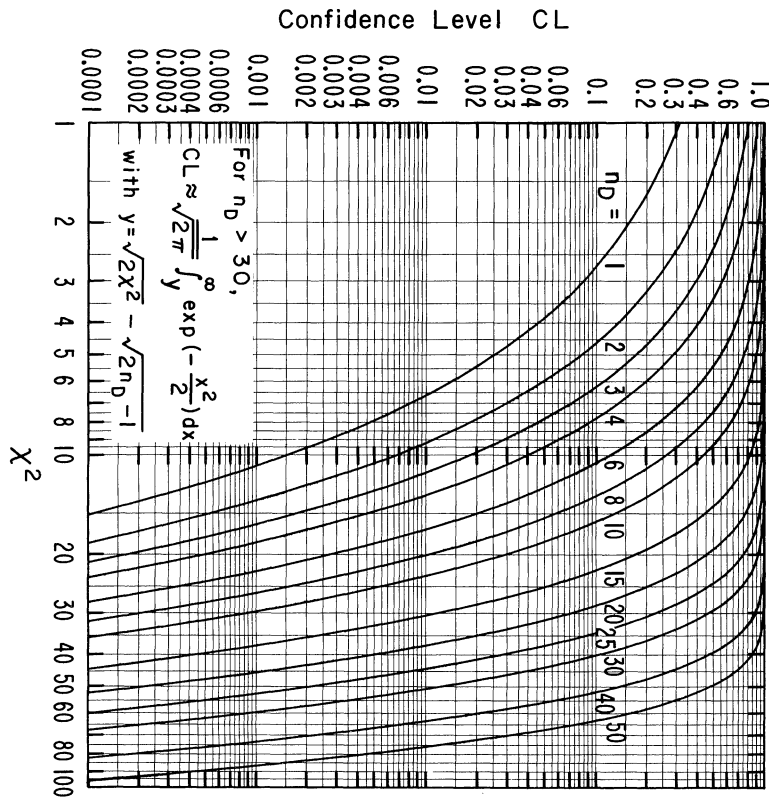
\* Four single coefficient tables are omitted; only the {27} is -1; the three with {35} are +1.

$\xi_{1 \rightarrow 27}$	$\xi_{1 \rightarrow 10}$	$\xi_{1 \rightarrow 10^*}$	$\xi_{1 \rightarrow 8_a}$	$\xi_{1 \rightarrow 8_s}$	$\xi_{1 \rightarrow 10}$
$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$	$\begin{matrix} \pi\Delta \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \\ \pi\Sigma \\ \pi\Lambda \\ \pi N \end{matrix}$
$\begin{matrix} \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \\ -\sqrt{10/5} \\ \sqrt{10/5} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$	$\begin{matrix} \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \\ -\sqrt{10/10} \\ \sqrt{10/10} \end{matrix}$

Multiplicity of 35:  $\bullet = 1, \times = 2$



# CONFIDENCE LEVEL VS. $\chi^2$ FOR $n_D$ DEGREES OF FREEDOM



## OTHER ATOMIC AND PHYSICAL CONSTANTS (cont'd)

Hydrogen-like atom (nonrelativistic,  $\mu$  = reduced mass)

$$\frac{v}{c} = \frac{ze^2}{n\hbar c}; E_n = \frac{\mu}{2} v^2 = \frac{\mu z^2 e^4}{2(n\hbar)^2}; a_n = \frac{2a_0}{n}; \mu = \frac{m_e m_p}{m_e + m_p}$$

$$R_\infty = m_e e^4 / 2\hbar^2 = m_e c^2 \alpha^2 / 2 = 13.60535 \text{ eV (Rydberg)}$$

pc = 0.3 Hp (MeV, kilogauss, cm); 0.3 (which is  $10^{-11}$ c) enters because there are  $\approx 300$  "volts"/esu volt.

- 1 year =  $3.1557 \times 10^7$  sec ( $\approx \pi \times 10^7$  sec)
- density of air =  $1.205 \text{ mg cm}^{-3}$  (at  $20^\circ\text{C}$ , 760 mm)
- acceleration by gravity =  $980.62 \text{ cm sec}^{-2}$  (sea level,  $45^\circ$ )
- gravitational constant =  $6.670 \times 10^{-8} \text{ cm}^3 \text{ g}^{-1} \text{ sec}^{-2}$
- 1 calorie = 4.184 joules
- 1 atmosphere = 1033.2 g  $\text{cm}^{-2}$

## NUMERICAL CONSTANTS

- 1 rad = 57.29578 deg
- ln 2 = 0.69315
- ln 10 = 2.30259
- e = 2.71828
- 1/c = 0.367879
- log<sub>10</sub>e = 0.43429
- log<sub>10</sub>2 = 0.30103

## GAUSSIANLIKE DISTRIBUTIONS

For  $n > -1$  but not necessarily integral:

$$\int_0^\infty x^{2n+1} \exp\left[-\frac{x^2}{2\sigma^2}\right] dx = 2^n n! \sigma^{2n+2}; \left(\frac{1}{2}\right)! = \sqrt{\pi}/2$$

Relation between standard deviation  $\sigma$  and mean deviation  $\sigma$ :

$$2\sigma^2 = \pi\sigma^2; \sigma = 1.4826 \text{ probable error}$$

Odds against exceeding one standard deviation = 2.15:1; two, 21:1; three, 370:1; four, 16,000:1; five, 1,700,000:1.

## MULTIPLE COULOMB SCATTERING<sup>b</sup>

The rms projected angle  $\theta$  due to multiple Coulomb scattering (only) of a particle of charge  $z$ , momentum  $P$ , velocity  $V$  is

$$\theta_{\text{proj}} = z \frac{15(\text{MeV})}{PV(\text{MeV})} \sqrt{\frac{L}{L(\text{rad})}} (1 + \epsilon) \text{ radians};$$

where  $L$  = length in scatterer.

For  $L \geq 1/10 L(\text{rad})$   $\epsilon$  is generally  $< 1/10$ . The distribution of  $\theta$  is not truly Gaussian.<sup>c</sup> The rms projected displacement  $y$  on traversing an absorber of thickness  $L$  is  $y_{\text{rms}} = L\theta_{\text{proj}}/\sqrt{3}$ .

## RADIOACTIVITY

- 1 curie =  $3.7 \times 10^{10}$  disintegrations/sec
- 1 R = 87.8 ergs/g air =  $5.49 \times 10^7$  MeV/g air
- Fluxes (per  $\text{cm}^2$ ) to liberate 1 R in carbon:  $3 \times 10^7$  minimum ionizing singly charged particles
- $0.9 \times 10^9$  photons of 1 MeV energy.
- (These fluxes are actually correct to within a factor of two for all materials.)

1 R of radiation, particularly for neutrons, may produce up to  $\sim 10$  "rem" (R equivalent for man), even 20 "rem" for  $\alpha$  and other heavy ions.

Natural background: 120  $\rightarrow$  130 milirem/year

- divided as follows:
- cosmic radiation - charged part. + neutrons  $\sim 25$  milirem/yr
- " " -  $\gamma$   $\sim 25$  " "
- Rock and air -  $\gamma$   $\sim 73$  " "

The permissible occupational dose for the whole body: 400 milirem/week, but 1.25 rem per calendar quarter.

b. Mainly from G. Z. Molière, Naturforsch. 3 (a), 78 (1948).

c. See, for example, the experimental work of A. D. Hansen, L. H. Lanzl, E. M. Lyman and M. B. Scott, Phys. Rev. 84, 634 (1951).

## Atomic and Nuclear Properties of Materials

Material	Z	A	Cross Section $\sigma^a$ barns	Collision Length $L_{\text{coll}}^b$ g $\text{cm}^{-2}$	Length $L_{\text{coll}}^b$ cm	Minimum $-dE/dx^c$ MeV $\text{g}^{-1} \text{cm}^2$	MeV $\text{cm}^{-1}$	Radiation Length $L_{\text{rad}}^d$ g $\text{cm}^{-2}$	Length $L_{\text{rad}}^d$ cm	Density $\rho$ g $\text{cm}^{-3}$
H <sub>2</sub>	1	1.01	0.063	26.5	374	4.13	0.292	58.0	819	0.0708 <sup>e</sup>
D <sub>2</sub>	1	2.01	0.100	33.4	202	2.07	0.342	116	703	0.165 <sup>e</sup>
He	2	4.00	0.16	42.0	336	1.94	0.242	85.4	683	0.125 <sup>e</sup>
Li	3	6.94	0.23	50.1	94.3	1.69	0.902	78.7	148	0.534
Be	4	9.01	0.28	55.0	29.9	1.60	2.96	63.7	34.7	1.848
C	6	12.01	0.33	60.4	f	1.78	f	42.4	f	$\approx 1.55^f$
N <sub>2</sub>	7	14.01	0.36	63.6	78.9	1.84	1.46	37.8	46.7	0.808 <sup>e</sup>
Ne	10	20.18	0.465	72.1	60.1	1.73	2.08	29.1 <sup>i</sup>	24.2 <sup>i</sup>	1.200 <sup>e</sup>
Al	13	26.98	0.57	79.2	29.3	1.62	4.37	24.0	8.9	2.70
Fe	26	55.85	0.92	104.2	12.8	1.48	11.6	13.9	1.8	7.87
Cu	29	63.54	1.00	105.4	11.8	1.44	42.9	12.0	1.34	8.96
Sn	50	118.69	1.55	129.7	17.8	1.28	9.4	8.89	1.22	7.31
W	74	183.85	2.02	150.8	7.80	1.17	22.6	6.89	0.36	19.3
Pb	82	207.19	2.20	156.2	13.8	1.13	12.8	6.52	0.58	11.35
U	92	238.03	2.42	163.6	$\approx 8.63$	1.09	$\approx 20.6$	6.13	$\approx 0.32$	$\approx 18.95$
Air				64.6	53620	1.81	0.0022	36.5	30290	0.001205 <sup>g</sup>
Freon (CF <sub>3</sub> Br)				87.1	$\approx 58.0$	1.52	$\approx 2.3$	16.6	$\approx 11$	$\approx 1.5$
H <sub>2</sub> (bubble chamber, 27°K)				26.5	442	4.13	0.248	58.0	970	$\approx 0.060^h$
H-Ne mixture (bubble chamber) <sup>j</sup>				67.3	96.1	1.83	1.28	29.8 <sup>i</sup>	42.5 <sup>i</sup>	.70
H <sub>2</sub> O				57.2	57.2	2.03	2.03	35.7	35.7	1.00
Ilford Emulsion				103.0	27.0	5.49	5.49	11.2	2.91	3.815
LiF				63.8	24.2	1.69	4.46	39.0	14.8	2.64
Mylar (C <sub>5</sub> H <sub>4</sub> O <sub>2</sub> )				59.1	42.8	1.91	2.64	39.6	28.7	1.38
NaI				119.0	32.4	1.32	4.84	9.58	2.61	3.67
Polyethylene (CH <sub>2</sub> )				51.0	$\approx 55.5$	2.09	$\approx 1.92$	44.1	$\approx 48$	$\approx 0.92$
Polystyrene (CH) [°]				54.9	$\approx 52.3$	2.03	$\approx 2.14$	43.4	$\approx 41.3$	$\approx 1.05$
Propane (C <sub>3</sub> H <sub>8</sub> bubble chamber)				48.9	119.3	2.28	0.935	44.6	109	0.44

$$a. \sigma = \sigma_{\text{natural}} = \pi (h/m_p c)^2 \times A^{2/3} = 62.8 \text{ mb} \times A^{2/3}$$

$$b. L_{\text{coll}} = A / (N\sigma_{\text{natural}}) = 26.5 \text{ g cm}^{-2} \times A^{1/3}$$

c. From W. H. Barkas and M. J. Berger, Tables of Energy Losses and Ranges of Heavy Charged Particles, NASA SP-3013 (1964)

d. Mainly from High Energy and Nuclear Physics Data Handbook, W. Galbraith and W. S. C. Williams, Ed. (N. I. R. N. S., Rutherford Lab., Chilton, Didcot, Berks.) 1964

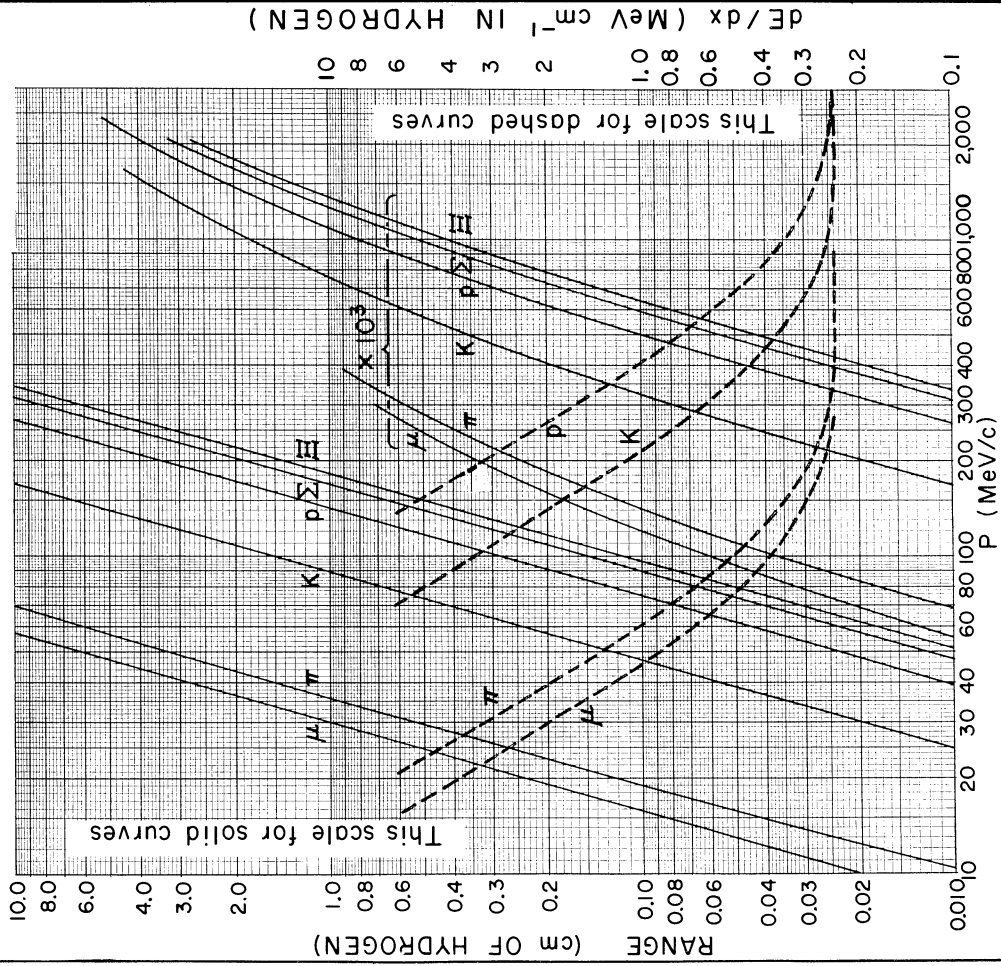
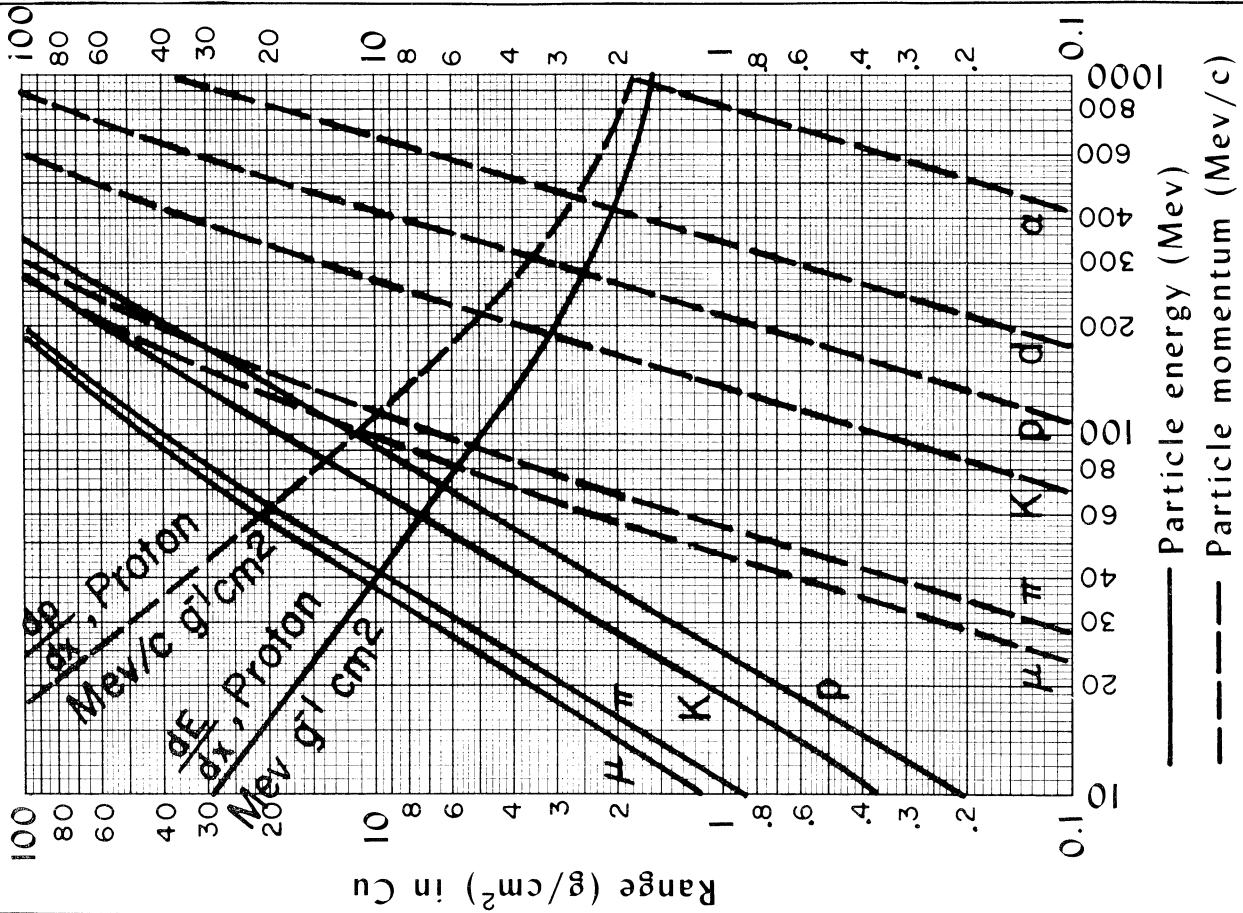
e. boiling at 1 atmosphere f, density variable g. at 20°C

h. May vary by about  $\pm 3\%$ , depending on operation conditions

i. From F. R. Huson, Ionization Loss, Range, Straggling and Multiple Scattering, BNL 11386 (1967)

j. 53.7 atomic percent Ne.

[\*] Typical scintillator, e.g. PILOT B has H/C = 1.1



Ranges in liquid hydrogen bubble chamber, determined by a  $\mu^+$  range of  $4.403 \pm 0.003$  cm from the  $\pi^+ \rightarrow \mu^+ \nu$  decay. Liquid hydrogen conditions:  $T = 27.6 \pm 0.1$  K;  $P = 48 \pm 5$  psia;  $\rho = (5.86 \pm 0.06) 10^{-2}$  g/cm<sup>3</sup>. (Data by Clark and Diehl.)