

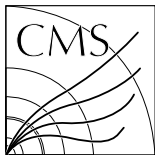
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Correlations between azimuthal anisotropy and mean transverse momentum  
in pp, pPb, and peripheral PbPb collisions

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# Correlations between azimuthal anisotropy and mean transverse momentum in pp, pPb, and peripheral PbPb collisions

The CMS Collaboration\*

## Abstract

Correlations between azimuthal anisotropy and mean transverse momentum of charged particles in proton-proton (pp), proton-lead (pPb), and peripheral lead-lead (PbPb) collisions are presented as a function of charged particle multiplicity. The pp, pPb and PbPb collision data were collected using the CMS detector at the LHC with a center-of-mass energy per nucleon pair of 13, 8.16 and 5.02 TeV, respectively. The two- and four-particle cumulants for the second- and third-order Fourier anisotropy harmonics are correlated with the mean transverse momentum of charged particles on an event-by-event basis. In pp and pPb systems, the observed correlation coefficients based on two-particle cumulants are found to change from negative to positive values as the charged particle multiplicity decreases. The sign changes disappear when the correlated particles are required to be further apart in pseudorapidity. Additionally, no sign changes in correlation coefficients are observed when employing four-particle cumulants. Models incorporating initial-state gluon saturation and final-state hydrodynamic evolutions are compared to pPb data and the predicted sign changes are not observed.

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A hot and dense medium known as the quark-gluon plasma (QGP) has been extensively studied using heavy ion collisions at the BNL RHIC [1–4] and CERN LHC [5–12]. The azimuthal anisotropy of the produced particles in these collisions is a powerful tool to study the collective dynamics and transport properties of the QGP. This anisotropy is characterized by the Fourier coefficients ( $v_n$ ) of the particle azimuthal angle ( $\phi$ ) distribution  $dN/d\phi \propto 1 + 2 \sum_n v_n \cos[n(\phi - \Psi_n)]$ , where  $\Psi_n$  represents the phase of the  $n^{\text{th}}$ -order azimuthal flow vector (also referred to as the angle of the  $n^{\text{th}}$ -order symmetry plane). In a hydrodynamic picture, the Fourier anisotropy harmonics are the result of strong partonic rescatterings in the final state, responding to the initial geometry of the colliding system [13–15]. In the past decade, a remarkable similarity in the azimuthal anisotropy signatures has been observed between heavy ion collisions and smaller collision systems, such as proton-proton (pp), proton-lead (pPb), and proton-gold (pAu) [16–26]. The similarity holds even for multiparticle correlations, which can suppress “nonflow” effects that do not result from correlations related to the bulk properties of the medium [23, 27]. These effects can include back-to-back jet correlations and resonance decays, which usually produce particles in localized pseudorapidity ( $\eta$ ) regions [28–30]. Apart from the final-state effects, the observed anisotropy in small systems can also originate from initial-state effects, for example, in the color-glass condensate (CGC) effective theory [31–33]. The dominant origin of the azimuthal anisotropy in small systems is still under active discussion [34, 35], because no observable has been found to unambiguously distinguish between the final- and initial-state effects.

In addition to generating the final-state azimuthal anisotropy the hydrodynamic response also results in a radial flow, which contributes to the mean transverse momentum  $[p_T]$  on an event-by-event basis. The correlations between radial and anisotropic flow can be quantified using a modified linear correlator [36, 37],

$$\rho(v_n^2, [p_T]) = \frac{\text{cov}(v_n^2, [p_T])}{\sqrt{\text{Var}(v_n^2)_{\text{dyn}}} \sqrt{\text{Var}([p_T])_{\text{dyn}}}}, \quad (1)$$

where  $\text{cov}(v_n^2, [p_T])$  is the covariance between  $v_n^2$  and  $[p_T]$ , and  $\text{Var}(v_n^2)_{\text{dyn}}$  and  $\text{Var}([p_T])_{\text{dyn}}$  are the dynamical variances of the  $v_n^2$  and  $[p_T]$  distributions, respectively. Dynamical variances remove auto-correlation effects, when compared with variances of  $v_n^2$  and  $[p_T]$  distributions, and better capture intrinsic initial-state fluctuations. This correlator defined by Eq. (1) is sensitive to the details of the initial conditions, and its magnitude can be traced back to the initial density profile of the nuclear overlap [38, 39].

Recently it was suggested that this correlator might be able to distinguish between initial- and final-state effects [40]. Specifically, the CGC model predicts a sign change of the  $\rho(v_n^2, [p_T])$  correlator going from high-multiplicity (dominated by final-state effects) to very low-multiplicity (dominated by initial-state effects) events, in small collision systems. A prior ATLAS measurement using two-particle cumulants shows no sign change across the explored multiplicity range in pPb collisions [41]. However, these results do not cover the multiplicity range where the initial-state sign change is expected. It was realized that a sign change could exist due to nonflow effects as shown by the PYTHIA8 event generator with the A2 tune [42, 43]. Measurements of the correlator with proper treatment of nonflow effects, and the searches for correlator sign changes in low-multiplicity pp, pPb, and peripheral PbPb collisions can provide insights into the origin of the azimuthal correlations in small systems.

In this Letter, these correlations are measured to very low multiplicity regions in three collision systems where sign changes are predicted due to the dominance of initial-state effects. Measurements using four-particle cumulants are also presented for the first time to further

suppress nonflow effects. The results in pPb collisions are compared directly with CGC predictions. In addition, the correlators for the third Fourier harmonic are also presented as a function of charged particle multiplicity. Tabulated results are provided in the HEPData record for this analysis [44].

The CMS apparatus [45, 46] is a multipurpose, nearly hermetic detector, designed to trigger on and identify electrons, muons, photons, and (charged and neutral) hadrons [47–49]. A global “particle-flow” (PF) algorithm [50] aims to reconstruct all individual particles in an event, combining information provided by the all-silicon inner tracker and by the crystal electromagnetic and brass-scintillator hadron calorimeters, operating inside a 3.8 T superconducting solenoid, with data from the gas-ionization muon detectors embedded in the flux-return yoke outside the solenoid. Forward hadronic calorimeters (HF), made of steel and quartz fibers, extend the  $\eta$  coverage provided by the barrel and endcap detectors up to  $|\eta| \sim 5$ . The silicon tracker used in 2016 measured charged particles within the range  $|\eta| < 2.5$ . For nonisolated particles of  $1 < p_T < 10$  GeV and  $|\eta| < 1.4$ , the track resolutions were typically 1.5% in  $p_T$  and 25–90 (45–150)  $\mu\text{m}$  in the transverse (longitudinal) impact parameter [49]. At the start of 2017, a new pixel detector was installed [51]; the upgraded tracker measured particles up to  $|\eta| < 3.0$  with typical resolutions of 1.5% in  $p_T$  and 20–75  $\mu\text{m}$  in the transverse impact parameter [52] for nonisolated particles of  $1 < p_T < 10$  GeV. The data sample is collected with a two-level trigger system: at level-1 events are selected by custom hardware processors and the high-level trigger uses fast versions of the offline software [53, 54].

The measurements presented in this Letter use pp, pPb and PbPb collisions with center-of-mass energy per nucleon pairs of 13, 8.16 and 5.02 TeV, taken in 2018, 2016, and 2018, and with integrated luminosities of  $28.6 \text{ pb}^{-1}$ ,  $186 \text{ nb}^{-1}$  and  $0.607 \text{ nb}^{-1}$ , respectively [55–58]. When measuring the correlations in pp and pPb collisions, the same beam crossing may contain multiple independent interactions, which constitute a background for the analysis of high-multiplicity events. The average number of collisions per bunch crossing (pileup) in pp (pPb) data varied between 0.1 and 1.3 (0.10 and 0.25). A procedure similar to that described in Ref. [27] is used for identifying and rejecting pileup. Minimum bias (MB) pp and pPb events are triggered by energy deposits in at least one of the two HFs above a threshold of approximately 1 GeV, and the presence of at least one track with  $p_T > 0.4$  GeV in the pixel tracker for pPb collisions. Events are also required to contain a primary vertex within 15 cm of the nominal interaction point along the beam axis and 0.2 cm in the transverse direction. The primary vertex is chosen as the reconstructed vertex with the largest number of associated tracks. At least two reconstructed tracks were required to be associated with the primary vertex. The MB PbPb events are selected to have signals above readout thresholds in the range of  $\sim 6$ –12 GeV on both sides of the HF calorimeters. The PbPb events are further filtered to have a primary vertex within 15 cm of the nominal interaction point along the beam axis and 0.2 cm in the transverse direction, and at least two towers with total deposited energy larger or equal to 4 GeV in each of the HF detectors. The trigger, event reconstruction, and selections are described in previous correlation analyses [27–30]. For all data sets analyzed, primary tracks, i.e., tracks that originate from the primary vertex and satisfy the high-purity criteria of Ref. [49], are used to perform the correlation measurements. In addition, the impact parameter significance of the tracks with respect to the primary vertex in the longitudinal and transverse direction is required to be less than 3 standard deviations. The relative  $p_T$  uncertainty must be less than 10% for the  $p_T$  range used. To ensure high tracking efficiency and minimal background contamination only tracks with  $|\eta| < 2.4$  and  $p_T > 0.5$  GeV are used, as in Ref. [49]. The selected tracks are corrected for tracking inefficiency and acceptance found using simulated Monte Carlo samples from PYTHIA 8.212 [59] tune CP5 [60], HIJING v1.35 [61], and HYDJET 1.9 [62] for pp, pPb, and PbPb, respec-

tively. If not specified otherwise, PYTHIA8 refers to this CP5 tune throughout the paper. The response of the CMS detector to these simulated events is based on GEANT4 [63].

All previous studies [40–43, 64] of the modified linear correlator use  $v_n^2$  from two-particle correlations in Eq. (1). The  $v_n^2$  term can be found from the Q-cumulant method described in Ref. [65]. The two- and four-particle  $v_n$  in this method can be written as,

$$v_n\{2\} = \sqrt{c_n\{2\}}, \quad v_n\{4\} = \sqrt[4]{-c_n\{4\}}. \quad (2)$$

The two- and four-particle cumulants  $c_n\{2\}$  and  $c_n\{4\}$  are

$$c_n\{2\} = \langle\langle 2 \rangle\rangle, \quad c_n\{4\} = \langle\langle 4 \rangle\rangle - 2\langle\langle 2 \rangle\rangle^2, \quad (3)$$

where the two- and four-particle correlations  $\langle\langle 2 \rangle\rangle$  and  $\langle\langle 4 \rangle\rangle$  are from

$$\langle\langle 2 \rangle\rangle = \langle\langle e^{in(\phi_1 - \phi_2)} \rangle\rangle, \quad \langle\langle 4 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle\rangle. \quad (4)$$

Here,  $\phi_j$  ( $j = 1, \dots, 4$ ) are the azimuthal angles of four different particles in an event, and the double average symbol  $\langle\langle \dots \rangle\rangle$  indicates that the average is taken over all particles from all events.

To suppress nonflow effects in the  $v_n\{2\}$ , all previous measurements applied the ‘‘subevent method’’, which separates each event into different subsets in separated  $\eta$  ranges [41, 64]. With two subevents for  $c_n\{2\}$ , the covariance of the correlator in Eq. (1) is

$$\text{cov}(c_n\{2\}, [p_T]) = \Re \left\langle \sum_{a,b} \exp^{in(\phi_a - \phi_b)} ([p_T] - \langle [p_T] \rangle) \right\rangle, \quad (5)$$

where  $\phi_a$  and  $\phi_b$  are the azimuthal angles of particles  $a$  and  $b$  in subevents A and B, respectively. The  $\langle [p_T] \rangle$  is the average  $[p_T]$  in all the events in a certain multiplicity range. We select tracks with  $-2.4 < \eta < -0.75$  to be subevent A, whereas tracks with  $0.75 < \eta < 2.4$  belong to subevent B. Tracks from the third subevent in the central region  $|\eta| < 0.5$  are used to obtain  $[p_T]$  in each event. These selections ensure that subevents A and B are symmetric in  $\eta$ , and that there is an  $\eta$  gap between each subevent to reduce nonflow effects.

The dynamical variance of  $c_n\{2\}$  and  $[p_T]$  are derived in Ref. [37]:

$$\text{Var}(c_n\{2\})_{\text{dyn}} = \langle\langle 4 \rangle\rangle - \langle\langle 2 \rangle\rangle^2, \quad (6)$$

$$\text{Var}([p_T])_{\text{dyn}} = \left\langle [(p_{Ti} - \langle [p_T] \rangle)(p_{Tj} - \langle [p_T] \rangle)] \right\rangle, \quad (7)$$

where the indices  $i \neq j$  and both run over all charged particles within  $|\eta| < 0.5$ .

The correlator  $\rho(c_n\{2\}, [p_T])$  is extracted using particles in the range  $0.5 < p_T < 5.0$  GeV. Particles with  $p_T > 5$  GeV are excluded to reduce nonflow effects. The analysis is repeated in different ranges of the number of reconstructed charged particles  $N_{\text{ch}}^{\text{rec}}$ , which is obtained with the same  $p_T$  range as the correlator within  $|\eta| < 2.4$ . The results are presented as a function of  $N_{\text{ch}}$ , which is found by applying acceptance and efficiency corrections to  $N_{\text{ch}}^{\text{rec}}$ . For each  $N_{\text{ch}}^{\text{rec}}$  range used in this analysis, Table 1 of the Appendix reports the mean values of  $N_{\text{ch}}$  and also of  $N_{\text{trk}}^{\text{offline}}$ , the number of offline-reconstructed charged particles with  $p_T > 0.4$  GeV and  $|\eta| < 2.4$  used in previous CMS measurements [16, 19, 27]. Since we are interested in small collision systems, only peripheral events with  $N_{\text{ch}} < 400$  are presented for PbPb collisions.

Studies using A2 tune of PYTHIA8 [42, 43] show that the remaining nonflow contributions to  $c_2\{2\}$  are not negligible when going to very low-multiplicity regions, even with the subevent

methods. In these studies, we improve the suppression of nonflow effects with two approaches. In the first approach, we increase the minimum  $\eta$  gap between subevents A and B from 1.5 to 2.0, by changing the  $c_2\{2\}$  analysis using particles from  $|\eta| > 0.75$  to  $|\eta| > 1.00$ . Nonflow correlations between particles inside back-to-back jets in the forward  $\eta$  regions still remain with this larger  $\eta$  gap [27].

In the second approach, we extend the current observable by replacing  $c_2\{2\}$  with four-particle cumulant  $c_2\{4\}$ . Although nonflow contributions to multiparticle cumulants are less dominant when compared with two-particle correlations, using the subevent method is important to remove remaining nonflow contributions in multiparticle correlations [66]. Particles with  $0.75 < |\eta| < 2.4$  are divided into three equal sized  $\eta$  intervals to obtain  $c_2\{4\}$  for each event. These intervals are  $-2.4 < \eta < -1.3$ ,  $0.75 < |\eta| < 1.3$ , and  $1.3 < \eta < 2.4$ . The event-by-event  $c_2\{4\}$  is then correlated with  $[p_T]$  in the same event. When nonflow contributions are not negligible, the method using  $c_2\{4\}$  suppresses them much more than the  $c_2\{2\}$  method. The correlator using  $c_2\{4\}$  is more sensitive to the presence of a sign change in the data. Corrections for tracking efficiency and misreconstructed tracks are done at particle level for both the cumulants and  $[p_T]$ .

The systematic uncertainties in the experimental procedure are evaluated as a function of  $N_{\text{ch}}$  by using alternative procedures for extracting the correlators. Systematic uncertainties due to tracking inefficiency and misreconstructed track rate are studied by varying the track quality requirements. The selection thresholds on the significances of the transverse and longitudinal track impact parameters divided by their uncertainties are varied from 2 to 5. In addition, the upper limit on the relative  $p_T$  uncertainty is varied from 5 to 10%. The resulting systematic uncertainty is from  $2 \times 10^{-7}$  to  $5 \times 10^{-6}$  for the covariances  $\text{cov}(c_n\{2\}, [p_T])$  and  $\text{cov}(c_n\{4\}, [p_T])$ , and from 0.001 to 0.008 for the correlator  $\rho(c_n\{2\}, [p_T])$ , depending on multiplicity and collision systems. The sensitivity of the results to the primary vertex position along the beam axis ( $z_{\text{vtx}}$ ) is quantified by comparing events with different  $z_{\text{vtx}}$  locations from  $-15$  to  $+15$  cm. The magnitude of this uncertainty is estimated to be from  $1 \times 10^{-7}$  to  $2 \times 10^{-5}$  for the covariances, and from 0.001 to 0.018 for  $\rho(c_n\{2\}, [p_T])$ , depending on multiplicity and collision system. Systematic effects from event selections are explored by comparing results obtained with and without the requirement of the coincidence of HF calorimeter tower signals above the threshold. The pileup effect is studied by requiring the presence of only a single reconstructed vertex. Systematic uncertainties from event selections and pileup are both found to be less than  $6 \times 10^{-6}$  for the covariances, and less than 0.006 for the correlators in all collision systems. The tracking and vertex position selections are the dominant sources of uncertainty. Systematic uncertainties originating from different sources are added in quadrature to obtain the total systematic uncertainty.

The measurements of covariances from two- and four-particle correlations for the second- and third-order Fourier harmonics in 13 TeV pp, 8.16 TeV pPb, and 5.02 TeV PbPb collisions are presented in Fig. 1. To compare  $\text{cov}(c_2\{2\}, [p_T])$  and  $\text{cov}(c_2\{4\}, [p_T])$  in the same scale, the values of  $\text{cov}(c_2\{4\}, [p_T])$  are multiplied by 4 in all of the panels. In both pp and pPb collisions,  $\text{cov}(c_2\{2\}, [p_T])$  for  $|\eta| > 0.75$  exhibits a sign change from positive to negative as  $N_{\text{ch}}$  increases. This trend is qualitatively consistent with the expectation of a sign change from the CGC model. However, no clear sign change is observed in pp and pPb collisions for  $\text{cov}(c_2\{4\}, [p_T])$ , which mitigates nonflow contributions compared to  $\text{cov}(c_2\{2\}, [p_T])$ . The  $\text{cov}(c_2\{4\}, [p_T])$  values are consistent with 0 in pp collisions with the current statistical precision. As  $N_{\text{ch}}$  becomes smaller in PbPb collisions, the values of  $\text{cov}(c_2\{2\}, [p_T])$  change from positive to negative below  $N_{\text{ch}} \approx 180$ , reach a minimum at  $N_{\text{ch}} \approx 60$ , and then approach zero at the lowest  $N_{\text{ch}}$  range.

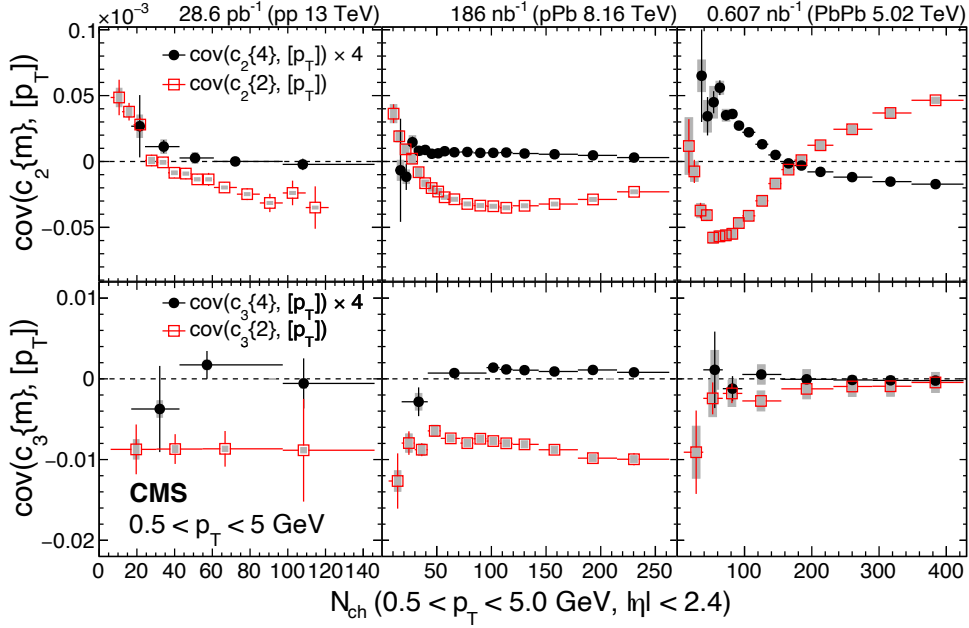


Figure 1: The covariances of cumulants from two- and four-particle correlations and  $[p_T]$  as a function of charged particle multiplicity ( $N_{\text{ch}}$ ) in 13 TeV pp (left), 8.16 TeV pPb (middle), and 5.02 TeV PbPb (right) collisions. The two-particle cumulants are obtained with  $|\eta| > 0.75$ . The upper (lower) panels show the second (third) harmonics. The error bars correspond to statistical uncertainties, while the shaded areas denote the systematic uncertainties.

One feature of the observable with  $c_2\{4\}$  is that, when nonflow contributions are not dominant, the variables using  $c_2\{4\}$  and  $c_2\{2\}$  have opposite signs and trends. This is because the relations between cumulants and  $v_n$  are  $c_2\{2\} = (v_2\{2\})^2$  and  $c_2\{4\} = -(v_2\{4\})^4$  from Eq. (2). As  $v_n$  increases, both  $v_2\{2\}$  and  $v_2\{4\}$  become larger if nonflow and flow fluctuations do not dominate the flow signal. Therefore, the values increase for  $c_2\{2\}$  but decrease for  $c_2\{4\}$  because of the minus sign. This opposing trends for  $\text{cov}(c_2\{2\}, [p_T])$  and  $\text{cov}(c_2\{4\}, [p_T])$  are clear for high-multiplicity PbPb results in Fig. 1.

Regarding the third harmonic,  $\text{cov}(c_3\{2\}, [p_T])$  is negative for the full  $N_{\text{ch}}$  range in the three collision systems. The signs of  $\text{cov}(c_3\{2\}, [p_T])$  and  $\text{cov}(c_3\{4\}, [p_T])$  are opposite in high-multiplicity pPb events, similar to  $\text{cov}(c_2\{2\}, [p_T])$  and  $\text{cov}(c_2\{4\}, [p_T])$ . In both pp and PbPb collisions, the  $\text{cov}(c_3\{4\}, [p_T])$  values are consistent with 0 with the current statistical precision.

The correlator with a larger  $\eta$  gap ( $|\eta| > 1.0$  corresponding to a minimum  $\eta$  gap of 2.0) for the cumulants is shown in Fig. 2. In both pp and pPb collisions, the sign change at low  $N_{\text{ch}}$  disappears with the larger  $\eta$  gap between the two subevents, which is also observed in calculations using PYTHIA8. The predictions in pPb collisions at 5.02 TeV from the IP-Glasma+MUSIC+UrQMD model [40] with  $0.5 < p_T < 5.0$  GeV are compared with the data in Fig. 2. This model includes gluon saturation in the initial state followed by hydrodynamic evolution and hadronic interactions. The characteristic sign change of the correlator predicted by this model is observed in data at the same  $N_{\text{ch}}$  location for  $|\eta| > 0.75$ , but it disappears when using  $|\eta| > 1.0$ , which leads to less nonflow effects. The results indicate that after removing more nonflow effects, the CGC signal is not observed in the data. More nonflow studies incorporated in the CGC model are needed to understand the origin of azimuthal anisotropy in small systems. Results from the ATLAS Collaboration in pPb (PbPb) collisions at 5.02 TeV with  $0.3 < p_T < 2.0$  GeV ( $0.5 < p_T < 5.0$  GeV) [41] are also compared and found to be consistent with results obtained from this analysis.



The  $n = 3$  correlators measured with the larger  $\eta$  gap are negative in all three collision systems. The results are compared with predictions from a hydrodynamic simulation of pPb collisions [67] using  $p_T > 0.5$  GeV and having an average root-mean-square (RMS) transverse radius of the initial fireball of either 0.9 or 1.5 fm. These sizes correspond to two versions of the Glauber model for the initial state [68]: the smaller fireball with 0.9 fm deposits entropy in the nucleon overlap region between the participant nucleons, while the larger fireball with 1.5 fm deposits entropy at both positions of the participant nucleons. Qualitatively, the data are better described by the smaller initial fireball, suggesting that entropy deposition should occur in the nucleon overlap region rather than at the center of both nucleons. The predictions from IP-Glasma+MUSIC+UrQMD for  $n = 3$  also indicate a sign change, although it is not observed in data.

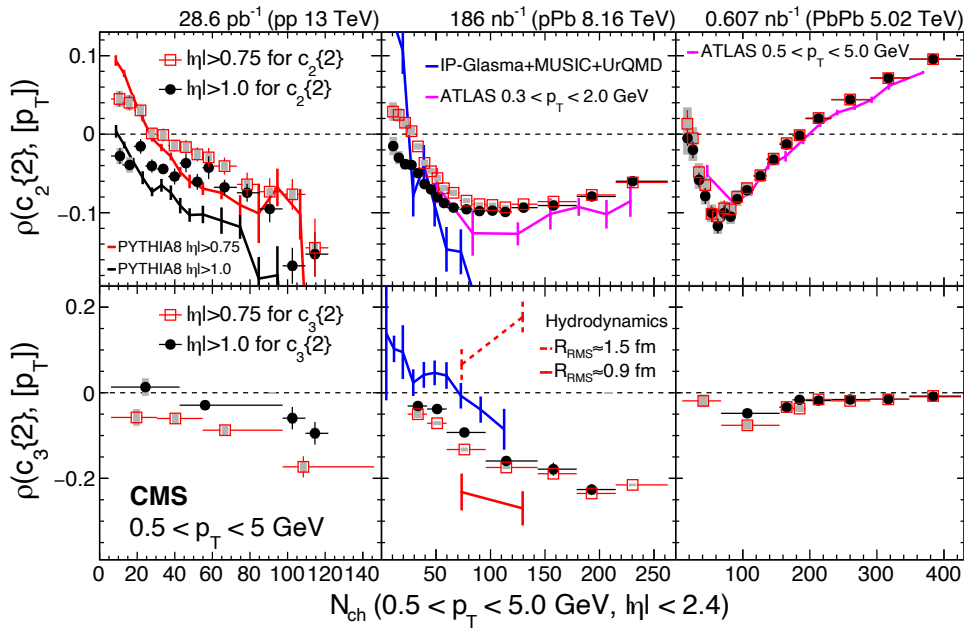


Figure 2: The correlator using two-particle cumulant from  $|\eta| > 0.75$  and  $|\eta| > 1.0$  as a function of  $N_{\text{ch}}$  in 13 TeV pp (left), 8.16 TeV pPb (middle), and 5.02 TeV PbPb (right) collisions. The upper (lower) panels show the second (third) harmonics. The error bars correspond to the statistical uncertainties, and the shaded areas denote the systematic uncertainties. Calculations from PYTHIA8 (upper left panel, red and black lines) and IP-Glasma+MUSIC+UrQMD (middle panels, blue lines) [40] are compared with the data. Hydrodynamic predictions (lower middle panel, red lines) [67] with an average RMS transverse radius of the initial fireball of either 0.9 or 1.5 fm are also included in the comparison. The magenta lines in the middle and right upper panels represent the ATLAS results [41] in pPb and PbPb collisions at 5.02 TeV using  $|\eta| > 0.75$  with the error bars denoting the statistical and systematic uncertainties added in quadrature.

In summary, correlations between mean transverse momentum  $[p_T]$  and multiparticle cumulants from two- and four-particle correlations for Fourier harmonics  $n = 2$  and  $n = 3$  are presented in proton-proton (pp) collisions at  $\sqrt{s} = 13$  TeV, proton-lead (pPb) collisions at  $\sqrt{s_{\text{NN}}} = 8.16$  TeV and peripheral lead-lead (PbPb) collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. Sign changes in the modified linear correlators are observed as a function of charged particle multiplicity when using two-particle cumulants with a minimum  $\eta$  gap of 1.5, in the pp and pPb systems. These sign changes disappear when nonflow effects are suppressed using a minimum  $\eta$  gap of 2.0. To further reduce nonflow contribution, four-particle cumulants  $c_2\{4\}$  are also correlated with  $[p_T]$  which shows no sign change in pp and pPb collisions, similar to the two-particle cor-

relation results with a larger  $\eta$  gap. This indicates that after removing nonflow effects, the sign change predicted at low multiplicity by initial-state saturation is not observed.

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## A Mapping between $N_{\text{ch}}$ and $N_{\text{trk}}^{\text{offline}}$ in pp, pPb, and PbPb collisions

Table 1: Average multiplicity of reconstructed tracks per  $N_{\text{ch}}^{\text{rec}}$  bin for  $N_{\text{ch}}$  in this analysis and  $N_{\text{trk}}^{\text{offline}}$  in previous CMS measurements [16, 19, 27] in pp, pPb, and peripheral PbPb collisions. Uncertainties in the tracking efficiency corrected  $N_{\text{ch}}$  are included.


$N_{\text{ch}}^{\text{rec}}$ range	pp		pPb		PbPb	
	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$	$\langle N_{\text{ch}} \rangle$	$\langle N_{\text{trk}}^{\text{offline}} \rangle$
[0, 20)	$8 \pm 0$	9	$11 \pm 0$	12	$16 \pm 1$	14
[20, 40)	$34 \pm 1$	34	$36 \pm 1$	36	$57 \pm 2$	48
[40, 60)	$58 \pm 2$	56	$60 \pm 2$	60	$96 \pm 4$	80
[60, 80)	$82 \pm 3$	78	$83 \pm 3$	82	$135 \pm 5$	112
[80, 100)	$106 \pm 4$	101	$107 \pm 4$	105	$175 \pm 7$	144
[100, 150)	$132 \pm 5$	125	$140 \pm 6$	137	$240 \pm 10$	197
[150, 200)			$198 \pm 8$	191	$335 \pm 13$	276
[200, 250)			$256 \pm 10$	246	$434 \pm 17$	353



























## B The CMS Collaboration

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


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













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




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J. Mejia Guisao , F. Ramirez , M. Rodriguez , J.D. Ruiz Alvarez 








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
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Z. Antunovic, M. Kovac , T. Sculac 


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V. Brigljevic , B.K. Chitroda , D. Ferencek , S. Mishra , M. Roguljic , A. Starodumov<sup>11</sup> , T. Susa 

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

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

























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























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


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

























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C. Dziwok , G. Flügge , W. Haj Ahmad<sup>21</sup> , O. Hlushchenko, T. Kress , A. Nowack , O. Pooth , A. Stahl , T. Ziemons , A. Zotz 



























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











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H. Bakhshiansohi<sup>44,45</sup> , E. Khazaie<sup>45</sup> , M. Zeinali<sup>46</sup> 
































**Institute for Research in Fundamental Sciences (IPM), Tehran, Iran**

S. Chenarani<sup>47</sup> , S.M. Etesami , M. Khakzad , M. Mohammadi Najafabadi 





























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














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






**INFN Laboratori Nazionali di Frascati, Frascati, Italy**












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



A. Benaglia<sup>a</sup> , G. Boldrini<sup>a</sup> , F. Brivio<sup>a,b</sup> , F. Cetorelli<sup>a,b</sup> , F. De Guio<sup>a,b</sup> , M.E. Dinardo<sup>a,b</sup> , P. Dini<sup>a</sup> , S. Gennai<sup>a</sup> , A. Ghezzi<sup>a,b</sup> , P. Govoni<sup>a,b</sup> , L. Guzzi<sup>a,b</sup> 

M.T. Lucchini<sup>a,b</sup> , M. Malberti<sup>a</sup> , S. Malvezzi<sup>a</sup> , A. Massironi<sup>a</sup> , D. Menasce<sup>a</sup> , L. Moroni<sup>a</sup> , M. Paganoni<sup>a,b</sup> , D. Pedrini<sup>a</sup> , B.S. Pinolini<sup>a</sup>, S. Ragazzi<sup>a,b</sup> , N. Redaelli<sup>a</sup> , T. Tabarelli de Fatis<sup>a,b</sup> , D. Zuolo<sup>a,b</sup> 












**INFN Sezione di Napoli<sup>a</sup>, Università di Napoli 'Federico II'<sup>b</sup>, Napoli, Italy; Università della Basilicata<sup>c</sup>, Potenza, Italy; Scuola Superiore Meridionale (SSM)<sup>d</sup>, Napoli, Italy**

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












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




























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
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



































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

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




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


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







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



























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





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





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















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











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









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
































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H. Flacher , M. Glowacki, J. Goldstein , H.F. Heath , L. Kreczko , B. Krikler ,  
S. Paramesvaran , S. Seif El Nasr-Storey, V.J. Smith , N. Stylianou<sup>79</sup> , K. Walkingshaw Pass,  
R. White 




**Rutherford Appleton Laboratory, Didcot, United Kingdom**

A.H. Ball, K.W. Bell , A. Belyaev<sup>80</sup> , C. Brew , R.M. Brown , D.J.A. Cockerill ,  
C. Cooke , K.V. Ellis, K. Harder , S. Harper , M.-L. Holmberg<sup>81</sup> , Sh. Jain , J. Linacre ,  
K. Manolopoulos, D.M. Newbold , E. Olaiya, D. Petyt , T. Reis , G. Salvi , T. Schuh,  
C.H. Shepherd-Themistocleous , I.R. Tomalin , T. Williams 

**Imperial College, London, United Kingdom**







R. Bainbridge , P. Bloch , S. Bonomally, J. Borg , C.E. Brown , O. Buchmuller, V. Cacchio,  
V. Cepaitis , G.S. Chahal<sup>82</sup> , D. Colling , J.S. Dancu, P. Dauncey , G. Davies , J. Davies,  
M. Della Negra , S. Fayer, G. Fedi , G. Hall , M.H. Hassanshahi , A. Howard, G. Iles ,  
J. Langford , L. Lyons , A.-M. Magnan , S. Malik, A. Martelli , M. Mieskolainen ,  
D.G. Monk , J. Nash<sup>83</sup> , M. Pesaresi , B.C. Radburn-Smith , D.M. Raymond, A. Richards,  
A. Rose , E. Scott , C. Seez , R. Shukla , A. Tapper , K. Uchida , G.P. Uttley ,  
L.H. Vage, T. Virdee<sup>28</sup> , M. Vojinovic , N. Wardle , S.N. Webb , D. Winterbottom 

**Brunel University, Uxbridge, United Kingdom**

K. Coldham, J.E. Cole , A. Khan, P. Kyberd , I.D. Reid 

**Baylor University, Waco, Texas, USA**

S. Abdullin , A. Brinkerhoff , B. Caraway , J. Dittmann , K. Hatakeyama , A.R. Kanu-

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









**Catholic University of America, Washington, DC, USA**

R. Bartek , A. Dominguez , R. Uniyal , A.M. Vargas Hernandez 





















**The University of Alabama, Tuscaloosa, Alabama, USA**

S.I. Cooper , D. Di Croce , S.V. Gleyzer , C. Henderson , C.U. Perez , P. Rumerio<sup>84</sup> ,  
C. West 


















**Boston University, Boston, Massachusetts, USA**

A. Akpınar , A. Albert , D. Arcaro , C. Cosby , Z. Demiragli , C. Erice ,  
E. Fontanesi , D. Gastler , S. May , J. Rohlf , K. Salyer , D. Sperka , D. Spitzbart ,  
I. Suarez , A. Tsatsos , S. Yuan 















**Brown University, Providence, Rhode Island, USA**

G. Benelli , B. Burkle , X. Coubez<sup>23</sup> , D. Cutts , M. Hadley , U. Heintz , J.M. Hogan<sup>85</sup> ,  
T. Kwon , G. Landsberg , K.T. Lau , D. Li , J. Luo , M. Narain , N. Pervan ,  
S. Sagir<sup>86</sup> , F. Simpson , E. Usai , W.Y. Wong , X. Yan , D. Yu , W. Zhang 










**University of California, Davis, Davis, California, USA**

J. Bonilla , C. Brainerd , R. Breedon , M. Calderon De La Barca Sanchez , M. Chertok ,  
J. Conway , P.T. Cox , R. Erbacher , G. Haza , F. Jensen , O. Kukral , G. Mocellin ,  
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















**University of California, Los Angeles, California, USA**

M. Bachtis , R. Cousins , A. Datta , D. Hamilton , J. Hauser , M. Ignatenko ,  
M.A. Iqbal , T. Lam , E. Manca , W.A. Nash , S. Regnard , D. Saltzberg , B. Stone ,  
V. Valuev 























**University of California, Riverside, Riverside, California, USA**

R. Clare , J.W. Gary , M. Gordon , G. Hanson , G. Karapostoli , O.R. Long ,  
N. Manganeli , W. Si , S. Wimpenny 














**University of California, San Diego, La Jolla, California, USA**

J.G. Branson , P. Chang , S. Cittolin , S. Cooperstein , D. Diaz , J. Duarte ,  
R. Gerosa , L. Giannini , J. Guiang , R. Kansal , V. Krutelyov , R. Lee , J. Letts ,  
M. Masciovecchio , F. Mokhtar , M. Pieri , B.V. Sathia Narayanan , V. Sharma ,  
M. Tadel , E. Vourliotis , F. Würthwein , Y. Xiang , A. Yagil 












**University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA**

N. Amin , C. Campagnari , M. Citron , G. Collura , A. Dorsett , V. Dutta ,  
J. Incandela , M. Kilpatrick , J. Kim , A.J. Li , P. Masterson , H. Mei , M. Oshiro ,  
M. Quinnan , J. Richman , U. Sarica , R. Schmitz , F. Setti , J. Sheplock , P. Siddireddy ,  
D. Stuart , S. Wang 

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A. Bornheim , O. Cerri , I. Dutta , A. Latorre , J.M. Lawhorn , J. Mao , H.B. Newman ,  
T. Q. Nguyen , M. Spiropulu , J.R. Vlimant , C. Wang , S. Xie , R.Y. Zhu 

**Carnegie Mellon University, Pittsburgh, Pennsylvania, USA**


















J. Alison , S. An , M.B. Andrews , P. Bryant , T. Ferguson , A. Harilal , C. Liu ,  
T. Mudholkar , S. Murthy , M. Paulini , A. Roberts , A. Sanchez , W. Terrill 

**University of Colorado Boulder, Boulder, Colorado, USA**







































































J.P. Cumalat , W.T. Ford , A. Hassani , G. Karathanasis , E. MacDonald, F. Marini , A. Perloff , C. Savard , N. Schonbeck , K. Stenson , K.A. Ulmer , S.R. Wagner , N. Zipper 



















#### **Cornell University, Ithaca, New York, USA**

J. Alexander , S. Bright-Thonney , X. Chen , D.J. Cranshaw , J. Fan , X. Fan , D. Gadkari , S. Hogan , J. Monroy , J.R. Patterson , D. Quach , J. Reichert , M. Reid , A. Ryd , J. Thom , P. Wittich , R. Zou 











#### **Fermi National Accelerator Laboratory, Batavia, Illinois, USA**

M. Albrow , M. Alyari , G. Apollinari , A. Apresyan , L.A.T. Bauerdick , D. Berry , J. Berryhill , P.C. Bhat , K. Burkett , J.N. Butler , A. Canepa , G.B. Cerati , H.W.K. Cheung , F. Chlebana , K.F. Di Petrillo , J. Dickinson , V.D. Elvira , Y. Feng , J. Freeman , A. Gandrakota , Z. Gecse , L. Gray , D. Green, S. Grünendahl , D. Guerrero , O. Gutsche , R.M. Harris , R. Heller , T.C. Herwig , J. Hirschauer , L. Horyn , B. Jayatilaka , S. Jindariani , M. Johnson , U. Joshi , T. Klijsma , B. Klima , K.H.M. Kwok , S. Lammel , D. Lincoln , R. Lipton , T. Liu , C. Madrid , K. Maeshima , C. Mantilla , D. Mason , P. McBride , P. Merkel , S. Mrenna , S. Nahn , J. Ngadiuba , D. Noonan , V. Papadimitriou , N. Pastika , K. Pedro , C. Pena<sup>87</sup> , F. Ravera , A. Reinsvold Hall<sup>88</sup> , L. Ristori , E. Sexton-Kennedy , N. Smith , A. Soha , L. Spiegel , J. Strait , L. Taylor , S. Tkaczyk , N.V. Tran , L. Uplegger , E.W. Vaandering , I. Zoi 







#### **University of Florida, Gainesville, Florida, USA**

P. Avery , D. Bourilkov , L. Cadamuro , V. Cherepanov , R.D. Field, M. Kim, E. Koenig , J. Konigsberg , A. Korytov , E. Kuznetsova<sup>89</sup> , K.H. Lo, K. Matchev , N. Menendez , G. Mitselmakher , A. Muthirakalayil Madhu , N. Rawal , D. Rosenzweig , S. Rosenzweig , K. Shi , J. Wang , Z. Wu 



















#### **Florida State University, Tallahassee, Florida, USA**

T. Adams , A. Askew , N. Bower , R. Habibullah , V. Hagopian , T. Kolberg , G. Martinez, H. Prosper , O. Viazlo , M. Wulansatiti , R. Yohay , J. Zhang












#### **Florida Institute of Technology, Melbourne, Florida, USA**

M.M. Baarmand , S. Butalla , T. Elkafrawy<sup>15</sup> , M. Hohlmann , R. Kumar Verma , M. Rahmani, F. Yumiceva 












#### **University of Illinois Chicago, Chicago, Illinois, USA**

M.R. Adams , H. Becerril Gonzalez , R. Cavanaugh , S. Dittmer , O. Evdokimov , C.E. Gerber , D.J. Hofman , D. S. Lemos , A.H. Merrit , C. Mills , G. Oh , T. Roy , S. Rudrabhatla , M.B. Tonjes , N. Varelas , X. Wang , Z. Ye , J. Yoo 












#### **The University of Iowa, Iowa City, Iowa, USA**

M. Alhousseini , K. Dilsiz<sup>90</sup> , L. Emediato , G. Karaman , O.K. Köseyan , J.-P. Merlo, A. Mestvirishvili<sup>91</sup> , J. Nachtman , O. Neogi, H. Ogul<sup>92</sup> , Y. Onel , A. Penzo , C. Snyder, E. Tiras<sup>93</sup> 

#### **Johns Hopkins University, Baltimore, Maryland, USA**







O. Amram , B. Blumenfeld , L. Corcodilos , J. Davis , A.V. Gritsan , S. Kyriacou , P. Maksimovic , J. Roskes , S. Sekhar , M. Swartz , T.Á. Vámi 

#### **The University of Kansas, Lawrence, Kansas, USA**

A. Abreu , L.F. Alcerro Alcerro , J. Anguiano , P. Baringer , A. Bean , Z. Flowers , T. Isidori , J. King , G. Krintiras , M. Lazarovits , C. Le Mahieu , C. Lindsey,

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






**Kansas State University, Manhattan, Kansas, USA**

B. Allmond , S. Duric, A. Ivanov , K. Kaadze , A. Kalogeropoulos , D. Kim, Y. Maravin , T. Mitchell, A. Modak, K. Nam, D. Roy 























**Lawrence Livermore National Laboratory, Livermore, California, USA**

F. Rebassoo , D. Wright 















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
















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















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












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




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








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














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