



# Search for singly produced vector-like top partners in multilepton final states with $139 \text{ fb}^{-1}$ of $pp$ collision data at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector

The ATLAS Collaboration

A search for the single production of a vector-like top partner ( $T$ ) decaying into a  $Z$  boson and a top quark is presented, using the full Run 2 dataset corresponding to  $139 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 13 \text{ TeV}$ , collected in 2015–2018 with the ATLAS detector at the Large Hadron Collider. The targeted final state is characterized by the presence of a pair of electrons or muons with opposite-sign charges which form a  $Z$ -boson candidate, as well as by the presence of  $b$ -tagged jets and forward jets. Events with exactly two or at least three leptons are categorized into two independently optimized analysis channels. No significant excess above the background expectation is observed and the results from the two channels are statistically combined to set exclusion limits at 95% confidence level on the masses and couplings of  $T$  for different benchmark models and generalized branching ratio scenarios.

# 1 Introduction

The discovery of the Higgs boson by the ATLAS and CMS experiments at the Large Hadron Collider (LHC) [1, 2] was a major milestone for the Standard Model (SM), the framework for understanding the fundamental constituents of our universe and their interactions. While finding the Higgs boson completed the SM, this model leaves many questions unaddressed about the physical nature of our universe. One of these questions is about the measured Higgs boson mass—specifically, the many orders of magnitude that separate it from the Planck scale, requiring quadratically divergent finely tuned corrections [3]. These large corrections are mitigated in some beyond Standard Model (BSM) theories. Vector-like Quarks (VLQs), colored spin- $\frac{1}{2}$  fermions with identical electroweak (EW) representation for both left and right chiralities, are predicted by a large subset of these theories [4–13]. The usual four VLQ species are denoted by  $X_{+\frac{5}{3}}, T_{+\frac{2}{3}}, B_{-\frac{1}{3}}$  and  $Y_{-\frac{4}{3}}$ , where the subscript indicates the electric charge of the corresponding particle. Their renormalizable EW representation consists of  $(T)$  or  $(B)$  singlets,  $(X, T)$ ,  $(T, B)$ , or  $(B, Y)$  doublets and  $(X, T, B)$  or  $(T, B, Y)$  triplets. In most representations, they couple to the SM quarks via an exchange of charged ( $W^\pm$ ) or neutral ( $Z, H$ ) bosons. The interaction of the VLQs with the SM quarks can be summarized with the following simplified Lagrangian [14]:

$$\mathcal{L} = \sum_{\zeta, q, Q} \left[ \frac{g_w}{2c_w} \tilde{\kappa}_\zeta^{Qq} \bar{Q} \not{Z} P_\zeta q + \frac{g_w}{\sqrt{2}} \kappa_\zeta^{Qq} \bar{Q} \not{W} P_\zeta q - \hat{\kappa}_\zeta^{Qq} H \bar{Q} P_\zeta q \right] + \text{h.c.}, \quad (1)$$

where  $Q$  represents a VLQ,  $\zeta$  represents the chirality with  $P_\zeta$  being the corresponding projection operator,  $q$  represents a SM quark of up or down type, and the electroweak couplings  $\tilde{\kappa}_\zeta^{Qq}$ ,  $\kappa_\zeta^{Qq}$ , and  $\hat{\kappa}_\zeta^{Qq}$  determine the coupling strengths between  $Q$  and  $q$  when mediated by the  $Z$ ,  $W$ , and  $H$  bosons respectively. The mass hierarchy of the SM quarks suggests that VLQs interact predominantly with the third-generation SM quarks [15, 16]. Hence, VLQ interactions with lighter generations are set to zero in the simplified representation of Eq. (1) and the VLQ species  $T$  and  $B$  are referred to as top and bottom partners respectively in the following discussion.

At the LHC, the dominant modes of VLQ production are pair and single production. During Run 1 (2010–2012) of the LHC, the VLQ searches at center-of-mass energies of  $\sqrt{s} = 7$  and 8 TeV typically probed VLQ masses lower than 1 TeV where VLQ pair production dominates [16]. Since the pair production of VLQs proceeds through QCD interactions, the cross-section under the narrow width approximation only depends on the VLQ mass and not on the species nor its EW representation. No significant excess was observed by the Run 1 searches looking for pair production of VLQs and lower exclusion limits on VLQ mass were reported in the range of approximately 600–1000 GeV, depending on the search [17–22]. The searches for single production of  $T$  or  $B$  also reported similar limits in Run 1 for a range of couplings probed in the singlet and doublet benchmarks [19, 23].

Run 2 (2015–2018) searches at  $\sqrt{s} = 13$  TeV for pair production of top and bottom partners from ATLAS and CMS have set lower limits on the VLQ masses in the range of 800–1400 GeV [24–34]. Since the single production of VLQs can have a larger cross-section at high masses for reasonably large couplings [16], the Run 2 searches are increasingly focusing on this production mode [35–45]. However, unlike pair production, single production of VLQs is mediated by electroweak processes. Both the kinematics and cross-sections of the single VLQ processes depend on the representation and the choice of couplings that determine the relative strength of these particles interacting with SM quarks and vector/Higgs bosons.

This paper focuses on the search for a singly produced  $T$  decaying to a  $Z$  boson and a top quark using the complete Run 2 dataset collected by the ATLAS detector. As shown in Figure 1, the electroweak interaction can produce a  $T$  in the  $s$ - and  $t$ -channel topologies. These processes are referred to as  $WTZt$  and  $ZTZt$  processes, where the first  $W$  or  $Z$  refers to the vector boson scattered off of the incoming parton,  $T$  refers to the VLQ exchanged in the intermediate stage, and  $Zt$  refers to the production of the  $Z$  boson and  $t$  quark. The  $WTZt$  mode dominates in the singlet representation of  $T$  since the smaller coupling between the  $T$  and the  $Z$  boson and the kinematically disfavored  $g \rightarrow t\bar{t}$  splitting suppress the contribution from  $ZTZt$ . On the other hand, the latter is the dominant mode in the doublet representation since the  $T$ 's coupling with the  $W$  boson is suppressed by the mass diagonalization matrix [16]. A model-independent approach [46–48] that allows probing a wide range of couplings for the  $W$ ,  $Z$ , and  $H$ -mediated interactions of these heavy quarks is adopted in this search.

The analysis is performed across two channels. The dilepton ( $2\ell$ ) channel selects events with exactly one pair of opposite-sign leptons, i.e. electrons or muons and a hadronically decaying top quark in the final state. The trilepton ( $3\ell$ ) channel requires, in addition to the pair of opposite-sign leptons, at least a third lepton from a leptonically decaying top quark. These results improve on a previous ATLAS search [35] with a data sample corresponding to an integrated luminosity of  $36.1 \text{ fb}^{-1}$ , by benefiting from the larger data sample of  $139 \text{ fb}^{-1}$ , improved kinematic selections and introducing a method to identify hadronically decaying top quarks.

The paper is organized as follows. Section 2 provides a brief overview of the ATLAS detector. The data and simulated background and signal samples are described in Section 3. Section 4 describes the object definition while event selection, categorization, and the analysis strategy are outlined in Section 5. The systematic uncertainties are discussed in Section 6. Results obtained from this search are reported in Section 7 and conclusions are presented in Section 8.

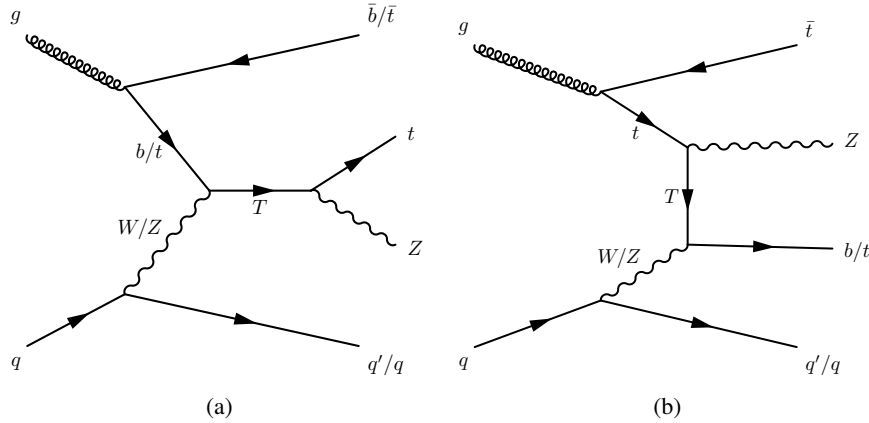


Figure 1: Illustrative (a)  $s$ -channel and (b)  $t$ -channel Feynman diagrams for  $WTZt$  and  $ZTZt$  processes.

## 2 ATLAS detector

The ATLAS detector [49] at the LHC covers nearly the entire solid angle around the collision point.<sup>1</sup> It consists of an inner tracking detector surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters, and a muon spectrometer incorporating three large superconducting air-core toroidal magnets.

The inner-detector system (ID) is immersed in a 2 T axial magnetic field and provides charged-particle tracking in the range of  $|\eta| < 2.5$ . The high-granularity silicon pixel detector covers the vertex region and typically provides four measurements per track, the first hit normally being in the insertable B-layer (IBL) installed before Run 2 [50]. It is followed by the silicon microstrip tracker (SCT), which usually provides eight measurements per track. These silicon detectors are complemented by the transition radiation tracker (TRT), which enables radially extended track reconstruction up to  $|\eta| = 2.0$ . The TRT also provides electron identification information based on the fraction of hits (typically 30 in total) above a higher energy-deposit threshold corresponding to transition radiation.

The calorimeter system covers the pseudorapidity range of  $|\eta| < 4.9$ . Within the region  $|\eta| = 3.2$ , electromagnetic calorimetry is provided by barrel and endcap high-granularity lead/liquid-argon (LAr) calorimeters, with an additional thin LAr presampler covering  $|\eta| < 1.8$  to correct for energy loss in material upstream of the calorimeters. Hadron calorimetry is provided by the steel/scintillator-tile calorimeter, segmented into three barrel structures within  $|\eta| = 1.7$ , and two copper/LAr hadron endcap calorimeters. The solid angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeter modules optimised for electromagnetic and hadronic energy measurements respectively.

The muon spectrometer (MS) comprises separate trigger and high-precision tracking chambers measuring the deflection of muons in a magnetic field generated by the superconducting air-core toroidal magnets. The field integral of the toroids ranges between 2.0 and 6.0 T m across most of the detector. Three layers of precision chambers, each consisting of layers of monitored drift tubes, cover the region  $|\eta| < 2.7$ , complemented by cathode-strip chambers in the forward region, where the background is highest. The muon trigger system covers the range of  $|\eta| < 2.4$  with resistive-plate chambers in the barrel, and thin-gap chambers in the endcap regions.

Interesting events are selected by the first-level trigger system implemented in custom hardware, followed by selections made by algorithms implemented in the software-based high-level trigger (HLT) [51]. The first-level trigger reduces the incoming data rate from the 40 MHz bunch crossings to a design value of 100 kHz, which is further reduced by the HLT to record events to disk at a rate of about 1 kHz.

An extensive software suite [52] is used in data simulation, in the reconstruction and analysis of real and simulated data, in detector operations, and in the trigger and data acquisition systems of the experiment.

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<sup>1</sup> ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the  $z$ -axis along the beam pipe. The  $x$ -axis points from the IP to the center of the LHC ring, while the  $y$ -axis points upwards. Cylindrical coordinates  $(r, \phi)$  are used in the transverse plane,  $\phi$  being the azimuthal angle around the  $z$ -axis. The pseudorapidity is defined in terms of the polar angle  $\theta$  as  $\eta = -\ln \tan(\theta/2)$ . Angular distance  $(\Delta R)$  is defined as  $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ .

### 3 Data and simulated event samples

The data sample was collected by the ATLAS detector in proton–proton ( $pp$ ) collisions at a center-of-mass energy of  $\sqrt{s} = 13$  TeV during Run 2 with all detector subsystems operational and with the LHC operating in stable beam conditions with 25 ns bunch spacing. The dataset corresponds to an integrated luminosity of  $139 \text{ fb}^{-1}$  with an average of about 34 simultaneous interactions per bunch crossing (pileup).

All of the nominal Monte Carlo (MC) simulation samples used in the analysis were processed with the ATLAS simulation framework [53], using a detailed simulation based on GEANT4 [54]. The effect of pileup was modeled by overlaying the simulated hard-scattering event with inelastic  $pp$  events generated with PYTHIA 8.186 [55] using the NNPDF2.3LO set of parton distribution functions (PDF) [56] and the A3 set of tuned parameters (tune) [57].

The nominal MC sample for  $Z$  boson production with jets ( $Z$ +jets) was generated with SHERPA 2.2.1 [58–61] and the nominal diboson ( $VV$ ) sample was generated with SHERPA 2.2.2. Both the samples were generated using the NNPDF3.0 [62] next-to-next-to-leading-order (NNLO) PDF set. The  $Z$ +jets sample includes events generated with up to two partons at next-to-leading-order (NLO) and up to four partons at leading-order (LO) and is normalized to the NNLO cross-section [63]. The  $VV$  sample is normalized to the SHERPA NLO cross-section and includes  $q\bar{q}$ -initiated events with up to one parton at NLO and up to three partons at LO and  $gg$ -initiated processes simulated using LO matrix elements for up to one additional jet. For both the samples, COMIX [60] and OPENLOOPS [64–66] were used and the matrix element (ME) was merged with the SHERPA parton shower [61] according to the MEPS@NLO prescription [59, 67–69]. To estimate modeling uncertainties for these backgrounds, additional samples were produced with MADGRAPH5\_AMC@NLO 2.2.3 [70], using the NNPDF3.0NNLO PDF set and interfaced to PYTHIA 8.210 [71] with the A14 tune [72] and the NNPDF2.3LO PDF for showering. In the  $2\ell$  channel, the simulated  $Z$ +jets events are categorized into  $Z$ +Light Flavor (LF) and  $Z$ +Heavy Flavor (HF) events in accordance with the absence or presence of heavy flavor jets in the MC generator record of these simulated events in the kinematic acceptance defined in Sections 4 and 5. An additional set of  $Z$ +jets events was generated with SHERPA 2.2.11 to compare the modeling of the  $Z$ +jets background in the  $2\ell$  channel.

The nominal  $t\bar{t}$  background sample uses the POWHEG method [73, 74] implemented in POWHEG BOX v2 [75, 76] with the NNPDF3.0NNLO PDF set. POWHEG BOX was interfaced with PYTHIA 8.230 with the A14 tune for showering and hadronization. The sample is normalized to the NNLO cross-section in QCD including resummation of next-to-next-to-leading-logarithmic (NNLL) soft gluon terms calculated with TOP++ [77–83]. For the evaluation of modeling uncertainties, samples were produced with the same matrix element generator as the nominal sample, but HERWIG 7 was used with the H7-UE-MMHT tune [84] for showering and hadronization. To assess the uncertainty in the matching of NLO matrix elements to the parton shower, additional samples were generated with MADGRAPH5\_AMC@NLO 2.6.0 and the NNPDF3.0NNLO PDF set, interfaced with PYTHIA 8.230 using the same showering configuration as the nominal sample.

The nominal sample for  $t\bar{t}$  production with a vector boson ( $t\bar{t} + W$  and  $t\bar{t} + Z$ ) was generated with MADGRAPH5\_AMC@NLO 2.3.3 interfaced with PYTHIA 8.210 for showering and hadronization, using the NNPDF2.3LO PDF set and the A14 tune. These events were normalized to the NLO cross-sections calculated with MADGRAPH5\_AMC@NLO [85, 86]. To evaluate modeling uncertainties, alternative samples were produced where either the A14 tune was varied or HERWIG 7 was used with the H7-UE-MMHT tune for the showering. Additional samples were produced using SHERPA 2.2.1 to evaluate the uncertainty due to the choice of matrix element generator. The production of  $t\bar{t}t\bar{t}$  and  $t\bar{t}WW$  events was modeled using

the MADGRAPH5\_AMC@NLO generator at LO with the NNPDF3.1<sub>NLO</sub> PDF, interfaced with PYTHIA 8.230 using the A14 tune and the NNPDF2.3<sub>LO</sub> PDF set. These samples were normalized to cross-sections calculated with NLO QCD and EW corrections [85]. These four processes are merged into a single process labeled  $t\bar{t} + X$  in which the contribution from  $t\bar{t} + Z$  dominates due to the requirement of an opposite-sign same-flavor (OS-SF) lepton pair with an invariant mass close to that of a  $Z$  boson in the final state.

The single-top-quark processes were simulated with POWHEG BOX [87, 88] using the NNPDF3.0<sub>NLO</sub> PDF set and interfaced to PYTHIA 8.234 with the A14 tune. The samples are normalized to their respective NLO QCD cross-sections [89, 90] for the  $t$ -channel and  $s$ -channel, and with additional NNLL soft gluon terms for  $Wt$  production [91–93]. The production of  $tZq$  and  $tWZ$  events was modeled using the MADGRAPH5\_AMC@NLO generator at NLO with the NNPDF3.0<sub>NLO</sub> PDF and interfaced with PYTHIA 8.212 using the A14 tune and the NNPDF2.3<sub>LO</sub> PDF set. The diagram-removal scheme [94] was used in the generation of  $Wt$  and  $tWZ$  events to address their overlaps with  $t\bar{t}$  and  $t\bar{t} + Z$  samples. These samples were also merged into a single process labeled as single-top.

Signal samples for the  $WTZt$  and  $ZTZt$  processes were generated at LO using MADGRAPH5\_AMC@NLO with the Universal FeynRules Output (UFO) model [95] from Ref. [14], interfaced with PYTHIA 8.244 and using the NNPDF3.0<sub>LO</sub> PDF set and the A14 tune. This model uses the four-flavor scheme. Top partner samples were obtained for various masses, ranging from 1.0 TeV and 2.7 TeV in steps of 100 GeV. For each mass point, the coupling strength  $\kappa$  was varied between 0.1 and 1.0 with a step of 0.05 (0.1) for couplings less (greater) than 0.5. The relative strength of the three coupling modes were obtained from  $\vec{\xi} = (\xi_W, \xi_Z, \xi_H) = (0.5, 0.25, 0.25)$ .<sup>2</sup> The MC samples were produced for a subset of the mass and coupling points. Those samples were used for other masses and couplings by implementing an event-by-event matrix element reweighting [96]. While the left- and right-handed couplings dominate in the singlet and doublet representations of  $T$ , the current analysis was found to be insensitive to the choice of chirality for the signal MC generation and all signal samples were generated using the left-handed couplings. The signal samples were normalized to the NLO cross-section calculated with narrow-width approximation [97]. However, since the top partner can have large width based on the choice of coupling, additional correction factors were applied to evaluate the cross-section at finite width [98]. The contribution from the non-resonant  $t$ -channel diagram was also taken into account [99].

## 4 Object reconstruction

Events are required to have at least one vertex with at least two tracks with transverse momentum  $p_T > 0.5$  GeV. The primary vertex (PV) is defined as the one with the largest  $\Sigma p_T^2$ , where the sum is performed over all the tracks matched to the vertex.

Electrons are reconstructed from energy clusters in the EM calorimeter matched with ID tracks and must fulfill the *tight likelihood* identification criteria [100]. Electrons are calibrated [100] and required to have  $p_T > 28$  GeV and to be reconstructed within  $|\eta| = 2.47$ , excluding the calorimeter barrel–endcap transition regions ( $1.37 < |\eta| < 1.52$ ). A high acceptance for the expected signal events is maintained by not applying an isolation requirement to electron candidates beyond those implicit in the trigger requirements, which are

<sup>2</sup> The parameterization of VLQ Lagrangian in terms of the  $\kappa$  and  $\vec{\xi}$  parameters was introduced in Ref. [47] and is used for the interpretation of the search presented in this paper. Its conversion to the coupling convention in Eq. (1) is obtained by doing a one-to-one mapping of the tree level couplings in the Lagrangian. The  $\xi_W, \xi_Z, \xi_H$  parameters satisfy the constraint  $\xi_W + \xi_Z + \xi_H = 1$  and asymptotically represent the branching ratios of  $T$  decaying into  $Wb$ ,  $Zt$ , and  $Ht$  in the narrow width limit.

explained in Section 5. Furthermore, the track matched with the candidate electron is required to have a longitudinal impact parameter with respect to the PV which satisfies  $|z_0 \cdot \sin \theta| < 0.5$  mm and a transverse impact parameter with respect to the beamline ( $d_0$ ) with a significance  $|d_0|/\sigma(d_0) < 5$ .

Muons are reconstructed [101] from combined tracks in the MS and the ID and must fulfill the *medium* identification criteria [101]. Muons are calibrated and required to have  $p_T > 28$  GeV and to be reconstructed within  $|\eta| = 2.5$ . Muon candidates must also satisfy the track-based isolation requirements defined by the *FixedCutTightTrackOnly* working point [101]. This working point uses the scalar sum of the  $p_T$  of all tracks that are inside of a cone of size  $\Delta R = \min \{0.3, 10 \text{ GeV}/p_T(\mu)\}$  around the muon candidate, where  $p_T(\mu)$  is the candidate muon  $p_T$ . The track matched with the muon candidate under consideration is excluded from the sum. The muon is selected if this sum is less than 15% of  $p_T(\mu)$ . Finally, muon candidates are required to have a matched track with  $|z_0 \cdot \sin \theta| < 0.5$  mm and a significance  $|d_0|/\sigma(d_0) < 3$ .

Jets are reconstructed from particle-flow objects [102] using the anti- $k_r$  algorithm [103, 104] with a radius parameter of  $R = 0.4$ . Jets are calibrated to the particle level from a combination of simulation-based corrections and measurements in data [105] and are required to fulfill  $p_T > 25$  GeV for  $|\eta| < 2.5$  and  $p_T > 35$  GeV for  $2.5 < |\eta| < 4.5$ . These two jet categories are respectively called central and forward jets. To reduce jet contributions from pileup, a ‘jet vertex tagger’ algorithm using a two-dimensional likelihood discriminant [106] is applied to jets with  $|\eta| < 2.4$  and  $p_T < 60$  GeV. The DL1r algorithm [107] is used to identify jets in the central region ( $|\eta| < 2.5$ ) containing  $b$ -hadrons ( $b$ -tagging) with a working point corresponding to a  $b$ -tagging efficiency in simulated  $t\bar{t}$  events of 77%, a  $c$ -jet rejection factor of  $\sim 6$ , and a light-jet rejection factor of  $\sim 192$ .

The missing transverse momentum [108], with magnitude  $E_T^{\text{miss}}$ , is defined as the negative vectorial sum of the transverse momenta of all the calibrated reconstructed lepton and jet candidates in the event and includes a ‘soft term’ with contributions from tracks emanating from the PV but not matched with any of the reconstructed objects.

A procedure to remove potential overlaps between reconstructed lepton and jet candidates is performed sequentially as follows. First, any electron sharing an ID track with one of the muons is removed. Next, any jet within  $\Delta R = 0.2$  of an electron is removed, followed by the removal of electrons within  $\Delta R = 0.4$  of any remaining jet. Then any jet with at most two tracks with  $p_T > 0.5$  GeV within  $\Delta R = 0.2$  of a muon is removed, unless it is  $b$ -tagged. At the end of the procedure, any muon within  $\Delta R = \min \{0.4, 0.04 + 10 \text{ GeV}/p_T(\mu)\}$  of any remaining jet is removed.

To identify hadronically decaying boosted top-quark jets, variable-radius reclustered jets (vRC jets) are reconstructed using the variable-radius jet algorithm [109] from calibrated jets with  $R = 0.4$  where the effective radius of the jet cone is chosen as  $R_0 = \frac{2m_t}{p_T}$ ,  $m_t$  and  $p_T$  representing the mass of the top quark and the jet transverse momentum respectively. These vRC jets are known to provide stable top-tagging performance for a wide range of jet  $p_T$  [110]. The vRC jets are called top-tagged (top-vetoed) if the vRC jet mass is greater (less) than 140 GeV. Top-tagged and top-vetoed jets are required to have  $p_T > 200$  GeV and to contain at least two  $R = 0.4$  jets as jet constituents when their transverse momentum is less than 700 GeV.

## 5 Event selection

A common initial event selection is performed for both the  $2\ell$  and  $3\ell$  channels where events are required to have passed the single lepton trigger selections [51, 111, 112]. For muons, triggers with a  $p_T$  threshold

of 20 GeV (26 GeV) in 2015 (2016–2018) and isolation requirements, are combined in logical OR with triggers with a 50 GeV  $p_T$  threshold with no isolation requirement. A trigger with a 60 GeV  $p_T$  threshold is added for the 2017–2018 data taking period. Similarly, electron triggers with isolation and identification requirements (tight likelihood identification in 2016–18, with a less restrictive requirement in 2015) and  $p_T$  thresholds of 24–26 GeV are combined with triggers with higher  $p_T$  thresholds between 120 GeV and 140 GeV, less restrictive identification criteria, and no isolation requirements.

Events are additionally required to have at least one OS-SF lepton pair with the transverse momentum of each lepton being greater than 28 GeV. The pair of OS-SF leptons with invariant mass  $m(\ell\ell)$  closest to the mass of the  $Z$ -boson ( $m_Z$ ) is identified as the  $Z$ -boson candidate. Based on the expected hadronic activity of the signal topologies, events are also required to have at least one vRC jet (in the  $2\ell$  channel) or at least two central jets (in the  $3\ell$  channel).

## 5.1 Dilepton channel

The preselection of events for the  $2\ell$  channel ( $2\ell$ PS) requires the  $Z$ -boson candidate to have an invariant mass satisfying  $|m(\ell\ell) - m_Z| < 10$  GeV. Given that the search is sensitive to energetic final-state objects, additional background reduction is achieved by requiring the  $Z$ -boson candidate to have a transverse momentum ( $p_T(\ell\ell)$ ) of at least 200 GeV and the scalar sum of transverse momenta of jets in the event, denoted by  $H_T$ , to be at least 300 GeV. Since signal events are expected to contain a hadronically decaying boosted top quark, events without any vRC jets are rejected. The four momentum of the leading vRC jet ( $J$ ), defined as the vRC jet with the highest  $p_T$ , is combined with that of the  $Z$ -boson candidate to obtain an estimate of the mass of  $T$ ,  $m_{\ell\ell J}$ . To reduce sensitivity to VLQ pair-production and to allow for a consistent phenomenological interpretation of observed data counts, a requirement of  $H_T + E_T^{\text{miss}} < m_{\ell\ell J}$  is also placed on preselection events [35].

The main background for the  $2\ell$  channel is comprised of  $Z$ +jets events, with minor contributions from  $VV$  and  $t\bar{t}$  processes. To improve the background modeling, a reweighting factor is extracted for the  $Z$ +jets samples by comparing data with simulated events in the  $2\ell$ PS region for the jet multiplicity and  $H_T + E_T^{\text{miss}}$  distributions. These reweighting factors are separately derived for  $2\ell$ PS events with and without  $b$ -tagged jets. Since the shapes of the  $Z$ +LF and  $Z$ +HF samples were found to be similar in these regions, these factors are applied on the  $Z$ +jets template obtained by combining these samples together. The reweighting function is defined bin-by-bin by the formula

$$R_{Z+\text{jets}}(x) = \frac{\text{Data}(x) - \text{MC}^{\text{non-}Z+\text{jets}}(x)}{\text{MC}^{Z+\text{jets}}(x)}, \quad (2)$$

where  $x$  represents the binning in each kinematic variable. After reweighting in jet multiplicity, a second reweighting is applied on  $H_T + E_T^{\text{miss}}$  using Eq. (2). The joint reweighting factors are applied on an event-by-event basis to estimate the MC contribution of  $Z$ +jets events for this channel. The jet multiplicity and  $H_T + E_T^{\text{miss}}$  distributions in the  $2\ell$ PS region are shown in Figure 2. In the bins of jet multiplicity and  $H_T + E_T^{\text{miss}}$  distributions, the contribution of these minor processes is less than 15%. The predicted signal contribution is also quite small. For a benchmark  $WTZt$  process with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in the singlet representation, the bin-by-bin signal-over-background ratio ( $S/B$ ) varies between 4–7% for most bins and is always less than 9%.

Figure 3 shows the distributions of the key kinematic variables for the expected background and benchmark signal processes in the  $2\ell$ PS region. Events with a singly produced  $T$  have a distinctive signature with



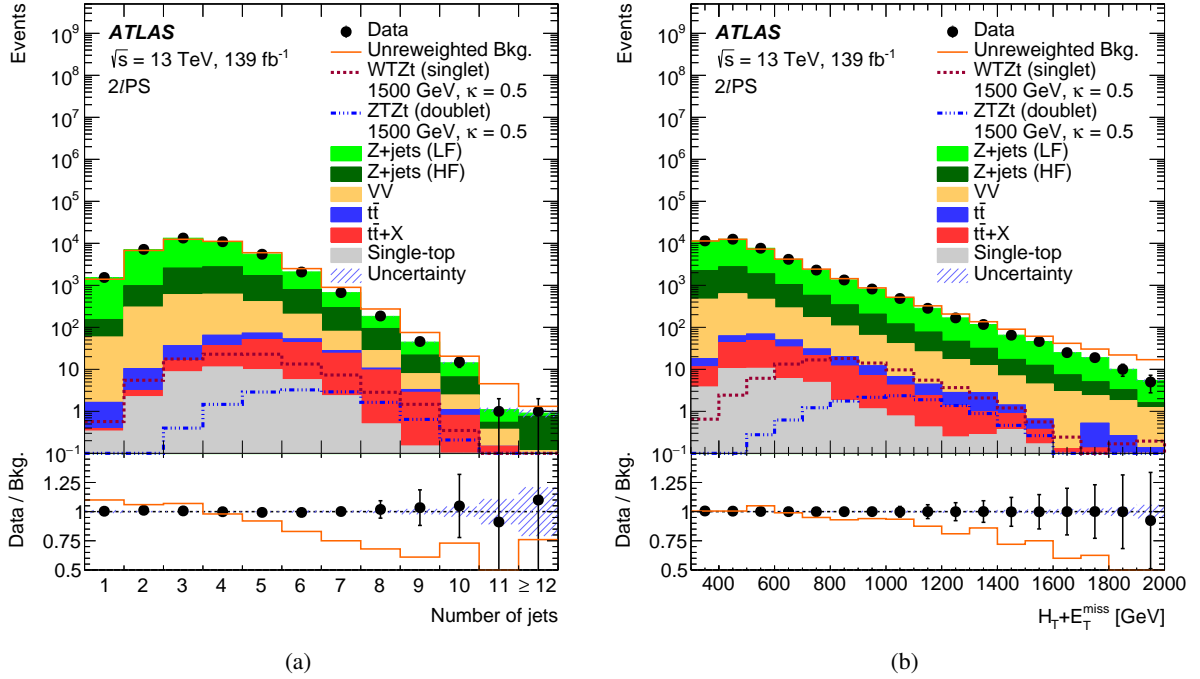


Figure 2: Distributions of (a) jet multiplicity and (b)  $H_T + E_T^{\text{miss}}$  in the  $2\ell\text{PS}$  region after reweighting. The overlaid solid line shows the total background prediction before the reweighting and the corresponding data-over-background ratio in the bottom panel. The red and blue dotted lines show the signal contributions from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in the singlet and doublet representations respectively. The last bin in each distribution contains the overflow.

a forward jet scattering off of a vector boson from one of the incoming partons. Events are additionally expected to have  $b$ -tagged jets. Hence, in addition to the preselection, the signal-enriched kinematic phase space, called the *signal region* ( $2\ell\text{SR}$ ), requires events with at least one forward jet, at least one  $b$ -tagged jet, and at least one top-tagged jet. The distribution of the transverse momentum of the  $Z$ -boson candidate,  $p_T(\ell\ell)$ , was the variable that offered maximum sensitivity and, hence, is used as the final discriminant in this channel to perform the statistical fit described in Section 7.

Three control regions (CRs) are defined to improve the background modeling and reduce its uncertainties. Two validation regions (VRs) are defined to validate the background prediction in the distribution of the final discriminant. The event selection criteria for these regions are summarized in Table 1. Each of the control regions is obtained by inverting exactly two of the three jet cuts of the  $2\ell\text{SR}$  shown in Table 1. The regions  $2\ell\text{CR1}$ ,  $2\ell\text{CR2}$ , and  $2\ell\text{CR3}$  are defined to have the same selection as  $2\ell\text{SR}$  in multiplicities of forward jets,  $b$ -tagged jets, and top-tagged jets respectively. The inversion of the requirement of the presence of top-tagged jets requires the presence of at least one top-vetoed tagged jet. Both the validation regions are required to have at least one top-tagged jet to maintain their kinematic similarity with the signal region. Identified as  $2\ell\text{VR1}$  and  $2\ell\text{VR2}$ , they are made orthogonal to the  $2\ell\text{SR}$  by requiring zero  $b$ -tagged jets and zero forward jets respectively.

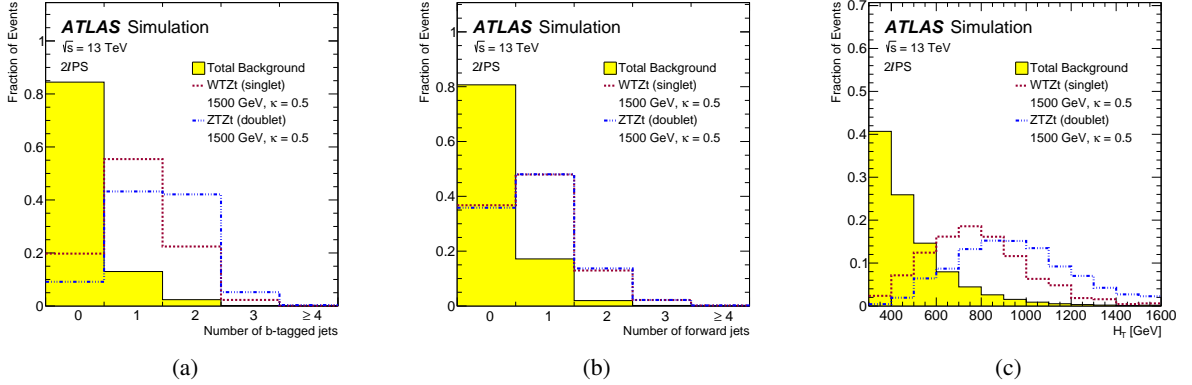


Figure 3: Simulated distributions of (a)  $b$ -tagged jet multiplicity, (b) forward jet multiplicity, and (c)  $H_T$  for total background (solid area) and benchmark signal processes (dashed lines) in the  $2\ell$ PS region before reweighting is performed in jet multiplicity and  $H_T + E_T^{\text{miss}}$  distributions. Each distribution is separately normalized to unit area. The signal contributions are from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in the singlet and doublet representations respectively. The last bin in each distribution contains the overflow.

Table 1: Summary of selections applied to define the control, validation, and signal regions for the  $2\ell$  channel.

	$2\ell$ CR1	$2\ell$ CR2	$2\ell$ CR3	$2\ell$ VVR1	$2\ell$ VVR2	$2\ell$ SR
Preselection	$= 2$ OS-SF leptons with $ m(\ell\ell) - m_Z  < 10$ GeV $p_T(\ell\ell) > 200$ GeV, $H_T > 300$ GeV $\geq 1$ vRC jet $H_T + E_T^{\text{miss}} < m_{\ell\ell J}$					
forward jets	$\geq 1$	0	0	$\geq 1$	0	$\geq 1$
$b$ -tagged jets	0	$\geq 1$	0	0	$\geq 1$	$\geq 1$
top-tagged jets	–	–	$\geq 1$	$\geq 1$	$\geq 1$	$\geq 1$
top-vetoed jets	$\geq 1$	$\geq 1$	–	–	–	–

## 5.2 Trilepton channel

As the name of the channel suggests, events with at least three leptons passing the preselection are selected. Although the selection is inclusive of events with four or more leptons, their overall contribution is quite small ( $< 5\%$ ). The lepton with highest  $p_T$  that does not comprise the  $Z$ -boson candidate is labeled as the third lepton ( $\ell_3$ ) in this channel. The main sources of background are diboson processes and  $t\bar{t}+X$ , the latter being primarily composed of  $t\bar{t} + Z$  and other small contributions from  $t\bar{t} + W$  and  $t\bar{t}t\bar{t}$  processes. Minor contributions result from background processes like single-top and  $Z$ +jets. To improve the quality of MC background modeling, a data-driven reweighting factor is determined for the  $VV$  and  $t\bar{t} + X$  samples in this channel. This reweighting is performed for the central jet multiplicity distribution for events passing the  $3\ell$ PS criteria. These events are classified into separate regions according to the presence or absence of  $b$ -tagged jets in the event. A pair of simultaneous linear equations can be formulated for these two sets of data and MC distributions:

$$\begin{aligned}
 n_{0,i}^{\text{data}} &= \alpha_i^{VV} n_{0,i}^{VV} + \alpha_i^{t\bar{t}X} n_{0,i}^{t\bar{t}X} + n_{0,i}^{\text{others}} \\
 n_{1,i}^{\text{data}} &= \alpha_i^{VV} n_{1,i}^{VV} + \alpha_i^{t\bar{t}X} n_{1,i}^{t\bar{t}X} + n_{1,i}^{\text{others}}
 \end{aligned} \tag{3}$$

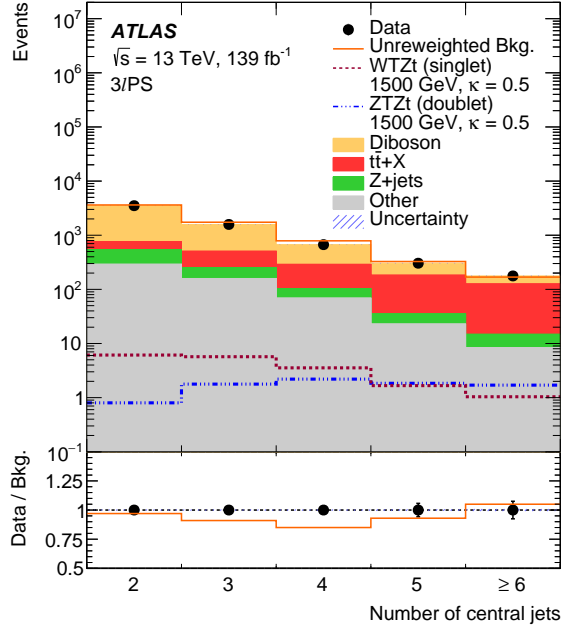


Figure 4: Distribution of central jet multiplicity in the  $3\ell$ PS region after reweighting. The overlaid solid line shows the total background prediction before the reweighting and the corresponding data-over-background ratio in the bottom panel. The red and blue dotted lines show the signal contributions from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in singlet and doublet representations respectively.

where  $n^{\text{data}}$  denotes the data yield and  $n^{VV}$ ,  $n^{t\bar{t}X}$ , and  $n^{\text{others}}$  denote the predicted yields from  $VV$ ,  $t\bar{t} + X$ , and other backgrounds respectively. The subscripts 0/1 refer to events with  $0/\geq 1$   $b$ -tagged jets and the index  $i$  runs over the bins of the central jet multiplicity distribution. The  $(\alpha_i^{VV}, \alpha_i^{t\bar{t}X})$  factors for each bin are obtained by simultaneously solving the pair of equations in Eq. (3) and applied to estimate the MC contributions of  $VV$  and  $t\bar{t} + X$  processes. The central jet multiplicity distribution in the  $3\ell$ PS is shown in Figure 4. In the regions where the reweighting factors are derived, the overall contribution of the minor backgrounds is up to 29%, reaching a bin-by-bin maximum of 37%. The predicted signal contributions in these regions are small. For a benchmark  $WTZt$  process with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in the singlet representation, the bin-by-bin  $S/B$  is less than 2%.

Figure 5 shows distributions of some of the kinematic variables for the expected background and benchmark signal processes in the  $3\ell$ PS region. Exploiting the distinctive signature of signal events, which are expected to have  $b$ -tagged and forward jets, the signal region in this channel ( $3\ell$ SR) requires events to have at least one of each of these jet types. To further increase the signal-to-background ratio, additional requirements are applied on the  $Z$ -boson candidate  $p_T$ ,  $p_T(\ell\ell) > 300$  GeV and the leading lepton  $p_T$ ,  $\max(p_T(\ell)) > 200$  GeV. The decay products of  $T$ , the  $Z$  boson and the top-quark, are expected to have large angular separation between them. Requiring the azimuthal separation of the  $Z$  candidate from the third lepton ( $\Delta\phi(Z, \ell_3)$ ) to be at least  $\frac{\pi}{2}$  significantly improves the signal purity in the signal region (Figure 5(c)). For similar reasons, the azimuthal separation between the  $Z$  candidate and the leading  $b$ -tagged jet, defined as the  $b$ -tagged jet with the largest  $p_T$  which is expected to emerge from the decay of the top quark, is required to satisfy  $\Delta\phi(Z, b_{\text{lead}}) > \frac{\pi}{2}$ .

Finally, to allow for a consistent phenomenological interpretation, the product of  $H_T$  and jet multiplicity,

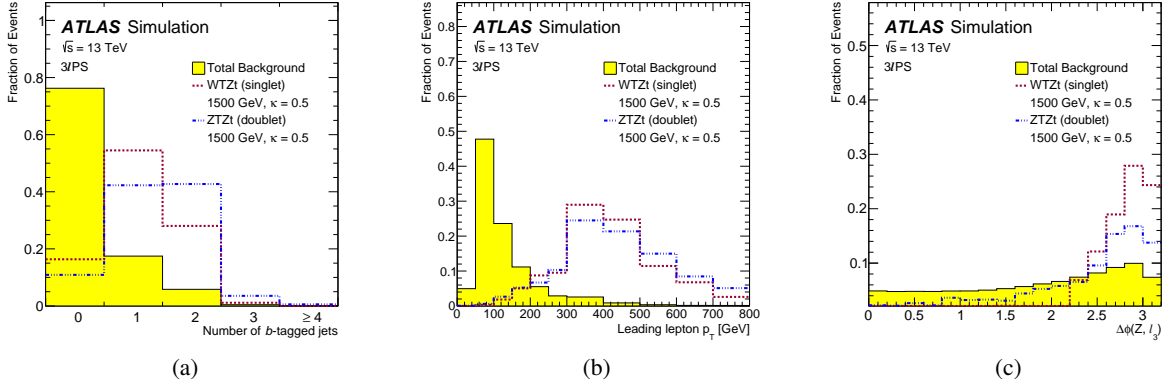


Figure 5: Simulated distributions of (a)  $b$ -tagged jet multiplicity, (b)  $p_T$  of the leading lepton, and (c)  $\Delta\phi$  between the  $Z$  candidate and the third lepton for total background (solid area) and benchmark signal processes (dashed lines) in the  $3\ell$ PS region before reweighting is performed in the central jet multiplicity distribution. Each distribution is separately normalized to unit area. The signal contributions are from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in the singlet and doublet representations respectively. The last bin in each distribution contains the overflow.

$H_T \cdot n(\text{jets})$ , is required to be less than 6 TeV in the signal region to reduce potential contamination from pair production of VLQs [35].

The statistical analysis in this channel also uses  $p_T(\ell\ell)$  as the final discriminant. Three control regions and one validation region are defined in this channel. The diboson control region ( $3\ell VV$ ) is defined by rejecting events with any  $b$ -tagged jet. The mixed control region ( $3\ell$ Mixed) and the  $t\bar{t} + X$  control region ( $3\ell t\bar{t}X$ ), are defined by accepting events with exactly one  $b$ -tagged jet and at least two  $b$ -tagged jets respectively. To maintain orthogonality with the signal region, both regions reject events with forward jets and require  $\Delta\phi(Z, \ell_3) < 2.6$ . Finally, a validation region ( $3\ell VVR$ ) is defined to validate the background modeling by requiring the same set of selections on the  $b$ -tagged jet and forward jet multiplicities as the signal region, but made orthogonal to it by requiring that at least one of the  $\Delta\phi$  requirements is reversed, i.e.,  $\Delta\phi(Z, \ell_3) < \frac{\pi}{2}$  or  $\Delta\phi(Z, b_{\text{lead}}) < \frac{\pi}{2}$ . The definitions of these regions are summarized in Table 2.

Table 2: Summary of selections applied to define the control, validation, and signal regions for the  $3\ell$  channel.

	$3\ell VV$	$3\ell$ Mixed	$3\ell t\bar{t}X$	$3\ell VVR$	$3\ell SR$
Preselection	$\geq 3$ leptons $\geq 1$ pair of OS-SF leptons with $ m(\ell\ell) - m_Z  < 10$ GeV $\geq 2$ central jets				
$b$ -tagged jets	0	1	$\geq 2$	$\geq 1$	$\geq 1$
forward jets	–	0	0	$\geq 1$	$\geq 1$
$\Delta\phi$ selections	–	$\Delta\phi(Z, \ell_3) < 2.6$	$\Delta\phi(Z, \ell_3) < 2.6$	$\Delta\phi(Z, \ell_3) < \frac{\pi}{2}$ <b>OR</b> $\Delta\phi(Z, b_{\text{lead}}) < \frac{\pi}{2}$	$\Delta\phi(Z, \ell_3) > \frac{\pi}{2}$ <b>AND</b> $\Delta\phi(Z, b_{\text{lead}}) > \frac{\pi}{2}$
other selections	–	–	–	–	$\max(p_T(\ell)) > 200$ GeV $p_T(\ell\ell) > 300$ GeV $H_T \cdot n(\text{jets}) < 6$ TeV

## 6 Systematic uncertainties

Uncertainties from experimental and theoretical sources are introduced in the normalization or shape of the final discriminant in the two analysis channels. Experimental uncertainties include effects on the electron energy scale and energy resolution, the muon momentum scale and resolution, and uncertainties in the data-to-MC correction factors for the electron and muon trigger, reconstruction, identification, and isolation efficiencies [100, 101]. Jet energy scale and resolution uncertainties are also included, as obtained from studies in data and simulation [105]. Flavor-tagging uncertainties include uncertainties in the  $b$ -jet tagging,  $c$ -jet mis-tagging, and light-jet mis-tagging efficiencies, and uncertainties due to extrapolations to regions not covered by the data used for the efficiency measurements [113–115]. Subdominant uncertainties include uncertainties related to the soft term in the  $E_T^{\text{miss}}$  calculation [116] and to the  $E_T^{\text{miss}}$  energy scale and resolution, uncertainties in the reweighting of the MC event samples to match the pileup conditions in data and a 1.7% [117] uncertainty in the integrated luminosity of the 2015–2018 data sample, obtained using the LUCID-2 detector [118] for the primary luminosity measurements.

Theoretical uncertainties include cross-section and other modeling uncertainties for all background samples. The cross-section uncertainties considered are between 5 – 6% for the  $Z$ +jets,  $t\bar{t}$ , and  $VV$  samples [77, 119, 120] and between 10 – 12% for the  $t\bar{t} + Z$  sample [121]. Uncertainties associated with missing higher orders are estimated by varying the renormalization and factorization scales by factors of 0.5 and 2 and then, constructing upward and downward varying envelopes by considering the largest positive and negative fluctuation on a bin-by-bin basis. PDF uncertainties when comparing distributions obtained with events generated with different PDF sets are obtained using the PDF4LHC recommendations [122]. For the  $2\ell$  channel, the  $Z$ +jets scale and PDF variations are decorrelated between  $Z$ +LF and  $Z$ +HF events. For the  $3\ell$  channel, the  $VV$  scale-variation uncertainties are broken down into their normalization and shape components, and treated as uncorrelated across regions with different the  $b$ -jet multiplicities.

Uncertainties due to the choice of generator or showering algorithm are estimated by comparing additional samples from alternative generators. Generator uncertainties for  $t\bar{t} + X$  are obtained from comparing the nominal MADGRAPH5\_AMC@NLO samples with the alternative SHERPA samples. For the  $3\ell$  channel where  $Z$ +jets is a minor background and mainly contributes via jets mis-identified as leptons, generator uncertainties are estimated by comparing nominal samples with alternative MADGRAPH5\_AMC@NLO samples and only taking the shape difference into account. For the  $2\ell$  channel where  $Z$ +jets is the main background, generator uncertainties are taken into account by comparing the SHERPA 2.2.1 sample with the SHERPA 2.2.11 sample where an independent reweighting based on jet multiplicity and  $H_T + E_T^{\text{miss}}$  distribution, as explained in Section 5.1, is applied on the latter. To estimate the uncertainties due to the modeling of initial-state radiation, final-state radiation, and the choice of parton shower algorithm, either additional event weights for the nominal samples or alternative MC samples were produced. These uncertainties were calculated by comparing the nominal MC distribution with the distributions obtained from these additional event weights or the alternative samples.

A conservative uncertainty of 30% on the fraction of  $VV$  and  $Z$ +jets events with heavy-flavor jets is applied in both the channels. This uncertainty is based on variations of the factorization and renormalization scales and the multijet merging scale [123]. Also, a shape uncertainty is introduced in the  $3\ell$  channel for  $VV$  and  $t\bar{t} + X$  by comparing the distributions before and after implementing jet multiplicity based reweighting explained in Section 5.2. In the  $2\ell$  channel, bin-by-bin variations in the reweighting factors due to the statistical fluctuations in the MC samples are calculated and these variations are used to obtain upward and downward fluctuating envelopes for the fit distributions. Shape uncertainties extracted from these fluctuations are applied on the  $Z$ +jets samples for the light and heavy-flavor components.

Events with jets misidentified as leptons can potentially arise in the  $2\ell$  channel mainly from the  $t\bar{t}$  process and in the  $3\ell$  channel from the  $t\bar{t}$  or  $Z$ +jets processes. After the  $m(\ell\ell)$  and  $p_T(\ell\ell)$  cuts, these contributions are estimated to be less than 2% in the signal regions. Following the prescription in Ref. [35], a conservative 25% uncertainty is applied on the normalization of the  $t\bar{t}$  ( $t\bar{t}$  and  $Z$ +jets) events in the  $2\ell$  ( $3\ell$ ) channel. This uncertainty is decorrelated across the control and signal regions in the  $2\ell$  and  $3\ell$  channels.

## 7 Statistical analysis and results

A binned profile likelihood fit in the discriminating variable  $p_T(\ell\ell)$  is performed in both channels to test if the data is compatible with the background-only hypothesis. The uncertainties are introduced in the likelihood as nuisance parameters (NP), which are constrained using Gaussian priors. Additional NPs are included to take into account the statistical uncertainties in each bin for each event category due to the limited size of the simulated samples [124]. The likelihood function  $L(\mu, \vec{\theta})$  is constructed as a product of Poisson probabilities for each bin of the discriminating variable in the control and signal regions,

$$L(\mu, \vec{\theta}) = \prod_{i=1}^{N_{\text{bins}}} \text{Pois}(n_i^{\text{data}} | b_i + \mu s_i) \times P(\vec{\theta}), \quad (4)$$

where  $N_{\text{bins}}$  is the total number of bins in control and signal regions. The parameters,  $n_i^{\text{data}}$ ,  $b_i$ ,  $s_i$ , represent the number of observed data, expected background and signal events respectively in the  $i$ -th bin. The product of prior distributions for the NPs,  $P(\vec{\theta})$ , is given by  $P(\vec{\theta})$ . The expected number of events in each bin depends on these NPs as well as the signal strength,  $\mu$ , which acts as a scaling factor for the signal cross-section.

The background and signal templates are fitted to the observed data to obtain the best estimate of signal and background contributions in the control and signal regions. The fitted values of  $\mu, \vec{\theta}$  are propagated to the validation regions to test the quality of the fit and the corresponding background modeling. The background-only *post-fit* distributions, obtained by maximizing the likelihood function in Eq. (4) with a fixed choice of  $\mu = 0$ , are shown for the  $2\ell$  and  $3\ell$  control, validation, and signal regions in Figures 6 and 7 respectively. The expected pre-fit signal yields for the  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  for the singlet and doublet representations of  $T$  are also shown in these figures. The contribution of different background processes, expected signal yields, and observed data in different regions for the  $2\ell$  and  $3\ell$  channels are summarized in Tables 3 and 4 respectively. The post-fit background contribution shows reasonable agreement in all regions with the observed data within the uncertainties.

To test whether the observed data is compatible with the background-only hypothesis, tests are performed with RooStats [125] with statistical models implemented using RooFit [126] and HistFactory [127]. No significant excess over the background expectation is observed. Hence, the results were used to set upper limits on the top partner total production cross-section  $\sigma(WTZt + ZTZt)$  at 95% confidence level (CL) using the  $\text{CL}_s$  method [128, 129]. These limits were calculated with the asymptotic approximation [130] of the test statistic and validated against the limits obtained using pseudoexperiments for benchmark signal samples with  $M_T = 1200, 1500, \text{ and } 2100$  GeV. For all these benchmarks, the expected and observed limits obtained from pseudoexperiments at 95% CL agreed with the limits obtained using the asymptotic approximation within 5%. The largest difference, close to 30%, in the results obtained with the two methods was observed in the  $-1\sigma$  and  $-2\sigma$  bands for the 2100 GeV mass point.

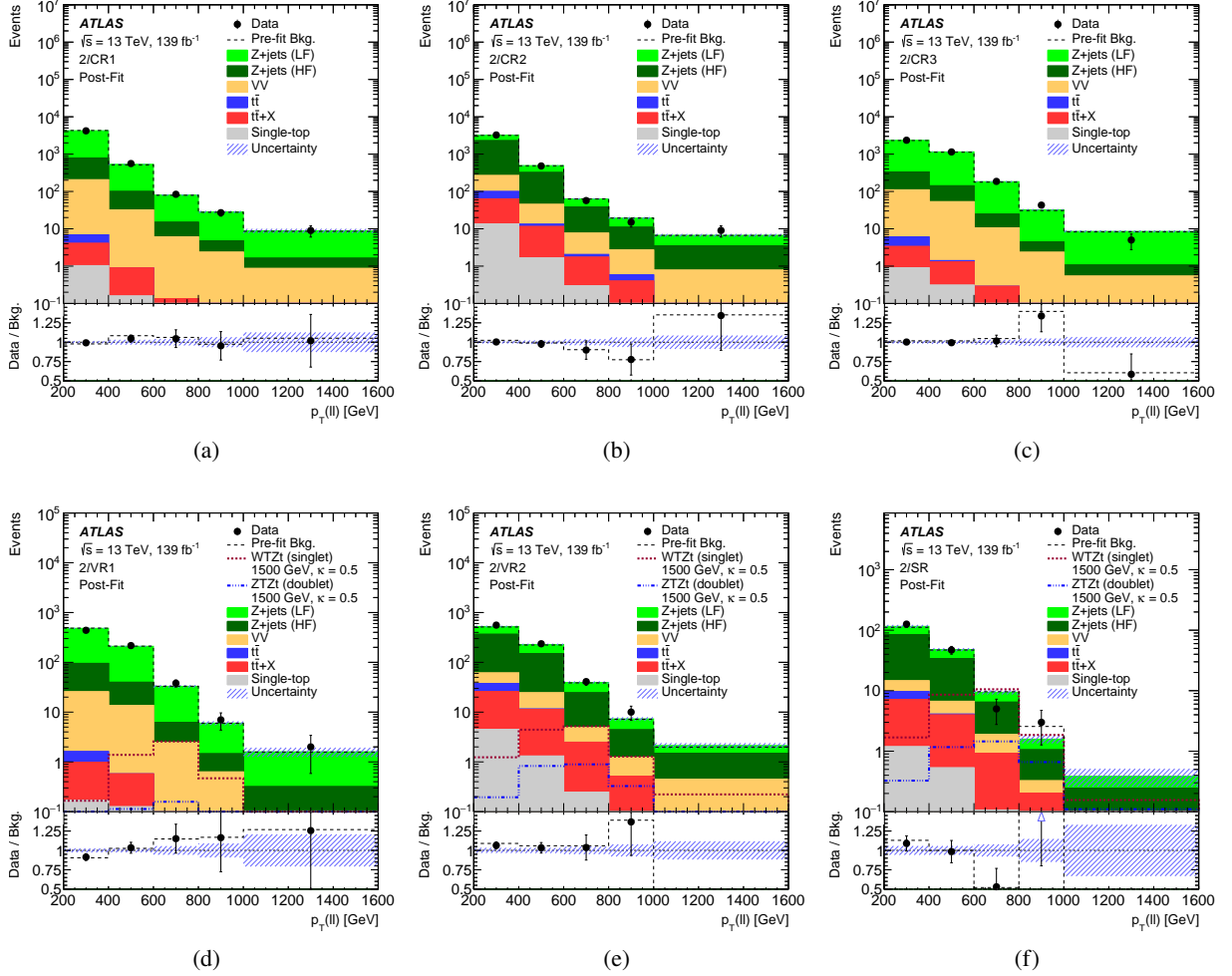


Figure 6: Distribution of the final discriminant,  $p_T(\ell\ell)$ , for the  $2\ell$  channel in the control regions (a)  $2\ell\text{CR1}$ , (b)  $2\ell\text{CR2}$ , (c)  $2\ell\text{CR3}$ , the validation regions (d)  $2\ell\text{VR1}$ , (e)  $2\ell\text{VR2}$ , and the signal region (f)  $2\ell\text{SR}$ . The distributions are shown after the background-only fit in the  $2\ell$  channel. The black, dashed line shows the total pre-fit background distribution and the corresponding data-over-background ratio in the bottom panel. The blue arrows in the bottom panel indicate points that are outside the vertical range of the figure. The overlaid red and blue dotted lines in (d)-(f) show the expected signal contributions in the corresponding regions from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in singlet and doublet representations respectively. Expected signal contributions in the control regions are negligible. The last bin in each distribution contains the overflow.

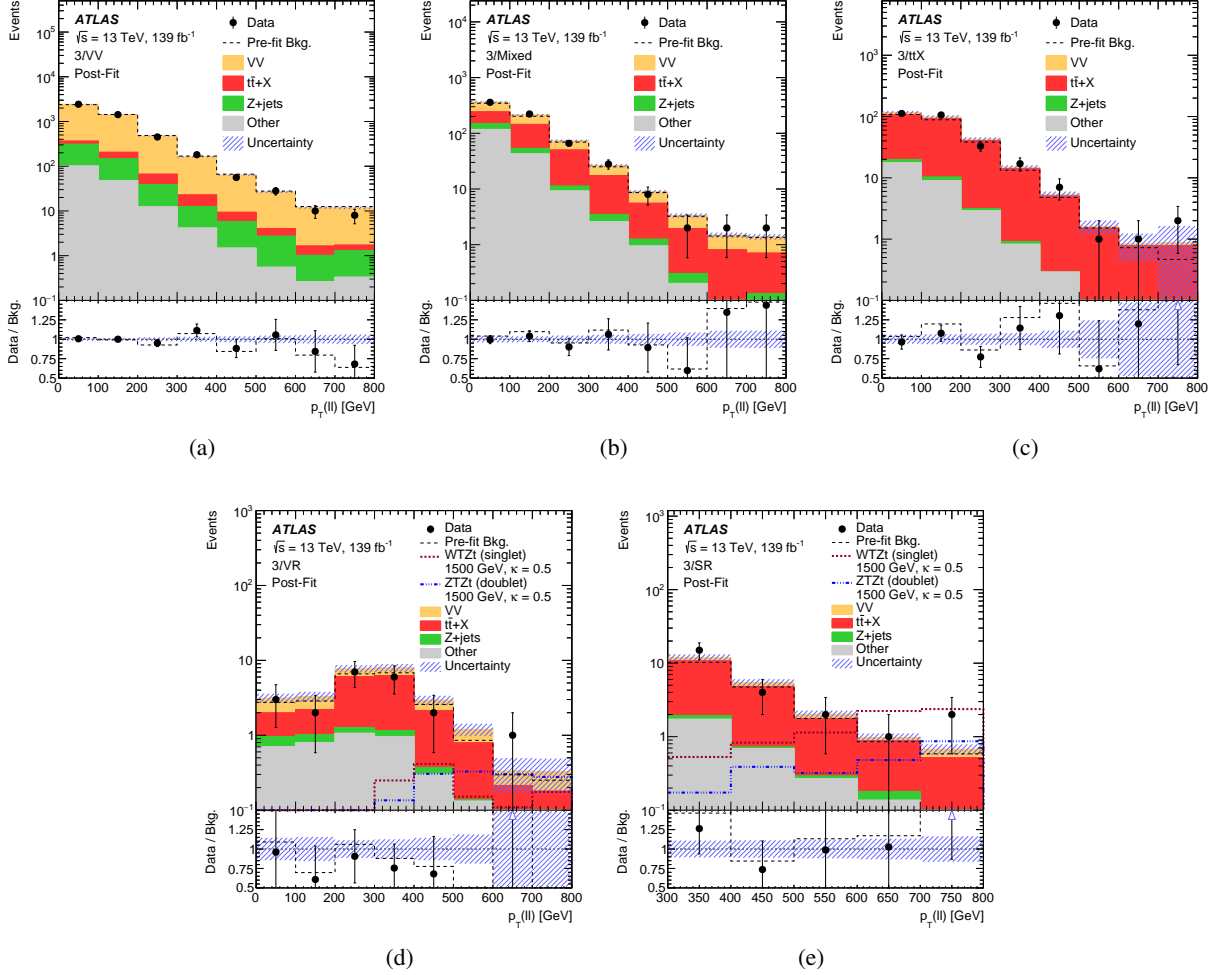


Figure 7: Distribution of the final discriminant,  $p_T(\ell\ell)$ , for the  $3\ell$  channel in the control regions (a)  $3\ell$ VV, (b)  $3\ell$ Mixed, (c)  $3\ell$ ttX, the validation region (d)  $3\ell$ VR, and the signal region (e)  $3\ell$ SR. The distributions are shown after the background-only fit in the  $3\ell$  channel. The black, dashed line shows the total pre-fit background distribution and the corresponding data-over-background ratio in the bottom panel. The blue arrows in the bottom panel indicate points that are outside the vertical range of the figure. The overlaid red and blue dotted lines in (d)-(e) show the expected signal contributions in the corresponding regions from  $WTZt$  and  $ZTZt$  processes with  $M_T = 1.5$  TeV and  $\kappa = 0.5$  in singlet and doublet representations respectively. Expected signal contributions in the control regions are negligible. The last bin in each distribution contains the overflow.



Table 3: Observed number of events in data and post-fit number of background events in the control, validation, and signal regions for the  $2\ell$  channel. The fit was performed under a background-only hypothesis. Statistical uncertainties from the limited size of MC samples and systematic uncertainties are added in quadrature. Systematic uncertainties take the correlations among nuisance parameters into account. The  $WTZt$  and  $ZTZt$  signal yields correspond to pre-fit yields.

	$2\ell\text{CR1}$	$2\ell\text{CR2}$	$2\ell\text{CR3}$	$2\ell\text{VR1}$	$2\ell\text{VR2}$	$2\ell\text{SR}$
$WTZt$ (singlet) $M_T = 1.5 \text{ TeV}, \kappa = 0.5$	$3.9 \pm 0.5$	$9.8 \pm 0.9$	$3.6 \pm 0.5$	$4.7 \pm 0.6$	$12.4 \pm 1.1$	$22.8 \pm 1.7$
$ZTZt$ (doublet) $M_T = 1.5 \text{ TeV}, \kappa = 0.5$	$0.34 \pm 0.04$	$1.64 \pm 0.16$	$0.19 \pm 0.04$	$0.32 \pm 0.05$	$2.28 \pm 0.22$	$3.72 \pm 0.24$
Z+jets (LF)	$3980 \pm 90$	$1170 \pm 110$	$3220 \pm 80$	$589 \pm 16$	$258 \pm 27$	$48 \pm 6$
Z+jets (HF)	$660 \pm 60$	$2330 \pm 140$	$330 \pm 50$	$101 \pm 9$	$452 \pm 28$	$103 \pm 10$
VV	$238 \pm 35$	$206 \pm 25$	$169 \pm 26$	$40 \pm 6$	$40 \pm 7$	$8.7 \pm 1.6$
$t\bar{t}$	$2.8 \pm 3.4$	$38 \pm 20$	$2.8 \pm 1.9$	$0.6 \pm 0.6$	$12 \pm 5$	$2.6 \pm 1.6$
$t\bar{t} + X$	$3.9 \pm 1.0$	$61 \pm 23$	$3.8 \pm 0.8$	$1.34 \pm 0.28$	$34 \pm 12$	$10 \pm 4$
Single-top	$1.23 \pm 0.17$	$15.6 \pm 1.8$	$1.27 \pm 0.15$	$0.33 \pm 0.09$	$6.1 \pm 0.8$	$1.85 \pm 0.26$
Total background	$4886 \pm 70$	$3821 \pm 60$	$3727 \pm 60$	$732 \pm 16$	$802 \pm 24$	$174 \pm 9$
Data	4887	3818	3735	704	846	181

Table 4: Observed number of events in data and post-fit number of background events in the control, validation, and signal regions for the  $3\ell$  channel. The fit was performed under a background-only hypothesis. Statistical uncertainties from the limited size of MC samples and systematic uncertainties are added in quadrature. Systematic uncertainties take the correlations among nuisance parameters into account. The  $WTZt$  and  $ZTZt$  signal yields correspond to pre-fit yields.

	$3\ell\text{VV}$	$3\ell\text{Mixed}$	$3\ell\text{ttX}$	$3\ell\text{VR}$	$3\ell\text{SR}$
$WTZt$ (singlet) $M_T = 1.5 \text{ TeV}, \kappa = 0.5$	$3.2 \pm 0.4$	$0.85 \pm 0.19$	$0.58 \pm 0.14$	$1.21 \pm 0.24$	$8.3 \pm 0.8$
$ZTZt$ (doublet) $M_T = 1.5 \text{ TeV}, \kappa = 0.5$	$0.84 \pm 0.24$	$0.68 \pm 0.12$	$0.40 \pm 0.12$	$2.0 \pm 0.8$	$2.04 \pm 0.32$
VV	$3920 \pm 160$	$230 \pm 40$	$21 \pm 4$	$7.1 \pm 1.5$	$2.4 \pm 0.5$
$t\bar{t} + X$	$156 \pm 27$	$245 \pm 28$	$225 \pm 20$	$14.9 \pm 2.5$	$15.4 \pm 2.2$
Z+jets	$350 \pm 120$	$45 \pm 13$	$3.8 \pm 1.5$	$1.0 \pm 0.4$	$0.31 \pm 0.11$
$t\bar{t}$	$82 \pm 30$	$93 \pm 16$	$5 \pm 6$	$0.9 \pm 1.1$	$0.005 \pm 0.007$
single-top	$63.1 \pm 2.6$	$77.9 \pm 1.6$	$25.8 \pm 0.8$	$2.99 \pm 0.17$	$2.88 \pm 0.17$
VVV	$22.6 \pm 1.2$	$1.22 \pm 0.09$	$0.087 \pm 0.012$	$0.078 \pm 0.015$	$0.023 \pm 0.004$
Total Bkg.	$4594 \pm 70$	$692 \pm 26$	$281 \pm 16$	$27.0 \pm 3.0$	$21.0 \pm 2.2$
Data	4590	690	279	21	24

The overall impact of the systematic uncertainties introduced in Section 6 on the final result is found to be small. Considering only statistical uncertainties, the limits on the total cross-section of the  $WTZt$  and  $ZTZt$  processes are reduced by less than 14%.

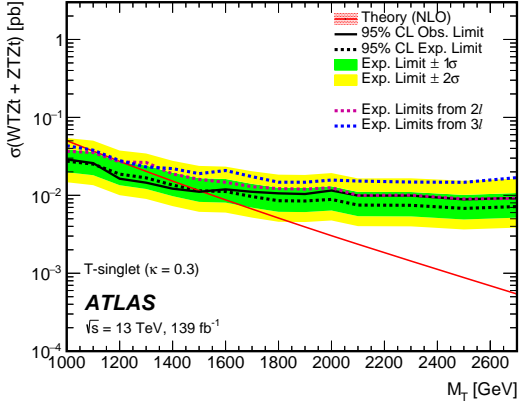
Given that the signal efficiencies for the  $WTZt$  and  $ZTZt$  signal modes can be different, limits were independently calculated for different combinations of top partner mass, coupling, and branching ratios, under the assumption that  $\xi_Z = \xi_H$ . The latter is motivated by the Goldstone equivalence theorem [131], which states that the asymptotic branching ratios of  $T$  decaying into  $Zt$  and  $Ht$  become similar in the large

$M_T$  limit under the narrow-width approximation.

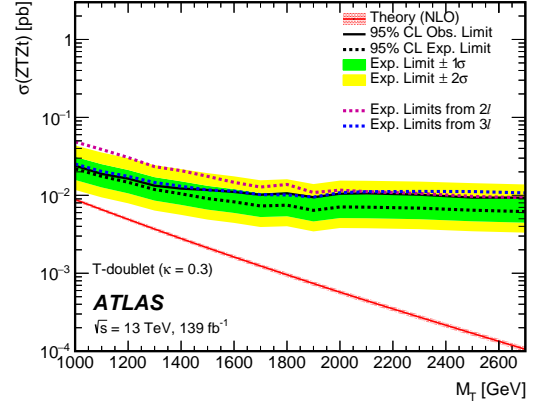
To maximize the sensitivity of the search, the results from both channels were statistically combined. All experimental uncertainties are treated as fully correlated across the channels while the theoretical and modeling uncertainties are taken to be uncorrelated. The corresponding limits on the signal strength translate into limits on the total cross-section of the  $WTZt$  and  $ZTZt$  processes for the corresponding choice of coupling and branching ratio. Limits corresponding to  $\kappa = 0.3, 0.5$ , and  $0.7$  for the singlet and doublet representations are shown in Figure 8. The combined expected limit is significantly stronger than limits obtained independently from the  $2\ell$  and  $3\ell$  channels. For the singlet representation with  $\kappa = 0.5$ , top partner masses less than 1825 GeV are excluded. The doublet representation only receives contribution from the  $ZTZt$  process that has a much smaller cross-section because of the suppressed matrix element contribution from the gluon splitting into  $t\bar{t}$  (Figure 1). Hence, the doublet exclusion limits are found to be considerably weaker than the singlet limits and none of the considered  $T$  masses are excluded for  $\kappa \leq 0.5$ .

The cross-section limits calculated for different choices of coupling can be reinterpreted as limits on the plane of  $(M_T, \kappa)$ , as shown in Figures 9(a) and 9(b) for the singlet and doublet representations only. Following the prescription of Ref. [132], this interpretation is only applied for the parametric subspace where the top partner's relative decay width,  $\frac{\Gamma_T}{M_T}$  is smaller than 50%. The correction factors for finite width and non-resonant contributions [98, 99] are also valid for the same parametric subspace. Hence, the limits in Figure 9 beyond  $\frac{\Gamma_T}{M_T} > 50\%$  are omitted for phenomenological consistency.

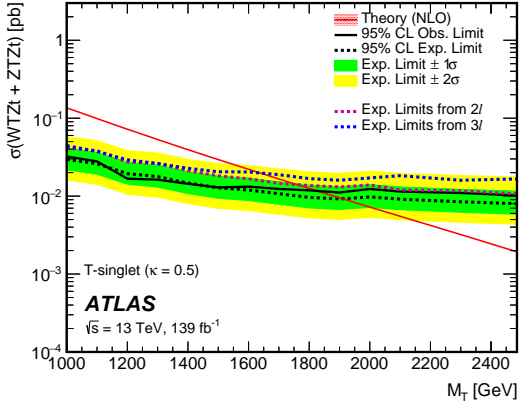
To further generalize the interpretation of the search results for any possible combination of top partner branching ratios, the semi-analytical interpretation strategy presented in Ref. [98] was adopted in Figure 10 to represent the excluded top partner mass as a function of its relative decay width and  $\xi_W$ , the asymptotic branching ratio for the  $T \rightarrow Wb$  decay mode. The  $\xi$  parameters for the other two decay modes are determined by the constraint  $\xi_Z = \xi_H = \frac{1-\xi_W}{2}$ .



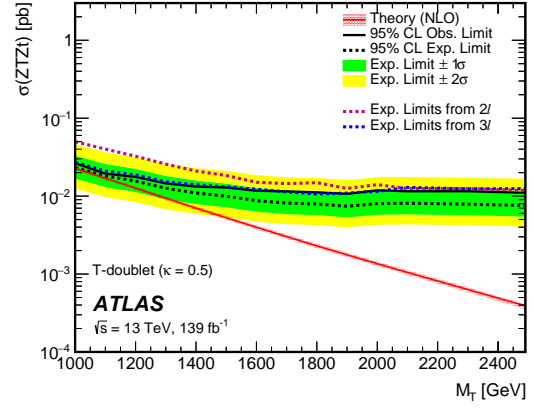
(a)



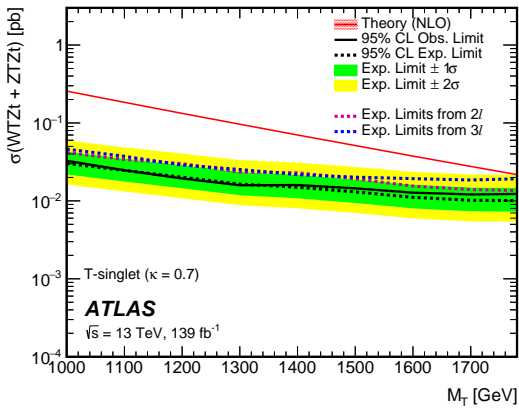
(b)



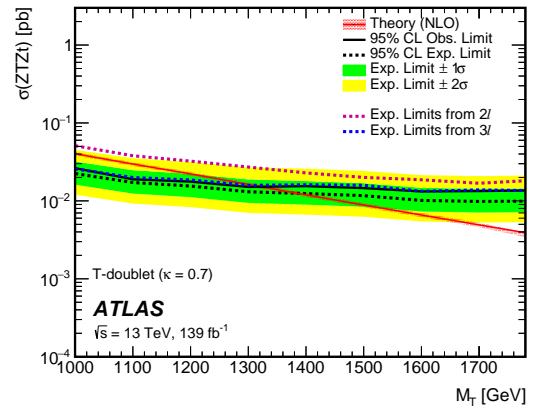
(c)



(d)



(e)



(f)

Figure 8: Observed (solid black line) and expected (dashed black line) limits calculated at 95% CL on total cross-section of the  $WTZt$  and  $ZTZt$  processes as a function of the  $T$  mass for a choice of coupling  $\kappa = 0.3$  (top row),  $\kappa = 0.5$  (middle row), and  $\kappa = 0.7$  (bottom row) for singlet (left column) and doublet (right column) representations. Expected limits calculated at 95% CL by independently fitting the  $2\ell$  and  $3\ell$  channels are shown as overlaid blue and red dotted lines respectively. The green (yellow) band is the 68% (95%) confidence interval around the median expected limit. The solid red line shows the theory prediction for the NLO cross-section, with the surrounding shaded band representing the corresponding uncertainty. The mass range in these figures is restricted to ensure that the relative decay width  $\frac{\Gamma_T}{M_T}$  is limited to less than 50%.

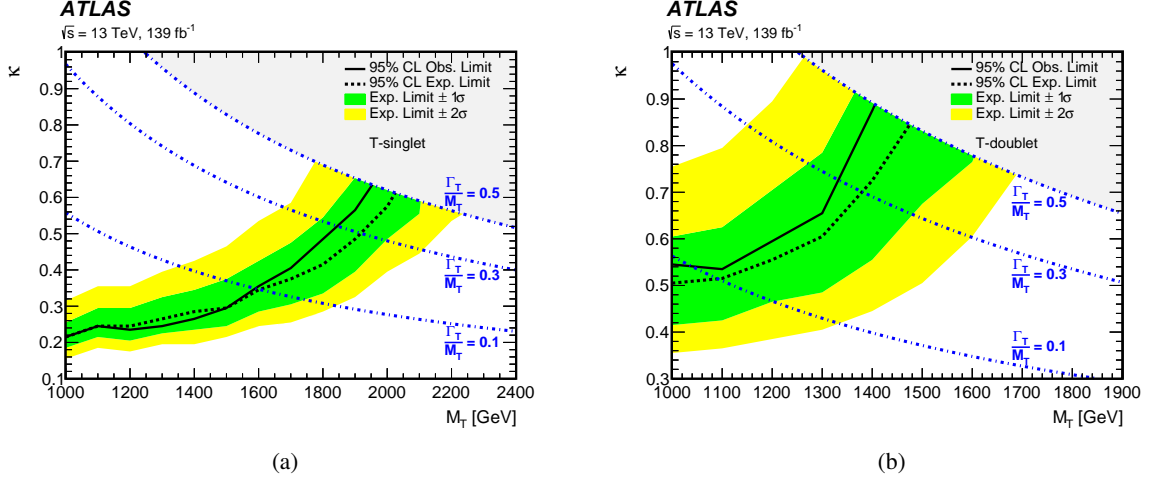


Figure 9: Observed (solid black line) and expected (dashed black line) limits calculated at 95% CL on the top partner coupling as a function of the  $T$  mass for (a) singlet and (b) doublet representations. The shaded regions represent the parametric space with  $\frac{\Gamma_T}{M_T} > 50\%$ . The green (yellow) band is the 68% (95%) confidence interval around the median expected limit. The dotted lines represent contours of equal values in  $\frac{\Gamma_T}{M_T}$ .

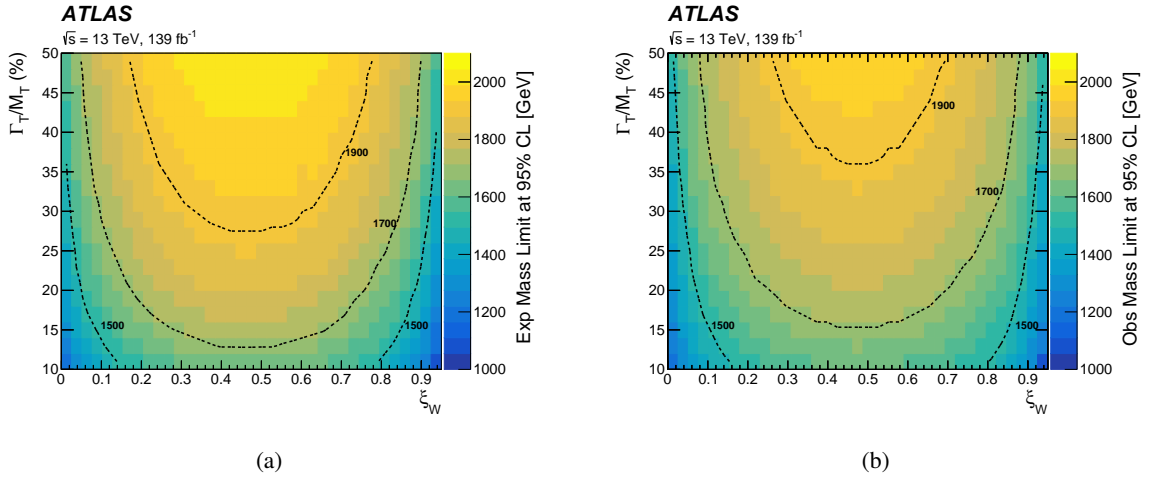


Figure 10: (a) Expected and (b) observed upper limits calculated at 95% CL on the excluded top partner mass as a function of its relative decay width ( $\frac{\Gamma_T}{M_T}$ ) and the relative coupling parameter  $\xi_W$  with the assumption of  $\xi_Z = \xi_H = \frac{1-\xi_W}{2}$ . The dotted contour lines represent exclusion limits of equal mass in units of GeV.

## 8 Conclusion

A search for the production of a single vector-like quark  $T$  with electric charge  $\frac{2}{3}e$  is presented in a final state with a leptonically decaying  $Z$  boson and a top quark. Two orthogonal channels with two or at least three leptons are separately optimized and their results are statistically combined to obtain the final result. No significant excess over the background expectation is observed, and therefore 95% CL upper limits on the inclusive single- $T$  cross-section are derived for different choices of  $T$  mass and couplings. These results are interpreted in terms of limits on the  $T$  mass and coupling for different electroweak representations and generalized branching ratio scenarios. For the singlet representation,  $\kappa$  values are excluded between 0.22 and 0.64 for masses of  $T$  between 1000 and 1975 GeV. The exclusion range for the doublet representation spans  $\kappa$  values between 0.54 and 0.88 for masses of  $T$  between 1000 and 1425 GeV. The strongest exclusion is observed for singlet representation with  $\xi_W \approx 0.5$  where masses up to 1975 GeV are excluded at relative decay width of  $\frac{\Gamma_T}{M_T} = 0.5$  for the top partner.

## Acknowledgments

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; ANID, Chile; CAS, MOST and NSFC, China; Minciencias, Colombia; MEYS CR, Czech Republic; DNRF and DNSRC, Denmark; IN2P3-CNRS and CEA-DRF/IRFU, France; SRNSFG, Georgia; BMBF, HGF and MPG, Germany; GSRI, Greece; RGC and Hong Kong SAR, China; ISF and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; NWO, Netherlands; RCN, Norway; MEiN, Poland; FCT, Portugal; MNE/IFA, Romania; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DSI/NRF, South Africa; MICINN, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TENMAK, Türkiye; STFC, United Kingdom; DOE and NSF, United States of America. In addition, individual groups and members have received support from BCKDF, CANARIE, Compute Canada and CRC, Canada; PRIMUS 21/SCI/017 and UNCE SCI/013, Czech Republic; COST, ERC, ERDF, Horizon 2020 and Marie Skłodowska-Curie Actions, European Union; Investissements d’Avenir Labex, Investissements d’Avenir IDEX and ANR, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF, Greece; BSF-NSF and MINERVA, Israel; Norwegian Financial Mechanism 2014-2021, Norway; NCN and NAWA, Poland; La Caixa Banking Foundation, CERCA Programme Generalitat de Catalunya and PROMETEO and GenT Programmes Generalitat Valenciana, Spain; Göran Gustafssons Stiftelse, Sweden; The Royal Society and Leverhulme Trust, United Kingdom.

The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN, the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA), the Tier-2 facilities worldwide and large non-WLCG resource providers. Major contributors of computing resources are listed in Ref. [133].

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J.D. Lomas <sup>20</sup>, J.D. Long <sup>162</sup>, I. Longarini <sup>160</sup>, L. Longo <sup>70a,70b</sup>, R. Longo <sup>162</sup>, I. Lopez Paz <sup>67</sup>,  
A. Lopez Solis <sup>48</sup>, J. Lorenz <sup>109</sup>, N. Lorenzo Martinez <sup>4</sup>, A.M. Lory <sup>109</sup>, O. Loseva <sup>37</sup>,  
X. Lou <sup>47a,47b</sup>, X. Lou <sup>14a,14e</sup>, A. Lounis <sup>66</sup>, J. Love <sup>6</sup>, P.A. Love <sup>91</sup>, G. Lu <sup>14a,14e</sup>, M. Lu <sup>80</sup>,  
S. Lu <sup>128</sup>, Y.J. Lu <sup>65</sup>, H.J. Lubatti <sup>138</sup>, C. Luci <sup>75a,75b</sup>, F.L. Lucio Alves <sup>14c</sup>, A. Lucotte <sup>60</sup>,  
F. Luehring <sup>68</sup>, I. Luise <sup>145</sup>, O. Lukianchuk <sup>66</sup>, O. Lundberg <sup>144</sup>, B. Lund-Jensen <sup>144</sup>,  
N.A. Luongo <sup>123</sup>, M.S. Lutz <sup>151</sup>, D. Lynn <sup>29</sup>, H. Lyons <sup>92</sup>, R. Lysak <sup>131</sup>, E. Lytken <sup>98</sup>,  
V. Lyubushkin <sup>38</sup>, T. Lyubushkina <sup>38</sup>, M.M. Lyukova <sup>145</sup>, H. Ma <sup>29</sup>, K. Ma <sup>62a</sup>, L.L. Ma <sup>62b</sup>,  
Y. Ma <sup>121</sup>, D.M. Mac Donell <sup>165</sup>, G. Maccarrone <sup>53</sup>, J.C. MacDonald <sup>100</sup>, R. Madar <sup>40</sup>,  
W.F. Mader <sup>50</sup>, J. Maeda <sup>85</sup>, T. Maeno <sup>29</sup>, M. Maerker <sup>50</sup>, H. Maguire <sup>139</sup>, V. Maiboroda <sup>135</sup>,  
A. Maio <sup>130a,130b,130d</sup>, K. Maj <sup>86a</sup>, O. Majersky <sup>48</sup>, S. Majewski <sup>123</sup>, N. Makovec <sup>66</sup>,  
V. Maksimovic <sup>15</sup>, B. Malaescu <sup>127</sup>, Pa. Malecki <sup>87</sup>, V.P. Maleev <sup>37</sup>, F. Malek <sup>60</sup>, M. Mali <sup>93</sup>,  
D. Malito <sup>95,r</sup>, U. Mallik <sup>80</sup>, S. Maltezos <sup>10</sup>, S. Malyukov <sup>38</sup>, J. Mamuzic <sup>13</sup>, G. Mancini <sup>53</sup>,  
G. Manco <sup>73a,73b</sup>, J.P. Mandalia <sup>94</sup>, I. Mandić <sup>93</sup>, L. Manhaes de Andrade Filho <sup>83a</sup>,  
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 J.L. Oliver <sup>160</sup>, A. Olszewski <sup>87</sup>, Ö.O. Öncel <sup>54</sup>, D.C. O'Neil <sup>142</sup>, A.P. O'Neill <sup>19</sup>,  
 A. Onofre <sup>130a,130e</sup>, P.U.E. Onyisi <sup>11</sup>, M.J. Oreglia <sup>39</sup>, G.E. Orellana <sup>90</sup>, D. Orestano <sup>77a,77b</sup>,  
 N. Orlando <sup>13</sup>, R.S. Orr <sup>155</sup>, V. O'Shea <sup>59</sup>, L.M. Osojnak <sup>128</sup>, R. Ospanov <sup>62a</sup>,  
 G. Otero y Garzon <sup>30</sup>, H. Otono <sup>89</sup>, P.S. Ott <sup>63a</sup>, G.J. Ottino <sup>17a</sup>, M. Ouchrif <sup>35d</sup>, J. Ouellette <sup>29</sup>,  
 F. Ould-Saada <sup>125</sup>, M. Owen <sup>59</sup>, R.E. Owen <sup>134</sup>, K.Y. Oyulmaz <sup>21a</sup>, V.E. Ozcan <sup>21a</sup>, N. Ozturk <sup>8</sup>,  
 S. Ozturk <sup>82</sup>, H.A. Pacey <sup>32</sup>, A. Pacheco Pages <sup>13</sup>, C. Padilla Aranda <sup>13</sup>, G. Padovano <sup>75a,75b</sup>,  
 S. Pagan Griso <sup>17a</sup>, G. Palacino <sup>68</sup>, A. Palazzo <sup>70a,70b</sup>, S. Palestini <sup>36</sup>, J. Pan <sup>172</sup>, T. Pan <sup>64a</sup>,  
 D.K. Panchal <sup>11</sup>, C.E. Pandini <sup>114</sup>, J.G. Panduro Vazquez <sup>95</sup>, H. Pang <sup>14b</sup>, P. Pani <sup>48</sup>,  
 G. Panizzo <sup>69a,69c</sup>, L. Paolozzi <sup>56</sup>, C. Papadatos <sup>108</sup>, S. Parajuli <sup>44</sup>, A. Paramonov <sup>6</sup>,  
 C. Paraskevopoulos <sup>10</sup>, D. Paredes Hernandez <sup>64b</sup>, T.H. Park <sup>155</sup>, M.A. Parker <sup>32</sup>, F. Parodi <sup>57b,57a</sup>,  
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 P. Pasuwan <sup>47a,47b</sup>, P. Patel <sup>87</sup>, U.M. Patel <sup>51</sup>, J.R. Pater <sup>101</sup>, T. Pauly <sup>36</sup>, J. Pearkes <sup>143</sup>,  
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