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Abstract

Cerium-doped Lutetium-Yttrium Oxyorthosilicate (LYSO:Ce) is one of the most widely used Cerium-doped Lutetium based scintillation crystals. Initially developed for medical detectors it rapidly became attractive for High Energy Particle Physics (HEP) applications, especially in the frame of high luminosity particle colliders. In this paper, a comprehensive and systematic study of LYSO:Ce $([Lu_{(1-x)}Y_x]_2SiO_5:Ce)$ crystals is presented. It involves for the first time a large number of crystal samples (180) of the same size from a dozen of producers. The study consists of a comparative characterization of LYSO:Ce crystal products available on the market by mechanical, optical and scintillation measurements and aims specifically, to investigate key parameters of timing applications for HEP.

Keywords: Scintillating crystals, LYSO, timing detectors, ...

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1. Introduction

 Cerium-doped Lutetium-Yttrium Oxyorthosilicate, commonly known as LYSO:Ce, is one of the most widely used Cerium-doped Lutetium based scin- tillation crystals. Initially developed for medical applications [1, 2], in par- ticular for Positron Emission Tomography (PET), its characteristics in terms of high mass density (twice the density of NaI(Tl)), fast scintillation kinetics $7 \left(6 \text{ times faster decay time than BGO}\right)$ and high light yield (40000 ph/MeV) attracted also the interest of the High Energy Physics (HEP) community.

 In the last decade, LYSO:Ce was employed to prototype and realize high precision electromagnetic calorimeters such as the one designed for the Mu2e experiment [3] and the CCALT forward calorimeter of the KLOE-2 experi-ment [4].

 More recently, a new crystal R&D effort driven by the requirement for high time resolution of second generation PET (Time of Flight PET) further improved the performance of LYSO leading to the industrial production of faster crystals (decay time < 40 ns) and with higher light yield than in the 17 past [5, 6]. The latter, together with the excellent resistance to γ radiation [7], neutrons [8] and charged hadrons [9], makes LYSO appealing for tim- ing applications in the harsh environment of future high-luminosity particle colliders. Here, the high rate of simultaneous interactions per bunch cross- $_{21}$ ing (*pileup*) will produce spatial overlap of tracks and energy deposits. This will affect the capability to disentangle physics events through the tradi- tional detector layers. A picosecond timing layer dedicated to time of arrival measurement of charged particles can help to associate tracks to the correct vertex, mitigating the pileup effect. In this context, the CMS experiment at the Large Hadron collider (LHC) chose LYSO:Ce crystals coupled to Silicon Photomultipliers (SiPMs) to design the sensor unit for the barrel part (BTL) of its timing layer, the MIP Timing Detector [10]. With this layout, the BTL will be able to provide precision timing of minimum ionizing particles with a resolution of 30 - 60 ps [11] restoring the event reconstruction performance of the pre high luminosity era.

 In this paper, a comparative and systematic study of LYSO:Ce crystal properties is carried out, for the first time for a wide number of crystal samples and crystal manufacturers and with particular attention to the key features responsible for the timing performance of the crystals. The aim is to offer a comprehensive review of the state of the art of LYSO:Ce crystal products currently available on the market and identify the best producers for the BTL project. The performance of LYSO:Ce crystals have been eval- uated using bare crystal samples (without wrapping) studying the following properties:

- mass density and correlation with Yttrium content;
- optical transmission characteristics and evaluation of the Ce^{3+} relative concentration;
- photoluminescence characteristics;
- light output and decay time;
- light yield and decay time temperature dependency (low temperature range);
- 48 γ radiation resistance.

2. Samples

 LYSO:Ce crystal samples from 12 manufacturers were studied and com- pared. A list of the manufacturers is provided below in alphabetic order. Each one is randomly associated with a number from 1 to 12 which identifies the producer's crystals throughout this work. Therefore the id number does not match with the order in the following list.

- Crystal Photonics, USA
- EPIC Crystal, China
- Hamamatsu Photonics, Japan
- Hypercrystal NSYSU, Taiwan
- Saint-Gobain, France
- Shanghai EBO Optoelectronics, China
- Shanghai Institute of Ceramics, China
- Simcrystals Technology, China

• SIPAT, China

• Suzhou JT Crystal Technology, China

- Tianle Photonics, China
- Zecotek Imaging System, Singapore

67 The LYSO:Ce $([Lu_{(1-x)}Y_x]_2SiO_5:Ce)$ crystals analyzed have a variable Yt- trium and Cerium content depending on the manufacturer. Both are related to fundamental properties of the crystals. The Yttrium content correlates with the mass density and consequently with the MIP deposited energy, while $_{71}$ the Cerium content is related to light yield and decay time. Dedicated mea- σ surements were performed to determine Yttrium and Ce^{3+} concentrations and are described in the following paragraphs.

2.1. Sample description

 The LYSO:Ce samples studied are 57 mm long crystal bars. The section is rectangular with 3 different thicknesses. The nominal dimensions are re- π ported in Table 1 for the 3 geometries. All manufacturers provided 15 crystal bars, 5 for each geometry and all cut from the same ingot. Samples were pro-⁷⁹ vided with an optical surface quality of $Ra < 15$ nm for all six faces. Most of the crystal properties were measured for all the samples of a manufacturer. When the set of crystals analyzed is smaller or with different characteristics, $\frac{1}{82}$ it is reported in detail. Fig. 1 shows an example of a crystal bar sample (left) and a cross-sectional view of the 3 different available geometries (right).

geometry	bar dimension (mm)			$\#$ of samples	
type	w			per producer	
	3.12	3.75	57.00		
	3.12	3.00	57.00		
	3.12		57.00		

Table 1: Nominal dimensions of the crystal bars. The bar width, thickness and length are labeled respectively as w, t, and L.

Figure 1: Example of a crystal bar sample (left). Cross-sectional view of the 3 types of bars used in this study (right).

2.2. Density of LYSO:Ce samples

 The density of the LYSO:Ce bars is calculated based on the measurements of dimensions and mass. The crystal density is an indicator of the Yttrium percentage in the crystal composition, as shown later in this section, and is directly related to the amount of deposited energy by a Minimum Ionizing Particle (MIP) crossing the crystal.

Dimensions measurement

 μ A high performance measurement system Mitutoyo LH-600 (Fig. 2(a)) was used to measure the three dimensions of the crystal bars. The digital 93 resolution of the instrument is $1 \mu m$ and the observed reproducibility of the 94 measurement is 2-3 μ m. The measurements were carried out on a flat granite ₉₅ table in a temperature controlled environment at T~ 20[°]C (±1[°]C).

 ϵ_{96} For each single bar, width (w) and thickness (t) are defined as the average σ of 16 measurements in different positions along the crystal axis (Fig. 2(b)), ⁹⁸ while the length (L) as the average of 8 measurements made in the 4×2 go corners of the ends of the bars, as shown in Fig. $2(c)$. The black 3D printed holder, with holes of different transverse size, was used to support vertically the bar without any pressure on it and to avoid accidental falls. In Fig. 3 (left), L is shown for all the crystal elements of producer 9. The data points and the error bars correspond to the average and the standard deviation of the 8 measurements performed to determine L respectively. With the purpose to add information related to the compliance with the dimension specifications of the samples provided by each producer, the results are given showing the

Figure 2: Mitutoyo LH-600 measurement system to measure the crystal bar dimensions (a); measurement of w and $t(b)$ and L dimension of the crystal bar with the Mitutoyo probe (c); Mettler Toledo XP105 balance for the mass measurements of the crystal bars (d).

Figure 3: Measured L for all crystal samples of producer 9. The red line corresponds to the nominal crystal length (left). Measured offsets with respect to the nominal dimensions of Table 1. The mean and the standard deviation of bar dimension offsets are shown for each producer. L data point for producer 9 corresponds to the average of the L values shown in the left plot subtracted by the nominal L value (right).

 offsets defined as the difference between the measured values and the nominal values reported in Table 1. Moreover, to characterize the uniformity of the samples, the offset values are averaged over the 15 crystals of each producer and the errors bars are the related standard deviations (Fig. 3, right). All the producers show a good mastering of the cutting technology. For almost 112 all of them, the standard deviation of each dimension is within $5 \mu m$ while 113 the tolerance with respect to the nominal dimensions is within 30 μ m.

Mass measurements

 The mass measurement of the crystal bars was performed with the high- precision Mettler Toledo XP105 balance (0.1 mg digital resolution). The reproducibility of the measurements is better than 0.5 mg; the balance is ¹¹⁸ provided with a glass enclosure for protection against drafts (Fig. 2(d)). The measurements were carried out in a temperature controlled environment at 120 T \sim 20 $^{\circ}$ C (\pm 1 $^{\circ}$ C). The mean values and the standard deviation of the mea- sured bar masses are given in Table 2 for each crystal geometry and for all the producers.

Producer	Type 1	Type 2	Type 3
	(mg)	(mg)	(mg)
1	4713 ± 8	3749 ± 11	2995 ± 3
$\overline{2}$	4760 ± 2	3805 ± 1	3043 ± 3
3	4795 ± 3	3833 ± 11	3060 ± 1
4	4622 ± 6	3730 ± 3	2943 ± 5
5	4721 ± 2	3778 ± 4	3017 ± 3
6	4765 ± 4	3800 ± 13	3014 ± 14
7	4906 ± 14	3921 ± 9	3148 ± 4
8	4782 ± 7	3816 ± 13	3055 ± 6
9	4738 ± 4	3777 ± 15	3041 ± 4
10	4935 ± 11	3938 ± 11	3169 ± 9
11	4734 ± 6	3771 ± 7	3024 ± 4
12	4765 ± 36	3839 ± 8	3079 ± 5

Table 2: LYSO bar mass per crystal geometry and producer. The mean and the standard deviation of the mass values are reported for each geometry.

Density measurements

 The density value is calculated by dividing the mass of the bar by its volume as calculated from the measured dimensions. The density uncertainty is obtained by the corresponding uncertainties on dimensions and mass (the latter being negligible). Results are summarized in Fig. 4 (left) where the $\frac{1}{228}$ density, between 7.0 and 7.4 g/cm³, is shown as the mean over the 15 crystals of the same producer. The error bar corresponds to the relative standard 130 deviation (standard deviation over the mean) which is well below 1% for all the producers.

Yttrium fraction with ICP-MS measurements and density correlation

 The chemical formula of the Cerium-doped LYSO crystals of this study ¹³⁴ is $[Lu_{(1-x)}Y_x]_2SiO_5:Ce$. The stoichiometry of $([Lu_{(1-x)}Y_x])$ group is not fixed and depends on the crystal growth recipe of each manufacturer (expected ¹³⁶ values for x are below 10 %). The large difference between the atomic mass of Lutetium (174.967 amu) and Yttrium (88.906 amu) leads to significant differences between the densities of LYSO crystals having different Yttrium content. The Yttrium molar fraction for at least one crystal bar from each producer was measured by Inductively Coupled Plasma Mass Spectrometry (ICP-MS), at the Gran Sasso National Laboratory (LNGS, Aquila, Italy). For one of the producers, a set of 6 crystals were measured in order to check the consistency of the measurement within the same producer. In total, 31 crystal bars were measured by the ICP-MS technique.

 The results showing the Yttrium content and its linear correlation with the measured mass density are reported in Fig. 4 (right). Measurements from all the crystals of the subsample analyzed are shown and correspond to a data point. The linear correlation of the Yttrium fraction of a crystal bar with its density is clearly demonstrated and the linear regression coefficient ¹⁵⁰ is $R = 0.95$. In addition, a linear fit with χ^2 minimization has been applied to the data. The linear fit parameters correspond, within the error, to the empirical linear relation of the Yttrium content and the density of the crystal ¹⁵³ which can be determined by the densities of pure LSO $(x = 0, \text{ density}$ ¹⁵⁴ 7.4 g/cm³) and pure YSO ($x = 1$, density = 4.5 g/cm³) crystals.

Figure 4: The mean density of the 15 bars is shown for each producer. The relative standard deviation of the bar density is well below 1% for all producers (left). Crystal density as a function of the Yttrium molar fraction (x) for a subsample of the crystal bars analyzed in this study. The linear correlation is clearly visible (right).

3. Optical properties

3.1. Transmission

 Transmission spectra were measured along the three directions of the ¹⁵⁸ crystal samples, one longitudinal (L) and two transversal along width (w) and thickness (t) . The measurements were performed at room temperature in the range 300 - 800 nm. Fig. 5 gives an example of transmission spectra measured for one crystal in all three directions. The figure shows on one side the reproducibility of the transmission measurement and on the other hand it illustrates the nature of the transmission threshold in the UV region. The transmission threshold is caused by the Cerium doping and not by the fundamental absorption of LYSO. Undoped LYSO crystals are indeed trans- parent in a wider range, with the fundamental absorption at 200 nm at room $_{167}$ temperature [12].

 The transverse dimensions of the samples (2 - 3 mm) did not allow for the study of the region below 300 nm due to saturating absorption on color cen- ters induced by dopants (mainly Ce). Double beam spectrophotometers were used: P.E. Lambda 950 at CERN and UV–Vis–NIR CARY 5000 (Varian, Ag-ilent Technologies Deutschland GmbH) at NIMP Bucharest.

 Although the sample dimensions were unsuitable (too thick) for a detailed analysis of the optical absorption, the spectra measured in the transverse

Figure 5: Optical transmission spectra measured in all three directions: longitudinal (black dots), and transversal along w (yellow dots and violet triangles) and t (green dots and blue triangles). Two measurements for each transverse direction are displayed.

¹⁷⁵ directions (w and t) allowed for the visualization of the $5d¹$ absorption band of Ce^{3+} and implicitly the evaluation of the relative Cerium concentration in the measured crystals. To this purpose, the absorbance spectra in the region of interest (ROI) from 440 nm down to 300 nm (2.8-4.1 eV) were obtained from the transmission spectra (Fig. 5, zoom) and fitted using a function which takes into consideration the main absorption centers acting in that ROI. In the considered ROI, the absorbance is found to be proportional to 182 the absorption coefficient (α) and the sample transverse size (d) :

$$
A \sim \alpha \cdot d \tag{1}
$$

¹⁸³ Details about how Eq. 1 was analytically obtained by the transmission ex-¹⁸⁴ pression are provided in Appendix A.

¹⁸⁵ The absorption coefficient can be decomposed into the sum of the con-186 tributions from different absorption centers j , each one described by an ab-187 sorption coefficient α_j , which is proportional to the concentration N_j of the ¹⁸⁸ respective absorption center:

$$
\alpha = \sum_{j} \alpha_{j}
$$

where: $\alpha_{j} = k_{j} \cdot N_{j}$ (2)

 $\frac{189}{189}$ In the defined ROI, α can be written as:

$$
\alpha = \alpha_{Ce^{3+}} + \alpha_{other} \tag{3}
$$

where $\alpha_{Ce^{3+}}$ represents the contribution of Ce^{3+} absorption centers while $191 \alpha_{other}$ takes into account the contribution of all the other absorption centers. The parameter $\alpha_{Ce^{3+}}$ is described by a Gaussian function of the energy ¹⁹³ The amplitude of the Gaussian function is proportional to the concentration ¹⁹⁴ of Cerium in the sample. The absorption due to all the other absorbing 195 centers (α_{other}) can be described by an empirical exponential function, similar ¹⁹⁶ to that applied in the Urbach approximation [13, 14]:

¹⁹⁷ The ratio between the amplitude of the Gaussian function and the sample 198 width can be used for a relative estimation of the concentration of Ce^{3+} centers in the sample $(N_{Ce^{3+}})$. The fit function is effective for all the spectra, ²⁰⁰ regardless of the Cerium doping and possible co-doping used by different ²⁰¹ crystal producers, as illustrated in Fig. 6.

Figure 6: Absorbance spectra for different crystals with the applied fit (blue line). The contributes to the fitting function coming from Ce^{3+} (red dotted line) and all the other (green dotted line) absorbing centers are also shown.

 Transmission spectra were measured for 39 crystals from different pro- ducers with at least two crystals from each producer. For producer 4, 5 and 6, samples from different ingots and with different declared Cerium concentration were studied. The corresponding $N_{Ce^{3+}}$ value are reported in Tab. 3. 206 A total of 23 crystals were measured in both transversal directions, w and t , and often more than one measurement was taken for a given direction, thus

Prod.	Sample $#$	$N_{Ce^{3+}}$	Lab.	Prod.	Sample $#$	$N_{Ce^{3+}}$	Lab.
1	1	1.7540	CERN	6	3	0.8835	NIMP
1	$\overline{2}$	1.4990	CERN	$\boldsymbol{6}$	$\overline{4}$	0.5733	NIMP
1	3	1.2230	NIMP	$\boldsymbol{6}$	5	0.4932	NIMP
1	4	1.2450	NIMP	7	1	0.5195	CERN
$\overline{2}$	$\mathbf{1}$	2.1010	CERN	$\overline{7}$	$\overline{2}$	0.5799	CERN
$\overline{2}$	$\overline{2}$	1.4590	CERN	7	3	0.5386	NIMP
$\overline{2}$	3	1.5520	NIMP	8	1	0.8030	CERN
3	1	0.3244	CERN	8	$\overline{2}$	0.5434	CERN
3	$\overline{2}$	0.3231	CERN	8	3	0.4948	NIMP
3	3	0.3240	NIMP	8	$\overline{4}$	0.5140	NIMP
4	$\mathbf{1}$	1.9800	CERN	$\overline{9}$	$\mathbf 1$	0.9132	CERN
4	$\overline{2}$	1.2480	NIMP	$\overline{9}$	$\overline{2}$	1.0730	CERN
4	3	1.5990	CERN	9	3	0.6914	NIMP
$\overline{4}$	$\overline{4}$	0.6741	NIMP	9	$\overline{4}$	0.7214	NIMP
$\overline{5}$	$\mathbf{1}$	0.3481	CERN	10	1	0.4885	NIMP
$\overline{5}$	$\overline{2}$	0.2560	NIMP	11	1	1.0490	NIMP
$\overline{5}$	3	0.3779	CERN	11	$\overline{2}$	0.8990	NIMP
$\overline{5}$	4	0.4304	NIMP	12	1	0.8548	NIMP
6	1	1.2850	CERN	12	$\overline{2}$	0.9264	NIMP
6	$\overline{2}$	1.1040	CERN				

Table 3: Ce^{3+} relative concentration $(N_{Ce^{3+}})$ reported per crystal sample. The uncertainty of the $N_{Ce^{3+}}$ corresponds to the stability of the fit procedure (6%). In the last column of the table, information about the laboratory in which the measurement was performed is also given.

 having a total of 75 optical transmission spectra analyzed. This was made in order to check both the reproducibility of the transmission spectrum measurement and the overall stability of the $(N_{Ce^{3+}})$ measurement procedure.

 The reproducibility of the transmission spectrum measurement was eval-212 uated repeating the measurement of the same kind of spectrum (along w or $_2$ ¹³ t) several times and it was found to be within 1%. The overall measure- ment process stability, depending on the reliability of the fit function, was ²¹⁵ evaluated at the level of 6% using the $N_{Ce^{3+}w,t}$ values obtained for crystals for which both the transverse spectra were available. In particular, it corre- ϵ_{217} sponds to the standard deviation of the distribution of $N_{Ce^{3+}w,t}$ divided by the corresponding average value over the two transverse spectra $\langle N_{Ce^{3+}} \rangle$.

 $N_{Ce^{3+}}$, as calculated from the absorption spectra, is expected to be cor- related with the light yield and the scintillation kinetics expressing the char- acteristic decay time of the crystals. However, these parameters depend on many other factors that may alter their direct correlation with the concen- tration of Cerium in the crystal. Possible correlations between the relative concentration of Ce^{3+} absorbing centers, $N_{Ce^{3+}}$, calculated from the absorp- tion spectra and scintillation parameters have been studied and the results are discussed in Sec. 4.3.

3.2. Photoluminescence

 Photoluminescence (PL) measurements were performed for crystals of different producers using an Edimburgh Instruments FS5 Spectrofluorom- eter at ENEA Casaccia R.C.(Calliope facility lab) in the excitation range 240 - 390 nm and emission range 370 - 550 nm. For the topics of interest in this article, only the emission spectra recorded in the range 370 - 550 nm by 233 exciting the crystals with $\lambda_{ex} = 358 \text{ nm}$ are reported. All the measurements were performed with 2 nm steps. The emission spectrum measurement re-235 producibility was found to be 1% . The emission spectra for crystals from different producers, normalized to the maximum intensity value, are given in Fig. 7.

Figure 7: Emission spectra for different crystals of different producers ($\lambda_{ex} = 358 \text{ nm}$).

Figure 8: Crystal #66, emission spectrum weighted for the transmittance (blue dots) and emission spectrum (red line).

 The emission spectra have the same characteristic shape with two peaks at 420 nm and 396 nm for all the crystals but the ratio of the two peaks is quite different from one producer to another.

 In Tab. 4, the relative intensity defined as the ratio of the intensity of ²⁴² the two peaks I_{420}/I_{396} for each producer is reported. Crystals that exhibit a ²⁴³ higher I_{420}/I_{396} ratio have intrinsically a better light collection efficiency (for the same optical quality of the crystal surface and bulk purity). This is due to the smaller presence of the self-absorption mechanism in correspondence of the 420 nm peak with respect to the 396 nm peak region [15, 16].

²⁴⁷ As example, the emission spectrum weighted for the transmittance is ²⁴⁸ shown in Fig. 8 for the same crystal $#66$. The resulting spectrum provides the information necessary to optimize the coupling of the crystals with the light detection sensor.

4. Scintillation properties

252 The light output (LO) and the decay time (τ) of the crystal samples from each producer were measured with dedicated setup and methods at the the INFN - Sezione di Roma and Sapienza University laboratory (Roma, Italy). The results are shown as the average values over the 15 samples of

Producer	I_{420}/I_{396}
1	0.98 ± 0.01
2	0.98 ± 0.01
3	0.97 ± 0.01
4	$0.93 + 0.01$
5	0.88 ± 0.01
6	0.94 ± 0.01
7	0.97 ± 0.01
8	0.96 ± 0.01
9	0.85 ± 0.01
10	0.98 ± 0.01
11	0.89 ± 0.01
12	0.86 ± 0.01

Table 4: Relative emission intensity defined as I_{420}/I_{396} for the crystals studied in the present work.

 each producer. Details about the reproducibility of the measurements are provided.

 LO and τ are key parameters for LYSO: Ce crystal timing applications. 259 The highest possible LO in the shortest possible time frame leads to the best timing performance for which a figure of merit can be defined as the ratio LO/ τ . Results for the figure of merit are also shown for all the producers.

 F_{262} Finally, the dependency of LO and τ on the relative Ce^{3+} concentration has been investigated in Sec. 4.3 with the aim to explore the possibility to 264 use Ce^{3+} concentration as a quality indicator of the scintillation and timing performance of the crystals.

4.1. Experimental setup, methods and tools

Setup description

 The experimental setup used for the measurement of the scintillation properties is shown in Fig. 9. It consists of a 51 mm diameter end window PMT (ET Enterprised model 9256B) placed inside a cylindrical box with a rectangular frame. The frame works as a guide to insert the bar holder which keeps the crystal bar vertical on the PMT photocathode window and is equipped with different transverse section holes for the housing of the 3 bar types. The crystal bars are inserted into the holder without any wrapping.

Figure 9: Experimental setup used to measure the single crystal bar LO and τ , showing a crystal placed in the plastic holder over the PMT window. The optical fiber coupled to a blue LED and the support for the ²²Na source are visible too.

 One crystal end face is in contact with the PMT window while the other one is free and in contact with air. No grease is applied to enhance the PMT-crystal optical contact. This precaution was taken to optimize the reproducibility of the measurement. The setup is enclosed in a black painted box whose temperature is kept stable at 20° C (within 0.1-0.2 $^{\circ}$ C over 24 h) by the use of a chiller. The PMT signal is readout by the DRS4 evaluation board [17], working at a sampling rate of 2 GS/s; this allows an integration window for the PMT signal extending up to 500 ns. The single photoelectron (SPE) response is calibrated using a pulsed, fast, blue LED. The LED light is brought inside the box using an optical fiber.

Light output measurement

286 The absolute LO measurement is performed using one of the annihila- $_{287}$ tion photons emitted by a 22 Na radioactive source placed beside the bar and evaluating the position of the 511 keV photoelectron peak in the crystal signal. The charge of the photoelectron peak is then divided by the SPE charge and by the energy of the photon to obtain the LO value expressed in photoelectrons per MeV of deposited energy.

SPE. The SPE charge value is extracted by fitting the charge spectra

 obtained with the LED with the convolution of a Poisson (accounting for gamma conversion process and first dynode photoelectron collection) and a Gauss distribution (accounting for multiplicative dynode system response), as shown in Fig. 10, top. In order to improve the fit stability, 5 charge spectra obtained with increasing LED pulses of different intensities are collected and simultaneously fitted leaving the SPE charge as common free parameter.

 The PMT signal acquisition is triggered by the coincidence signal pro- vided by the LED driver and the charge is integrated in a 30 ns window after the baseline subtraction.

 511 keV photo-peak. The charge associated to the 511 keV photo-peak is obtained using a 17 parameter fit which fully describes the energy deposit of $_{304}$ both the 511 keV and the 1275 keV photons emitted by 22 Na, including the contributions due to Compton, photo-electric and back-scatter interactions. A turn-on function is also used to describe the trigger behavior.

 In this case the PMT signal acquisition is triggered on the PMT sig- nal itself using an optimal threshold. The charge is integrated in a 450 ns time window after the baseline subtraction. An example of charge spectra used to extract the 511 keV photo-peak values is presented together with the corresponding fitting functions in Fig. 10, bottom.

Decay time measurement

 The acquisition with a fast sampling digitizer allows the extraction of the $_{314}$ scintillation τ directly from the acquired waveform of the PMT signal. An average over all PMT signals with an associated total charge above roughly $_{316}$ 100 keV in the 22 Na runs is performed. The average waveform is passed through a Butterworth filter with a cut-off frequency of 20 MHz to reduce oscillations due the imperfect impedance matching between the PMT anode 319 output and the DRS4 buffer input. τ is extracted from a fit which includes a single exponential decay function and a Gaussian turn-on. An example of this fit is shown in Fig. 11. From the average waveform it is also possible to estimate the amount of light emitted in a time window smaller than 450 ns, integrating the waveform in different time windows.

 The reproducibility of the LO and τ measurements was estimated repeat- ing them daily over one month using a reference crystal and it was found to be 4% and better than 1% , respectively.

Figure 10: Example of charge spectrum used to extract the charge associated to the SPE. The contributes to the fitting function (orange line) from the pedestal (black dotted line), 1 (red line), 2 (green line), 3 (blue line), 4 (yellow line) and 5 (magenta line) photoelectrons are visible (top). Example of a charge spectrum obtained with the 22 Na radioactive source. The total charge fitting function (red line) used to extract the 511 keV photo-peak and the background events fitting function (green line) are shown (bottom).

³²⁷ 4.2. Measurement results

 328 The LO and τ measurement results are averaged over the 15 crystals 329 provided by each producer and are displayed in Fig. 12. The LO (Fig. 12,

Figure 11: Average waveform from a 22 Na run with the superimposed fit (red line) performed on filtered data (black dots) to estimate the τ of the crystal.

 top) is expressed in photons/MeV and represents the number of scintillation photons produced per MeV of energy deposit which impinge on the photosen- sor and are successfully detected. It is corrected for the quantum efficiency of the sensor and corresponds to the intrinsic crystal light yield (LY) times the light collection efficiency (LCE). The latter depends on the optical sur- face quality of the crystal and the transparency of the bulk as well as the crystal-sensor coupling (which is, however, the same for all the crystals). The quantum efficiency correction factor is obtained by the quantum efficiency of the PMT, as provided by the producer, weighted over the LYSO spectrum 339 and corresponds to about 25% .

³⁴⁰ The relative standard deviation of the LO values for different producers is about 8%. The LO standard deviation (error bars in Fig. 12, top) for ³⁴² samples of the same producer is mostly comparable with the reproducibility 343 of the measurement (4%) , although some show higher values revealing a less ³⁴⁴ uniform LO among the provided samples. The standard deviation value of producer 1 can be explained by 2 outlier crystals.

 The τ value ranges from 45 down to 38 ns for the slowest to the fastest crystal, as illustrated in Fig. 12, bottom. The relative standard deviation of $\frac{348}{100}$ the τ values for different producers is about 5% while the relative standard

Figure 12: LO (top) and τ (bottom) results for the 12 producers.

Figure 13: Figure of merit for the timing performance of the LYSO:Ce crystals defined as the ratio of LO over τ .

 349 deviation for crystals from the same producer is around 1% and thus com-³⁵⁰ parable with the reproducibility of the measurement. Finally, Fig. 13 shows $_{351}$ the figure of merit LO/τ for the timing performance for each producer. The ³⁵² relative standard deviation of the values from different producers is within $353 \quad 5\,\%$.

354 4.3. Study of the main scintillation parameters as a function of Ce^{3+} relative ³⁵⁵ concentration

 F_{356} Figures 14 and 15 show τ and LO as a function of the Ce^{3+} relative concentration calculated as described in Sec. 3.1. The linear dependence 358 between τ and LO and the calculated Ce^{3+} relative concentration is too weak to recommend the use of the latter for an indirect assessment of the first two parameters i.e. the scintillation performance of crystals. In particular, 361 while the τ trend with respect to Ce^{3+} relative concentration is close to the expectation of linear correlation, for the LO the linear dependency hypothesis is weaker.

 The scintillation performance of LYSO:Ce crystals depends on several factors, not only the intensity of the optical absorption peak at 360 nm, which ³⁶⁶ represents the standard indicator for the concentration of Ce^{3+} absorbing centers (i.e. concentration of Cerium used as scintillation activator). First of

Figure 14: τ as a function of the relative concentration of Ce^{3+} absorbing centers.

Figure 15: LO as a function of the relative concentration of Ce^{3+} absorbing centers. The LO value is normalized to the LO of a reference crystal.

³⁶⁸ all, the intensity of the absorption peak at 360 nm reflects only the content of Ce^{3+} [18, 19] while the content of Ce^{4+} , which has an important contribution σ to the scintillation LO, remains unknown. The ratio between the Ce^{3+} and Ce^{4+} concentrations depends on possible co-doping applied by each crystal producer and also on unintentional impurities and defects.

 As mentioned above, the samples are too thick for measuring the broad $_{374}$ band in the UV region which may possibly give a hint on the Ce^{3+}/Ce^{4+} concentrations ratio. Furthermore, the light yield depends on the competi- tion between radiative and non-radiative recombination. This competition might be strongly affected by co-doping and unintentional impurities due to different raw materials used by different crystal producers.

 On the basis of these arguments, the results in Fig. 14 and 15 can be 380 explained as a milder sensitivity of τ to the presence of other dopants and $_{381}$ impurities or defects with respect to the one exhibited by the LO. Co-doping, defects and impurities depend indeed by the specific LYSO:Ce recipe and growing process chosen by each manufacturer.

5. γ radiation hardness

 \mathcal{L}_{385} Radiation hardness of the crystal samples against ionizing radiation by γ rays was studied at the Calliope facility of ENEA-Casaccia Research Centre (Rome, Italy). Calliope is a pool-type facility equipped with a 60° Co radio- isotopic source array in a large volume shielded cell [20]. The irradiation tests involved at least one crystal bar of type 2 for each producer. All the samples were irradiated at the same dose rate of 9 kGy/h and received a total integrated absorbed dose of 50 kGy. The dose rate value is experimen- tally measured by an alanine-ESR dosimetric system mapping the Calliope 393 irradiation area. The dose rate uncertainty is 5% .

5.1. Scintillation properties

 All the irradiated samples were measured before and after the irradiation with the setup described in Sec. 4.1.

 After irradiation, all the crystals exhibited phosphorescence light with an approximate decay time of 2 - 3 h as estimated from the presence of a transient noise in the baseline of the PMT signal acquired ∼ every hour for 12 h and displayed in Fig. 16. For this reason, the samples were measured again at $_{401}$ least 16 h after the irradiation to evaluate the ratio of the LO and the τ after and before irradiation. The results are shown in Fig. 17. The average light 403 output loss amounts to 9% with a relative standard deviation of 3% among the different producers (Fig. 17, top). The scintillation τ (Fig. 17, bottom) after irradiation remains unchanged within the measurement uncertainties

Figure 16: Impact of the LYSO phosphorescence light on the standard deviation of the PMT signal baseline as a function of the time from the end of the irradiation. The standard deviation of the PMT signal charge is calculated in a 20 ns time window before the scintillation signal and averaged over the events of a source run.

⁴⁰⁶ compared to the pre-irradiation value for most of the producers. The average $\frac{407}{407}$ ratio of τ after and before the irradiation is 1% with a standard deviation of 408 2% .

⁴⁰⁹ In general, the scintillation mechanism of LYSO:Ce is not damaged by 410 γ-ray irradiation [21]. The LO decrease depends on the γ-induced trans-⁴¹¹ parency loss which is due to the creation of absorbing centers. The LO can 412 be further recovered through a air annealing of the crystal at \sim 300°C for ⁴¹³ some hours. Slow (few days) spontaneous recovery can also be observed at ⁴¹⁴ room temperature [18].

Figure 17: Ratio of LO (top) and τ (bottom) after and before γ irradiation for the 12 producers. The error bars are determined by propagation of the measurement uncertainties corresponding in this case to the reproducibility of the LO and τ measurements.

6. Scintillation properties at low temperature

 Due to its radiation hardness against photons and hadrons, LYSO:Ce can be employed for timing purposes in the harsh environment of the new gen-eration particle colliders such as the HL-LHC. Here, to mitigate the impact of the radiation damage on the performance of the detector components, es- pecially the silicon ones, the operating temperature is usually lowered below 421 O°C by some tens of degrees. This will be, for example, the case of the barrel part of the timing detector of CMS-phase II. In BTL, LYSO crystals are coupled to Silicon PhotoMultipliers (SiPM). Radiation exposure increases the noise due to the SiPM dark count rate and lowers the LO of the crys- tals deteriorating the time resolution. For this reason the detector will be $_{426}$ operated at low temperature, between -45° C and -35° C.

⁴²⁷ With the aim to extend and complete the set of information collected in 428 this paper, additional measurements of LO and τ in this range of temper- atures for crystal bars from each of the 12 producers were performed. The experimental setup and the results are presented in this section.

6.1. Experimental setup

 The experimental bench used for cold measurements of LYSO:Ce crystal features the same concept of the PMT bench used in the crystal characteri- $_{434}$ zation campaign at room temperature $(20^{\circ}C)$ and described in Sec. 4. Also ⁴³⁵ the methods and the analysis tools to obtain the values of LO (expressed in 436 photoelectrons per MeV of deposited energy) and τ are the same. The LO value is corrected for the temperature dependency of the PMT gain using the charge of the SPE measured at the same temperature with the LED. The same LED runs have been used to exclude a non-negligible dependency of the PMT quantum efficiency (QE) on temperature. This was obtained verifying that the average number of photoelectrons in a LED run (LED 442 intensity set to give an average number of photoelectrons \simeq 1) remains con- stant with the temperature. To reach and stabilize the temperature down $_{444}$ to -30° C, the setup was enclosed into a thermostatic chamber (Angelantoni TY110) and equipped with a temperature monitor. The temperature fluctu- ations during a standard data taking have been measured and found to be $_{447}$ \pm 0.2°C. In a preliminary study, the response of the PMT used (Hamamatsu 448 R7378) was measured and proved to be linear down to -30° C. The PMT signal is brought outside the chamber through a circular feedthrough and readout by a 12 bit 3.2 GS/s digitizer (CAEN DT5743). The reproducibility $_{451}$ of the LO and the τ measurement was evaluated repeating several times the corresponding measurements using a reference crystal and it was found to be 453 2% and $\langle 1\% \rangle$, respectively. The better performance in term of LO measure-ment reproducibility of this test bench with respect to the one used for the ⁴⁵⁵ measurements described in Sec. 4.1 is probably due to the better tempera-⁴⁵⁶ ture stabilization provided by the high performance thermostatic chamber in which the setup was enclosed.

Figure 18: (left) A picture of the experimental bench used for the characterization of single crystal bars at cold temperatures. The bench is inserted in a thermostatic chamber able to provide stable temperatures down to −40◦C. (right) A detailed picture of the setup components.

457

⁴⁵⁸ 6.2. Results

⁴⁵⁹ At least one crystal bar of the smallest geometry for each of the 12 pro-⁴⁶⁰ ducers was measured. Six measurement points have been acquired with tem- $_{461}$ peratures ranging from 20 $^{\circ}$ C down to -30° C. Lowering the temperature, 462 both the LO and τ increase slowly. In Fig. 19 (top) an example of LO as ⁴⁶³ a function of the temperature and normalized to the corresponding value at $_{464}$ T= 20° C is shown.

⁴⁶⁵ The LO is linear with the temperature for all producers. The tempera-⁴⁶⁶ ture coefficient is on average -0.15% /°C ranging between -0.28% /°C and $_{467}$ -0.08 %/°C as shown in Fig. 19 (bottom). The LO relative variation as a ⁴⁶⁸ function of the temperature is equal to the light yield (LY) relative variation 469 because the LO can be factorized as $LY \times LCE \times QE$ and the LCE and the ⁴⁷⁰ QE can be assumed constant with the temperature and therefore cancel out ⁴⁷¹ in the ratio.

 τ dependency on the temperature is linear down to -30° C only for 6 473 producers over 12 (regression coefficient R > 0.85) and in general the variation

Figure 19: (top) LY normalized to the corresponding value at $T=20\degree C$ as a function of the temperature. The normalized LY error bars were determined by propagation of the measurement uncertainties. From the linear fit, the LY temperature coefficient is obtained. (bottom) Light yield temperature coefficient for the 12 producers. The error bars correspond to the fit uncertainties.

 474 with temperature is smaller than for the LO. In Fig. 20 (top) the linear 475 dependency of τ for producer 5 is shown as an example. For the other 476 producers, no linear relation between the temperature and τ can be assumed

Figure 20: (top) τ normalized to the corresponding value at T=20[°]C as a function of the temperature. The normalized τ error bars correspond to the reproducibility of the τ measurement. For this producer, τ has a linear behavior and from a linear fit the τ temperature coefficient can be obtained. (bottom) τ normalized to the corresponding value at $T=20^{\circ}$ C as a function of the temperature for producer 2, 4 and 7. For this producers, τ does not feature a linear dependency on the temperature.

 477 (R<0.75). In Fig. 20 (bottom), τ vs. T is shown for crystals from this 478 subset of producers; in particular for producer 2 (R=0.41), 4 (R=0.73) and ⁴⁷⁹ 7 (R=0.76). For these producers, additional measurement points at low

- 480 temperature would be needed for a more rigorous description of τ dependency $_{481}$ down to -30° C.
- I_{482} In Fig. 21 the ratio of the figure of merit (LO/τ) measured at -30°C and
- 483 at 20° C is also shown. Its average value and standard deviation are 1.05 and
- 484 0.02 respectively. For all producers the ratio is >1 . This demonstrates that
- ⁴⁸⁵ lowering the operating temperature of the crystal can help to improve their timing performance.

Figure 21: Ratio of the figure of merit (here expressed as LY/τ) for timing performance measured at −30◦C and at 20◦C for the 12 producers. The error bars were determined by propagation of the measurement uncertainties.

486

7. Discussion

⁴⁸⁸ A set of 15 small crystal bars $(3 \, mm \times 3 \, mm \times 57 \, mm)$ from 12 different producers were studied and compared with respect to a set of properties and performance fundamental for HEP applications with a special focus on timing applications.

 All producers are shown to have mastered the cutting technology pro- ducing samples with uniform dimensions at the level of per mille, and within the requested specifications at a level better than 1% . From the dimensions and the mass measurement, the crystal density value was derived for every 496 sample. It ranges from 7.1 to 7.4 g/cm^3 and its relative standard deviation $_{497}$ among the samples of the same producer is well below 1%.

 The mass density study is complemented, for at least one crystal per pro- ducer, by inductively coupled plasma mass spectrometry (ICP-MS) measure- ments from which the Yttrium molar fraction was evaluated. The Yttrium fraction is indeed expected to linearly correlate with the mass density. The expectation has been confirmed by data $(R = 0.95)$ and the spread of the Yttrium fraction among the different producers is about 30 %.

 Optical transmission spectra and photoluminescence properties were also studied for all producers. In particular, the evaluation of the relative con- $\frac{1}{506}$ centration of the main crystal luminescence center (Ce^{3+}) was obtained from the transmission spectra. Its correlation with the light output (LO) and $\frac{1}{508}$ decay time (τ) of the crystals has been investigated in the attempt to estab- lish a method to characterize the timing performance of the crystals. The data do not match the expectations showing a poor linear correlation of the $_{511}$ (Ce^{3+}) relative concentration with both scintillation parameters. This has been mainly ascribed to the possible presence of different co-dopants, im- purities and defects which may have an important role in the scintillation dynamics.

 LO and τ were measured for all the crystal samples, together with the $_{516}$ figure of merit for timing application defined as LO/τ . all producers' sam- ples show similar scintillation properties. The spread of the LO value for $_{518}$ different producers is at the level of 8% while for τ , ranging from 38 to 45 ns, it is within 5 %. The uniformity of the crystal samples provided by each producer with respect to these scintillation parameters is comparable with ϵ_{521} the reproducibility of the measurements: 4% for the LO and 1% for τ .

 $\frac{522}{10}$ In order to test the radiation hardness of the crystal samples against γ , LO and τ were also measured after irradiation with 50 kGy at a dose rate of

 $524 \text{ } 9 \text{ kGy/h}$ for a subsample of crystals from all 12 producers. While τ remains essentially unchanged for all producers, the LO loss is on average at the level of 10 %. The study did not include a thermal annealing campaign. Nevertheless it is a well established concept that the LO damage is not permanent and it can be fully recovered by thermal annealing.

 $\frac{529}{20}$ Finally, the LO and τ dependency on temperature was analyzed for a sso subsample of crystals down to -30° C. The LO exhibits a linear dependency μ_{531} on temperature with a temperature coefficient ranging between $-0.28\,\%$ / °C μ ₅₃₂ and -0.08% /°C. Only 6 producer over 12 shows a linear τ dependency on $_{533}$ the temperature down to -30° C. More data points at low temperature would $_{534}$ be needed to study the non-linearity of τ for the other producers.

Nevertheless, the figure of merit at -30° C compared with the results ob- $_{536}$ tained at 20° C shows that lowering the operating temperature of the crystals ⁵³⁷ can help to improve their timing performance. This holds true for all the 538 producers and with an relative standard deviation of $\simeq 2\%$.

 The most important crystal features measured in this study are summa- rized in Tab. 5 and Tab. 6 for each producer. All producers showed similar ⁵⁴¹ characteristics within \simeq 10 %, except for the Ce^{3+} relative concentration and the LY temperature coefficient. For these crystal properties the spread among the producers is at the level of 50 %. Despite this, their impact on the key performance for HEP and especially for timing application is limited. The Ce^{3+} relative concentration has shown a poor correlation with LO and τ while the spread in the LY temperature coefficients does not reflect in the 547 figure of merit LY/τ .

Crystal	M. density	$N_{Ce^{3+}}$	LO	τ	LO/τ
Prod.	(g/cm^3)	(a.u)	(ph./MeV)	(ns)	$(\text{ph/MeV}\,\text{ns})$
	7.088 ± 0.020	1.430 ± 0.216	5164 ± 580	45.13 ± 1.43	115 ± 15
$\overline{2}$	7.093 ± 0.008	1.704 ± 0.283	5261 ± 186	45.03 ± 0.49	117 ± 4
3	7.250 ± 0.005	0.324 ± 0.001	4708 ± 156	39.99 ± 0.63	118 ± 4
$\overline{4}$	7.137 ± 0.006	1.609 ± 0.299	4688 ± 273	41.60 ± 0.52	113 ± 7
$\overline{5}$	7.103 ± 0.008	0.327 ± 0.052	4847 ± 169	41.21 ± 0.93	118 ± 5
6	7.109 ± 0.011	1.091 ± 0.164	4216 ± 116	38.02 ± 0.41	111 ± 3
$\overline{7}$	7.313 ± 0.009	0.546 ± 0.025	5381 ± 190	42.05 ± 0.31	128 ± 5
8	7.175 ± 0.008	0.589 ± 0.125	4662 ± 313	39.76 ± 0.30	117 ± 8
9	7.078 ± 0.016	0.850 ± 0.155	4852 ± 141	40.30 ± 0.18	120 ± 3
10	7.334 ± 0.009	0.488 ± 0.029	5274 ± 89	42.50 ± 0.31	124 ± 2
11	7.116 ± 0.006	0.974 ± 0.075	4740 ± 116	40.96 ± 0.30	116 ± 2
12	7.110 ± 0.008	0.891 ± 0.036	5061 ± 111	42.29 ± 0.31	120 ± 2

Table 5: Average and standard deviation values for Mass density, $N_{Ce^{3+}}$, LO , τ and LO/τ measured for the crystals from each producers.

Crystal	LO_{irr}/LO	τ_{irr}/τ	dLY/dT	$T_{-30^{\circ}C}$	'LO (τ) $-30^{\circ}C$
Prod.			$(\%LY/°C)$	(ns)	$(LO/\tau)_{20^{\circ}C}$
1	0.893 ± 0.050	1.009 ± 0.019	-0.079 ± 0.017	46.28 ± 1.12	1.034 ± 0.012
$\overline{2}$	0.868 ± 0.049	0.984 ± 0.016	-0.141 ± 0.017	45.83 ± 0.11	1.065 ± 0.011
3	0.900 ± 0.051	1.031 ± 0.018	-0.131 ± 0.018	43.86 ± 0.11	1.025 ± 0.011
4	0.891 ± 0.050	0.992 ± 0.015	-0.131 ± 0.018	42.18 ± 0.10	1.054 ± 0.011
5	0.897 ± 0.051	0.987 ± 0.017	-0.097 ± 0.017	42.29 ± 0.11	1.026 ± 0.011
6	0.861 ± 0.049	1.002 ± 0.019	-0.279 ± 0.018	38.67 ± 0.10	1.090 ± 0.012
7	0.933 ± 0.053	0.991 ± 0.018	-0.094 ± 0.017	43.35 ± 0.11	1.038 ± 0.011
8	0.954 ± 0.054	1.000 ± 0.019	-0.277 ± 0.020	42.44 ± 0.11	1.093 ± 0.013
9	0.866 ± 0.049	0.985 ± 0.018	-0.179 ± 0.018	41.59 ± 0.10	1.075 ± 0.011
10	0.897 ± 0.051	0.981 ± 0.017	-0.127 ± 0.018	43.49 ± 0.11	1.060 ± 0.011
11	0.978 ± 0.055	0.992 ± 0.018	-0.197 ± 0.017	42.23 ± 0.11	1.061 ± 0.011
12	0.897 ± 0.051	0.961 ± 0.017	-0.063 ± 0.018	43.44 ± 0.11	1.027 ± 0.011

Table 6: Summary of the crystal scintillation properties measured after γ irradiation and at low temperatures down to $-30^{\circ}\mathrm{C}$ for at least a crystal per producer.

8. Conclusions

549 A comprehensive and systematic study of LYSO:Ce $([Lu_{(1-x)}Y_x]_2SiO_5:Ce)$ crystals is presented. It involves for the first time a large number of crystal samples (180) of the same size from several producers. The study consists of a comparative characterization of LYSO:Ce crystal products available on the market and aims, in particular, to investigate key parameters of timing applications for HEP.

 λ set of 15 small crystal bars $(3\,mm \times 3\,mm \times 57\,mm)$ for each producer were measured with respect to mechanical, optical and scintillation proper- ties. The latter were studied before and after the irradiation of the crystals 558 with a 50 kGy integrated dose of γ -ray and at temperatures down to -30° C. The timing performance of the crystals was evaluated by a figure of merit $_{560}$ defined as LO/τ . Finally, the number of the samples provided by each pro- ducer allowed for the study of the uniformity of the crystal properties within a producer batch.

 The LYSO:Ce products considered in this study fully qualify for tim- ing applications at future HEP colliders. LYSO:Ce crystals of all producers show in general similar properties and an excellent uniformity of the samples. The spread of the crystal characteristics with a direct impact on the timing $_{567}$ performance is within 10% among the different producers.

 This review of LYSO:Ce crystals does not identify a single producer or a set of producers with globally superior performance. The detected differ- ences in the crystal products, although limited, could however be used to guide the selection process of the LYSO:Ce crystals best suited for a specific application.

Appendix A. Appendix A: Absorbance analytical expression in the approximation of multiple reflection between parallel crystal faces

The absorbance is defined as:

$$
A = 2 - \log_{10} T(\%) \tag{A.1}
$$

 where T corresponds, in the present study, to the measured optical transmis-sion (transmittance).

 The transmittance is defined as the ratio I/I_0 of light intensities at the 580 exit (I) and the entrance (I_0) of the measured sample. When accounting for

Figure A.22: Sketch of the multiple reflection of the light between parallel crystal faces at a distance d. The analytical expressions for the contributions of the reflected, absorbed and transmitted light are also reported along the light path.

⁵⁸¹ multiple reflections on the crystal faces, the numerator is given by the sum 582 of the I_j contributions exiting the crystal and displayed in Fig. A.22:

$$
T = \frac{I}{I_0} = \lim_{n \to \infty} \frac{\sum_{j=1}^n I_j}{I_0} = \lim_{n \to \infty} (1 - R)^2 e^{-\alpha d} [1 + R^2 e^{-2\alpha d} + R^4 e^{-4\alpha d} + ... R^{2n} e^{-2n\alpha d}]
$$
\n(A.2)

 583 where R and α are the reflection and the absorption coefficients, respectively, \mathfrak{so}_4 and d is the sample transverse size $(w \text{ or } t)$.

⁵⁸⁵ Since the term in square bracket corresponds to a geometric progression 586 with common ratio $R^2e^{-2\alpha d}$, T can be written as:

$$
T = (1 - R)^{2} e^{-\alpha d} \lim_{n \to \infty} \frac{1 - (R^{2} e^{-2\alpha d})^{n}}{1 - R^{2} e^{-2\alpha d}}
$$
(A.3)

 587 Considering the value of α and the value of the refraction index of LYSO σ_{188} ($n_r \simeq 1.7$) in the analyzed ROI as reported in [22] and the expression of R ⁵⁸⁹ at normal incidence:

$$
R = \left(\frac{n-1}{n+1}\right)^2 \tag{A.4}
$$

⁵⁹⁰ $R^2e^{-2\alpha d}$ < 1 holds true and T converge to:

$$
T = (1 - R)^{2} e^{-\alpha d} \frac{1}{1 - R^{2} e^{-2\alpha d}} \simeq (1 - R)^{2} e^{-\alpha d}
$$
 (A.5)

 $_{591}$ Using Eq. A.5 in Eq. A.1 and R being constant in the considered ROI, A is found to be proportional to d:

$$
A \sim \alpha \cdot d \tag{A.6}
$$

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