

Searches for Higgs boson pair production with the f ull LHC Run 2 dataset in ATLAS

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10th International Conference on New Frontiers in Physics (ICNFP 2021) 24/08/21

Why study Higgs pair production?

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$$
L = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu}
$$

+ $i \overline{\psi} \psi + 4.c.$
+ $\overline{\psi} : y_{ij} \psi + 4.c.$
+ $\overline{\psi} \psi^2 - V(\phi)$

Why study Higgs pair production?

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BSM resonances predicted to decay to HH...

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BSM physics that can alter non-resonant HH production...

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Modifications to the Higgs self-coupling

> **BSM physics that can alter non-resonant HH production...**

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Several others...

Modifications to the Higgs self-coupling

New couplings of Higgs to SM particles

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> **New particles contributing in loops**

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Modifications to the Higgs self-coupling

New couplings of Higgs to SM particles

BSM physics that can alter non-resonant HH production...

> **New particles contributing in**

loops

Rich theoretical basis to search for BSM effects in the HH production!

14 **So far, no such BSM effects have been observed experimentally.** 24/08/21

HH Production at the LHC

 Total HH production σ = 34.4 fb[†], which is ~1000 times smaller than $\sigma_{\sf H}^{\sf H}$!

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[†]for √s = 13 TeV with m_H = 125 GeV 15

ATLAS HH Analyses

HH decay branching ratios HH analyses in ATLAS focus primarily on setting \lim its on σ _{HH}, couplings involving HH and BSM **resonance models.**

There are several HH analyses targeting different decay channels → overall aim: combine results to maximise sensitivity.

Today, will cover the latest results from:

[HH → bb](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-016/)γγ [HH → bb](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-030/)ττ [HH → bbbb](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-035/)

All use the full ATLAS Run 2 (2015-18) dataset (ℒ **= 139 fb-1) at √s = 13 TeV.**

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Publication: [ATLAS-CONF-2021-016](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-016/) Physics briefing:<https://atlas.cern/updates/briefing/twice-higgs-twice-challenge>

HH → bbγγ

Non-resonant: search for HH production for κ_{λ} hypotheses in range -10 to 10 (SM = 1). **Resonant: search for a narrow-width scalar resonance with mass between 251-1000 GeV.**

Low HH BR at ~0.3%, but clean signal due to H→γγ.

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Main bkg is yy+jets, with some ttyy, fakes and single H

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Non-Resonant Results

No significant excesses over background observed \rightarrow **set limits¹ on** $\sigma_{_{\rm HH}}$ **and** $\kappa_{_{\lambda}}.$

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No significant excesses over background observed → set upper limits¹ on σ for narrow width scalar resonances at the various mass hypotheses tested.

For 251 \leq m_x \leq 1000 GeV, obs' (exp') **varies between 610–47 fb (360–43 fb)**

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1at the 95% CL_s. 11

²**HH→bbγγ, ℒ = 36fb⁻¹.** 29

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2 to 3-fold improvement on previous analysis² depending on m_χ

~⅓ due to analysis improvements!

24/08/21 1at the 95% CL_s ²HH→bbγγ, ℒ = 36fb⁻¹ . 30

Publication: [ATLAS-CONF-2021-030](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-030/) Physics briefing:<https://atlas.cern/updates/briefing/two-Higgs-better-one>

HH → bbττ

Non-resonant: search for SM non-resonant HH production.

Resonant: search for a narrow-width scalar resonance with mass in range 251 to 1600 GeV.

Balance of moderate HH BR of ~7.3% and relatively clean ττ signature.

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τs can decay hadronically or leptonically... e, μ, c W^{-} $\bar{\nu_e}, \bar{\nu_\mu}, \bar{u}$

Non-resonant: search for SM non-resonant HH production.

Resonant: search for a narrow-width scalar resonance with mass in range 251 to 1600 GeV.

Balance of moderate HH BR of ~7.3% and relatively clean ττ signature.

Main background is top, Z+jets and multijet.

Modelled true τ bkg using MC. Modelled fake contribution using data.

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Modelled true τ bkg using MC. Modelled fake contribution using data.

Used suite of MVAs to separate sig from bkg

Analysis Overview

Non-Resonant Results

No significant excesses over background observed → set upper limits¹ on σ_{HH}.

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No significant excesses over background observed → set upper limits¹ on σ_{HH}.

4-fold improvement on previous SM limits² .

~2-fold due to analysis improvements e.g. MVA event selection, improved fake estimate, object reco/calibration.

Dominant uncertainty is statistical. Largest systematic is background modelling.

Resonant Results

At mX = 1 TeV, a small excess of local significance 3.0σ and a global significance 2.0 (+0.4, −0.2) σ was observed.

Publication: [ATLAS-CONF-2021-035](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2021-035/) Physics briefing:<https://atlas.cern/updates/briefing/double-Higgs-to-bottoms>

$HH \rightarrow bbbb$

Analysis Overview

Resonant: search for narrow-width scalar resonance and spin-2 Kaluza-Klein Graviton† resonance with mass in range 251-3000 GeV.

Largest HH BR (~34%) but dominant multijet background → data-driven background estimate.

24/08/21 **tfrom the Randall-Sundrum model where k /** $\overline{\mathsf{M}}_{\mathsf{pl}}$ **= 1 42** $\overline{\mathsf{M}}$

Resolved Analysis Overview

Resolved Analysis Overview

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Resolved Analysis Overview

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Boosted Analysis Overview

Boosted Analysis Overview

Boosted Analysis Overview

$HH \rightarrow bbbb$ Results

No significant excesses above background observed → set upper limits† on σ for two models at various mass hypotheses.

Results Summary

Non-Resonant SM HH production SCALAR Scalar resonance

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Looking Forward...

~21 x as much as full Run 2 dataset!

Projections for High Luminosity LHC (HL-LHC) with 3ab-1 of data based on 36fb-1 HH analyses (not those shown today!):

Looking Forward...

~21 x as much as full Run 2 dataset!

Projections for High Luminosity LHC (HL-LHC) with 3ab-1 of data based on 36fb-1 HH analyses (not those shown today!):

Future is bright for HH searches in ATLAS in the near future and beyond!

Back Up Slides

Primary text

Side note

HH → bbγγ

Figure 4: Reconstructed four-body mass for m_X = 300 GeV and m_X = 500 GeV resonant signal benchmarks and for the $\gamma\gamma$ +jets background. Dashed lines represent the distribution of $m_{b\bar{b}\gamma\gamma}$ while solid lines represent the distribution of $m^*_{b\bar{b}\gamma\gamma}$, defined in Section 4.2.1. Distributions are normalized to unit area.

(a) $ggFHH$ production mode

(b) VBF HH production mode

Figure 5: The $m^*_{b\bar{b}\gamma\gamma}$ distributions after the common preselection for (a) non-resonant ggF HH and (b) VBF HH signals with several K_{λ} values. $m_{b\bar{b}\gamma\gamma}^{*} = 350$ GeV is chosen as the separating boundary between categories targeting the SM and BSM κ_{λ} signals.

Table 2: Variables used in the BDT for the non-resonant analysis. The b -tag status identifies the highest fixed b -tag working point (60%, 70%, 77%) that the jet passes. All vectors in the event are rotated so that the leading photon ϕ is equal to zero.

 \cdot

Table 3: Definition of the categories used in the HH non-resonant search. Before entering the BDT-based categories, events are required to satisfy the common preselection.

Figure 6: The BDT distribution of the di-Higgs ggF signal for two different values of k_{λ} and the main backgrounds in the (a) low and (b) high mass region. Distributions are normalized to unit area. The dotted lines denote the category boundaries. Events with a BDT score below 0.881 in the low mass region or below 0.857 in the high mass region are discarded.

Variable	Definition
Photon-related kinematic variables	
$p_{\rm T}^{\gamma\gamma}$, $y^{\gamma\gamma}$	Transverse momentum and rapidity of the di-photon system
$\Delta\phi_{\gamma\gamma}$ and $\Delta R_{\gamma\gamma}$	Azimuthal angular distance and ΔR between the two photons
Jet-related kinematic variables	
$m_{h\bar{b}}$, $p_{\rm T}^{bb}$ and $y_{b\bar{b}}$	Invariant mass, transverse momentum and rapidity of the <i>b</i> -tagged jets system
$\Delta\phi_{b\bar{b}}$ and $\Delta R_{b\bar{b}}$	Azimuthal angular distance and ΔR between the two <i>b</i> -tagged jets
N_{jets} and $N_{b-\text{jets}}$	Number of jets and number of b -tagged jets
$H_{\rm T}$	Scalar sum of the p_T of the jets in the event
Photons and jets-related kinematic variables	
$m_{b\bar{b}\gamma\gamma}$	Invariant mass built with the di-photon and b-tagged jets system
$\Delta y_{\gamma\gamma,b\bar{b}}, \Delta \phi_{\gamma\gamma,b\bar{b}}$ and $\Delta R_{\gamma\gamma,b\bar{b}}$	Distance in rapidity, azimuthal angle and ΔR between the di-photon and the b -tagged jets system

Table 4: Variables used in the BDT for the resonant analysis. For variables depending on b-tagged jets, only jets b-tagged using the 77% working point are considered as described in Section 4.1.

Figure 7: The BDT score for the benchmark signals ((a) $m_X = 300$ GeV and (b) $m_X = 500$ GeV) and for the main backgrounds. Distributions are normalized to unit area. The dotted lines denote the event selection thresholds. Events with a BDT score below 0.85 for $m_X = 300$ GeV or below 0.75 for $m_X = 500$ GeV are discarded.

Figure 8: Distributions of (a) $m_{\gamma\gamma}$ and (b) $m_{b\bar{b}\gamma\gamma}^*$ for events passing the common preselection criteria. The continuum background is scaled by the $\gamma\gamma$, γ -jet or jet- γ , and di-jet fractions and normalized to the data sideband.

Figure 9: Distributions of $m_{\gamma\gamma}$ in all signal categories for the non-resonant HH search: (a) high mass BDT tight, (b) high mass BDT loose, (c) low mass BDT tight, (d) low mass BDT loose. The continuum background is scaled by the $\gamma\gamma$, γ -jet, and di-jet fractions and normalized to the data sideband.

Figure 10: Distributions of $m_{\gamma\gamma}$ for the selections used for the resonance mass points (a) $m_X = 300$ GeV and (b) m_X = 500 GeV for the resonant search. The non-resonant background is scaled by the $\gamma\gamma$, γ -jet, and di-jet fractions and normalized to the data sideband. The scalar resonance signal is scaled to an arbitrary cross section value.

Likelihood function for bbγγ:

$$
\mathcal{L} = \prod_c \left(\text{Pois}(n_c | N_c(\boldsymbol{\theta})) \cdot \prod_{i=1}^{n_c} f_c(m_{\gamma\gamma}^i, \boldsymbol{\theta}) \cdot G(\boldsymbol{\theta}) \right)
$$

$$
N_c(\boldsymbol{\theta}) = \mu \cdot N_{HH,c}(\boldsymbol{\theta}_{HH}^{\text{yield}}) + N_{\text{bkg},c}^{\text{res}}(\boldsymbol{\theta}_{\text{res}}^{\text{yield}}) + N_{SS,c} \cdot \boldsymbol{\theta}^{\text{SS},c} + N_{\text{bkg},c}^{\text{non-res}}
$$

$$
f_c(m_{\gamma\gamma}, \theta) = [\mu \cdot N_{HH,c}(\theta_{HH}^{\text{yield}}) \cdot f_{HH,c}(m_{\gamma\gamma}, \theta_{HH}^{\text{shape}}) + N_{bkg,c}^{\text{res}}(\theta_{\text{res}}^{\text{yield}}) \cdot f_{bkg,c}^{\text{res}}(m_{\gamma\gamma}, \theta_{\text{res}}^{\text{shape}})
$$

+ $N_{SS,c} \cdot \theta_{HH}^{SS,c} \cdot f_{HH,c}(m_{\gamma\gamma}, \theta_{HH}^{\text{shape}}) + N_{bkg,c}^{\text{non-res}} \cdot f_{bkg,c}^{\text{non-res}}(m_{\gamma\gamma}, \theta_{\text{non-res}}^{\text{shape}})]/N_c(\theta_{\text{non-res}}^{\text{yield}})$

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Figure 11: Data are compared to the background-only fit for the four categories of the non-resonant search. Both the continuum background and the background from single Higgs boson production are considered.

Table 6: Expected and observed numbers of events in the categories of the non-resonant search. An additional requirement of 120 GeV $\langle m_{\gamma\gamma} \rangle$ < 130 GeV is applied. The uncertainties on the continuum background are those arising from the fitting procedure. The uncertainties on the single Higgs boson and Higgs boson pair productions are from MC statistical error.

Table 8: Breakdown of the dominant systematic uncertainties. The impact of the uncertainties is defined according to the statistical analysis described in Section 7. It corresponds to the variation on the upper limit on the signal strength when re-evaluating the profile likelihood ratio after fixing the nuisance parameter in question to its best-fit value increased or decreased by one standard deviation, while all remaining nuisance parameters remain free to float. The impact is shown in $\%$. Only systematic uncertainties with an impact of at least 0.5% are shown. Uncertainties of Norm. + Shape type have effects on both the normalization and the parameters of the functional form, the rest of uncertainties affects only the yields.

Figure 12: Observed and expected limits at 95% CL on the cross section of non-resonant Higgs boson pair production as a function of the Higgs boson self-coupling modifier $\kappa_{\lambda} = \lambda_{HHH}/\lambda_{HHH}^{SM}$. The constraints on κ_{λ} are obtained over an expected hypothesis excluding $pp \rightarrow HH$ production. The $\pm 1\sigma$ and $\pm 2\sigma$ variations about the expected limit due to statistical and systematic uncertainties are also shown. The theory prediction curve represents the scenario where all parameters and couplings are set to their SM values except for κ_{λ} . The uncertainty band of the theory prediction curve shows the cross section uncertainty.

Table 7: Expected and observed numbers of events of the resonant HH search. An additional requirement of 120 GeV $\langle m_{\gamma\gamma} \rangle$ < 130 GeV is applied. The event numbers quoted for the scalar resonance signal assume an arbitrary total production cross section $\sigma(pp \to X \to HH)$ equal to the observed exclusion limits of Figure 14. The uncertainties on the continuum background are those arising from the fitting procedure. The uncertainties on the single Higgs boson, Higgs boson pair and scalar resonance production are from the MC statistical error.

Figure 13: Data are compared to the background-only fit for the resonant search for the (a) $m_X = 300$ GeV and (b) m_X = 500 GeV mass hypotheses. The continuum background, as well as the background from single Higgs boson production and from the SM HH production are considered.

Figure 14: Observed and expected limits at 95% CL on the production cross section of a narrow width scalar resonance X as a function of the mass m_X of the hypothetical scalar particle. The black solid line represents the observed upper limits. The dashed line represents the expected upper limits. The $\pm 1\sigma$ and $\pm 2\sigma$ variations about the expected limit due to statistical and systematic uncertainties are also shown.

HH → bbττ

Table 2: Summary of the event selections, shown separately in the different trigger categories. In cases where pairs of reconstructed objects of the same type are required, thresholds on the (sub-)leading p_T object are given outside (within) parentheses. When the selection depends on the year of data-taking, the possible values of the requirements are separated by commas, except for the jet selection in the lepton-plus- τ_{had} -vis trigger and di- τ_{had} -vis triggers, which use multiple selection criteria as described in Section 5.1. The trigger p_T thresholds shown correspond to the offline requirements.

Figure 2: Acceptance times efficiency for the full analysis selection as a function of the resonance mass m_X in the τ_{had} , τ_{lep} τ_{had} single-lepton trigger and τ_{lep} , τ_{had} lepton-plus- $\tau_{\text{had-vis}}$ trigger categories, shown in solid line with square markers, dashed and dotted lines, respectively. The solid line with circle markers is the acceptance times efficiency curve for the combined $\tau_{\rm lep} \tau_{\rm had}$ category. The acceptance times efficiency is evaluated for $X \to HH \to b\bar{b}\tau^+\tau^$ decays, with respect to the targeted τ -lepton decay modes ($\tau_{\rm lep} \tau_{\rm had}$ or $\tau_{\rm had} \tau_{\rm had}$). James Grundy 24/08/21 77

Table 3: Variables used as inputs to the MVAs in the three analysis categories. The same choice of input variables is used for the resonant and non-resonant production modes. The variables are defined in the main text.

 (a)

 (b)

 (c)

 (d)

 (e)

 (f)

Figure 3: Signal (solid lines), post-fit background (filled histograms) and data (dots with error bars) distributions of m_{HH} (top), $m_{\tau\tau}^{MMC}$ (middle row) and m_{bb} (bottom) for events in the $\tau_{had}\tau_{had}$ (left), $\tau_{lep}\tau_{had}$ single-lepton trigger (middle column) and $\tau_{\rm lep} \tau_{\rm had}$ lepton-plus- $\tau_{\rm had\text{-}vis}$ trigger (right) categories. The normalisation and shape of the backgrounds and the uncertainty on the total background shown are determined from the likelihood fit to data in the non-resonant HH search. The expected non-resonant signal is overlaid with its normalisation scaled by a factor of 100, and the m_X = 500 GeV and m_X = 1000 GeV resonant signals are overlaid in the m_{HH} distributions with their cross-section set to 1 pb. The dashed histogram shows the total pre-fit background. The size of the combined statistical and systematic uncertainty of the background is indicated by the hatched band. The ratio of the data to the sum of the backgrounds is shown in the lower panels.

Figure 4: Schematic depiction of the combined fake-factor method used to estimate multi-jet and $t\bar{t}$ backgrounds with fake- $\tau_{\text{had-vis}}$ in the $\tau_{\text{lep}}\tau_{\text{had}}$ channel. Backgrounds which are not from events with fake- $\tau_{\text{had-vis}}$ originating from jets are estimated from simulation and are subtracted from data in all control regions. Events in which an electron or a muon is misidentified as a $\tau_{\text{had-vis}}$ are also subtracted, but their contribution is very small. Both sources are indicated by "True- $\tau_{\text{had-vis}}$ subtracted" in the legend.

Figure 5: Schematic depiction of the combined fake-factor method to estimate the multi-jet background with fake- $\tau_{\text{had-vis}}$ in the $\tau_{\text{had}}\tau_{\text{had}}$ channel. Backgrounds with true- $\tau_{\text{had-vis}}$ that are not from multi-jet events are simulated and subtracted from data in all the control regions. This is indicated by "Non-multi-jet subtracted" in the legend.

Table 4: Breakdown of the relative contributions to the uncertainty in the extracted signal cross-sections, as determined in the likelihood fit to data. These are obtained by fixing the relevant nuisance parameters in the likelihood fit, and subtracting the obtained uncertainty on the fitted signal cross-sections in quadrature from the total uncertainty, and then dividing the result by the total uncertainty. The sum in quadrature of the individual components differs from the total uncertainty due to correlations between the groups of uncertainties.

Figure 7: The MVA output distributions in the search for non-resonant HH signal (top) and in the search for resonant HH signal with $m_X = 500$ GeV (middle row) and $m_X = 1000$ GeV (bottom), in the $\tau_{had}\tau_{had}$ (left), $\tau_{len}\tau_{had}$ single-lepton trigger (middle column) and $\tau_{\text{len}}\tau_{\text{had}}$ lepton-plus- $\tau_{\text{had-vis}}$ trigger (right) categories. The distributions are shown after the fit to the background-only hypothesis. The signal is overlaid and scaled to the combined expected limit. The dashed histogram shows the total pre-fit background. The lower panels show the ratio between data and the total post-fit background, where the hatched band shows the statistical and systematic uncertainties on that background. For visualisation purposes, these histograms are displayed using uniform bin widths instead of the bin edges used in the fit, though the bin contents correspond to those used in the fit.

Table 5: Observed and expected upper limits at 95% CL on the cross-section of non-resonant HH production according to SM-like kinematics, and on the cross-section of non-resonant HH production divided by the SM prediction. The $\pm 1 \sigma$ and $\pm 2 \sigma$ variations around the expected limit are also shown.

Figure 8: Observed and expected limits at 95% CL on the cross-section of the resonant HH production as a function of the scalar resonance mass m_X . The dashed lines show the expected limits while the solid lines show the observed limits. The blue and red lines are the limits for the $\tau_{had}\tau_{had}$ channel and $\tau_{lep}\tau_{had}$ channel, respectively. The black lines are the combined limits of the two channels. The $\pm 1\sigma$ and $\pm 2\sigma$ variations around the expected combined limit are indicated by the turquoise and yellow bands, respectively. The limits are obtained using the profile-likelihood test statistic and the modified frequentist CL_s technique.

$HH \rightarrow bbbb$

$$
X_{HH} = \sqrt{\left(\frac{m(H_1) - 120 \text{ GeV}}{0.1 \times m(H_1)}\right)^2 + \left(\frac{m(H_2) - 110 \text{ GeV}}{0.1 \times m(H_2)}\right)^2}.
$$

$$
R_{HH}^{\text{VR}} \equiv \sqrt{\left(m(H_1) - 1.03 \times 120 \,\text{GeV}\right)^2 + \left(m(H_2) - 1.03 \times 110 \,\text{GeV}\right)^2} < 30 \,\text{GeV}.
$$

$$
R_{HH}^{CR} \equiv \sqrt{\left(m(H_1) - 1.05 \times 120 \,\text{GeV}\right)^2 + \left(m(H_2) - 1.05 \times 110 \,\text{GeV}\right)^2} < 45 \,\text{GeV}.
$$

Figure 3: Cumulative acceptance times efficiency as a function of resonance mass for each event selection step in the resolved channel for (a) the spin-0 and (b) the spin-2 signal models.

Figure 4: Corrected $m(HH)$ distributions for the 2b control region (teal histogram) and 4b control region (dots) in the resolved channel. The statistical uncertainty in the $2b$ control region is represented by the grey band. The error bars on the 4b points represent the Poisson uncertainties corresponding to their event yields. The 2b data are shown (a) before and (b) after the kinematic reweighting procedure. In both cases the $2b$ distributions are normalized to the 4b event yields for a pure shape comparison. The final bin of each distribution includes overflow. The bottom panel shows the difference between the $4b$ and $2b$ distributions, relative to the $2b$ distribution.

Figure 5: Corrected $m(HH)$ distribution in the resolved 4b validation region (dots), compared to the reweighted distribution in 2b validation region (teal histogram). The error bars on the $4b$ points represent the Poisson uncertainties corresponding to their event yields. The final bin includes overflow. The background uncertainty (grey band) is computed by adding all individual sources in quadrature. The bottom panel shows the difference between the 4b and reweighted 2b distributions, relative to the 2b distribution.

Table 2: Resolved 4b signal region data, estimated background, and signal event yields in corrected $m(HH)$ windows containing roughly 90% of each signal, for representative spin-2 mass hypotheses. The signal is normalized to the overall expected limit on its cross-section; its uncertainties are evaluated by adding all individual components in quadrature. The background yields and uncertainties are evaluated after a background-only fit to the data.

Figure 9: Cumulative signal acceptance times efficiency as a function of the resonance mass for various selection steps in the boosted channel. The steps up to the b -tag categorization are shown for (a) the spin-0 and (b) the spin-2 signal models. The efficiencies of the three b -tag categories are shown for (c) the spin-0 and (d) the spin-2 scenarios; this efficiency is obtained after the other selection steps including the SR definition. The signal efficiency in the $4b$ region has a maximum around 1.5 TeV. Above that value the track jets starts to merge together, and for the highest resonance masses the $3b$ and $2b$ categories become the most efficient.

Figure 10: Reconstructed mass distributions of the leading H candidate for the data (dots) and the background model (stacked histograms) in the (a) 2b, (b) 3b and (c) 4b control regions. The error bars on the data points represent the Poisson uncertainties corresponding to their event yields. The statistical uncertainty in the background model is represented by the grey band. This distribution is used to normalize the multijet and $t\bar{t}$ background components. The enhanced event rates at low and high masses are due to the geometry of the CR. The bottom panel shows the difference between the data and the background model, normalized to the background model.