Elastic and inelastic cross section measurements with the ATLAS detector

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Outline:

- Physics motivation.
- Inelastic cross section at $\sqrt{s} = 13$ TeV.
- Elastic, inelastic and total cross sections at $\sqrt{s} = 8$ TeV.
- Future elastic measurements.

Physics motivation

- The elastic (σ_{el}), inelastic (σ_{inel}) and total (σ_{tot}) *pp* cross sections are fundamental quantities which cannot be calculated with perturbative QCD.
- Regge theory provides general description but data is needed to constrain models.
- σ_{tot} gives the upper bound on any *pp* process and is observed to rise with \sqrt{s} .
- A substantial fraction of σ_{inel} is diffractive processes and a measurement of σ_{inel} gives better background determination for high p_T processes.



Inelastic cross section measurement at $\sqrt{s} =$ 13 TeV

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Inelastic cross section at 13 TeV - Strategy

- Measure the σ_{inel} in a fiducial region and extrapolate to full phase-space using simulations.
 - The better detector coverage, the smaller extrapolation uncertainty.
- ATLAS uses the MBTS plastic scintillator discs at z = ±3.6 m covering 2.07 < |η| < 3.86
 - Corresponds to

$$\xi \equiv M_x^2/s > 10^{-6}$$

- 8 inner and 4 outer counters in each disc.
- \sim 99 % efficiency for charged particles.
- The fiducial σ_{inel} is the number of observed events with ≥ 2 hits corrected for background, pile-up, efficiencies and luminosity.
- Use 60 μ b⁻¹ of *pp* collisions at $\sqrt{s} =$ 13 TeV with $\mu \sim$ 0.002



Inelastic cross section at 13 TeV - Tuning models

- The inelastic cross section is the sum of the non-diffractive and the diffractive cross section.
- The ratio $f_D = (\sigma_{SD} + \sigma_{DD})/\sigma_{inel}$ is poorly known and differs between models.
- The fraction of single-sided events, R_{SS}, is related to f_D and used to tune f_D in the models.
- Using the *f*_D-tuned models, the hit multiplicity in the MBTS for the models are compared to data:
 - The DL (Donnachie-Landshoff) pomeron flux model is best.
 - ε is a free parameter in the pomeron Regge trajectory.
 - The EPOS LHC and QGSJET-II models (developed mostly for cosmic-ray showering) are worst.



Inelastic cross section at 13 TeV - Results

• The fiducial inelastic cross section in $\xi > 10^{-6}$ is measured to be

$$\sigma_{ ext{inel}}^{ ext{fid}} = \!\!68.1 \pm \! 0.6_{(ext{exp.})} \pm \! 1.3_{(ext{lum.})}$$
 mb

• This is in good agreement with the Pythia8 DL model.



(inner error bars are without luminosity uncertainty)

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• The fiducial cross section is extrapolated to the full phase space using the Pythia8 DL and MBR models:

$$\sigma_{ ext{inel}} =$$
 78.1 \pm 0.6 $_{ ext{(exp.)}} \pm$ 1.3 $_{ ext{(lum.)}} \pm$ 2.6 $_{ ext{(extr.)}}$ mb

- The inelastic cross section is still increasing with √s.
- This is in agreement with model predictions.
- The CMS result is:

 $\sigma_{inel}=\!71.3\pm3.5~\text{mb}$

which is a 1.5 σ difference.



Elastic, inelastic and total cross sections measurement at $\sqrt{s} = 8$ TeV

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• From the optical theorem we get:

$$\sigma_{\text{tot}}^2 = \frac{16\pi(\hbar c)^2}{1+\rho^2} \frac{1}{L} \frac{dN_{\text{el}}}{dt}\Big|_{t=0} \quad \text{with} \quad \rho \equiv \frac{\text{Re}\left[F_{\text{el}}(t=0)\right]}{\text{Im}\left[F_{\text{el}}(t=0)\right]}$$

• The four-momentum transfer t is calculated as

$$-t=(p\times\theta^*)^2.$$

where the scattering angle θ^* is calculated from the proton trajectories and *p* is the beam momentum.

• Minimum accessible *t*-value is given by

$$-t_{\rm min} \propto \frac{d^2}{\beta^*}$$

- Use 500 μ b⁻¹ with a $\beta^* = 90$ m collision optics and low pile-up ($\mu \approx 0.1$).
 - Standard LHC running uses $\beta^* \sim$ 0.5 m.

Elastic analysis at 8 TeV - The ALFA detector

- Built to measure elastically scattered protons at µrad angles.
- Located 240 m from the ATLAS interaction point inside Roman Pots.
- The main detector is built of scintillating fibers.
 - The fiber width of 500 μm and layer staggering gives \approx 30 μm tracking resolution.
- The scattering angle is reconstructed from the tracks using knowledge about the LHC magnet strengths.





- Events are selected when all four detectors in an arm have a track.
- Momentum conservation provides strong selection cuts.
- Total number of elastic candidates is 3.8 M.
- Background level is 0.12 % estimated from the antigolden topology.
 - Primarily Double Pomeron Exchange (DPE).
- Observed t-spectrum is corrected for acceptance, detector resolution, luminosity, and efficiencies.





 The differential elastic cross section is a superposition of the strong interaction amplitude F_N and the Coulomb amplitude F_C added in quadrature giving

$$\frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}t} \propto \frac{G^4(t)}{|t|^2} + \sigma_{\mathrm{tot}}^2(1+\rho^2) \cdot \exp(-\mathbf{B}|t|) - \frac{\sigma_{\mathrm{tot}}G^2(t)}{|t|} \left[\sin(\phi(t)) + \rho\cos(\phi(t))\right] \cdot \exp\left(\frac{-\mathbf{B}|t|}{2}\right)$$

- The fit range is $0.014 \le |t| \le 0.1 \text{ GeV}^2$.

No sensitivity to:

- The ρ-parameter fixed to 0.1362 from COMPETE predictions.
- The interference phase taken as:

$$\phi(t) = -\ln\left(\frac{B|t|}{2}\right) - \gamma_E$$

• The proton form factor taken as:

$$G(t) = \left(rac{\Lambda}{\Lambda+|t|}
ight)^2 \,, \ \ \Lambda = 0.71 \; {
m GeV}^2$$



• The result for the total cross section is:

$$\sigma_{tot} = 96.07 \pm 0.18_{(stat)} \pm 0.85_{(exp.)} \pm 0.31_{(extr.)} \text{ mb}$$
 .



• The elastic cross section σ_{el} is the integral of the nuclear part:

 $\sigma_{el} = 24.33 \pm 0.04_{(stat.)} \pm 0.39_{(syst.)} \text{ mb}$.

• The inelastic cross section is the difference between the total and the elastic:

 $\sigma_{inel} = 71.73 \pm 0.15_{(stat.)} \pm 0.69_{(syst.)} \text{ mb}$.

• This is about 4 times more precise than the direct measurement at 13 TeV (page 7):

$$\sigma_{\mathsf{inel}}^{\mathsf{13 \ TeV}} =$$
 78.1 \pm 0.6 $_{(\mathsf{exp.})}$ \pm 1.3 $_{(\mathsf{lum.})}$ \pm 2.6 $_{(\mathsf{extr.})}$ mb

The extrapolation uncertainty on σ_{inel} from elastic scattering comes from a simple fit range variation.

- Data in the Coulomb-Nuclear-Interference region at $-t \approx 10^{-3} \text{ GeV}^2$ allows a measurement of the p-parameter.
 - Provides insight to the understanding of elastic scattering.
 - Dispersion relations derived from analyticity relates energy evolution of the σ_{tot} and $\rho.$
 - $\bullet~$ The simultaneous measurement of σ_{tot} and ρ therefore tests a very basic assumption.
 - High energy predictions of σ_{tot} will be possible.

- Data has been collected at $\sqrt{s} = 8$ TeV with $\beta^* = 1$ km optics.
 - The analysis is in review.
- Data has been collected at $\sqrt{s} = 13$ TeV with $\beta^* = 2.5$ km optics.
 - The analysis is ongoing.



Summary

Inelastic cross section at $\sqrt{s} = 13$ TeV:

- The diffractive fraction in several models are tuned to match an observable in data.
- The fiducial cross section agrees well with Pythia8 DL model.
- The full inelastic cross section is in agreement with models and is still observed to rise with $\sqrt{s}.$

Total cross section at $\sqrt{s} = 8$ TeV:

- The differential elastic cross section is measured with tracking detectors in Roman pots.
- The total cross section is inferred using the optical theorem and is still observed to rise with $\sqrt{s}.$
- Future measurements of the ρ-parameter will provide further insight to the non-perturbative regime of QCD.

Backup slides

Backup - Investigations of nuclear slope parametrization

- Different parametrizations for the nuclear slope have been investigated.
- The upper limit of the fit range was increased to $|t| = 0.3 \text{ GeV}^2$ in order to increase the sensitivity of additional parameters.
- The quality of the fit is increased due to the higher number of free parameters.

| | $\sigma_{ m tot}[{ m mb}]$ | Model |
|--------------|----------------------------|---|
| Nominal | 96.07 ± 0.86 | $f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2}$ |
| Ct^2 | 96.16 ± 0.80 | $f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2 - Ct^2/2} $ |
| $c\sqrt{-t}$ | 96.40 ± 0.80 | $f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2 - c/2(\sqrt{4\mu^2 - t} - 2\mu)} , \mu = m_{\pi}$ |
| SVM | 96.16 ± 0.80 | $f_{\rm N}(t) = \rho \frac{\sigma_{\rm tot}}{\hbar c} e^{-B_R t/2} + i \frac{\sigma_{\rm tot}}{\hbar c} e^{-B_I t/2}$ |
| BP | 96.81 ± 0.95 | $f_{\rm el}(t) = i \left[G^2(t) \sqrt{A} e^{-Bt/2} + e^{i\phi} \sqrt{C} e^{-Dt/2} \right]$ |
| BSW | 96.67 ± 0.99 | $\operatorname{Re} f_{\mathrm{el}}(t) = c_1(t_1 + t)e^{-b_1t/2}$, $\operatorname{Im} f_{\mathrm{el}}(t) = c_2(t_2 + t)e^{-b_1t/2}$ |

Backup - Results for the nuclear B-slope

- ATLAS measurement: $B = 19.73 \pm 0.24 \text{ GeV}^{-2}$
- TOTEM measurement: $B = 19.9 \pm 0.3 \text{ GeV}^{-2}$
- Pre-LHC expectations was a linear evolution of the B-slope with ln(s)
- LHC measurements of the *B*-slope favours a second $\ln^2(s)$ term.



(fit not updated with latest ATLAS result)

- **Detector acceptance** is highly dependent on detector distance to the beam and beam divergence.
 - Found from simulation tuned to data.
- t-resolution is influenced by detector resolution and beam divergence.
 - Relative *t*-resolution is better than 10 % and corrected for by unfolding.





- The beam optics has direct influence on the *t*-reconstruction.
- Different *t*-reconstructions gives different results ⇒ the initial **design** optics needs modifications.
- Elastic data is used to constrain an optics fit including magnet strengths whereby an **effective** optics in obtained.

