

Top Physics at the electron proton colliders

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In this talk we present some top physics program at future electron proton colliders. Mentioned topics involve but not limited to top structure function, top parton distribution functions, top spin polarization, top electric charge, measurement of V_{tb} , V_{ts} , V_{td} CKM matrix elements, anomalous $t\bar{t}\gamma$, $t\bar{t}Z$, $t\bar{t}W$, $tq\gamma$, tuh couplings and CP phase of tth coupling. The study of flavor changing neutral current tuh couplings and the measurement of $V_{td(s)}$ CKM matrix elements are highlight to show current progress in the top sector.

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1. Introduction

Electron-proton (ep) colliders are hybrids between electron-positron (e^+e^-) and proton-proton (pp) colliders, which consist of a hadron beam with an electron beam. They provide a cleaner environment compared to the pp colliders and higher center-of-mass (c.m.s.) energies to the e^+e^- ones. The scenarios studied here involve an electron beam energy of 60 GeV and a LHC proton beam of 7 TeV (LHeC) leading to a c.m.s. energy of 1.3 TeV, and a FCC-hh proton beam of 50 TeV (FCC-eh) leading to a c.m.s. energy of 3.5 TeV, respectively. Both projects can deliver up to 2 ab^{-1} integrated luminosity over their lifetime. Such colliders would become top factories allowing to analyze the electroweak(EW) interactions of the top quark particularly well, and would allow to perform sensitive searches for new physics up to TeV scale.

This talk focuses on top physics program at the LHeC and FCC-eh which are summarized and quickly mentioned in Section.2. Two selected topics, the study of flavor changing neutral current (FCNC) t - H couplings and the measurement of $V_{td(s)}$ CKM matrix elements are highlight to show current progress in the top sector. Finally, we end this talk with a short summary.

2. Top quark production at the ep colliders

By ep collisions, top quark can be produced in pair or singly, through both neutral current (NC) and charged current(CC) productions. The NC top productions include deep inelastic scattering (DIS) or photoproduction(γp) modes, while the CC ones are through W boson exchange. Some production rates depending on the collision energy are list in Fig.1. We see that standard

		top pair			single top		
		NC		CC	NC		CC
		DIS	rp	w-exch.	DIS	rp	w-exch.
30GeV	7TeV	0.0040pb	0.0091pb	-	4.653fb	12.54fb	0.7599pb
40GeV	7TeV	0.0090pb	0.0205pb		9.193fb	24.16fb	1.1850pb
50GeV	7TeV	0.0165pb	0.0354pb		14.85fb	38.27fb	1.6270pb
60GeV	7TeV	0.0253pb	0.0531pb		21.37fb	54.31fb	2.0835pb
60GeV	50TeV	0.6268pb	1.1660pb		40.29fb	942.8fb	16.701pb

Figure 1: Top production rate at ep colliders.

model(SM) top quark production at a future ep collider is dominated by single top quark production, mainly via CC mode. The total cross section is $\sim 2 \text{ pb}$ at the LHeC and $\sim 16 \text{ pb}$ at the FCC-eh. The other important top quark production mode is $t\bar{t}$ photoproduction with a total cross section of 53 fb at the LHeC and $\sim 1 \text{ pb}$ at the FCC-eh. This makes a future ep collider an ideal tool to study in particular the EW interactions of the top quark.

3. Selected topics in top sector at the ep colliders

3.1 Topics in top sector summarize

DIS fermion pair production is sensitive to the gluon density in proton. By using DIS $t\bar{t}$ production one can study the top component of the structure function at small x region[1]. Top

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pair photoproduction can be used to measure the $t\bar{t}\gamma$ vertex, thus the electric charge of the top quark. When the anomalous $t\bar{t}\gamma$ vertex is introduced, top pair photoproduction can also be used to set constraints on the Magnetic and Electric Dipole moment[2]. Similarly but through DIS top pair production from Gluon-Z fusion, one can study the anomalous $t\bar{t}Z$ couplings, though with less sensitivity. For the anomalous $tq\gamma$ interaction, it is found that ep based γp collider can provide a nice place to probe it[3]. With 1ab^{-1} integrated luminosity $\text{Br}(t \rightarrow q\gamma)$ can be probed up to order of 10^{-6} , which shows quite nice feature[4].

The CC single top production can be precisely measured through W boson leptonic or hadronic decay modes, and can be used to measure the CKM matrix element V_{tb} . It is found that at the LHeC, with luminosity of 100fb^{-1} , V_{tb} can be measured with a precision of 0.5%. In addition, the top quark spin polarization can also be precisely measured through this production[5]. Furthermore, the CC single top production is also a good way to measure the anomalous $t\bar{b}W$ couplings. Study was performed in [6] by means of the following effective CP conserving lagrangian:

$$\mathcal{L}_{Wtb} = \frac{g}{\sqrt{2}} \left[W_\mu \bar{t} \gamma^\mu (V_{tb} f_1^L P_L + f_1^R P_R) b - \frac{1}{2m_W} W_{\mu\nu} \bar{t} \sigma^{\mu\nu} (f_2^L P_L + f_2^R P_R) b \right] + \text{h.c.} \quad (3.1)$$

where $f_1^L (\equiv 1 + \Delta f_1^L)$ and f_1^R are left- and right-handed vector couplings, $f_2^{L,R}$ are left- and right-handed tensor couplings. The analysis shows that the $t\bar{b}W$ vertex can be probed at the LHeC to a very high accuracy. There is another CC top production mode making it possible to consider quark density for the top quark, since at very high scales the top may be considered "light". A six-flavor scheme has been proposed in [4], thus the LHeC offers new field of research for the top quark pdfs. The CP-nature of the $t\bar{t}H$ coupling can also be studied through top-Higgs associated production[7]. The CP-phase dependent lagrangian can be written here,

$$\mathcal{L} = -i \frac{m_t}{v} \bar{t} [\kappa \cos \zeta_t + i \gamma_5 \sin \zeta_t] t H. \quad (3.2)$$

Here ζ_t is the phases of the top-Higgs couplings. The case $\kappa = 1$, $\zeta_t = 0$ corresponds to the SM. $\zeta_t = 0$ or $\zeta_t = \pi$ correspond to a pure scalar state while $\zeta_t = \pi/2$ to a pure pseudo scalar state. The ranges $0 < \zeta_t < \pi/2$ or $\pi/2 < \zeta_t < \pi$ represent a mixture of the different CP-states. The study was performed by considering H to $b\bar{b}$ and top quark leptonic decay modes. It was found the LHeC provides a better environment to test the CP nature of Higgs boson couplings compare to LHC.

3.2 Progress in top sector: V_{ts} and V_{td} CKM matrix elements

The signals that involve at least one V_{tx} vertex we studied are

$$\begin{aligned} \text{Signal.1} : & \quad p e^- \rightarrow \nu_e \bar{t} \rightarrow \nu_e W^- \bar{b} \rightarrow \nu_e \ell^- \nu_\ell \bar{b}, \\ \text{Signal.2} : & \quad p e^- \rightarrow \nu_e W^- b \rightarrow \nu_e \ell^- \nu_\ell b, \\ \text{Signal.3} : & \quad p e^- \rightarrow \nu_e \bar{t} \rightarrow \nu_e W^- j \rightarrow \nu_e \ell^- \nu_\ell j, \\ \text{Signal.4} : & \quad p e^- \rightarrow \nu_e W^- j \rightarrow \nu_e \ell^- \nu_\ell j. \end{aligned} \quad (3.3)$$

We present the measurement potential by using different channels separately, depending on the possibility to distinguish them kinematically. But the same final state contributions are fully considered when considering different channels. Our conclusion show that Signal.1, 2 and 3 show very

$\mathcal{P}_e = 0.8$ 5% syst.	LHeC SS	$1ab^{-1}$		$2ab^{-1}$		$1ab^{-1}$		$2ab^{-1}$	
		R_d	V_{td}	R_d	V_{td}	R_s	V_{ts}	R_s	V_{ts}
signal.1	2σ	8.16	0.0699	6.86	0.0588	2.39	0.0980	2.01	0.0824
	3σ	9.99	0.0857	8.40	0.0720	2.92	0.1200	2.46	0.1009
	5σ	12.90	0.1106	10.84	0.0930	3.77	0.1550	3.17	0.1303
signal.2	2σ	10.53	0.0903	8.85	0.0759	5.33	0.2192	4.49	0.1843
	3σ	12.90	0.1106	10.84	0.0930	6.53	0.2684	5.49	0.2257
	5σ	16.65	0.1428	14.00	0.1201	8.44	0.3465	7.09	0.2914
signal.3	2σ	6.90	0.0592	5.83	0.0500	1.56	0.0641	1.32	0.0541
	3σ	8.43	0.0723	7.13	0.0611	1.91	0.0804	1.61	0.0662
	5σ	10.82	0.0928	9.16	0.0785	2.45	0.1008	2.07	0.0852

Table 1: The 2, 3, and 5 σ limits for V_{td} and V_{ts} at the LHeC.

$\mathcal{P}_e = 0.8$ 5% syst.	FCC-eh SS	$1ab^{-1}$		$2ab^{-1}$		$1ab^{-1}$		$2ab^{-1}$	
		R_d	V_{td}	R_d	V_{td}	R_s	V_{ts}	R_s	V_{ts}
signal.1	2σ	5.46	0.0468	4.60	0.0394	1.74	0.0716	1.47	0.0602
	3σ	6.69	0.0573	5.63	0.0483	2.13	0.0876	1.80	0.0738
	5σ	8.63	0.0740	7.27	0.0623	2.75	0.1131	2.32	0.0952
signal.2	2σ	6.46	0.0554	5.44	0.0466	2.75	0.1131	2.32	0.0952
	3σ	7.91	0.0679	6.66	0.0571	3.37	0.1386	2.84	0.1166
	5σ	10.22	0.0876	8.60	0.0737	4.35	0.1789	3.67	0.1506
signal.3	2σ	5.01	0.0430	4.27	0.0366	1.05	0.0433	0.90	0.0369
	3σ	6.13	0.0526	5.22	0.0448	1.29	0.0530	1.10	0.0452
	5σ	7.90	0.0677	6.73	0.0577	1.66	0.0684	1.42	0.0582

Table 2: The 2, 3, and 5 σ limits for V_{td} and V_{ts} at the FCC-eh.

good potential for the measurement of both V_{td} and V_{ts} , and Signal.3 is the most prompt one. In contrary, Signal.4 is of lower quality to be recommended. We displayed the 2, 3, and 5 σ limits in Tab.1 for LHeC and 2 for FCC-eh. We find that considering the proposed $2ab^{-1}$ luminosity and 80% electron polarization, the 2σ limits for R_{td} parameters for Signal.1, 2, 3 are 6.86, 8.85, 5.83 at the LHeC and 4.6, 5.44, 4.27 at the FCC-eh, corresponding to the limits of V_{td} equal 0.0588, 0.0759, 0.05 at the LHeC and 0.0394, 0.0466, 0.0366 at the FCC-eh. Similarly, for R_{ts} , the 2σ R parameter limits are 2.01, 4.49, 1.32 at the LHeC and 1.47, 2.32, 0.90 at the FCC-eh, corresponding to the limits of V_{ts} equal 0.0824, 0.1843, 0.0541 at the LHeC and 0.0602, 0.0952, 0.0369 at the FCC-eh. In summary, the electron proton collider provide a nice platform in measuring the V_{td} and V_{ts} CKM matrix elements.

3.3 Progress in top sector: FCNC tuH couplings

We consider two signal productions that contain the top-Higgs FCNC couplings[8]. One is

$$\text{signal.I: } e^- p \rightarrow \nu_e \bar{t} \rightarrow \nu_e h \bar{q} \rightarrow \nu_e b \bar{b} \bar{q}, \quad (3.4)$$

where $q=u$ or c . In this case, the five flavor scheme should be applied and an initial state bottom quark will collide with a w boson to produce a single top, which decay anomalously to a Higgs and a light quark. The studied topology gives rise to the $E_T^{\text{miss}} + \text{jets}$ signature characterized by three (or more than three) jets and a missing transverse momentum (E_T^{miss}) from the undetected neutrino. Two of the jets should be tagged as B-jets. The combination of the two B-jets should appear as a narrow resonance centered around the SM Higgs boson mass. Together with the remaining light jet(s), they should be able to reconstruct a resonant top quark. The second channel we considered is

$$\text{singal.II} : e^- p \rightarrow \nu_e h b \rightarrow \nu_e b \bar{b} b, \quad (3.5)$$

where the FCNC tqh couplings are induced through light quarks that directly emitting from the proton. Considering the studied topology, three tagged B-jets should be required for signal.II. The Feynman diagrams are plotted in Fig.2.

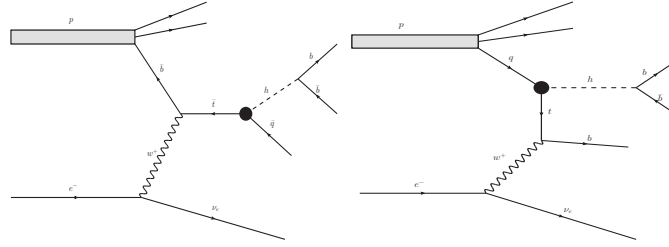


Figure 2: Illustrated Feynman diagrams for the processes $e^- p \rightarrow \nu_e \bar{t} \rightarrow \nu_e h \bar{q} \rightarrow \nu_e b \bar{b} \bar{q}$ (signal.I) and $e^- p \rightarrow \nu_e h b \rightarrow \nu_e b \bar{b} b$ (signal.II) at the ep colliders that contain flavor changing top-Higgs interactions.

In Fig.3, the upper limit on $\text{Br}(t \rightarrow uh)$ at 99.99, 99.73, 95.40, 68.27% C.L. as a function of the integrated luminosity are plotted. Our conclusion is that, for signal.I, at the high luminosity (up to 1ab^{-1}) ep colliders where the electrons have a polarisation of 80%, the 1σ , 2σ , 3σ and 5σ upper limit on $\text{Br}(t \rightarrow uh)$ are $0.075 \times 10^{-2} (0.14 \times 10^{-3})$, $0.15 \times 10^{-2} (0.29 \times 10^{-3})$, $0.22 \times 10^{-2} (0.43 \times 10^{-3})$ and $0.38 \times 10^{-2} (0.72 \times 10^{-3})$ at the LHeC(FCC-eh). For signal.II, the boundaries are becoming $0.064 \times 10^{-2} (0.097 \times 10^{-3})$, $0.15 \times 10^{-2} (0.22 \times 10^{-3})$, $0.26 \times 10^{-2} (0.35 \times 10^{-3})$ and $0.53 \times 10^{-2} (0.68 \times 10^{-3})$ at the LHeC(FCC-eh) respectively. We can see that signal.II can even have better potential than signal.I at the FCC-eh due to its clean environment. Notice here we use 5% systematic uncertainty for backgrounds yields only at both ep colliders.

4. SUMMARY

In this talk we present an overview of top physics at the ep colliders. Mentioned topics involve but not limited to top structure function, top parton distribution functions, top spin polarization, top electric charge, measurement of V_{tb} , V_{ts} , V_{td} CKM matrix elements, anomalous $t\bar{t}\gamma$, $t\bar{t}Z$, $t\bar{t}W$, $tq\gamma$, tuh couplings and CP phase of tth coupling, etc. The flavor changing neutral current tuh couplings, and the study of $V_{td(s)}$ CKM matrix elements are highlight to show current progress. New ideas or contributions are also welcomed in order to further emphasize the strength of the LHeC (FCC-eh) when it comes to the top sector.

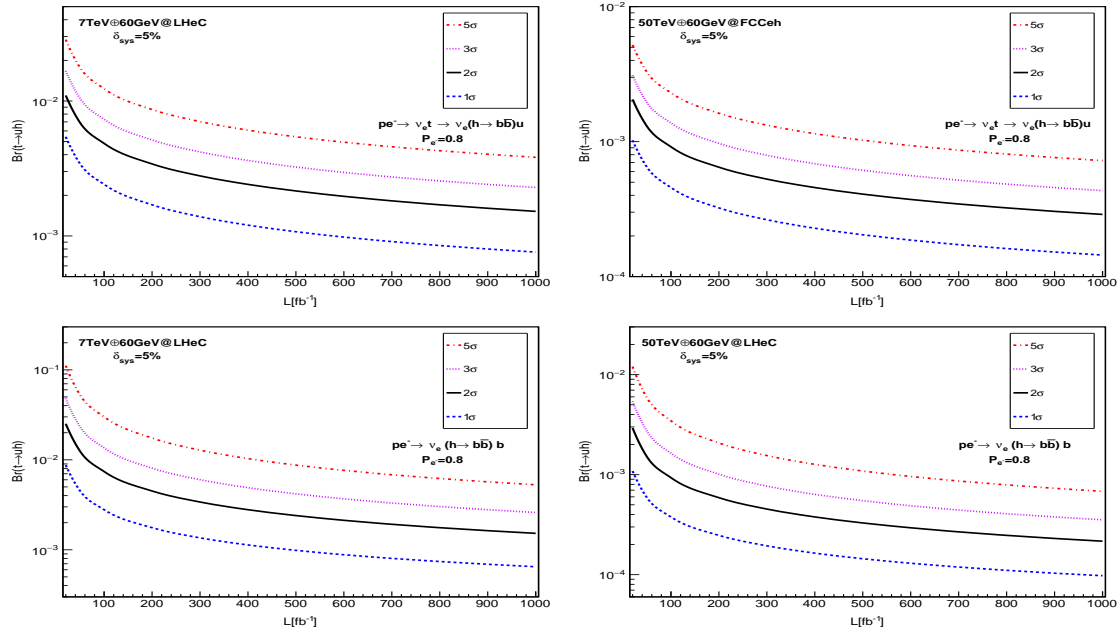


Figure 3: The upper limit on $\text{Br}(t \rightarrow uh)$ at 99.99, 99.73, 95.40, 68.27% C.L. as a function of the integrated luminosity at the LHeC(FCC-eh).

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