



Search for exclusive Higgs and Z boson decays to $\phi\gamma$ and $\rho\gamma$ with the ATLAS detector

The ATLAS Collaboration

A search for the exclusive decays of the Higgs and Z bosons to a ϕ or ρ meson and a photon is performed with a pp collision data sample corresponding to an integrated luminosity of up to 35.6 fb^{-1} collected at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector at the CERN Large Hadron Collider. These decays have been suggested as a probe of the Higgs boson couplings to light quarks. No significant excess of events is observed above the background, as expected from the Standard Model. Upper limits at 95% confidence level were obtained on the branching fractions of the Higgs boson decays to $\phi\gamma$ and $\rho\gamma$ of 5.0×10^{-4} and 10.4×10^{-4} , respectively. The corresponding 95% confidence level upper limits for the Z boson decays are 0.7×10^{-6} and 4.0×10^{-6} for $\phi\gamma$ and $\rho\gamma$, respectively.

1 Introduction

Following the observation [1, 2] of a Higgs boson, H , with a mass of approximately 125 GeV [3] by the ATLAS and CMS collaborations at the Large Hadron Collider (LHC), the properties of its interactions with the electroweak gauge bosons have been measured extensively [4–6]. The coupling of the Higgs boson to leptons has been established through the observation of the $H \rightarrow \tau^+\tau^-$ channel [4, 7, 8], while in the quark sector indirect evidence is available for the coupling of the Higgs boson to the top-quark [4] and evidence for the Higgs boson decays into $b\bar{b}$ has been recently presented [9, 10]. Despite this progress, the Higgs boson interaction with the fermions of the first and second generations is still to be confirmed experimentally. In the Standard Model (SM), Higgs boson interactions to fermions are implemented through Yukawa couplings, while a wealth of beyond-the-SM theories predict substantial modifications. Such scenarios include the Minimal Flavour Violation framework [11], the Froggatt–Nielsen mechanism [12], the Higgs-dependent Yukawa couplings model [13], the Randall–Sundrum family of models [14], and the possibility of the Higgs boson being a composite pseudo-Goldstone boson [15]. An overview of relevant models of new physics is provided in Ref. [16].

The rare decays of the Higgs boson into a heavy quarkonium state, J/ψ or $\Upsilon(nS)$ with $n = 1, 2, 3$, and a photon have been suggested for probing the charm- and bottom-quark couplings to the Higgs boson [17–20] and have already been searched for by the ATLAS Collaboration [21], resulting in 95% confidence level (CL) upper limits of 1.5×10^{-3} and $(1.3, 1.9, 1.3) \times 10^{-3}$ on the branching fractions, respectively. The $H \rightarrow J/\psi\gamma$ decay mode has also been searched for by the CMS Collaboration [22], yielding the same upper limit. The corresponding SM predictions for these branching fractions [23] are $\mathcal{B}(H \rightarrow J/\psi\gamma) = (2.95 \pm 0.17) \times 10^{-6}$ and $\mathcal{B}(H \rightarrow \Upsilon(nS)\gamma) = (4.6_{-1.2}^{+1.7}, 2.3_{-1.0}^{+0.8}, 2.1_{-1.1}^{+0.8}) \times 10^{-9}$. The prospects for observing and studying exclusive Higgs boson decays into a meson and a photon with an upgraded High Luminosity LHC [16] or a future hadron collider [24] have also been studied.

Currently, the light (u, d, s) quark couplings to the Higgs boson are loosely constrained by existing data on the total Higgs boson width, while the large multijet background at the LHC inhibits the study of such couplings with inclusive $H \rightarrow q\bar{q}$ decays. Rare exclusive decays of the Higgs boson into a light meson, M , and a photon, γ , have been suggested as a probe of the couplings of the Higgs boson to light quarks and would allow a search for potential deviations from the SM prediction [23, 25, 26]. Specifically, the observation of the Higgs boson decay to a ϕ or $\rho(770)$ (denoted as ρ in the following) meson and a photon would provide sensitivity to its couplings to the strange-quark, and the up- and down-quarks, respectively. The expected SM branching fractions are $\mathcal{B}(H \rightarrow \phi\gamma) = (2.31 \pm 0.11) \times 10^{-6}$ and $\mathcal{B}(H \rightarrow \rho\gamma) = (1.68 \pm 0.08) \times 10^{-5}$ [23]. The decay amplitude receives two main contributions that interfere destructively. The first is referred to as “direct” and proceeds through the $H \rightarrow q\bar{q}$ coupling, where subsequently a photon is emitted before the $q\bar{q}$ hadronises exclusively to M . The second is referred to as “indirect” and proceeds via the $H \rightarrow \gamma\gamma$ coupling followed by the fragmentation $\gamma^* \rightarrow M$. In the SM, owing to the smallness of the light-quark Yukawa couplings, the latter amplitude dominates, despite being loop induced. As a result, the expected branching fraction predominantly arises from the “indirect” process, while the Higgs boson couplings to the light quarks are probed by searching for modifications of this branching fraction due to changes in the “direct” amplitude.

This paper describes a search for Higgs boson decays into the exclusive final states $\phi\gamma$ and $\rho\gamma$. The decay $\phi \rightarrow K^+K^-$ is used to reconstruct the ϕ meson, and the decay $\rho \rightarrow \pi^+\pi^-$ is used to reconstruct the ρ meson. The branching fractions of the respective meson decays are well known and are accounted for when calculating the expected signal yields. The presented search uses approximately 13 times more

integrated luminosity than the first search for $H \rightarrow \phi\gamma$ decays [27], which led to a 95% CL upper limit of $\mathcal{B}(H \rightarrow \phi\gamma) < 1.4 \times 10^{-3}$, assuming SM production rates of the Higgs boson. Currently, no other experimental information about the $H \rightarrow \rho\gamma$ decay mode exists.

The searches for the analogous decays of the Z boson into a meson and a photon are also presented in this paper. These have been theoretically studied [28, 29] as a unique precision test of the SM and the factorisation approach in quantum chromodynamics (QCD), in an environment where the power corrections in terms of the QCD energy scale over the vector boson's mass are small [29]. The large Z boson production cross section at the LHC means that rare Z boson decays can be probed at branching fractions much smaller than for Higgs boson decays into the same final states. The SM branching fraction predictions for the decays considered in this paper are $\mathcal{B}(Z \rightarrow \phi\gamma) = (1.04 \pm 0.12) \times 10^{-8}$ [28, 29] and $\mathcal{B}(Z \rightarrow \rho\gamma) = (4.19 \pm 0.47) \times 10^{-8}$ [29]. The first search for $Z \rightarrow \phi\gamma$ decays by the ATLAS Collaboration was presented in Ref. [27] and a 95% CL upper limit of $\mathcal{B}(Z \rightarrow \phi\gamma) < 8.3 \times 10^{-6}$ was obtained. So far no direct experimental information about the decay $Z \rightarrow \rho\gamma$ exists.

2 ATLAS detector

ATLAS [30] is a multi-purpose particle physics detector with a forward-backward symmetric cylindrical geometry and near 4π coverage in solid angle.¹ It consists of an inner tracking detector surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters, and a muon spectrometer.

The inner tracking detector (ID) covers the pseudorapidity range $|\eta| < 2.5$, and is surrounded by a thin superconducting solenoid providing a 2 T magnetic field. At small radii, a high-granularity silicon pixel detector covers the vertex region and typically provides three measurements per track. A new innermost pixel-detector layer, the insertable B-layer, was added before 13 TeV data-taking began in 2015 and provides an additional measurement at a radius of about 33 mm around a new and thinner beam pipe [31]. The pixel detectors are followed by a silicon microstrip tracker, which typically provides four space-point measurements per track. The silicon detectors are complemented by a gas-filled straw-tube transition radiation tracker, which enables radially extended track reconstruction up to $|\eta| = 2.0$, with typically 35 measurements per track.

The calorimeter system covers the pseudorapidity range $|\eta| < 4.9$. A high-granularity lead/liquid-argon (LAr) sampling electromagnetic calorimeter covers the region $|\eta| < 3.2$, with an additional thin LAr presampler covering $|\eta| < 1.8$ to correct for energy losses upstream. The electromagnetic calorimeter is divided into a barrel section covering $|\eta| < 1.475$ and two endcap sections covering $1.375 < |\eta| < 3.2$. For $|\eta| < 2.5$ it is divided into three layers in depth, which are finely segmented in η and ϕ . A steel/scintillator-tile calorimeter provides hadronic calorimetry in the range $|\eta| < 1.7$. LAr technology, with copper as absorber, is used for the hadronic calorimeters in the endcap region, $1.5 < |\eta| < 3.2$. The solid-angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeter modules in $3.1 < |\eta| < 4.9$, optimised for electromagnetic and hadronic measurements, respectively.

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

The muon spectrometer surrounds the calorimeters and comprises separate trigger and high-precision tracking chambers measuring the deflection of muons in a magnetic field provided by three air-core superconducting toroids.

A two-level trigger and data acquisition system is used to provide an online selection and record events for offline analysis [32]. The level-1 trigger is implemented in hardware and uses a subset of detector information to reduce the event rate to 100 kHz or less from the maximum LHC collision rate of 40 MHz. It is followed by a software-based high-level trigger which filters events using the full detector information and records events for detailed offline analysis at an average rate of 1 kHz.

3 Data and Monte Carlo simulation

The search is performed with a sample of pp collision data recorded at a centre-of-mass energy $\sqrt{s} = 13$ TeV. Events are retained for further analysis only if they were collected under stable LHC beam conditions and the detector was operating normally. This results in an integrated luminosity of 35.6 and 32.3 fb⁻¹ for the $\phi\gamma$ and $\rho\gamma$ final states, respectively. The integrated luminosity of the data sample has an uncertainty of 3.4% derived using the method described in Ref. [33].

The $\phi\gamma$ and $\rho\gamma$ data samples used in this analysis were each collected with a specifically designed trigger. Both triggers require an isolated photon with a transverse momentum, p_T , greater than 35 GeV and an isolated pair of ID tracks, one of which must have a p_T greater than 15 GeV, associated with a topological cluster of calorimeter cells [34] with a transverse energy greater than 25 GeV. The photon part of the trigger follows the same process as the inclusive photon trigger requiring an electromagnetic cluster in the calorimeter consistent with a photon and is described with more detail in Ref. [32], while requirements on the ID tracks are applied in the high-level trigger through an appropriately modified version of the τ -lepton trigger algorithms which are described in more detail in Ref. [35]. The trigger for the $\phi\gamma$ final state was introduced in September 2015. This trigger requires that the invariant mass of the pair of tracks, under the charged-kaon hypothesis, is in the range 987–1060 MeV, consistent with the ϕ meson mass. The trigger efficiency for both the Higgs and Z boson signals is approximately 83% with respect to the offline selection, as described in Section 4. The corresponding trigger for the $\rho\gamma$ final state was introduced in May 2016. This trigger requires the invariant mass of the pair of tracks, under the charged-pion hypothesis, to be in the range 475–1075 MeV to include the bulk of the broad ρ meson mass distribution. The trigger efficiency for the Higgs boson signal is approximately 78% and for the Z boson signal is approximately 72% with respect to the offline selection.

Higgs boson production through the gluon–gluon fusion (ggH) and vector-boson fusion (VBF) processes was modelled up to next-to-leading order (NLO) in α_S using the POWHEG-Box v2 Monte Carlo (MC) event generator [36–40] with CT10 parton distribution functions [41]. POWHEG-Box was interfaced with the PYTHIA 8.186 MC event generator [42, 43] to model the parton shower, hadronisation and underlying event. The corresponding parameter values were set according to the AZNLO tune [44]. Additional contributions from the associated production of a Higgs boson and a W or Z boson (denoted by WH and ZH , respectively) are modelled by the PYTHIA 8.186 MC event generator with NNPDF23LO parton distribution functions [45] and the A14 tune for hadronisation and the underlying event [46]. The production rates and kinematic distributions for the SM Higgs boson with $m_H = 125$ GeV are assumed throughout. These were obtained from Ref. [16] and are summarised below. The ggH production rate is normalised such that it reproduces the total cross section predicted by a next-to-next-to-next-to-leading-order QCD calculation with NLO electroweak corrections applied [47–50]. The VBF production rate is

normalised to an approximate NNLO QCD cross section with NLO electroweak corrections applied [51–53]. The WH and ZH production rates are normalised to cross sections calculated at next-to-next-leading order (NNLO) in QCD with NLO electroweak corrections [54, 55] including the NLO QCD corrections [56] for $gg \rightarrow ZH$. The expected signal yield is corrected to include the 2% contribution from the production of a Higgs boson in association with a $t\bar{t}$ or a $b\bar{b}$ pair.

The POWHEG-BOX v2 MC event generator with CT10 parton distribution functions was also used to model inclusive Z boson production. PYTHIA 8.186 with CTEQ6L1 parton distribution functions [57] and the AZNLO parameter tune was used to simulate parton showering and hadronisation. The prediction is normalised to the total cross section obtained from the measurement in Ref. [58], which has an uncertainty of 2.9%. The Higgs and Z boson decays were simulated as a cascade of two-body decays, respecting angular momentum conservation. The meson line shapes were simulated by PYTHIA. The branching fraction for the decay $\phi \rightarrow K^+K^-$ is $(48.9 \pm 0.5)\%$ whereas the decay $\rho \rightarrow \pi^+\pi^-$ has a branching fraction close to 100% [59]. The simulated events were passed through the detailed GEANT 4 simulation of the ATLAS detector [60, 61] and processed with the same software used to reconstruct the data. Simulated pile-up events (additional pp collisions in the same or nearby bunch crossings) are also included and the distribution of these is matched to the conditions observed in the data.

4 Event selection for $\phi\gamma \rightarrow K^+K^-\gamma$ and $\rho\gamma \rightarrow \pi^+\pi^-\gamma$ final states

The $\phi\gamma$ and $\rho\gamma$ exclusive final states are very similar. Both final states consist of a pair of oppositely charged reconstructed ID tracks. The difference is that for the former the mass of the pair, under the charged-kaon hypothesis for the two tracks, is consistent with the ϕ meson mass, while for the later, under the charged-pion hypothesis for the tracks, it is consistent with the ρ meson mass. Events with a pp interaction vertex reconstructed from at least two ID tracks with $p_T > 400$ MeV are considered in the analysis. Within an event, the primary vertex is defined as the reconstructed vertex with the largest $\sum p_T^2$ of associated ID tracks.

Photons are reconstructed from clusters of energy in the electromagnetic calorimeter. Clusters without matching ID tracks are classified as unconverted photon candidates while clusters matched to ID tracks consistent with the hypothesis of a photon conversion into e^+e^- are classified as converted photon candidates [62]. Reconstructed photon candidates are required to have $p_T^\gamma > 35$ GeV, $|\eta^\gamma| < 2.37$, excluding the barrel/endcap calorimeter transition region $1.37 < |\eta^\gamma| < 1.52$, and to satisfy “tight” photon identification criteria [62]. An isolation requirement is imposed to further suppress contamination from jets. The sum of the transverse momenta of all tracks within $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = 0.2$ of the photon direction, excluding those associated with the reconstructed photon, is required to be less than 5% of p_T^γ . Moreover, the sum of the transverse momenta of all calorimeter energy deposits within $\Delta R = 0.4$ of the photon direction, excluding those associated with the reconstructed photon, is required to be less than 2.45 GeV $+ 0.022 \times p_T^\gamma$. To mitigate the effects of multiple pp interactions in the same or neighbouring bunch crossings, only ID tracks which originate from the primary vertex are considered in the photon track-based isolation. For the calorimeter-based isolation the effects of the underlying event and multiple pp interactions are also accounted for on an event by event basis using an average underlying event energy density determined from data, as described in Ref. [62].

Charged particles satisfying the requirements detailed below are assumed to be a K^\pm meson in the $\phi\gamma$ analysis and a π^\pm meson in the $\rho\gamma$ analysis. No further particle identification requirements are applied. In the following, when referring to charged particles collectively the term “charged-hadron candidates” is

used, while when referring to the charged particles relevant to the $\phi\gamma$ and the $\rho\gamma$ analyses the terms “kaon candidates” and “pion candidates” are used, respectively, along with the corresponding masses. A pair of oppositely-charged charged-hadron candidates is referred to collectively as M .

Charged-hadron candidates are reconstructed from ID tracks which are required to have $|\eta| < 2.5$, $p_T > 15$ GeV and to satisfy basic quality criteria, including a requirement on the number of hits in the silicon detectors [63]. The $\phi \rightarrow K^+K^-$ and $\rho \rightarrow \pi^+\pi^-$ decays are reconstructed from pairs of oppositely charged-hadron candidates; the candidate with the higher p_T , referred to as the leading charged-hadron candidate, is required to have $p_T > 20$ GeV.

Pairs of charged-hadron candidates are selected based on their invariant masses. Those with an invariant mass, under the charged-kaon hypothesis, $m_{K^+K^-}$ between 1012 MeV and 1028 MeV are selected as $\phi \rightarrow K^+K^-$ candidates. Pairs with an invariant mass, under the charged-pion hypothesis, $m_{\pi^+\pi^-}$ between 635 MeV and 915 MeV are selected as $\rho \rightarrow \pi^+\pi^-$ candidates. The candidates where $m_{K^+K^-}$ is consistent with the ϕ meson mass are rejected from the $\rho\gamma$ analysis. This requirement rejects a negligible fraction of the signal in the $\rho\gamma$ analysis. Selected M candidates are required to satisfy an isolation requirement: the sum of the p_T of the reconstructed ID tracks from the primary vertex within $\Delta R = 0.2$ of the leading charged hadron candidate (excluding the charged-hadron candidates defining the pair) is required to be less than 10% of the p_T of the M candidate.

The M candidates are combined with the photon candidates, to form $M\gamma$ candidates. When multiple combinations are possible, a situation that arises only in a few percent of the events, the combination of the highest- p_T photon and the M candidate with an invariant mass closest to the respective meson mass is selected. The event is retained for further analysis if the requirement $\Delta\phi(M, \gamma) > \pi/2$ is satisfied. The transverse momentum of the M candidates is required to be greater than a threshold that varies as a function of the invariant mass of the three-body system, $m_{M\gamma}$. Thresholds of 40 GeV and 47.2 GeV are imposed on p_T^M for the regions $m_{M\gamma} < 91$ GeV and $m_{M\gamma} \geq 140$ GeV, respectively. The threshold is varied from 40 GeV to 47.2 GeV as a linear function of $m_{M\gamma}$ in the region $91 \leq m_{M\gamma} < 140$ GeV. This approach ensures good sensitivity for both the Higgs and Z boson searches, while keeping a single kinematic selection.

For the $\phi(\rightarrow K^+K^-)\gamma$ final state, the total signal efficiencies (kinematic acceptance, trigger and reconstruction efficiencies) are 17% and 10% for the Higgs and Z boson decays, respectively. The corresponding efficiencies for the $\rho\gamma$ final state are 8% and 2.4%. The difference in efficiency between the Higgs and Z boson decays arises primarily from the softer p_T distributions of the photon and charged-hadron candidates associated with the $Z \rightarrow M\gamma$ production, as can be seen for the $\phi\gamma$ case by comparing Figures 1(a) and 1(b). The overall lower efficiency in the $\rho\gamma$ final state is a result of the lower efficiency of the m_M requirement due to the large ρ -meson natural width and the different kinematics of the ρ decay products, as presented in Figures 1(c) and 1(d). Meson helicity effects have a relatively small impact for the $\phi \rightarrow K^+K^-$ decays, where the kaons carry very little momentum in the ϕ rest frame. Specifically, the expected Higgs (Z) boson signal yield in the signal region is 2.4% larger (9% larger) than in the hypothetical scenario where the meson is unpolarised. For the $\rho \rightarrow \pi^+\pi^-$ decays the yields are increased by 12% (increased by 7%).

The average $m_{M\gamma}$ resolution is 1.8% for both the Higgs and Z boson decays. The Higgs boson signal $m_{M\gamma}$ distribution is modelled with a sum of two Gaussian probability density functions (pdf) with a common mean value, while the Z boson signal $m_{M\gamma}$ distribution is modelled with a double Voigtian pdf (a convolution of relativistic Breit–Wigner and Gaussian pdfs) corrected with a mass-dependent efficiency factor.

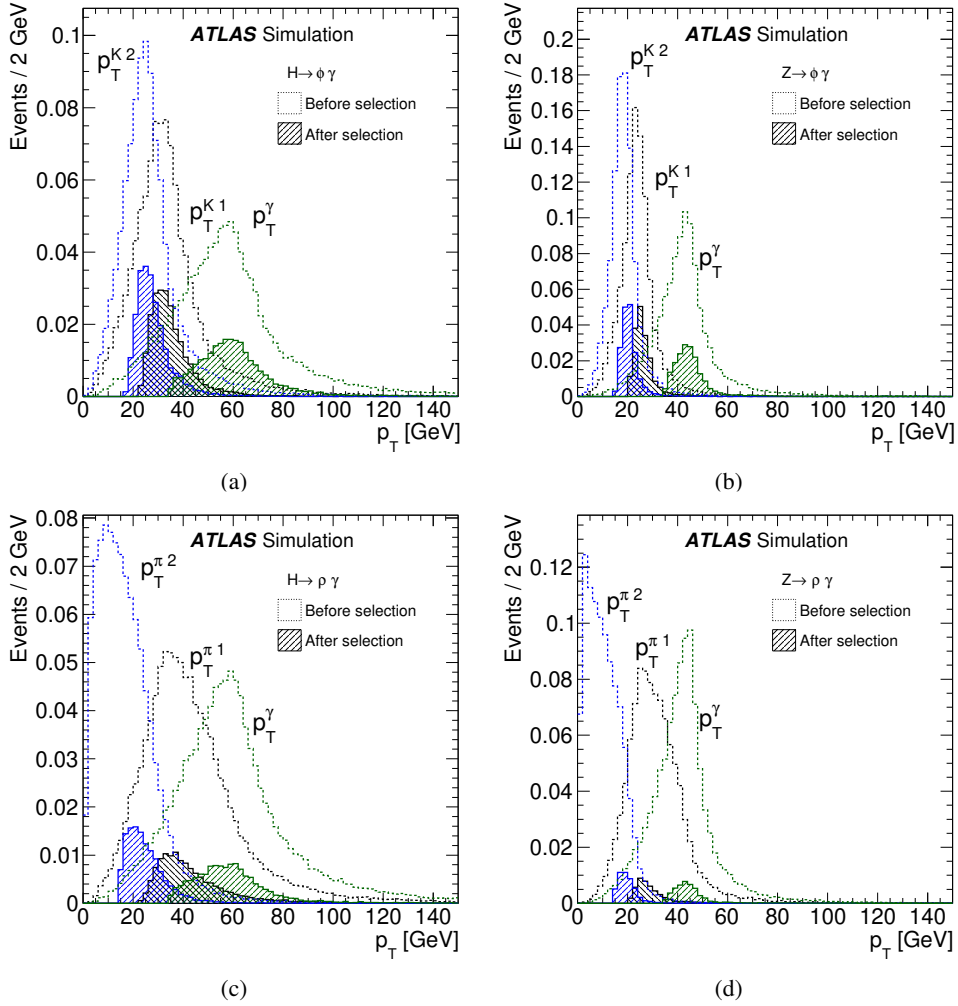


Figure 1: Generator-level transverse momentum (p_T) distributions of the photon and of the charged-hadrons, ordered in p_T , for (a) $H \rightarrow \phi \gamma$, (b) $Z \rightarrow \phi \gamma$, (c) $H \rightarrow \rho \gamma$ and (d) $Z \rightarrow \rho \gamma$ simulated signal events, respectively. The hatched histograms denote the full event selection while the dashed histograms show the events at generator level that fall within the analysis geometric acceptance (both charged-hadrons are required to have $|\eta| < 2.5$ while the photon is required to have $|\eta| < 2.37$, excluding the region $1.37 < |\eta| < 1.52$). The dashed histograms are normalised to unity, and the relative difference between the two sets of distributions corresponds to the effects of reconstruction, trigger, and event selection efficiencies. The leading charged-hadron candidate $h = K, \pi$ is denoted by p_T^{h1} and the sub-leading candidate by p_T^{h2} .

The $m_{K^+K^-}$ distribution for the selected $\phi \gamma$ candidates, with no $m_{K^+K^-}$ requirement applied, is shown in Figure 2(a) exhibiting a visible peak at the ϕ meson mass. The ϕ peak is fitted with a Voigtian pdf, while the background is modelled with a function typically used to describe kinematic thresholds [64]. The experimental resolution in $m_{K^+K^-}$ is approximately 4 MeV, comparable to the 4.3 MeV [59] width of the ϕ meson. In Figure 2(b), the corresponding distribution for the selected $\rho \gamma$ candidates is shown, where the ρ meson can also be observed. The ρ peak is fitted with a single Breit–Wigner pdf, modified by a mass-dependent width to match the distribution obtained from PYTHIA [42]. The background is fitted with the sum of a combinatoric background, estimated from events containing a same-sign di-track pair, and other backgrounds determined in the fit using a linear combination of Chebyshev polynomials up to

the second order. Figure 2 only qualitatively illustrates the meson selection in the studied final state, and is not used any further in this analysis.

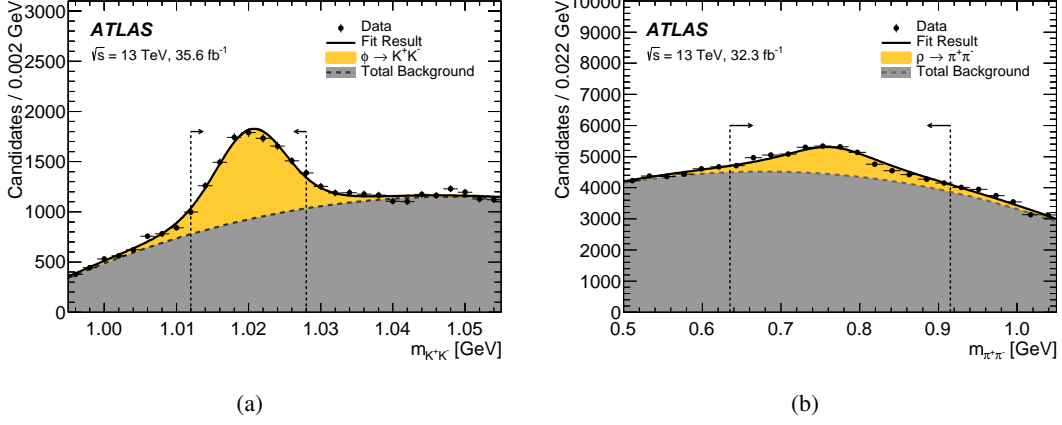


Figure 2: The (a) $m_{K^+K^-}$ and (b) $m_{\pi^+\pi^-}$ distributions for $\phi\gamma$ and $\rho\gamma$ candidates, respectively. The candidates fulfil the complete event selection (see text), apart from requirements on $m_{K^+K^-}$ or $m_{\pi^+\pi^-}$. These requirements are marked on the figures with dashed lines topped with arrows indicating the included area. The signal and background models are discussed in the text.

5 Background

For both the $\phi\gamma$ and $\rho\gamma$ final states, the main sources of background in the searches are events involving inclusive photon + jet or multijet processes where an M candidate is reconstructed from ID tracks originating from a jet.

From the selection criteria discussed earlier, the shape of this background exhibits a turn-on structure in the $m_{M\gamma}$ distribution around 100 GeV, in the region of the Z boson signal, and a smoothly falling background in the region of the Higgs boson signal. Given the complex shape of this background, these processes are modelled in an inclusive fashion with a non-parametric data-driven approach using templates to describe the relevant distributions. The background normalisation and shape are simultaneously extracted from a fit to the data. A similar procedure was used in the earlier search for Higgs and Z boson decays into $\phi\gamma$ [27] and the search for Higgs and Z boson decays into $J/\psi\gamma$ and $\Upsilon(nS)\gamma$ described in Ref. [21].

5.1 Background modelling

The background modelling procedure for each final state exploits a sample of approximately 54 000 $K^+K^-\gamma$ and 220 000 $\pi^+\pi^-\gamma$ candidate events in data. These events pass all the kinematic selection requirements described previously, except that the photon and M candidates are not required to satisfy the nominal isolation requirements, and a looser $p_T^M > 35$ GeV requirement is imposed. This selection defines the background-dominated “generation region” (GR). From these events, pdfs are constructed to describe the distributions of the relevant kinematic and isolation variables and their most important correlations. In

this way, in the absence of appropriate simulations, pseudocandidate events are generated, from which the background shape in the discriminating variable is derived.

This ensemble of pseudocandidate events is produced by randomly sampling the distributions of the relevant kinematic and isolation variables, which are estimated from the data in the GR. Each pseudocandidate event is described by M and γ four-momentum vectors and the associated M and photon isolation variables. The M four-momentum vector is constructed from sampled η_M , ϕ_M , m_M and p_T^M values. For the γ four-momentum vector, the η_γ and ϕ_γ are determined from the sampled $\Delta\phi(M, \gamma)$ and $\Delta\eta(M, \gamma)$ values whereas p_T^γ is sampled directly.

The most important correlations among these kinematic and isolation variables in background events are retained in the generation of the pseudocandidates through the following sampling scheme, where the steps are performed sequentially:

- i) Values for η_M , ϕ_M , m_M and p_T^M are drawn randomly and independently according to the corresponding pdfs.
- ii) The distribution of p_T^γ values is parameterised in bins of p_T^M , and values are drawn from the corresponding bins given the previously generated value of p_T^M . The M isolation variable is parameterised in bins of $p_T^M(p_T^\gamma)$ for the $\phi\gamma$ ($\rho\gamma$) model and sampled accordingly. The difference between the two approaches for the $\phi\gamma$ and $\rho\gamma$ accounts for the difference in the observed correlations arising in the different datasets.
- iii) The distributions of the values for $\Delta\eta(M, \gamma)$, photon calorimeter isolation, normalised to p_T^γ , and their correlations are parameterised in a two-dimensional distribution. For the $\phi\gamma$ analysis, several distributions are produced corresponding to the p_T^M bins used earlier to describe the p_T^γ and M isolation variables, whereas for the $\rho\gamma$ final state the two-dimensional distribution is produced inclusively for all p_T^M values.
- iv) The photon track isolation, normalised to p_T^γ , and the $\Delta\phi(M, \gamma)$ variables are sampled from pdfs generated in bins of relative photon calorimeter isolation and $\Delta\eta(M, \gamma)$, respectively, using the values drawn in step iii).

The nominal selection requirements are imposed on the ensemble, and the surviving pseudocandidates are used to construct templates for the $m_{M\gamma}$ distribution, which are then smoothed using Gaussian kernel density estimation [65]. It was verified through signal injection tests that the shape of the background model is not affected by potential signal contamination.

5.2 Background validation

To validate the background model, the $m_{M\gamma}$ distributions in several validation regions, defined by kinematic and isolation requirements looser than the nominal signal requirements, are used to compare the prediction of the background model with the data. Three validation regions are defined, each based on the GR selection and adding one of the following: the p_T^M requirement (VR1), the photon isolation requirements (VR2), or the meson isolation requirement (VR3). The $m_{M\gamma}$ distributions in these validation regions are shown in Figure 3. The background model is found to describe the data in all regions within uncertainties (see Section 6).

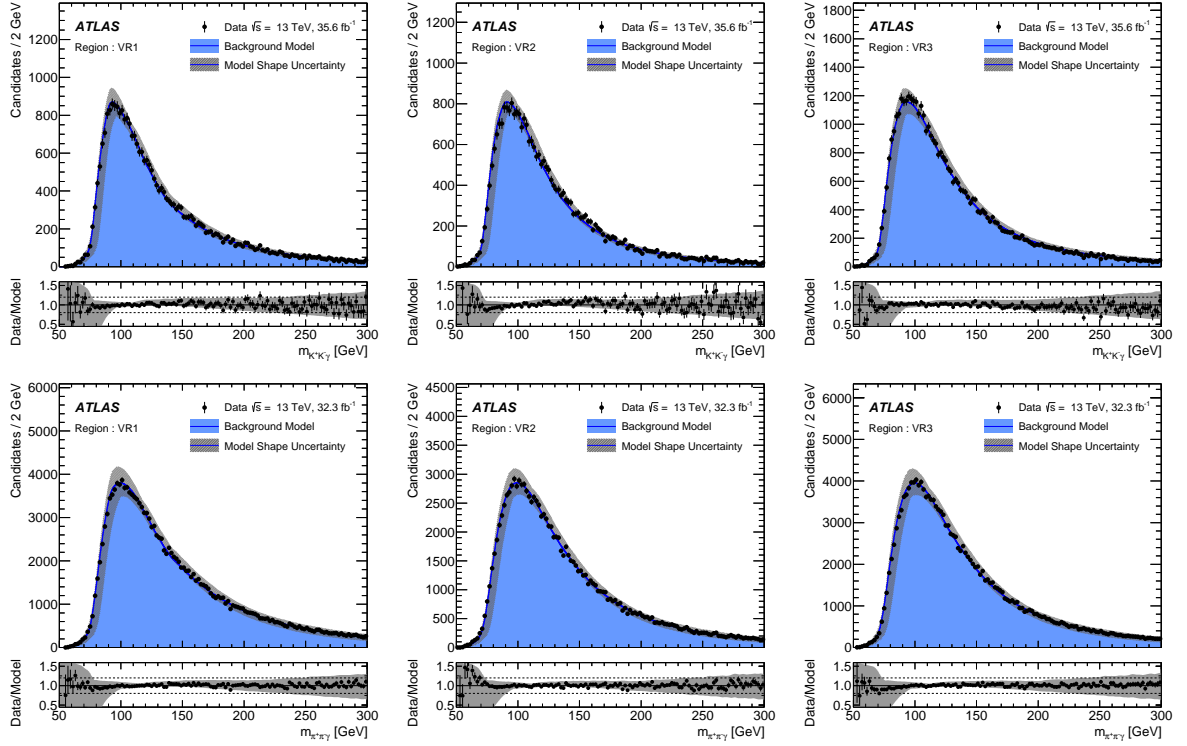


Figure 3: The distribution of $m_{K^+K^- \gamma}$ top ($m_{\pi^+ \pi^- \gamma}$ bottom) in data compared to the prediction of the background model for the VR1, VR2 and VR3 validation regions. The background model is normalised to the observed number of events within the region shown. The uncertainty band corresponds to the uncertainty envelope derived from variations in the background modelling procedure. The ratio of the data to the background model is shown below the distributions.

Potential background contributions from $Z \rightarrow \ell\ell\gamma$ decays and inclusive Higgs decays were studied and found to be negligible for the selection requirements and dataset used in this analysis.

A further validation of the background modelling is performed using events within a sideband of the M mass distribution. For the $\phi\gamma$ analysis the sideband region is defined by $1.035 \text{ GeV} < m_{K^+K^-} < 1.051 \text{ GeV}$. For the $\rho\gamma$ analysis the sideband region is defined by $950 \text{ MeV} < m_{\pi^+ \pi^-} < 1050 \text{ MeV}$. All other selection requirements and modelling procedures are identical to those used in the signal region. Figures 4(a) and 4(b) show the $m_{M\gamma}$ distributions for the sideband region. The background model is found to describe the data within the systematic uncertainties described in Section 6.

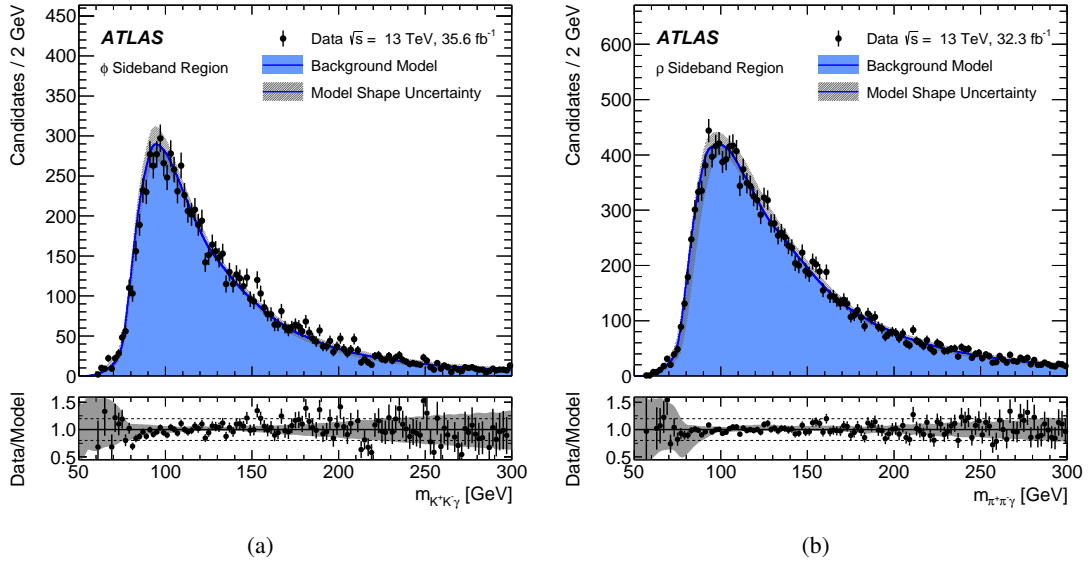


Figure 4: The distribution of $m_{M\gamma}$ for the (a) $\phi\gamma$ and (b) $\rho\gamma$ selections in the sideband control region. The background model is normalised to the observed number of events within the region shown. The uncertainty band corresponds to the uncertainty envelope derived from variations in the background modelling procedure. The ratio of the data to the background model is shown below the distributions.

6 Systematic uncertainties

Trigger and identification efficiencies for photons are determined from samples enriched with $Z \rightarrow e^+e^-$ events in data [32, 62]. The systematic uncertainty in the expected signal yield associated with the trigger efficiency is estimated to be 2.0%. The photon identification and isolation uncertainties, for both the converted and unconverted photons, are estimated to be 2.4% and 2.6% for the Higgs and Z boson signals, respectively. An uncertainty of 6.0% per M candidate is assigned to the track reconstruction efficiency and accounts for effects associated with the modelling of ID material and track reconstruction algorithms if a nearby charged particle is present. This uncertainty is derived conservatively by assuming a 3% uncertainty in the reconstruction efficiency of each track [66], and further assuming the uncertainty to be fully correlated between the two tracks of the M candidate.

The systematic uncertainties in the Higgs production cross section are obtained from Ref. [16] as described in Section 3. The Z boson production cross-section uncertainty is taken from the measurement in Ref. [58].

The photon energy scale uncertainty, determined from $Z \rightarrow e^+e^-$ events and validated using $Z \rightarrow \ell\ell\gamma$ events [67], is applied to the simulated signal samples as a function of η^γ and p_T^γ . The impact of the photon energy scale uncertainty on the Higgs and Z boson mass distributions does not exceed 0.2%. The uncertainty associated with the photon energy resolution is found to have a negligible impact. Similarly, the systematic uncertainty associated with the ID track momentum measurement is found to be negligible. The systematic uncertainties in the expected signal yields are summarised in Table 1.

The shape of the background model is allowed to vary around the nominal shape, and the parameters controlling these systematic variations are treated as nuisance parameters in the maximum-likelihood fit used to extract the signal and background yields. Three such shape variations are implemented through

varying p_T^γ , linear distortions of the shape of the $\Delta\phi(M, \gamma)$, and a global tilt of the three-body mass. The first two variations alter the kinematics of the pseudocandidates that are propagated to the three-body mass.

Table 1: Summary of the relative systematic uncertainties in the expected signal yields. The magnitude of the effects are the same for both the $\phi\gamma$ and $\rho\gamma$ selections.

Source of systematic uncertainty	Yield uncertainty
Total H cross section	6.3%
Total Z cross section	2.9%
Integrated luminosity	3.4%
Photon ID efficiency	2.5%
Trigger efficiency	2.0%
Tracking efficiency	6.0%

7 Results

The data are compared to background and signal predictions using an unbinned maximum-likelihood fit to the $m_{M\gamma}$ distribution. The parameters of interest are the Higgs and Z boson signal normalisations. Systematic uncertainties are modelled using additional nuisance parameters in the fit; in particular the background normalisation is a free parameter in the model. The fit uses the selected events with $m_{M\gamma} < 300$ GeV. The expected and observed numbers of background events within the $m_{M\gamma}$ ranges relevant to the Higgs and Z boson signals are shown in Table 2. The observed yields are consistent with the number of events expected from the background-only prediction within the systematic and statistical uncertainties. The results of the background-only fits for the $\phi\gamma$ and $\rho\gamma$ analyses are shown in Figures 5(a) and 5(b), respectively.

Table 2: The number of observed events and the mean expected background, estimated from the maximum-likelihood fit and shown with the associated total uncertainty, for the $m_{M\gamma}$ ranges of interest. The expected Higgs and Z boson signal yields, along with the total systematic uncertainty, for $\phi\gamma$ and $\rho\gamma$, estimated using simulations, are also shown in parentheses.

	Observed yields (Mean expected background)				Expected signal yields		
	Mass range [GeV]				H	Z	
	All	81–101		120–130	$[\mathcal{B} = 10^{-4}]$	$[\mathcal{B} = 10^{-6}]$	
$\phi\gamma$	12051	3364	(3500 \pm 30)	1076	(1038 \pm 9)	15.1 \pm 1.5	98 \pm 8
$\rho\gamma$	58702	12583	(12660 \pm 60)	5473	(5450 \pm 30)	14.3 \pm 1.4	47 \pm 4

Upper limits are set on the branching fractions for the Higgs and Z boson decays into $M\gamma$ using the CL_s modified frequentist formalism [68] with the profile-likelihood-ratio test statistic [69]. For the upper limits on the branching fractions, the SM production cross section is assumed for the Higgs boson [16], while the ATLAS measurement of the inclusive Z boson cross section is used for the Z boson signal [58], as discussed in Section 3. The results are summarised in Table 3. The observed 95% CL upper limits on the branching fractions for $H \rightarrow \phi\gamma$ and $Z \rightarrow \phi\gamma$ decays are 208 and 87 times the expected SM branching fractions, respectively. The corresponding values for the $\rho\gamma$ decays are 52 and 597 times the expected SM branching fractions, respectively. Upper limits at 95% CL on the production cross section times branching

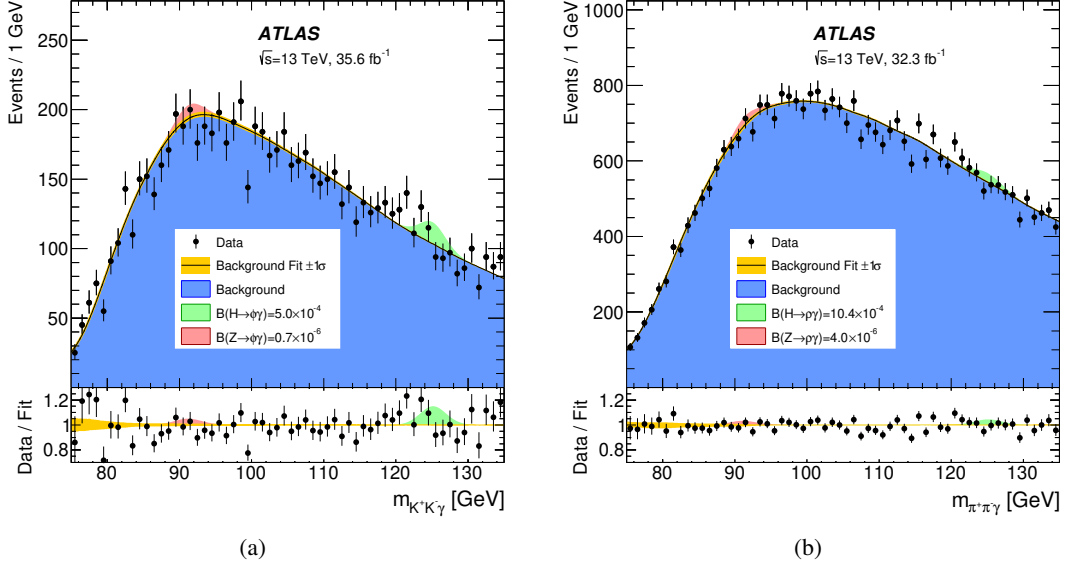


Figure 5: The (a) $m_{K^+K^- \gamma}$ and (b) $m_{\pi^+\pi^- \gamma}$ distributions of the selected $\phi\gamma$ and $\rho\gamma$ candidates, respectively, along with the results of the maximum-likelihood fits with a background-only model. The Higgs and Z boson contributions for the branching fraction values corresponding to the observed 95% CL upper limits are also shown. Below the figures the ratio of the data to the background-only fit is shown.

fraction are also estimated for the Higgs boson decays, yielding 26.1 fb for the $H \rightarrow \phi\gamma$ decay, and 54.8 fb for the $H \rightarrow \rho\gamma$ decay.

The systematic uncertainties described in Section 6 result in a 14% deterioration of the post-fit expected 95% CL upper limit on the branching fraction in the $H \rightarrow \phi\gamma$ and $Z \rightarrow \phi\gamma$ analyses, compared to the result including only statistical uncertainties. For the $\rho\gamma$ analysis the systematic uncertainties result in a 2.3% increase in the post-fit expected upper limit for the Higgs boson decay, while for the Z boson decay the upper limit deteriorates by 29%.

Table 3: Expected and observed branching fraction upper limits at 95% CL for the $\phi\gamma$ and $\rho\gamma$ analyses. The $\pm 1\sigma$ intervals of the expected limits are also given.

Branching Fraction Limit (95% CL)	Expected	Observed
$\mathcal{B}(H \rightarrow \phi\gamma) [10^{-4}]$	$4.2^{+1.8}_{-1.2}$	5.0
$\mathcal{B}(Z \rightarrow \phi\gamma) [10^{-6}]$	$1.1^{+0.5}_{-0.3}$	0.7
$\mathcal{B}(H \rightarrow \rho\gamma) [10^{-4}]$	$10.0^{+4.9}_{-2.8}$	10.4
$\mathcal{B}(Z \rightarrow \rho\gamma) [10^{-6}]$	$5.1^{+2.1}_{-1.4}$	4.0

8 Summary

A search for the decays of Higgs and Z bosons into $\phi\gamma$ and $\rho\gamma$ has been performed with $\sqrt{s} = 13$ TeV pp collision data samples collected with the ATLAS detector at the LHC corresponding to integrated luminosities of up to 35.6 fb^{-1} . The ϕ and ρ mesons are reconstructed via their dominant decays into

the K^+K^- and $\pi^+\pi^-$ final states, respectively. The background model is derived using a fully data driven approach and validated in a number of control regions including sidebands in the K^+K^- and $\pi^+\pi^-$ mass distributions.

No significant excess of events above the background expectations is observed, as expected from the SM. The obtained 95% CL upper limits are $\mathcal{B}(H \rightarrow \phi\gamma) < 5.0 \times 10^{-4}$, $\mathcal{B}(Z \rightarrow \phi\gamma) < 0.7 \times 10^{-6}$, $\mathcal{B}(H \rightarrow \rho\gamma) < 10.4 \times 10^{-4}$ and $\mathcal{B}(Z \rightarrow \rho\gamma) < 4.0 \times 10^{-6}$.

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The ATLAS Collaboration

M. Aaboud^{34d}, G. Aad⁹⁹, B. Abbott¹²⁴, O. Abdinov^{13,*}, B. Abeloos¹²⁸, S.H. Abidi¹⁶⁵, O.S. AbouZeid¹⁴³, N.L. Abraham¹⁵³, H. Abramowicz¹⁵⁹, H. Abreu¹⁵⁸, R. Abreu¹²⁷, Y. Abulaiti^{43a,43b}, B.S. Acharya^{64a,64b,o}, S. Adachi¹⁶¹, L. Adamczyk^{81a}, J. Adelman¹¹⁹, M. Adersberger¹¹², T. Adye¹⁴¹, A.A. Affolder¹⁴³, Y. Afik¹⁵⁸, T. Agatonovic-Jovin¹⁶, C. Agheorghiesei^{27c}, J.A. Aguilar-Saavedra^{136f,136a}, F. Ahmadov^{77,ag}, G. Aielli^{71a,71b}, S. Akatsuka⁸³, H. Akerstedt^{43a,43b}, T.P.A. Åkesson⁹⁴, E. Akilli⁵², A.V. Akimov¹⁰⁸, G.L. Alberghi^{23b,23a}, J. Albert¹⁷⁴, P. Albicocco⁴⁹, M.J. Alconada Verzini⁸⁶, S. Alderweireldt¹¹⁷, M. Aleksa³⁵, I.N. Aleksandrov⁷⁷, C. Alexa^{27b}, G. Alexander¹⁵⁹, T. Alexopoulos¹⁰, M. Alhroob¹²⁴, B. Ali¹³⁸, G. Alimonti^{66a}, J. Alison³⁶, S.P. Alkire³⁸, B.M.M. Allbrooke¹⁵³, B.W. Allen¹²⁷, P.P. Allport²¹, A. Aloisio^{67a,67b}, A. Alonso³⁹, F. Alonso⁸⁶, C. Alpigiani¹⁴⁵, A.A. Alshehri⁵⁵, M.I. Alstaty⁹⁹, B. Alvarez Gonzalez³⁵, D. Álvarez Piqueras¹⁷², M.G. Alviggi^{67a,67b}, B.T. Amadio¹⁸, Y. Amaral Coutinho^{78b}, C. Amelung²⁶, D. Amidei¹⁰³, S.P. Amor Dos Santos^{136a,136c}, S. Amoroso³⁵, G. Amundsen²⁶, C. Anastopoulos¹⁴⁶, L.S. Ancu⁵², N. Andari²¹, T. Andeen¹¹, C.F. Anders^{59b}, J.K. Anders⁸⁸, K.J. Anderson³⁶, A. Andreazza^{66a,66b}, V. Andrei^{59a}, S. Angelidakis³⁷, I. Angelozzi¹¹⁸, A. Angerami³⁸, A.V. Anisenkov^{120b,120a}, N. Anjos¹⁴, A. Annovi^{69a}, C. Antel^{59a}, M. Antonelli⁴⁹, A. Antonov^{110,*}, D.J.A. Antrim¹⁶⁹, F. Anulli^{70a}, M. Aoki⁷⁹, L. Aperio Bella³⁵, G. Arabidze¹⁰⁴, Y. Arai⁷⁹, J.P. Araque^{136a}, V. Araujo Ferraz^{78b}, A.T.H. Arce⁴⁷, R.E. Ardell⁹¹, F.A. Arduh⁸⁶, J-F. Arguin¹⁰⁷, S. Argyropoulos⁷⁵, M. Arik^{12c}, A.J. Armbruster³⁵, L.J. Armitage⁹⁰, O. Arnaez¹⁶⁵, H. Arnold⁵⁰, M. Arratia³¹, O. Arslan²⁴, A. Artamonov^{109,*}, G. Artoni¹³¹, S. Artz⁹⁷, S. Asai¹⁶¹, N. Asbah⁴⁴, A. Ashkenazi¹⁵⁹, L. Asquith¹⁵³, K. Assamagan²⁹, R. Astalos^{28a}, M. Atkinson¹⁷¹, N.B. Atlay¹⁴⁸, K. Augsten¹³⁸, G. Avolio³⁵, B. Axen¹⁸, M.K. Ayoub^{15a}, G. Azuelos^{107,av}, A.E. Baas^{59a}, M.J. Baca²¹, H. Bachacou¹⁴², K. Bachas^{65a,65b}, M. Backes¹³¹, P. Bagnaia^{70a,70b}, M. Bahmani⁸², H. Bahrasemani¹⁴⁹, J.T. Baines¹⁴¹, M. Bajic³⁹, O.K. Baker¹⁸¹, P.J. Bakker¹¹⁸, E.M. Baldin^{120b,120a}, P. Balek¹⁷⁸, F. Balli¹⁴², W.K. Balunas¹³³, E. Banas⁸², A. Bandyopadhyay²⁴, S. Banerjee^{179,k}, A.A.E. Bannoura¹⁸⁰, L. Barak¹⁵⁹, E.L. Barberio¹⁰², D. Barberis^{53b,53a}, M. Barbero⁹⁹, T. Barillari¹¹³, M-S. Barisits³⁵, J. Barkeloo¹²⁷, T. Barklow¹⁵⁰, N. Barlow³¹, S.L. Barnes^{58c}, B.M. Barnett¹⁴¹, R.M. Barnett¹⁸, Z. Barnovska-Blenessy^{58a}, A. Baroncelli^{72a}, G. Barone²⁶, A.J. Barr¹³¹, L. Barranco Navarro¹⁷², F. Barreiro⁹⁶, J. Barreiro Guimarães da Costa^{15a}, R. Bartoldus¹⁵⁰, A.E. Barton⁸⁷, P. Bartos^{28a}, A. Basalaev¹³⁴, A. Bassalat¹²⁸, R.L. Bates⁵⁵, S.J. Batista¹⁶⁵, J.R. Batley³¹, M. Battaglia¹⁴³, M. Baucé^{70a,70b}, F. Bauer¹⁴², H.S. Bawa^{150,m}, J.B. Beacham¹²², M.D. Beattie⁸⁷, T. Beau¹³², P.H. Beauchemin¹⁶⁸, P. Bechtel²⁴, H.C. Beck⁵¹, H.P. Beck^{20,s}, K. Becker¹³¹, M. Becker⁹⁷, C. Becot¹²¹, A. Beddall^{12d}, A.J. Beddall^{12a}, V.A. Bednyakov⁷⁷, M. Bedognetti¹¹⁸, C.P. Bee¹⁵², T.A. Beermann³⁵, M. Begalli^{78b}, M. Beger²⁹, J.K. Behr⁴⁴, A.S. Bell⁹², G. Bella¹⁵⁹, L. Bellagamba^{23b}, A. Bellerive³³, M. Bellomo¹⁵⁸, K. Belotskiy¹¹⁰, O. Beltramello³⁵, N.L. Belyaev¹¹⁰, O. Benary^{159,*}, D. Benchekroun^{34a}, M. Bender¹¹², N. Benekos¹⁰, Y. Benhammou¹⁵⁹, E. Benhar Noccioli¹⁸¹, J. Benitez⁷⁵, D.P. Benjamin⁴⁷, M. Benoit⁵², J.R. Bensinger²⁶, S. Bentvelsen¹¹⁸, L. Beresford¹³¹, M. Beretta⁴⁹, D. Berge¹¹⁸, E. Bergeaas Kuutmann¹⁷⁰, N. Berger⁵, L.J. Bergsten²⁶, J. Beringer¹⁸, S. Berlendis⁵⁶, N.R. Bernard¹⁰⁰, G. Bernardi¹³², C. Bernius¹⁵⁰, F.U. Bernlochner²⁴, T. Berry⁹¹, P. Berta⁹⁷, C. Bertella^{15a}, G. Bertoli^{43a,43b}, I.A. Bertram⁸⁷, C. Bertsche⁴⁴, G.J. Besjes³⁹, O. Bessidskaia Bylund^{43a,43b}, M. Bessner⁴⁴, N. Besson¹⁴², A. Bethani⁹⁸, S. Bethke¹¹³, A. Betti²⁴, A.J. Bevan⁹⁰, J. Beyer¹¹³, R.M.B. Bianchi¹³⁵, O. Biebel¹¹², D. Biedermann¹⁹, R. Bielski⁹⁸, K. Bierwagen⁹⁷, N.V. Biesuz^{69a,69b}, M. Biglietti^{72a}, T.R.V. Billoud¹⁰⁷, H. Bilokon⁴⁹, M. Bindi⁵¹, A. Bingul^{12d}, C. Bini^{70a,70b}, S. Biondi^{23b,23a}, T. Bisanz⁵¹, C. Bittrich⁴⁶, D.M. Bjergaard⁴⁷, J.E. Black¹⁵⁰, K.M. Black²⁵, R.E. Blair⁶, T. Blazek^{28a}, I. Bloch⁴⁴, C. Blocker²⁶, A. Blue⁵⁵, U. Blumenschein⁹⁰, Dr. Blunier^{144a}, G.J. Bobbink¹¹⁸, V.S. Bobrovnikov^{120b,120a}, S.S. Bocchetta⁹⁴, A. Bocci⁴⁷, C. Bock¹¹², M. Boehler⁵⁰, D. Boerner¹⁸⁰, D. Bogavac¹¹², A.G. Bogdanchikov^{120b,120a},

C. Bohm^{43a}, V. Boisvert⁹¹, P. Bokan^{170,y}, T. Bold^{81a}, A.S. Boldyrev¹¹¹, A.E. Bolz^{59b}, M. Bomben¹³²,
 M. Bona⁹⁰, M. Boonekamp¹⁴², A. Borisov¹⁴⁰, G. Borissov⁸⁷, J. Bortfeldt³⁵, D. Bortoletto¹³¹,
 V. Bortolotto^{61a,61b,61c}, D. Boscherini^{23b}, M. Bosman¹⁴, J.D. Bossio Sola³⁰, J. Boudreau¹³⁵,
 E.V. Bouhova-Thacker⁸⁷, D. Boumediene³⁷, C. Bourdarios¹²⁸, S.K. Boutle⁵⁵, A. Boveia¹²², J. Boyd³⁵,
 I.R. Boyko⁷⁷, A.J. Bozson⁹¹, J. Bracinik²¹, A. Brandt⁸, G. Brandt⁵¹, O. Brandt^{59a}, F. Braren⁴⁴,
 U. Bratzler¹⁶², B. Brau¹⁰⁰, J.E. Brau¹²⁷, W.D. Breaden Madden⁵⁵, K. Brendlinger⁴⁴, A.J. Brennan¹⁰²,
 L. Brenner¹¹⁸, R. Brenner¹⁷⁰, S. Bressler¹⁷⁸, D.L. Briglin²¹, T.M. Bristow⁴⁸, D. Britton⁵⁵, D. Britzger⁴⁴,
 I. Brock²⁴, R. Brock¹⁰⁴, G. Brooijmans³⁸, T. Brooks⁹¹, W.K. Brooks^{144b}, J. Brosamer¹⁸, E. Brost¹¹⁹,
 J.H. Broughton²¹, P.A. Bruckman de Renstrom⁸², D. Bruncko^{28b}, A. Bruni^{23b}, G. Bruni^{23b}, L.S. Bruni¹¹⁸,
 S. Bruno^{71a,71b}, B.H. Brunt³¹, M. Bruschi^{23b}, N. Bruscolo¹³⁵, P. Bryant³⁶, L. Bryngemark⁴⁴,
 T. Buanes¹⁷, Q. Buat¹⁴⁹, P. Buchholz¹⁴⁸, A.G. Buckley⁵⁵, I.A. Budagov⁷⁷, F. Buehrer⁵⁰, M.K. Bugge¹³⁰,
 O. Bulekov¹¹⁰, D. Bullock⁸, T.J. Burch¹¹⁹, S. Burdin⁸⁸, C.D. Burgard¹¹⁸, A.M. Burger⁵,
 B. Burghgrave¹¹⁹, K. Burka⁸², S. Burke¹⁴¹, I. Burmeister⁴⁵, J.T.P. Burr¹³¹, D. Büscher⁵⁰, V. Büscher⁹⁷,
 P. Bussey⁵⁵, J.M. Butler²⁵, C.M. Buttar⁵⁵, J.M. Butterworth⁹², P. Butti³⁵, W. Buttinger²⁹, A. Buzatu¹⁵⁵,
 A.R. Buzykaev^{120b,120a}, S. Cabrera Urbán¹⁷², D. Caforio¹³⁸, H. Cai¹⁷¹, V.M.M. Cairo^{40b,40a}, O. Cakir^{4a},
 N. Calace⁵², P. Calafiura¹⁸, A. Calandri⁹⁹, G. Calderini¹³², P. Calfayan⁶³, G. Callea^{40b,40a},
 L.P. Caloba^{78b}, S. Calvente Lopez⁹⁶, D. Calvet³⁷, S. Calvet³⁷, T.P. Calvet⁹⁹, R. Camacho Toro³⁶,
 S. Camarda³⁵, P. Camarri^{71a,71b}, D. Cameron¹³⁰, R. Caminal Armadans¹⁷¹, C. Camincher⁵⁶,
 S. Campana³⁵, M. Campanelli⁹², A. Camplani^{66a,66b}, A. Campoverde¹⁴⁸, V. Canale^{67a,67b},
 M. Cano Bret^{58c}, J. Cantero¹²⁵, T. Cao¹⁵⁹, M.D.M. Capeans Garrido³⁵, I. Caprini^{27b}, M. Caprini^{27b},
 M. Capua^{40b,40a}, R.M. Carbone³⁸, R. Cardarelli^{71a}, F.C. Cardillo⁵⁰, I. Carli¹³⁹, T. Carli³⁵, G. Carlino^{67a},
 B.T. Carlson¹³⁵, L. Carminati^{66a,66b}, R.M.D. Carney^{43a,43b}, S. Caron¹¹⁷, E. Carquin^{144b}, S. Carrá^{66a,66b},
 G.D. Carrillo-Montoya³⁵, D. Casadei²¹, M.P. Casado^{14,g}, A.F. Casha¹⁶⁵, M. Casolino¹⁴, D.W. Casper¹⁶⁹,
 R. Castelijin¹¹⁸, V. Castillo Gimenez¹⁷², N.F. Castro^{136a}, A. Catinaccio³⁵, J.R. Catmore¹³⁰, A. Cattai³⁵,
 J. Caudron²⁴, V. Cavaliere¹⁷¹, E. Cavallaro¹⁴, D. Cavalli^{66a}, M. Cavalli-Sforza¹⁴, V. Cavasinni^{69a,69b},
 E. Celebi^{12b}, F. Ceradini^{72a,72b}, L. Cerda Alberich¹⁷², A.S. Cerqueira^{78a}, A. Cerri¹⁵³, L. Cerrito^{71a,71b},
 F. Cerutti¹⁸, A. Cervelli^{23b,23a}, S.A. Cetin^{12b}, A. Chafaq^{34a}, D. Chakraborty¹¹⁹, S.K. Chan⁵⁷,
 W.S. Chan¹¹⁸, Y.L. Chan^{61a}, P. Chang¹⁷¹, J.D. Chapman³¹, D.G. Charlton²¹, C.C. Chau³³,
 C.A. Chavez Barajas¹⁵³, S. Che¹²², S. Cheatham^{64a,64c}, A. Chegwidan¹⁰⁴, S. Chekanov⁶,
 S.V. Chekulaev^{166a}, G.A. Chelkov^{77,au}, M.A. Chelstowska³⁵, C. Chen^{58a}, C.H. Chen⁷⁶, H. Chen²⁹,
 J. Chen^{58a}, S. Chen¹⁶¹, S.J. Chen^{15c}, X. Chen^{15b,at}, Y. Chen⁸⁰, H.C. Cheng¹⁰³, H.J. Cheng^{15d},
 A. Cheplakov⁷⁷, E. Cheremushkina¹⁴⁰, R. Cherkaoui El Moursli^{34e}, E. Cheu⁷, K. Cheung⁶²,
 L. Chevalier¹⁴², V. Chiarella⁴⁹, G. Chiarelli^{69a}, G. Chiodini^{65a}, A.S. Chisholm³⁵, A. Chitan^{27b},
 Y.H. Chiu¹⁷⁴, M.V. Chizhov⁷⁷, K. Choi⁶³, A.R. Chomont³⁷, S. Chouridou¹⁶⁰, Y.S. Chow^{61a},
 V. Christodoulou⁹², M.C. Chu^{61a}, J. Chudoba¹³⁷, A.J. Chuinard¹⁰¹, J.J. Chwastowski⁸², L. Chytka¹²⁶,
 A.K. Ciftci^{4a}, D. Cinca⁴⁵, V. Cindro⁸⁹, I.A. Cioară²⁴, A. Ciocio¹⁸, F. Ciotto^{67a,67b}, Z.H. Citron¹⁷⁸,
 M. Citterio^{66a}, M. Ciubancan^{27b}, A. Clark⁵², B.L. Clark⁵⁷, M.R. Clark³⁸, P.J. Clark⁴⁸, R.N. Clarke¹⁸,
 C. Clement^{43a,43b}, Y. Coadou⁹⁹, M. Coba^{64a,64c}, A. Coccaro⁵², J. Cochran⁷⁶, L. Colasurdo¹¹⁷, B. Cole³⁸,
 A.P. Colijn¹¹⁸, J. Collot⁵⁶, T. Colombo¹⁶⁹, P. Conde Muiño^{136a,136b}, E. Coniavitis⁵⁰, S.H. Connell^{32b},
 I.A. Connelly⁹⁸, S. Constantinescu^{27b}, G. Conti³⁵, F. Conventi^{67a,aw}, M. Cooke¹⁸,
 A.M. Cooper-Sarkar¹³¹, F. Cormier¹⁷³, K.J.R. Cormier¹⁶⁵, M. Corradi^{70a,70b}, F. Corriveau^{101,ae},
 A. Cortes-Gonzalez³⁵, G. Costa^{66a}, M.J. Costa¹⁷², D. Costanzo¹⁴⁶, G. Cottin³¹, G. Cowan⁹¹, B.E. Cox⁹⁸,
 K. Cranmer¹²¹, S.J. Crawley⁵⁵, R.A. Creager¹³³, G. Cree³³, S. Crépe-Renaudin⁵⁶, F. Crescioli¹³²,
 W.A. Cribbs^{43a,43b}, M. Cristinziani²⁴, V. Croft¹²¹, G. Crosetti^{40b,40a}, A. Cueto⁹⁶,
 T. Cuhadar Donszelmann¹⁴⁶, A.R. Cukierman¹⁵⁰, J. Cummings¹⁸¹, M. Curatolo⁴⁹, J. Cúth⁹⁷,
 S. Czekaierda⁸², P. Czodrowski³⁵, M.J. Da Cunha Sargedas De Sousa^{136a,136b}, C. Da Via⁹⁸,
 W. Dabrowski^{81a}, T. Dado^{28a,y}, T. Dai¹⁰³, O. Dale¹⁷, F. Dallaire¹⁰⁷, C. Dallapiccola¹⁰⁰, M. Dam³⁹,

G. D'amen^{23b,23a}, J.R. Dandoy¹³³, M.F. Daneri³⁰, N.P. Dang^{179,k}, A.C. Daniells²¹, N.D. Dann⁹⁸, M. Danning¹⁷³, M. Dano Hoffmann¹⁴², V. Dao¹⁵², G. Darbo^{53b}, S. Darmora⁸, J. Dassoulas³, A. Dattagupta¹²⁷, T. Daubney⁴⁴, S. D'Auria⁵⁵, W. Davey²⁴, C. David⁴⁴, T. Davidek¹³⁹, D.R. Davis⁴⁷, P. Davison⁹², E. Dawe¹⁰², I. Dawson¹⁴⁶, K. De⁸, R. De Asmundis^{67a}, A. De Benedetti¹²⁴, S. De Castro^{23b,23a}, S. De Cecco¹³², N. De Groot¹¹⁷, P. de Jong¹¹⁸, H. De la Torre¹⁰⁴, F. De Lorenzi⁷⁶, A. De Maria^{51,u}, D. De Pedis^{70a}, A. De Salvo^{70a}, U. De Sanctis^{71a,71b}, A. De Santo¹⁵³, K. De Vasconcelos Corga⁹⁹, J.B. De Vivie De Regie¹²⁸, R. Debbe²⁹, C. Debenedetti¹⁴³, D.V. Dedovich⁷⁷, N. Dehghanian³, I. Deigaard¹¹⁸, M. Del Gaudio^{40b,40a}, J. Del Peso⁹⁶, D. Delgove¹²⁸, F. Deliot¹⁴², C.M. Delitzsch⁷, M. Della Pietra^{67a,67b}, D. Della Volpe⁵², A. Dell'Acqua³⁵, L. Dell'Asta²⁵, M. Dell'Orso^{69a,69b}, M. Delmastro⁵, C. Delporte¹²⁸, P.A. Delsart⁵⁶, D.A. DeMarco¹⁶⁵, S. Demers¹⁸¹, M. Demichev⁷⁷, A. Demilly¹³², S.P. Denisov¹⁴⁰, D. Denysiuk¹⁴², L. D'Eramo¹³², D. Derendarz⁸², J.E. Derkaoui^{34d}, F. Derue¹³², P. Dervan⁸⁸, K. Desch²⁴, C. Deterre⁴⁴, K. Dette¹⁶⁵, M.R. Devesa³⁰, P.O. Deviveiros³⁵, A. Dewhurst¹⁴¹, S. Dhaliwal²⁶, F.A. Di Bello⁵², A. Di Ciaccio^{71a,71b}, L. Di Ciaccio⁵, W.K. Di Clemente¹³³, C. Di Donato^{67a,67b}, A. Di Girolamo³⁵, B. Di Girolamo³⁵, B. Di Micco^{72a,72b}, R. Di Nardo³⁵, K.F. Di Petrillo⁵⁷, A. Di Simone⁵⁰, R. Di Sipio¹⁶⁵, D. Di Valentino³³, C. Diaconu⁹⁹, M. Diamond¹⁶⁵, F.A. Dias³⁹, M.A. Diaz^{144a}, J. Dickinson¹⁸, E.B. Diehl¹⁰³, J. Dietrich¹⁹, S. Díez Cornell⁴⁴, A. Dimitrievska¹⁶, J. Dingfelder²⁴, P. Dita^{27b}, S. Dita^{27b}, F. Dittus³⁵, F. Djama⁹⁹, T. Djobava^{157b}, J.I. Djuvslund^{59a}, M.A.B. Do Vale^{78c}, D. Dobos³⁵, M. Dobre^{27b}, D. Dodsworth²⁶, C. Doglioni⁹⁴, J. Dolejsi¹³⁹, Z. Dolezal¹³⁹, M. Donadelli^{78d}, S. Donati^{69a,69b}, P. Dondero^{68a,68b}, J. Donini³⁷, M. D'Onofrio⁸⁸, J. Dopke¹⁴¹, A. Doria^{67a}, M.T. Dova⁸⁶, A.T. Doyle⁵⁵, E. Drechsler⁵¹, M. Dris¹⁰, Y. Du^{58b}, J. Duarte-Campderros¹⁵⁹, F. Dubinin¹⁰⁸, A. Dubreuil⁵², E. Duchovni¹⁷⁸, G. Duckeck¹¹², A. Ducourthial¹³², O.A. Ducu^{107,x}, D. Duda¹¹⁸, A. Dudarev³⁵, A.C. Dudder⁹⁷, E.M. Duffield¹⁸, L. Dufflot¹²⁸, M. Dührssen³⁵, C. Dülsen¹⁸⁰, M. Dumancic¹⁷⁸, A.E. Dumitriu^{27b,e}, A.K. Duncan⁵⁵, M. Dunford^{59a}, A. Duperrin⁹⁹, H. Duran Yildiz^{4a}, M. Dören⁵⁴, A. Durglishvili^{157b}, D. Duschinger⁴⁶, B. Dutta⁴⁴, D. Duvnjak¹, M. Dyndal⁴⁴, B.S. Dziedzic⁸², C. Eckardt⁴⁴, K.M. Ecker¹¹³, R.C. Edgar¹⁰³, T. Eifert³⁵, G. Eigen¹⁷, K. Einsweiler¹⁸, T. Ekelof¹⁷⁰, M. El Kacimi^{34c}, R. El Kosseifi⁹⁹, V. Ellajosyula⁹⁹, M. Ellert¹⁷⁰, S. Elles⁵, F. Ellinghaus¹⁸⁰, A.A. Elliot¹⁷⁴, N. Ellis³⁵, J. Elmsheuser²⁹, M. Elsing³⁵, D. Emelianov¹⁴¹, Y. Enari¹⁶¹, J.S. Ennis¹⁷⁶, M.B. Epland⁴⁷, J. Erdmann⁴⁵, A. Ereditato²⁰, M. Ernst²⁹, S. Errede¹⁷¹, M. Escalier¹²⁸, C. Escobar¹⁷², B. Esposito⁴⁹, O. Estrada Pastor¹⁷², A.I. Etienve¹⁴², E. Etzion¹⁵⁹, H. Evans⁶³, A. Ezhilov¹³⁴, M. Ezzi^{34e}, F. Fabbri^{23b,23a}, L. Fabbri^{23b,23a}, V. Fabiani¹¹⁷, G. Facini⁹², R.M. Fakhruddinov¹⁴⁰, S. Falciano^{70a}, R.J. Falla⁹², J. Faltova³⁵, Y. Fang^{15a}, M. Fanti^{66a,66b}, A. Farbin⁸, A. Farilla^{72a}, C. Farina¹³⁵, E.M. Farina^{68a,68b}, T. Farooque¹⁰⁴, S. Farrell¹⁸, S.M. Farrington¹⁷⁶, P. Farthouat³⁵, F. Fassi^{34e}, P. Fassnacht³⁵, D. Fassouliotis⁹, M. Fauci Giannelli⁴⁸, A. Favareto^{53b,53a}, W.J. Fawcett¹³¹, L. Fayard¹²⁸, O.L. Fedin^{134,q}, W. Fedorko¹⁷³, S. Feigl¹³⁰, L. Feligioni⁹⁹, C. Feng^{58b}, E.J. Feng³⁵, M.J. Fenton⁵⁵, A.B. Fenyuk¹⁴⁰, L. Feremenga⁸, P. Fernandez Martinez¹⁷², J. Ferrando⁴⁴, A. Ferrari¹⁷⁰, P. Ferrari¹¹⁸, R. Ferrari^{68a}, D.E. Ferreira de Lima^{59b}, A. Ferrer¹⁷², D. Ferrere⁵², C. Ferretti¹⁰³, F. Fiedler⁹⁷, M. Filipuzzi⁴⁴, A. Filipčić⁸⁹, F. Filthaut¹¹⁷, M. Fincke-Keeler¹⁷⁴, K.D. Finelli²⁵, M.C.N. Fiolhais^{136a,136c,b}, L. Fiorini¹⁷², A. Fischer², C. Fischer¹⁴, J. Fischer¹⁸⁰, W.C. Fisher¹⁰⁴, N. Flaschel⁴⁴, I. Fleck¹⁴⁸, P. Fleischmann¹⁰³, R.R.M. Fletcher¹³³, T. Flick¹⁸⁰, B.M. Flierl¹¹², L.R. Flores Castillo^{61a}, M.J. Flowerdew¹¹³, G.T. Forcolin⁹⁸, A. Formica¹⁴², F.A. Förster¹⁴, A.C. Forti⁹⁸, A.G. Foster²¹, D. Fournier¹²⁸, H. Fox⁸⁷, S. Fracchia¹⁴⁶, P. Francavilla^{69a,69b}, M. Franchini^{23b,23a}, S. Franchino^{59a}, D. Francis³⁵, L. Franconi¹³⁰, M. Franklin⁵⁷, M. Frate¹⁶⁹, M. Fraternali^{68a,68b}, D. Freeborn⁹², S.M. Fressard-Batraneanu³⁵, B. Freund¹⁰⁷, D. Froidevaux³⁵, J.A. Frost¹³¹, C. Fukunaga¹⁶², T. Fusayasu¹¹⁴, J. Fuster¹⁷², O. Gabizon¹⁵⁸, A. Gabrielli^{23b,23a}, A. Gabrielli¹⁸, G.P. Gach^{81a}, S. Gadatsch³⁵, S. Gadomski⁵², G. Gagliardi^{53b,53a}, L.G. Gagnon¹⁰⁷, C. Galea¹¹⁷, B. Galhardo^{136a,136c}, E.J. Gallas¹³¹, B.J. Gallop¹⁴¹, P. Gallus¹³⁸, G. Galster³⁹, K.K. Gan¹²², S. Ganguly³⁷, Y. Gao⁸⁸, Y.S. Gao^{150,m}, C. García¹⁷²,

J.E. García Navarro¹⁷², J.A. García Pascual^{15a}, M. Garcia-Sciveres¹⁸, R.W. Gardner³⁶, N. Garelli¹⁵⁰, V. Garonne¹³⁰, A. Gascon Bravo⁴⁴, K. Gasnikova⁴⁴, C. Gatti⁴⁹, A. Gaudiello^{53b,53a}, G. Gaudio^{68a}, I.L. Gavrilenko¹⁰⁸, C. Gay¹⁷³, G. Gaycken²⁴, E.N. Gazis¹⁰, C.N.P. Gee¹⁴¹, J. Geisen⁵¹, M. Geisen⁹⁷, M.P. Geisler^{59a}, K. Gellerstedt^{43a,43b}, C. Gemme^{53b}, M.H. Genest⁵⁶, C. Geng¹⁰³, S. Gentile^{70a,70b}, C. Gentsos¹⁶⁰, S. George⁹¹, D. Gerbaudo¹⁴, G. Gessner⁴⁵, S. Ghasemi¹⁴⁸, M. Ghneimat²⁴, B. Giacobbe^{23b}, S. Giagu^{70a,70b}, N. Giangiacomi^{23b,23a}, P. Giannetti^{69a}, S.M. Gibson⁹¹, M. Gignac¹⁷³, M. Gilchriese¹⁸, D. Gillberg³³, G. Gilles¹⁸⁰, D.M. Gingrich^{3,av}, M.P. Giordani^{64a,64c}, F.M. Giorgi^{23b}, P.F. Giraud¹⁴², P. Giromini⁵⁷, G. Giugliarelli^{64a,64c}, D. Giugni^{66a}, F. Giuli¹³¹, C. Giuliani¹¹³, M. Giulini^{59b}, B.K. Gjelsten¹³⁰, S. Gkaitatzis¹⁶⁰, I. Gkialas^{9j}, E.L. Gkoukousis¹⁴, P. Gkoutoumis¹⁰, L.K. Gladilin¹¹¹, C. Glasman⁹⁶, J. Glatzer¹⁴, P.C.F. Glaysher⁴⁴, A. Glazov⁴⁴, M. Goblirsch-Kolb²⁶, J. Godlewski⁸², S. Goldfarb¹⁰², T. Golling⁵², D. Golubkov¹⁴⁰, A. Gomes^{136a,136b,136d}, R. Goncalves Gama^{78b}, J. Goncalves Pinto Firmino Da Costa¹⁴², R. Gonçalves^{136a}, G. Gonella⁵⁰, L. Gonella²¹, A. Gongadze⁷⁷, J.L. Gonski⁵⁷, S. González de la Hoz¹⁷², S. Gonzalez-Sevilla⁵², L. Goossens³⁵, P.A. Gorbounov¹⁰⁹, H.A. Gordon²⁹, I. Gorelov¹¹⁶, B. Gorini³⁵, E. Gorini^{65a,65b}, A. Gorišek⁸⁹, A.T. Goshaw⁴⁷, C. Gössling⁴⁵, M.I. Gostkin⁷⁷, C.A. Gottardo²⁴, C.R. Goudet¹²⁸, D. Goujdami^{34c}, A.G. Goussiou¹⁴⁵, N. Govender^{32b,c}, E. Gozani¹⁵⁸, I. Grabowska-Bold^{81a}, P.O.J. Gradin¹⁷⁰, J. Gramling¹⁶⁹, E. Gramstad¹³⁰, S. Grancagnolo¹⁹, V. Gratchev¹³⁴, P.M. Gravila^{27f}, C. Gray⁵⁵, H.M. Gray¹⁸, Z.D. Greenwood^{93,aj}, C. Grefe²⁴, K. Gregersen⁹², I.M. Gregor⁴⁴, P. Grenier¹⁵⁰, K. Grevtsov⁵, J. Griffiths⁸, A.A. Grillo¹⁴³, K. Grimm⁸⁷, S. Grinstein^{14,z}, Ph. Gris³⁷, J.-F. Grivaz¹²⁸, S. Groh⁹⁷, E. Gross¹⁷⁸, J. Grosse-Knetter⁵¹, G.C. Grossi⁹³, Z.J. Grout⁹², A. Grummer¹¹⁶, L. Guan¹⁰³, W. Guan¹⁷⁹, J. Guenther³⁵, F. Guescini^{166a}, D. Guest¹⁶⁹, O. Gueta¹⁵⁹, B. Gui¹²², E. Guido^{53b,53a}, T. Guillemin⁵, S. Guindon³⁵, U. Gul⁵⁵, C. Gumpert³⁵, J. Guo^{58c}, W. Guo¹⁰³, Y. Guo^{58a,t}, R. Gupta⁴¹, S. Gurbuz^{12c}, G. Gustavino¹²⁴, B.J. Gutelman¹⁵⁸, P. Gutierrez¹²⁴, N.G. Gutierrez Ortiz⁹², C. Gutsche⁹², C. Guyot¹⁴², M.P. Guzik^{81a}, C. Gwenlan¹³¹, C.B. Gwilliam⁸⁸, A. Haas¹²¹, C. Haber¹⁸, H.K. Hadavand⁸, N. Haddad^{34e}, A. Hader⁹⁹, S. Hageböck²⁴, M. Hagihara¹⁶⁷, H. Hakobyan^{182,*}, M. Haleem⁴⁴, J. Haley¹²⁵, G. Halladjian¹⁰⁴, G.D. Hallowell⁹⁹, K. Hamacher¹⁸⁰, P. Hamal¹²⁶, K. Hamano¹⁷⁴, A. Hamilton^{32a}, G.N. Hamity¹⁴⁶, P.G. Hamnett⁴⁴, L. Han^{58a}, S. Han^{15d}, K. Hanagaki^{79,w}, K. Hanawa¹⁶¹, M. Hance¹⁴³, D.M. Handl¹¹², B. Haney¹³³, P. Hanke^{59a}, J.B. Hansen³⁹, J.D. Hansen³⁹, M.C. Hansen²⁴, P.H. Hansen³⁹, K. Hara¹⁶⁷, A.S. Hard¹⁷⁹, T. Harenberg¹⁸⁰, F. Hariri¹²⁸, S. Harkusha¹⁰⁵, P.F. Harrison¹⁷⁶, N.M. Hartmann¹¹², Y. Hasegawa¹⁴⁷, A. Hasib⁴⁸, S. Hassani¹⁴², S. Haug²⁰, R. Hauser¹⁰⁴, L. Hauswald⁴⁶, L.B. Havener³⁸, M. Havranek¹³⁸, C.M. Hawkes²¹, R.J. Hawkins³⁵, D. Hayakawa¹⁶³, D. Hayden¹⁰⁴, C.P. Hays¹³¹, J.M. Hays⁹⁰, H.S. Hayward⁸⁸, S.J. Haywood¹⁴¹, S.J. Head²¹, T. Heck⁹⁷, V. Hedberg⁹⁴, L. Heelan⁸, S. Heer²⁴, K.K. Heidegger⁵⁰, S. Heim⁴⁴, T. Heim¹⁸, B. Heinemann^{44,aq}, J.J. Heinrich¹¹², L. Heinrich¹²¹, C. Heinz⁵⁴, J. Hejbal¹³⁷, L. Helary³⁵, A. Held¹⁷³, S. Hellman^{43a,43b}, C. Hensens³⁵, R.C.W. Henderson⁸⁷, Y. Heng¹⁷⁹, S. Henkelmann¹⁷³, A.M. Henriques Correia³⁵, S. Henrot-Versille¹²⁸, G.H. Herbert¹⁹, H. Herde²⁶, V. Herget¹⁷⁵, Y. Hernández Jiménez^{32c}, H. Herr⁹⁷, G. Herten⁵⁰, R. Hertenberger¹¹², L. Hervas³⁵, T.C. Herwig¹³³, G.G. Hesketh⁹², N.P. Hessey^{166a}, J.W. Hetherly⁴¹, S. Higashino⁷⁹, E. Higón-Rodríguez¹⁷², K. Hildebrand³⁶, E. Hill¹⁷⁴, J.C. Hill³¹, K.H. Hiller⁴⁴, S.J. Hillier²¹, M. Hils⁴⁶, I. Hinchliffe¹⁸, M. Hirose⁵⁰, D. Hirschbuehl¹⁸⁰, B. Hiti⁸⁹, O. Hladik¹³⁷, D.R. Hlaluku^{32c}, X. Hoad⁴⁸, J. Hobbs¹⁵², N. Hod^{166a}, M.C. Hodgkinson¹⁴⁶, P. Hodgson¹⁴⁶, A. Hoecker³⁵, M.R. Hoferkamp¹¹⁶, F. Hoenig¹¹², D. Hohn²⁴, T.R. Holmes³⁶, M. Holzbock¹¹², M. Homann⁴⁵, S. Honda¹⁶⁷, T. Honda⁷⁹, T.M. Hong¹³⁵, B.H. Hooberman¹⁷¹, W.H. Hopkins¹²⁷, Y. Horii¹¹⁵, A.J. Horton¹⁴⁹, J.-Y. Hostachy⁵⁶, A. Hostiuc¹⁴⁵, S. Hou¹⁵⁵, A. Hoummada^{34a}, J. Howarth⁹⁸, J. Hoya⁸⁶, M. Hrabovsky¹²⁶, J. Hrdinka³⁵, I. Hristova¹⁹, J. Hrivnac¹²⁸, A. Hrynevich¹⁰⁶, T. Hryn'ova⁵, P.J. Hsu⁶², S.-C. Hsu¹⁴⁵, Q. Hu²⁹, S. Hu^{58c}, Y. Huang^{15a}, Z. Hubacek¹³⁸, F. Hubaut⁹⁹, F. Huegging²⁴, T.B. Huffman¹³¹, E.W. Hughes³⁸, M. Huhtinen³⁵, R.F.H. Hunter³³, P. Huo¹⁵², N. Huseynov^{77,ag}, J. Huston¹⁰⁴, J. Huth⁵⁷, R. Hyneman¹⁰³, G. Iacobucci⁵²,

G. Iakovidis²⁹, I. Ibragimov¹⁴⁸, L. Iconomidou-Fayard¹²⁸, Z. Idrissi^{34e}, P. Iengo³⁵, O. Igonkina^{118,ac}, T. Iizawa¹⁷⁷, Y. Ikegami⁷⁹, M. Ikeno⁷⁹, Y. Ilchenko¹¹, D. Iliadis¹⁶⁰, N. Ilic¹⁵⁰, F. Iltzsche⁴⁶, G. Introzzi^{68a,68b}, P. Ioannou^{9,*}, M. Iodice^{72a}, K. Iordanidou³⁸, V. Ippolito⁵⁷, M.F. Isacson¹⁷⁰, N. Ishijima¹²⁹, M. Ishino¹⁶¹, M. Ishitsuka¹⁶³, C. Issever¹³¹, S. Istin^{12c,an}, F. Ito¹⁶⁷, J.M. Iturbe Ponce^{61a}, R. Iuppa^{73a,73b}, H. Iwasaki⁷⁹, J.M. Izen⁴², V. Izzo^{67a}, S. Jabbar³, P. Jackson¹, R.M. Jacobs²⁴, V. Jain², G. Jäkel¹⁸⁰, K.B. Jakobi⁹⁷, K. Jakobs⁵⁰, S. Jakobsen⁷⁴, T. Jakoubek¹³⁷, D.O. Jamin¹²⁵, D.K. Jana⁹³, R. Jansky⁵², J. Janssen²⁴, M. Janus⁵¹, P.A. Janus^{81a}, G. Jarlskog⁹⁴, N. Javadov^{77,ag}, T. Javůrek⁵⁰, M. Javurkova⁵⁰, F. Jeanneau¹⁴², L. Jeanty¹⁸, J. Jejelava^{157a,ah}, A. Jelinskas¹⁷⁶, P. Jenni^{50,d}, C. Jeske¹⁷⁶, S. Jézéquel⁵, H. Ji¹⁷⁹, J. Jia¹⁵², H. Jiang⁷⁶, Y. Jiang^{58a}, Z. Jiang¹⁵⁰, S. Jiggins⁹², J. Jimenez Pena¹⁷², S. Jin^{15c}, A. Jinaru^{27b}, O. Jinnouchi¹⁶³, H. Jivan^{32c}, P. Johansson¹⁴⁶, K.A. Johns⁷, C.A. Johnson⁶³, W.J. Johnson¹⁴⁵, K. Jon-And^{43a,43b}, R.W.L. Jones⁸⁷, S.D. Jones¹⁵³, S. Jones⁷, T.J. Jones⁸⁸, J. Jongmanns^{59a}, P.M. Jorge^{136a,136b}, J. Jovicevic^{166a}, X. Ju¹⁷⁹, A. Juste Rozas^{14,z}, A. Kaczmarska⁸², M. Kado¹²⁸, H. Kagan¹²², M. Kagan¹⁵⁰, S.J. Kahn⁹⁹, T. Kaji¹⁷⁷, E. Kajomovitz¹⁵⁸, C.W. Kalderon⁹⁴, A. Kaluza⁹⁷, S. Kama⁴¹, A. Kamenshchikov¹⁴⁰, N. Kanaya¹⁶¹, L. Kanjir⁸⁹, V.A. Kantserov¹¹⁰, J. Kanzaki⁷⁹, B. Kaplan¹²¹, L.S. Kaplan¹⁷⁹, D. Kar^{32c}, K. Karakostas¹⁰, N. Karastathis¹⁰, M.J. Kareem^{166b}, E. Karentzos¹⁰, S.N. Karpov⁷⁷, Z.M. Karpova⁷⁷, K. Karthik¹²¹, V. Kartvelishvili⁸⁷, A.N. Karyukhin¹⁴⁰, K. Kasahara¹⁶⁷, L. Kashif¹⁷⁹, R.D. Kass¹²², A. Kastanas¹⁵¹, Y. Kataoka¹⁶¹, C. Kato¹⁶¹, A. Katre⁵², J. Katzy⁴⁴, K. Kawade⁸⁰, K. Kawagoe⁸⁵, T. Kawamoto¹⁶¹, G. Kawamura⁵¹, E.F. Kay⁸⁸, V.F. Kazanin^{120b,120a}, R. Keeler¹⁷⁴, R. Kehoe⁴¹, J.S. Keller³³, E. Kellermann⁹⁴, J.J. Kempster⁹¹, J. Kendrick²¹, H. Keoshkerian¹⁶⁵, O. Kepka¹³⁷, S. Kersten¹⁸⁰, B.P. Kerševan⁸⁹, R.A. Keyes¹⁰¹, M. Khader¹⁷¹, F. Khalil-Zada¹³, A. Khanov¹²⁵, A.G. Kharlamov^{120b,120a}, T. Kharlamova^{120b,120a}, A. Khodinov¹⁶⁴, T.J. Khoo⁵², V. Khovanskiy^{109,*}, E. Khramov⁷⁷, J. Khubua^{157b}, S. Kido⁸⁰, C.R. Kilby⁹¹, H.Y. Kim⁸, S.H. Kim¹⁶⁷, Y.K. Kim³⁶, N. Kimura¹⁶⁰, O.M. Kind¹⁹, B.T. King⁸⁸, D. Kirchmeier⁴⁶, J. Kirk¹⁴¹, A.E. Kiryunin¹¹³, T. Kishimoto¹⁶¹, D. Kisielewska^{81a}, V. Kitali⁴⁴, O. Kivernyk⁵, E. Kladiva^{28b}, T. Klapdor-Kleingrothaus⁵⁰, M.H. Klein¹⁰³, M. Klein⁸⁸, U. Klein⁸⁸, K. Kleinknecht⁹⁷, P. Klimek¹¹⁹, A. Klimentov²⁹, R. Klingenberg^{45,*}, T. Klingl²⁴, T. Klioutchnikova³⁵, F.F. Klitzner¹¹², P. Kluit¹¹⁸, S. Kluth¹¹³, E. Kneringer⁷⁴, E.B.F.G. Knoops⁹⁹, A. Knue¹¹³, A. Kobayashi¹⁶¹, D. Kobayashi⁸⁵, T. Kobayashi¹⁶¹, M. Kobel⁴⁶, M. Kocian¹⁵⁰, P. Kodys¹³⁹, T. Koffas³³, E. Koffeman¹¹⁸, M.K. Köhler¹⁷⁸, N.M. Köhler¹¹³, T. Koi¹⁵⁰, M. Kolb^{59b}, I. Koletsou⁵, A.A. Komar^{108,*}, T. Kondo⁷⁹, N. Kondrashova^{58c}, K. Köneke⁵⁰, A.C. König¹¹⁷, T. Kono^{79,ap}, R. Konoplich^{121,ak}, N. Konstantinidis⁹², B. Konya⁹⁴, R. Kopeliansky⁶³, S. Koperny^{81a}, A.K. Kopp⁵⁰, K. Korcyl⁸², K. Kordas¹⁶⁰, A. Korn⁹², A.A. Korol^{120b,120a,ao}, I. Korolkov¹⁴, E.V. Korolkova¹⁴⁶, O. Kortner¹¹³, S. Kortner¹¹³, T. Kosek¹³⁹, V.V. Kostyukhin²⁴, A. Kotwal⁴⁷, A. Koulouris¹⁰, A. Kourkoumeli-Charalampidi^{68a,68b}, C. Kourkoumelis⁹, E. Kourlitis¹⁴⁶, V. Kouskoura²⁹, A.B. Kowalewska⁸², R. Kowalewski¹⁷⁴, T.Z. Kowalski^{81a}, C. Kozakai¹⁶¹, W. Kozanecki¹⁴², A.S. Kozhin¹⁴⁰, V.A. Kramarenko¹¹¹, G. Kramberger⁸⁹, D. Krasnopevtsev¹¹⁰, M.W. Krasny¹³², A. Krasznahorkay³⁵, D. Krauss¹¹³, J.A. Kremer^{81a}, J. Kretzschmar⁸⁸, K. Kreutzfeldt⁵⁴, P. Krieger¹⁶⁵, K. Krizka¹⁸, K. Kroeninger⁴⁵, H. Kroha¹¹³, J. Kroll¹³⁷, J. Kroll¹³³, J. Kroseberg²⁴, J. Krstic¹⁶, U. Kruchonak⁷⁷, H. Krüger²⁴, N. Krumnack⁷⁶, M.C. Kruse⁴⁷, T. Kubota¹⁰², H. Kucuk⁹², S. Kudah^{4b}, J.T. Kuechler¹⁸⁰, S. Kuehn³⁵, A. Kugel^{59a}, F. Kuger¹⁷⁵, T. Kuhl⁴⁴, V. Kukhtin⁷⁷, R. Kukla⁹⁹, Y. Kulchitsky¹⁰⁵, S. Kuleshov^{144b}, Y.P. Kulinich¹⁷¹, M. Kuna^{70a,70b}, T. Kunigo⁸³, A. Kupco¹³⁷, T. Kupfer⁴⁵, O. Kuprash¹⁵⁹, H. Kurashige⁸⁰, L.L. Kurchaninov^{166a}, Y.A. Kurochkin¹⁰⁵, M.G. Kurth^{15d}, E.S. Kuwertz¹⁷⁴, M. Kuze¹⁶³, J. Kvita¹²⁶, T. Kwan¹⁷⁴, D. Kyriazopoulos¹⁴⁶, A. La Rosa¹¹³, J.L. La Rosa Navarro^{78d}, L. La Rotonda^{40b,40a}, F. La Ruffa^{40b,40a}, C. Lacasta¹⁷², F. Lacava^{70a,70b}, J. Lacey⁴⁴, D.P.J. Lack⁹⁸, H. Lacker¹⁹, D. Lacour¹³², E. Ladygin⁷⁷, R. Lafaye⁵, B. Laforge¹³², S. Lai⁵¹, S. Lammers⁶³, W. Lampl⁷, E. Lançon²⁹, U. Landgraf⁵⁰, M.P.J. Landon⁹⁰, M.C. Lanfermann⁵², V.S. Lang⁴⁴, J.C. Lange¹⁴, R.J. Langenberg³⁵, A.J. Lankford¹⁶⁹, F. Lanni²⁹, K. Lantzsich²⁴, A. Lanza^{68a},

A. Lapertosa^{53b,53a}, S. Laplace¹³², J.F. Laporte¹⁴², T. Lari^{66a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁵,
 T.S. Lau^{61a}, P. Laurelli⁴⁹, W. Lavrijsen¹⁸, A.T. Law¹⁴³, P. Laycock⁸⁸, T. Lazovich⁵⁷, M. Lazzaroni^{66a,66b},
 B. Le¹⁰², O. Le Dortz¹³², E. Le Guirriec⁹⁹, E.P. Le Quilleuc¹⁴², M. LeBlanc¹⁷⁴, T. LeCompte⁶,
 F. Ledroit-Guillon⁵⁶, C.A. Lee²⁹, G.R. Lee^{144a}, L. Lee⁵⁷, S.C. Lee¹⁵⁵, B. Lefebvre¹⁰¹, G. Lefebvre¹³²,
 M. Lefebvre¹⁷⁴, F. Legger¹¹², C. Leggett¹⁸, G. Lehmann Miotto³⁵, X. Lei⁷, W.A. Leight⁴⁴,
 M.A.L. Leite^{78d}, R. Leitner¹³⁹, D. Lellouch¹⁷⁸, B. Lemmer⁵¹, K.J.C. Leney⁹², T. Lenz²⁴, B. Lenzi³⁵,
 R. Leone⁷, S. Leone^{69a}, C. Leonidopoulos⁴⁸, G. Lerner¹⁵³, C. Leroy¹⁰⁷, R. Les¹⁶⁵, A.A.J. Lesage¹⁴²,
 C.G. Lester³¹, M. Levchenko¹³⁴, J. Levêque⁵, D. Levin¹⁰³, L.J. Levinson¹⁷⁸, M. Levy²¹, D. Lewis⁹⁰,
 B. Li^{58a,t}, C-Q. Li^{58a}, H. Li¹⁵², L. Li^{58c}, Q. Li^{15d}, Q.Y. Li^{58a}, S. Li⁴⁷, X. Li^{58c}, Y. Li¹⁴⁸, Z. Liang^{15a},
 B. Liberti^{71a}, A. Liblong¹⁶⁵, K. Lie^{61c}, J. Liebal²⁴, W. Liebig¹⁷, A. Limosani¹⁵⁴, C.Y. Lin³¹, K. Lin¹⁰⁴,
 S.C. Lin¹⁵⁶, T.H. Lin⁹⁷, R.A. Linck⁶³, B.E. Lindquist¹⁵², A.L. Lioni⁵², E. Lipeles¹³³, A. Lipniacka¹⁷,
 M. Lisovyi^{59b}, T.M. Liss^{171,as}, A. Lister¹⁷³, A.M. Litke¹⁴³, B. Liu⁷⁶, H.B. Liu²⁹, H. Liu¹⁰³, J.B. Liu^{58a},
 J.K.K. Liu¹³¹, J. Liu^{58b}, K. Liu⁹⁹, L. Liu¹⁷¹, M. Liu^{58a}, Y.L. Liu^{58a}, Y.W. Liu^{58a}, M. Livan^{68a,68b},
 A. Lleres⁵⁶, J. Llorente Merino^{15a}, S.L. Lloyd⁹⁰, C.Y. Lo^{61b}, F. Lo Sterzo⁴¹, E.M. Lobodzinska⁴⁴,
 P. Loch⁷, F.K. Loebinger⁹⁸, A. Loesle⁵⁰, K.M. Loew²⁶, T. Lohse¹⁹, K. Lohwasser¹⁴⁶, M. Lokajicek¹³⁷,
 B.A. Long²⁵, J.D. Long¹⁷¹, R.E. Long⁸⁷, L. Longo^{65a,65b}, K.A. Looper¹²², J.A. Lopez^{144b},
 I. Lopez Paz¹⁴, A. Lopez Solis¹³², J. Lorenz¹¹², N. Lorenzo Martinez⁵, M. Losada²², P.J. Lösel¹¹²,
 X. Lou^{15a}, A. Lounis¹²⁸, J. Love⁶, P.A. Love⁸⁷, H. Lu^{61a}, N. Lu¹⁰³, Y.J. Lu⁶², H.J. Lubatti¹⁴⁵,
 C. Luci^{70a,70b}, A. Lucotte⁵⁶, C. Luedtke⁵⁰, F. Luehring⁶³, W. Lukas⁷⁴, L. Luminari^{70a},
 O. Lundberg^{43a,43b}, B. Lund-Jensen¹⁵¹, M.S. Lutz¹⁰⁰, P.M. Luzi¹³², D. Lynn²⁹, R. Lysak¹³⁷, E. Lytken⁹⁴,
 F. Lyu^{15a}, V. Lyubushkin⁷⁷, H. Ma²⁹, L.L. Ma^{58b}, Y. Ma^{58b}, G. Maccarrone⁴⁹, A. Macchiolo¹¹³,
 C.M. Macdonald¹⁴⁶, J. Machado Miguens^{133,136b}, D. Madaffari¹⁷², R. Madar³⁷, W.F. Mader⁴⁶,
 A. Madsen⁴⁴, N. Madysa⁴⁶, J. Maeda⁸⁰, S. Maeland¹⁷, T. Maeno²⁹, A.S. Maevskiy¹¹¹, V. Magerl⁵⁰,
 C. Maiani¹²⁸, C. Maidantchik^{78b}, T. Maier¹¹², A. Maio^{136a,136b,136d}, O. Majersky^{28a}, S. Majewski¹²⁷,
 Y. Makida⁷⁹, N. Makovec¹²⁸, B. Malaescu¹³², Pa. Malecki⁸², V.P. Maleev¹³⁴, F. Malek⁵⁶, U. Mallik⁷⁵,
 D. Malon⁶, C. Malone³¹, S. Maltezos¹⁰, S. Malyukov³⁵, J. Mamuzic¹⁷², G. Mancini⁴⁹, I. Mandić⁸⁹,
 J. Maneira^{136a,136b}, L. Manhaes de Andrade Filho^{78a}, J. Manjarres Ramos⁴⁶, K.H. Mankinen⁹⁴,
 A. Mann¹¹², A. Manousos³⁵, B. Mansoulie¹⁴², J.D. Mansour^{15a}, R. Mantifel¹⁰¹, M. Mantoani⁵¹,
 S. Manzoni^{66a,66b}, L. Mapelli³⁵, G. Marceca³⁰, L. March⁵², L. Marchese¹³¹, G. Marchiori¹³²,
 M. Marcisovsky¹³⁷, C.A. Marin Tobon³⁵, M. Marjanovic³⁷, D.E. Marley¹⁰³, F. Marroquim^{78b},
 S.P. Marsden⁹⁸, Z. Marshall¹⁸, M.U.F. Martensson¹⁷⁰, S. Marti-Garcia¹⁷², C.B. Martin¹²²,
 T.A. Martin¹⁷⁶, V.J. Martin⁴⁸, B. Martin dit Latour¹⁷, M. Martinez^{14,z}, V.I. Martinez Outschoorn¹⁷¹,
 S. Martin-Haugh¹⁴¹, V.S. Martoiu^{27b}, A.C. Martyniuk⁹², A. Marzin³⁵, L. Masetti⁹⁷, T. Mashimo¹⁶¹,
 R. Mashinistov¹⁰⁸, J. Masik⁹⁸, A.L. Maslennikov^{120b,120a}, L.H. Mason¹⁰², L. Massa^{71a,71b},
 P. Mastrandrea⁵, A. Mastroberardino^{40b,40a}, T. Masubuchi¹⁶¹, P. Mättig¹⁸⁰, J. Maurer^{27b}, B. Maček⁸⁹,
 S.J. Maxfield⁸⁸, D.A. Maximov^{120b,120a}, R. Mazini¹⁵⁵, I. Maznas¹⁶⁰, S.M. Mazza^{66a,66b},
 N.C. Mc Fadden¹¹⁶, G. Mc Goldrick¹⁶⁵, S.P. Mc Kee¹⁰³, A. McCarn¹⁰³, R.L. McCarthy¹⁵²,
 T.G. McCarthy¹¹³, L.I. McClymont⁹², E.F. McDonald¹⁰², J.A. Mcfayden³⁵, G. Mchedlidze⁵¹,
 S.J. McMahon¹⁴¹, P.C. McNamara¹⁰², C.J. McNicol¹⁷⁶, R.A. McPherson^{174,ae}, S. Meehan¹⁴⁵, T. Megy⁵⁰,
 S. Mehlhase¹¹², A. Mehta⁸⁸, T. Meideck⁵⁶, B. Meirose⁴², D. Melini^{172,h}, B.R. Mellado Garcia^{32c},
 J.D. Mellenthin⁵¹, M. Melo^{28a}, F. Meloni²⁰, A. Melzer²⁴, S.B. Menary⁹⁸, L. Meng⁸⁸, X.T. Meng¹⁰³,
 A. Mengarelli^{23b,23a}, S. Menke¹¹³, E. Meoni^{40b,40a}, S. Mergelmeyer¹⁹, C. Merlassino²⁰, P. Mermod⁵²,
 L. Merola^{67a,67b}, C. Meroni^{66a}, F.S. Merritt³⁶, A. Messina^{70a,70b}, J. Metcalfe⁶, A.S. Mete¹⁶⁹,
 C. Meyer¹³³, J. Meyer¹¹⁸, J-P. Meyer¹⁴², H. Meyer Zu Theenhausen^{59a}, F. Miano¹⁵³, R.P. Middleton¹⁴¹,
 S. Miglioranzani^{53b,53a}, L. Mijović⁴⁸, G. Mikenberg¹⁷⁸, M. Mikesikova¹³⁷, M. Mikuž⁸⁹, M. Milesi¹⁰²,
 A. Milic¹⁶⁵, D.A. Millar⁹⁰, D.W. Miller³⁶, C. Mills⁴⁸, A. Milov¹⁷⁸, D.A. Milstead^{43a,43b},
 A.A. Minaenko¹⁴⁰, Y. Minami¹⁶¹, I.A. Minashvili^{157b}, A.I. Mincer¹²¹, B. Mindur^{81a}, M. Mineev⁷⁷,

Y. Minegishi¹⁶¹, Y. Ming¹⁷⁹, L.M. Mir¹⁴, A. Mirto^{65a,65b}, K.P. Mistry¹³³, T. Mitani¹⁷⁷, J. Mitrevski¹¹²,
 V.A. Mitsou¹⁷², A. Miucci²⁰, P.S. Miyagawa¹⁴⁶, A. Mizukami⁷⁹, J.U. Mjörnmark⁹⁴, T. Mkrtchyan¹⁸²,
 M. Mlynarikova¹³⁹, T. Moa^{43a,43b}, K. Mochizuki¹⁰⁷, P. Mogg⁵⁰, S. Mohapatra³⁸, S. Molander^{43a,43b},
 R. Moles-Valls²⁴, M.C. Mondragon¹⁰⁴, K. Mönig⁴⁴, J. Monk³⁹, E. Monnier⁹⁹, A. Montalbano¹⁵²,
 J. Montejo Berlingen³⁵, F. Monticelli⁸⁶, S. Monzani^{66a}, R.W. Moore³, N. Morange¹²⁸, D. Moreno²²,
 M. Moreno Llácer³⁵, P. Moretini^{53b}, S. Morgenstern³⁵, D. Mori¹⁴⁹, T. Mori¹⁶¹, M. Morii⁵⁷,
 M. Morinaga¹⁷⁷, V. Morisbak¹³⁰, A.K. Morley³⁵, G. Mornacchi³⁵, J.D. Morris⁹⁰, L. Morvaj¹⁵²,
 P. Moschovakos¹⁰, M. Mosidze^{157b}, H.J. Moss¹⁴⁶, J. Moss^{150,n}, K. Motohashi¹⁶³, R. Mount¹⁵⁰,
 E. Mountricha²⁹, E.J.W. Moyse¹⁰⁰, S. Muanza⁹⁹, F. Mueller¹¹³, J. Mueller¹³⁵, R.S.P. Mueller¹¹²,
 D. Muenstermann⁸⁷, P. Mullen⁵⁵, G.A. Mullier²⁰, F.J. Munoz Sanchez⁹⁸, W.J. Murray^{176,141},
 H. Musheghyan³⁵, M. Muškinja⁸⁹, A.G. Myagkov^{140,al}, M. Myska¹³⁸, B.P. Nachman¹⁸,
 O. Nackenhorst⁵², K. Nagai¹³¹, R. Nagai^{79,ap}, K. Nagano⁷⁹, Y. Nagasaka⁶⁰, K. Nagata¹⁶⁷, M. Nagel⁵⁰,
 E. Nagy⁹⁹, A.M. Nairz³⁵, Y. Nakahama¹¹⁵, K. Nakamura⁷⁹, T. Nakamura¹⁶¹, I. Nakano¹²³,
 R.F. Naranjo Garcia⁴⁴, R. Narayan¹¹, D.I. Narrias Villar^{59a}, I. Naryshkin¹³⁴, T. Naumann⁴⁴,
 G. Navarro²², R. Nayyar⁷, H.A. Neal¹⁰³, P.Y. Nechaeva¹⁰⁸, T.J. Neep¹⁴², A. Negri^{68a,68b}, M. Negrini^{23b},
 S. Nektarijevic¹¹⁷, C. Nellist⁵¹, A. Nelson¹⁶⁹, M.E. Nelson¹³¹, S. Nemecek¹³⁷, P. Nemethy¹²¹,
 M. Nessi^{35,f}, M.S. Neubauer¹⁷¹, M. Neumann¹⁸⁰, P.R. Newman²¹, T.Y. Ng^{61c}, Y.S. Ng¹⁹,
 T. Nguyen Manh¹⁰⁷, R.B. Nickerson¹³¹, R. Nicolaidou¹⁴², J. Nielsen¹⁴³, N. Nikiforou¹¹,
 V. Nikolaenko^{140,al}, I. Nikolic-Audit¹³², K. Nikolopoulos²¹, P. Nilsson²⁹, Y. Ninomiya⁷⁹, A. Nisati^{70a},
 N. Nishu^{58c}, R. Nisius¹¹³, I. Nitsche⁴⁵, T. Nitta¹⁷⁷, T. Nobe¹⁶¹, Y. Noguchi⁸³, M. Nomachi¹²⁹,
 I. Nomidis³³, M.A. Nomura²⁹, T. Nooney⁹⁰, M. Nordberg³⁵, N. Norjoharuddeen¹³¹, O. Novgorodova⁴⁶,
 M. Nozaki⁷⁹, L. Nozka¹²⁶, K. Ntekas¹⁶⁹, E. Nurse⁹², F. Nuti¹⁰², F.G. Oakham^{33,av}, H. Oberlack¹¹³,
 T. Obermann²⁴, J. Ocariz¹³², A. Ochi⁸⁰, I. Ochoa³⁸, J.P. Ochoa-Ricoux^{144a}, K. O'Connor²⁶, S. Oda⁸⁵,
 S. Odaka⁷⁹, A. Oh⁹⁸, S.H. Oh⁴⁷, C.C. Ohm¹⁵¹, H. Ohman¹⁷⁰, H. Oide^{53b,53a,*}, H. Okawa¹⁶⁷,
 Y. Okumura¹⁶¹, T. Okuyama⁷⁹, A. Olariu^{27b}, L.F. Oleiro Seabra^{136a}, S.A. Olivares Pino^{144a},
 D. Oliveira Damazio²⁹, M.J.R. Olsson³⁶, A. Olszewski⁸², J. Olszowska⁸², D.C. O'Neil¹⁴⁹,
 A. Onofre^{136a,136e}, K. Onogi¹¹⁵, P.U.E. Onyisi¹¹, H. Oppen¹³⁰, M.J. Oreglia³⁶, Y. Oren¹⁵⁹,
 D. Orestano^{72a,72b}, N. Orlando^{61b}, A.A. O'Rourke⁴⁴, R.S. Orr¹⁶⁵, B. Osculati^{53b,53a,*}, V. O'Shea⁵⁵,
 R. Ospanov^{58a}, G. Otero y Garzon³⁰, H. Otono⁸⁵, M. Ouchrif^{34d}, F. Ould-Saada¹³⁰, A. Ouraou¹⁴²,
 K.P. Oussoren¹¹⁸, Q. Ouyang^{15a}, M. Owen⁵⁵, R.E. Owen²¹, V.E. Ozcan^{12c}, N. Ozturk⁸, K. Pachal¹⁴⁹,
 A. Pacheco Pages¹⁴, L. Pacheco Rodriguez¹⁴², C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, M. Paganini¹⁸¹,
 F. Paige²⁹, G. Palacino⁶³, S. Palazzo^{40b,40a}, S. Palestini³⁵, M. Palka^{81b}, D. Pallin³⁷,
 E.St. Panagiotopoulou¹⁰, I. Panagoulas¹⁰, C.E. Pandini⁵², J.G. Panduro Vazquez⁹¹, P. Pani³⁵,
 S. Panitkin²⁹, D. Pantea^{27b}, L. Paolozzi⁵², T.D. Papadopoulou¹⁰, K. Papageorgiou^{9j}, A. Paramonov⁶,
 D. Paredes Hernandez¹⁸¹, A.J. Parker⁸⁷, K.A. Parker⁴⁴, M.A. Parker³¹, F. Parodi^{53b,53a,*}, J.A. Parsons³⁸,
 U. Parzefall⁵⁰, V.R. Pascuzzi¹⁶⁵, J.M.P. Pasner¹⁴³, E. Pasqualucci^{70a}, S. Passaggio^{53b}, F. Pastore⁹¹,
 S. Pataria⁹⁷, J.R. Pater⁹⁸, T. Pauly³⁵, B. Pearson¹¹³, S. Pedraza Lopez¹⁷², R. Pedro^{136a,136b},
 S.V. Peleganchuk^{120b,120a}, O. Penc¹³⁷, C. Peng^{15d}, H. Peng^{58a}, J. Penwell⁶³, B.S. Peralva^{78a},
 M.M. Perego¹⁴², D.V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{66a,66b}, H. Pernegger³⁵, S. Perrella^{67a,67b},
 R. Peschke⁴⁴, V.D. Peshekhonov^{77,*}, K. Peters⁴⁴, R.F.Y. Peters⁹⁸, B.A. Petersen³⁵, T.C. Petersen³⁹,
 E. Petit⁵⁶, A. Petridis¹, C. Petridou¹⁶⁰, P. Petroff¹²⁸, E. Petrolo^{70a}, M. Petrov¹³¹, F. Petrucci^{72a,72b},
 N.E. Pettersson¹⁰⁰, A. Peyaud¹⁴², R. Pezoa^{144b}, F.H. Phillips¹⁰⁴, P.W. Phillips¹⁴¹, G. Piacquadio¹⁵²,
 E. Pianori¹⁷⁶, A. Picazio¹⁰⁰, M.A. Pickering¹³¹, R. Piegai³⁰, J.E. Pilcher³⁶, A.D. Pilkington⁹⁸,
 M. Pinamonti^{71a,71b}, J.L. Pinfold³, H. Pirumov⁴⁴, M. Pitt¹⁷⁸, L. Plazak^{28a}, M-A. Pleier²⁹, V. Pleskot⁹⁷,
 E. Plotnikova⁷⁷, D. Pluth⁷⁶, P. Podberezko^{120b,120a}, R. Poettgen⁹⁴, R. Poggi^{68a,68b}, L. Poggioli¹²⁸,
 I. Pogrebnyak¹⁰⁴, D. Pohl²⁴, I. Pokharel⁵¹, G. Polesello^{68a}, A. Poley⁴⁴, A. Policicchio^{40b,40a}, R. Polifka³⁵,
 A. Polini^{23b}, C.S. Pollard⁵⁵, V. Polychronakos²⁹, K. Pommès³⁵, D. Ponomarenko¹¹⁰, L. Pontecorvo^{70a},

G.A. Popeneciu^{27d}, D.M. Portillo Quintero¹³², S. Pospisil¹³⁸, K. Potamianos⁴⁴, I.N. Potrap⁷⁷, C.J. Potter³¹, H. Potti¹¹, T. Poulsen⁹⁴, J. Poveda³⁵, M.E. Pozo Astigarraga³⁵, P. Pralavorio⁹⁹, A. Pranko¹⁸, S. Prell⁷⁶, D. Price⁹⁸, M. Primavera^{65a}, S. Prince¹⁰¹, N. Proklova¹¹⁰, K. Prokofiev^{61c}, F. Prokoshin^{144b}, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{81a}, A. Puri¹⁷¹, P. Puzo¹²⁸, J. Qian¹⁰³, G. Qin⁵⁵, Y. Qin⁹⁸, A. Quadt⁵¹, M. Queitsch-Maitland⁴⁴, D. Quilty⁵⁵, S. Raddum¹³⁰, V. Radeka²⁹, V. Radescu¹³¹, S.K. Radhakrishnan¹⁵², P. Radloff¹²⁷, P. Rados¹⁰², F. Ragusa^{66a,66b}, G. Rahal⁹⁵, J.A. Raine⁹⁸, S. Rajagopalan²⁹, C. Rangel-Smith¹⁷⁰, T. Rashid¹²⁸, S. Raspopov⁵, M.G. Ratti^{66a,66b}, D.M. Rauch⁴⁴, F. Rauscher¹¹², S. Rave⁹⁷, I. Ravinovich¹⁷⁸, J.H. Rawling⁹⁸, M. Raymond³⁵, A.L. Read¹³⁰, N.P. Readioff⁵⁶, M. Reale^{65a,65b}, D.M. Rebuzzi^{68a,68b}, A. Redelbach¹⁷⁵, G. Redlinger²⁹, R. Reece¹⁴³, R.G. Reed^{32c}, K. Reeves⁴², L. Rehnisch¹⁹, J. Reichert¹³³, A. Reiss⁹⁷, C. Rembser³⁵, H. Ren^{15d}, M. Rescigno^{70a}, S. Resconi^{66a}, E.D. Resseguie¹³³, S. Rettie¹⁷³, E. Reynolds²¹, O.L. Rezanova^{120b,120a}, P. Reznicek¹³⁹, R. Rezvani¹⁰⁷, R. Richter¹¹³, S. Richter⁹², E. Richter-Was^{81b}, O. Ricken²⁴, M. Ridel¹³², P. Rieck¹¹³, C.J. Riegel¹⁸⁰, J. Rieger⁵¹, O. Rifki¹²⁴, M. Rijssenbeek¹⁵², A. Rimoldi^{68a,68b}, M. Rimoldi²⁰, L. Rinaldi^{23b}, G. Ripellino¹⁵¹, B. Ristić³⁵, E. Ritsch³⁵, I. Riu¹⁴, F. Rizatdinova¹²⁵, E. Rizvi⁹⁰, C. Rizzi¹⁴, R.T. Roberts⁹⁸, S.H. Robertson^{101,ae}, A. Robichaud-Veronneau¹⁰¹, D. Robinson³¹, J.E.M. Robinson⁴⁴, A. Robson⁵⁵, E. Rocco⁹⁷, C. Roda^{69a,69b}, Y. Rodina^{99,aa}, S. Rodriguez Bosca¹⁷², A. Rodriguez Perez¹⁴, D. Rodriguez Rodriguez¹⁷², S. Roe³⁵, C.S. Rogan⁵⁷, O. Røhne¹³⁰, J. Roloff⁵⁷, A. Romaniouk¹¹⁰, M. Romano^{23b,23a}, S.M. Romano Saez³⁷, E. Romero Adam¹⁷², N. Rompotis⁸⁸, M. Ronzani⁵⁰, L. Roos¹³², S. Rosati^{70a}, K. Rosbach⁵⁰, P. Rose¹⁴³, N.-A. Rosien⁵¹, E. Rossi^{67a,67b}, L.P. Rossi^{53b}, J.H.N. Rosten³¹, R. Rosten¹⁴⁵, M. Rotaru^{27b}, J. Rothberg¹⁴⁵, D. Rousseau¹²⁸, D. Roy^{32c}, A. Rozanov⁹⁹, Y. Rozen¹⁵⁸, X. Ruan^{32c}, F. Rubbo¹⁵⁰, F. Rühr⁵⁰, A. Ruiz-Martinez³³, Z. Rurikova⁵⁰, N.A. Rusakovich⁷⁷, H.L. Russell¹⁰¹, J.P. Rutherford⁷, N. Ruthmann³⁵, E.M. Rüttinger^{44,1}, Y.F. Ryabov¹³⁴, M. Rybar¹⁷¹, G. Rybkin¹²⁸, S. Ryu⁶, A. Ryzhov¹⁴⁰, G.F. Rzehorz⁵¹, A.F. Saavedra¹⁵⁴, G. Sabato¹¹⁸, S. Sacerdoti³⁰, H.F.-W. Sadrozinski¹⁴³, R. Sadykov⁷⁷, F. Safai Tehrani^{70a}, P. Saha¹¹⁹, M. Sahinsoy^{59a}, M. Saimpert⁴⁴, M. Saito¹⁶¹, T. Saito¹⁶¹, H. Sakamoto¹⁶¹, Y. Sakurai¹⁷⁷, G. Salamanna^{72a,72b}, J.E. Salazar Loyola^{144b}, D. Salek¹¹⁸, P.H. Sales De Bruin¹⁷⁰, D. Salihagic¹¹³, A. Salnikov¹⁵⁰, J. Salt¹⁷², D. Salvatore^{40b,40a}, F. Salvatore¹⁵³, A. Salvucci^{61a,61b,61c}, A. Salzburger³⁵, D. Sammel⁵⁰, D. Sampsonidis¹⁶⁰, D. Sampsonidou¹⁶⁰, J. Sánchez¹⁷², V. Sanchez Martinez¹⁷², A. Sanchez Pineda^{64a,64c}, H. Sandaker¹³⁰, R.L. Sandbach⁹⁰, C.O. Sander⁴⁴, M. Sandhoff¹⁸⁰, C. Sandoval²², D.P.C. Sankey¹⁴¹, M. Sannino^{53b,53a}, Y. Sano¹¹⁵, A. Sansoni⁴⁹, C. Santoni³⁷, H. Santos^{136a}, I. Santoyo Castillo¹⁵³, A. Saponov⁷⁷, J.G. Saraiva^{136a,136d}, B. Sarrazin²⁴, O. Sasaki⁷⁹, K. Sato¹⁶⁷, E. Sauvan⁵, G. Savage⁹¹, P. Savard^{165,av}, N. Savic¹¹³, C. Sawyer¹⁴¹, L. Sawyer^{93,aj}, J. Saxon³⁶, C. Sbarra^{23b}, A. Sbrizzi^{23b,23a}, T. Scanlon⁹², D.A. Scannicchio¹⁶⁹, J. Schaarschmidt¹⁴⁵, P. Schacht¹¹³, B.M. Schachtner¹¹², D. Schaefer³⁶, L. Schaefer¹³³, R. Schaefer⁴⁴, J. Schaeffer⁹⁷, S. Schaepe³⁵, S. Schaezel^{59b}, U. Schäfer⁹⁷, A.C. Schaffer¹²⁸, D. Schaile¹¹², R.D. Schamberger¹⁵², V.A. Schegelsky¹³⁴, D. Scheirich¹³⁹, F. Schenck¹⁹, M. Schernau¹⁶⁹, C. Schiavi^{53b,53a}, S. Schier¹⁴³, L.K. Schildgen²⁴, C. Schillo⁵⁰, M. Schioppa^{40b,40a}, S. Schlenker³⁵, K.R. Schmidt-Sommerfeld¹¹³, K. Schmieden³⁵, C. Schmitt⁹⁷, S. Schmitt⁴⁴, S. Schmitz⁹⁷, U. Schnoor⁵⁰, L. Schoeffel¹⁴², A. Schoening^{59b}, B.D. Schoenrock¹⁰⁴, E. Schopf²⁴, M. Schott⁹⁷, J.F.P. Schouwenberg¹¹⁷, J. Schovancova³⁵, S. Schramm⁵², N. Schuh⁹⁷, A. Schulte⁹⁷, M.J. Schultens²⁴, H.-C. Schultz-Coulon^{59a}, H. Schulz¹⁹, M. Schumacher⁵⁰, B.A. Schumm¹⁴³, Ph. Schune¹⁴², A. Schwartzman¹⁵⁰, T.A. Schwarz¹⁰³, H. Schweiger⁹⁸, Ph. Schwemling¹⁴², R. Schwienhorst¹⁰⁴, A. Sciandra²⁴, G. Sciolla²⁶, M. Scornajenghi^{40b,40a}, F. Scuri^{69a}, F. Scutti¹⁰², J. Searcy¹⁰³, P. Seema²⁴, S.C. Seidel¹¹⁶, A. Seiden¹⁴³, J.M. Seixas^{78b}, G. Sekhniaidze^{67a}, K. Sekhon¹⁰³, S.J. Sekula⁴¹, N. Semprini-Cesari^{23b,23a}, S. Senkin³⁷, C. Serfon¹³⁰, L. Serin¹²⁸, L. Serkin^{64a,64b}, M. Sessa^{72a,72b}, R. Seuster¹⁷⁴, H. Severini¹²⁴, F. Sforza¹⁶⁸, A. Sfyrla⁵², E. Shabalina⁵¹, N.W. Shaikh^{43a,43b}, L.Y. Shan^{15a}, R. Shang¹⁷¹, J.T. Shank²⁵, M. Shapiro¹⁸, P.B. Shatalov¹⁰⁹, K. Shaw^{64a,64b}, S.M. Shaw⁹⁸, A. Shcherbakova^{43a,43b}, C.Y. Shehu¹⁵³, Y. Shen¹²⁴, N. Sherafati³³,

A.D. Sherman²⁵, P. Sherwood⁹², L. Shi^{155,ar}, S. Shimizu⁸⁰, C.O. Shimmin¹⁸¹, M. Shimojima¹¹⁴,
 I.P.J. Shipsey¹³¹, S. Shirabe⁸⁵, M. Shiyakova⁷⁷, J. Shlomi¹⁷⁸, A. Shmeleva¹⁰⁸, D. Shoaleh Saadi¹⁰⁷,
 M.J. Shochet³⁶, S. Shojaii¹⁰², D.R. Shope¹²⁴, S. Shrestha¹²², E. Shulga¹¹⁰, M.A. Shupe⁷, P. Sicho¹³⁷,
 A.M. Sickles¹⁷¹, P.E. Sidebo¹⁵¹, E. Sideras Haddad^{32c}, O. Sidiropoulou¹⁷⁵, A. Sidoti^{23b,23a}, F. Siegert⁴⁶,
 Dj. Sijacki¹⁶, J. Silva^{136a,136d}, S.B. Silverstein^{43a}, V. Simak¹³⁸, L. Simic⁷⁷, S. Simion¹²⁸, E. Simioni⁹⁷,
 B. Simmons⁹², M. Simon⁹⁷, P. Sinervo¹⁶⁵, N.B. Sinev¹²⁷, M. Sioli^{23b,23a}, G. Siragusa¹⁷⁵, I. Siral¹⁰³,
 S.Yu. Sivoklov¹¹¹, J. Sjölin^{43a,43b}, M.B. Skinner⁸⁷, P. Skubic¹²⁴, M. Slater²¹, T. Slavicek¹³⁸,
 M. Slawinska⁸², K. Sliwa¹⁶⁸, R. Slovak¹³⁹, V. Smakhtin¹⁷⁸, B.H. Smart⁵, J. Smiesko^{28a}, N. Smirnov¹¹⁰,
 S.Yu. Smirnov¹¹⁰, Y. Smirnov¹¹⁰, L.N. Smirnova¹¹¹, O. Smirnova⁹⁴, J.W. Smith⁵¹, M.N.K. Smith³⁸,
 R.W. Smith³⁸, M. Smizanska⁸⁷, K. Smolek¹³⁸, A.A. Snesarev¹⁰⁸, I.M. Snyder¹²⁷, S. Snyder²⁹,
 R. Sobie^{174,ae}, F. Socher⁴⁶, A. Soffer¹⁵⁹, A. Søggaard⁴⁸, D.A. Soh¹⁵⁵, G. Sokhrannyi⁸⁹,
 C.A. Solans Sanchez³⁵, M. Solar¹³⁸, E.Yu. Soldatov¹¹⁰, U. Soldevila¹⁷², A.A. Solodkov¹⁴⁰,
 A. Soloshenko⁷⁷, O.V. Solovyanov¹⁴⁰, V. Solovyev¹³⁴, P. Sommer¹⁴⁶, H. Son¹⁶⁸, A. Sopczak¹³⁸,
 D. Sosa^{59b}, C.L. Sotiropoulou^{69a,69b}, S. Sottocornola^{68a,68b}, R. Soualah^{64a,64c,i}, A.M. Soukharev^{120b,120a},
 D. South⁴⁴, B.C. Sowden⁹¹, S. Spagnolo^{65a,65b}, M. Spalla^{69a,69b}, M. Spangenberg¹⁷⁶, F. Spanò⁹¹,
 D. Sperlich¹⁹, F. Spettel¹¹³, T.M. Spieker^{59a}, R. Spighi^{23b}, G. Spigo³⁵, L.A. Spiller¹⁰², M. Spousta¹³⁹,
 R.D. St. Denis^{55,*}, A. Stabile^{66a,66b}, R. Stamen^{59a}, S. Stamm¹⁹, E. Stanecka⁸², R.W. Stanek⁶,
 C. Stanescu^{72a}, M.M. Stanitzki⁴⁴, B. Stapf¹¹⁸, S. Stapnes¹³⁰, E.A. Starchenko¹⁴⁰, G.H. Stark³⁶,
 J. Stark⁵⁶, S.H. Stark³⁹, P. Staroba¹³⁷, P. Starovoitov^{59a}, S. Stärz³⁵, R. Staszewski⁸², M. Stegler⁴⁴,
 P. Steinberg²⁹, B. Stelzer¹⁴⁹, H.J. Stelzer³⁵, O. Stelzer-Chilton^{166a}, H. Stenzel⁵⁴, T.J. Stevenson⁹⁰,
 G.A. Stewart⁵⁵, M.C. Stockton¹²⁷, M. Stoebe¹⁰¹, G. Stoicea^{27b}, P. Stolte⁵¹, S. Stonjek¹¹³,
 A.R. Stradling⁸, A. Straessner⁴⁶, M.E. Stramaglia²⁰, J. Strandberg¹⁵¹, S. Strandberg^{43a,43b}, M. Strauss¹²⁴,
 P. Strizenc^{28b}, R. Ströhmer¹⁷⁵, D.M. Strom¹²⁷, R. Stroynowski⁴¹, A. Strubig⁴⁸, S.A. Stucci²⁹,
 B. Stugu¹⁷, N.A. Styles⁴⁴, D. Su¹⁵⁰, J. Su¹³⁵, S. Sucheck^{59a}, Y. Sugaya¹²⁹, M. Suk¹³⁸, V.V. Sulin¹⁰⁸,
 D.M.S. Sultan^{73a,73b}, S. Sultansoy^{4c}, T. Sumida⁸³, S. Sun⁵⁷, X. Sun³, K. Suruliz¹⁵³, C.J.E. Suster¹⁵⁴,
 M.R. Sutton¹⁵³, S. Suzuki⁷⁹, M. Svatos¹³⁷, M. Swiatlowski³⁶, S.P. Swift², I. Sykora^{28a}, T. Sykora¹³⁹,
 D. Ta⁵⁰, K. Tackmann^{44,ab}, J. Taenzer¹⁵⁹, A. Taffard¹⁶⁹, R. Tafirout^{166a}, E. Tahirovic⁹⁰, N. Taiblum¹⁵⁹,
 H. Takai²⁹, R. Takashima⁸⁴, E.H. Takasugi¹¹³, K. Takeda⁸⁰, T. Takeshita¹⁴⁷, Y. Takubo⁷⁹, M. Talby⁹⁹,
 A.A. Talyshev^{120b,120a}, J. Tanaka¹⁶¹, M. Tanaka¹⁶³, R. Tanaka¹²⁸, S. Tanaka⁷⁹, R. Tanioka⁸⁰,
 B.B. Tannenwald¹²², S. Tapia Araya^{144b}, S. Tapprogge⁹⁷, S. Tarem¹⁵⁸, G.F. Tartarelli^{66a}, P. Tas¹³⁹,
 M. Tasevsky¹³⁷, T. Tashiro⁸³, E. Tassi^{40b,40a}, A. Tavares Delgado^{136a,136b}, Y. Tayalati^{34e}, A.C. Taylor¹¹⁶,
 A.J. Taylor⁴⁸, G.N. Taylor¹⁰², P.T.E. Taylor¹⁰², W. Taylor^{166b}, P. Teixeira-Dias⁹¹, D. Temple¹⁴⁹,
 H. Ten Kate³⁵, P.K. Teng¹⁵⁵, J.J. Teoh¹²⁹, F. Tepel¹⁸⁰, S. Terada⁷⁹, K. Terashi¹⁶¹, J. Terron⁹⁶, S. Terzo¹⁴,
 M. Testa⁴⁹, R.J. Teuscher^{165,ae}, S.J. Thais¹⁸¹, T. Thevenaux-Pelzer⁹⁹, F. Thiele³⁹, J.P. Thomas²¹,
 J. Thomas-Wilsker⁹¹, A.S. Thompson⁵⁵, P.D. Thompson²¹, L.A. Thomsen¹⁸¹, E. Thomson¹³³, Y. Tian³⁸,
 M.J. Tibbetts¹⁸, R.E. Ticse Torres⁵¹, V.O. Tikhomirov^{108,am}, Yu.A. Tikhonov^{120b,120a}, S. Timoshenko¹¹⁰,
 P. Tipton¹⁸¹, S. Tisserant⁹⁹, K. Todome¹⁶³, S. Todorova-Nova⁵, S. Todt⁴⁶, J. Tojo⁸⁵, S. Tokár^{28a},
 K. Tokushuku⁷⁹, E. Tolley¹²², L. Tomlinson⁹⁸, M. Tomoto¹¹⁵, L. Tompkins^{150,r}, K. Toms¹¹⁶, B. Tong⁵⁷,
 P. Tornambe⁵⁰, E. Torrence¹²⁷, H. Torres⁴⁶, E. Torró Pastor¹⁴⁵, J. Toth^{99,ad}, F. Touchard⁹⁹, D.R. Tovey¹⁴⁶,
 C.J. Treado¹²¹, T. Trefzger¹⁷⁵, F. Tresoldi¹⁵³, A. Tricoli²⁹, I.M. Trigger^{166a}, S. Trincaz-Duvoid¹³²,
 M.F. Tripiana¹⁴, W. Trischuk¹⁶⁵, B. Trocme⁵⁶, A. Trofymov⁴⁴, C. Troncon^{66a}, M. Trotter-McDonald¹⁸,
 M. Trovatelli¹⁷⁴, L. Truong^{32b}, M. Trzebinski⁸², A. Trzupek⁸², K.W. Tsang^{61a}, J.C.-L. Tseng¹³¹,
 P.V. Tsiarshka¹⁰⁵, G. Tsipolitis¹⁰, N. Tsirintanis⁹, S. Tsiskaridze¹⁴, V. Tsiskaridze⁵⁰,
 E.G. Tskhadadze^{157a}, I.I. Tsukerman¹⁰⁹, V. Tsulaia¹⁸, S. Tsuno⁷⁹, D. Tsybychev¹⁵², Y. Tu^{61b},
 A. Tudorache^{27b}, V. Tudorache^{27b}, T.T. Tulbure^{27a}, A.N. Tuna⁵⁷, S. Turchikhin⁷⁷, D. Turgeman¹⁷⁸,
 I. Turk Cakir^{4b,v}, R. Turra^{66a}, P.M. Tuts³⁸, G. Uccielli^{23b,23a}, I. Ueda⁷⁹, M. Ughetto^{43a,43b},
 F. Ukegawa¹⁶⁷, G. Unal³⁵, A. Undrus²⁹, G. Unel¹⁶⁹, F.C. Ungaro¹⁰², Y. Unno⁷⁹, K. Uno¹⁶¹,

C. Unverdorben¹¹², J. Urban^{28b}, P. Urquijo¹⁰², P. Urrejola⁹⁷, G. Usai⁸, J. Usui⁷⁹, L. Vacavant⁹⁹, V. Vacek¹³⁸, B. Vachon¹⁰¹, K.O.H. Vadla¹³⁰, A. Vaidya⁹², C. Valderanis¹¹², E. Valdes Santurio^{43a,43b}, M. Valente⁵², S. Valentineti^{23b,23a}, A. Valero¹⁷², L. Valéry¹⁴, S. Valkar¹³⁹, A. Vallier⁵, J.A. Valls Ferrer¹⁷², W. Van Den Wollenberg¹¹⁸, H. Van der Graaf¹¹⁸, P. Van Gemmeren⁶, J. Van Nieuwkoop¹⁴⁹, I. Van Vulpen¹¹⁸, M.C. van Woerden¹¹⁸, M. Vanadia^{71a,71b}, W. Vandelli³⁵, A. Vaniachine¹⁶⁴, P. Vankov¹¹⁸, G. Vardanyan¹⁸², R. Vari^{70a}, E.W. Varnes⁷, C. Varni^{53b,53a}, T. Varol⁴¹, D. Varouchas¹²⁸, A. Vartapetian⁸, K.E. Varvell¹⁵⁴, G.A. Vasquez^{144b}, J.G. Vasquez¹⁸¹, F. Vazeille³⁷, D. Vazquez Furelos¹⁴, T. Vazquez Schroeder¹⁰¹, J. Veatch⁵¹, V. Veeraraghavan⁷, L.M. Veloce¹⁶⁵, F. Veloso^{136a,136c}, S. Veneziano^{70a}, A. Ventura^{65a,65b}, M. Venturi¹⁷⁴, N. Venturi³⁵, A. Venturini²⁶, V. Vercesi^{68a}, M. Verducci^{72a,72b}, W. Verkerke¹¹⁸, A.T. Vermeulen¹¹⁸, J.C. Vermeulen¹¹⁸, M.C. Vetterli^{149,av}, N. Viaux Maira^{144b}, O. Viazlo⁹⁴, I. Vichou^{171,*}, T. Vickey¹⁴⁶, O.E. Vickey Boeriu¹⁴⁶, G.H.A. Viehhauser¹³¹, S. Viel¹⁸, L. Vigani¹³¹, M. Villa^{23b,23a}, M. Villaplana Perez^{66a,66b}, E. Vilucchi⁴⁹, M.G. Vinciter³³, V.B. Vinogradov⁷⁷, A. Vishwakarma⁴⁴, C. Vittori^{23b,23a}, I. Vivarelli¹⁵³, S. Vlachos¹⁰, M. Vogel¹⁸⁰, P. Vokac¹³⁸, G. Volpi¹⁴, H. von der Schmitt¹¹³, E. Von Toerne²⁴, V. Vorobel¹³⁹, K. Vorobev¹¹⁰, M. Vos¹⁷², R. Voss³⁵, J.H. Vosseveld⁸⁸, N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹³⁸, M. Vreeswijk¹¹⁸, T. Šfiligoj⁸⁹, R. Vuillermet³⁵, I. Vukotic³⁶, T. Ženiš^{28a}, L. Živkovic¹⁶, P. Wagner²⁴, W. Wagner¹⁸⁰, J. Wagner-Kuhr¹¹², H. Wahlberg⁸⁶, S. Wahrmund⁴⁶, K. Wakamiya⁸⁰, J. Walder⁸⁷, R. Walker¹¹², W. Walkowiak¹⁴⁸, V. Wallangen^{43a,43b}, C. Wang^{15c}, C. Wang^{58b,e}, F. Wang¹⁷⁹, H. Wang¹⁸, H. Wang³, J. Wang¹⁵⁴, J. Wang⁴⁴, Q. Wang¹²⁴, R.-J. Wang¹³², R. Wang⁶, S.M. Wang¹⁵⁵, T. Wang³⁸, W. Wang^{155,p}, W.X. Wang^{58a,af}, Z. Wang^{58c}, C. Wanotayaroj⁴⁴, A. Warburton¹⁰¹, C.P. Ward³¹, D.R. Wardrope⁹², A. Washbrook⁴⁸, P.M. Watkins²¹, A.T. Watson²¹, M.F. Watson²¹, G. Watts¹⁴⁵, S. Watts⁹⁸, B.M. Waugh⁹², A.F. Webb¹¹, S. Webb⁹⁷, M.S. Weber²⁰, S.A. Weber³³, S.M. Weber^{59a}, S.W. Weber¹⁷⁵, J.S. Webster⁶, A.R. Weidberg¹³¹, B. Weinert⁶³, J. Weingarten⁵¹, M. Weirich⁹⁷, C. Weiser⁵⁰, H. Weits¹¹⁸, P.S. Wells³⁵, T. Wenaus²⁹, T. Wengler³⁵, S. Wenig³⁵, N. Wermes²⁴, M.D. Werner⁷⁶, P. Werner³⁵, M. Wessels^{59a}, T.D. Weston²⁰, K. Whalen¹²⁷, N.L. Whallon¹⁴⁵, A.M. Wharton⁸⁷, A.S. White¹⁰³, A. White⁸, M.J. White¹, R. White^{144b}, D. Whiteson¹⁶⁹, B.W. Whitmore⁸⁷, F.J. Wickens¹⁴¹, W. Wiedenmann¹⁷⁹, M. Wielers¹⁴¹, C. Wiglesworth³⁹, L.A.M. Wiik-Fuchs⁵⁰, A. Wildauer¹¹³, F. Wilk⁹⁸, H.G. Wilkens³⁵, H.H. Williams¹³³, S. Williams³¹, C. Willis¹⁰⁴, S. Willocq¹⁰⁰, J.A. Wilson²¹, I. Wingerter-Seez⁵, E. Winkels¹⁵³, F. Winklmeier¹²⁷, O.J. Winston¹⁵³, B.T. Winter²⁴, M. Wittgen¹⁵⁰, M. Wobisch⁹³, A. Wolf⁹⁷, T.M.H. Wolf¹¹⁸, R. Wolff⁹⁹, M.W. Wolter⁸², H. Wolters^{136a,136c}, V.W.S. Wong¹⁷³, N.L. Woods¹⁴³, S.D. Worm²¹, B.K. Wosiek⁸², J. Wotschack³⁵, K.W. Woźniak⁸², M. Wu³⁶, S.L. Wu¹⁷⁹, X. Wu⁵², Y. Wu¹⁰³, T.R. Wyatt⁹⁸, B.M. Wynne⁴⁸, S. Xella³⁹, Z. Xi¹⁰³, L. Xia^{15b}, D. Xu^{15a}, L. Xu²⁹, T. Xu¹⁴², W. Xu¹⁰³, B. Yabsley¹⁵⁴, S. Yacoob^{32a}, D. Yamaguchi¹⁶³, Y. Yamaguchi¹⁶³, A. Yamamoto⁷⁹, S. Yamamoto¹⁶¹, T. Yamanaka¹⁶¹, F. Yamane⁸⁰, M. Yamatani¹⁶¹, T. Yamazaki¹⁶¹, Y. Yamazaki⁸⁰, Z. Yan²⁵, H.J. Yang^{58c,58d}, H.T. Yang¹⁸, Y. Yang¹⁵⁵, Z. Yang¹⁷, W.-M. Yao¹⁸, Y.C. Yap⁴⁴, Y. Yasu⁷⁹, E. Yatsenko⁵, K.H. Yau Wong²⁴, J. Ye⁴¹, S. Ye²⁹, I. Yeletskikh⁷⁷, E. Yigitbasi²⁵, E. Yildirim⁹⁷, K. Yorita¹⁷⁷, K. Yoshihara¹³³, C.J.S. Young³⁵, C. Young¹⁵⁰, J. Yu⁸, J. Yu⁷⁶, S.P.Y. Yuen²⁴, I. Yusuff^{31,a}, B. Zabinski⁸², G. Zacharis¹⁰, R. Zaidan¹⁴, A.M. Zaitsev^{140,al}, N. Zakharchuk⁴⁴, J. Zalieckas¹⁷, A. Zaman¹⁵², S. Zambito⁵⁷, D. Zanzi¹⁰², C. Zeitnitz¹⁸⁰, G. Zemaityte¹³¹, A. Zemla^{81a}, J.C. Zeng¹⁷¹, Q. Zeng¹⁵⁰, O. Zenin¹⁴⁰, D. Zerwas¹²⁸, D.F. Zhang^{58b}, D. Zhang¹⁰³, F. Zhang¹⁷⁹, G. Zhang^{58a,af}, H. Zhang¹²⁸, J. Zhang⁶, L. Zhang⁵⁰, L. Zhang^{58a}, M. Zhang¹⁷¹, P. Zhang^{15c}, R. Zhang^{58a,e}, R. Zhang²⁴, X. Zhang^{58b}, Y. Zhang^{15d}, Z. Zhang¹²⁸, X. Zhao⁴¹, Y. Zhao^{58b,128,ai}, Z. Zhao^{58a}, A. Zhemchugov⁷⁷, B. Zhou¹⁰³, C. Zhou¹⁷⁹, L. Zhou⁴¹, M.S. Zhou^{15d}, M. Zhou¹⁵², N. Zhou^{58c}, Y. Zhou⁷, C.G. Zhu^{58b}, H. Zhu^{15a}, J. Zhu¹⁰³, Y. Zhu^{58a}, X. Zhuang^{15a}, K. Zhukov¹⁰⁸, A. Zibell¹⁷⁵, D. Zieminska⁶³, N.I. Zimine⁷⁷, C. Zimmermann⁹⁷, S. Zimmermann⁵⁰, Z. Zinonos¹¹³, M. Zinser⁹⁷, M. Ziolkowski¹⁴⁸, G. Zobernig¹⁷⁹, A. Zoccoli^{23b,23a}, R. Zou³⁶, M. Zur Nedden¹⁹, L. Zwalinski³⁵.

- ¹Department of Physics, University of Adelaide, Adelaide; Australia.
- ²Physics Department, SUNY Albany, Albany NY; United States of America.
- ³Department of Physics, University of Alberta, Edmonton AB; Canada.
- ⁴(^a)Department of Physics, Ankara University, Ankara; (^b)Istanbul Aydin University, Istanbul; (^c)Division of Physics, TOBB University of Economics and Technology, Ankara; Turkey.
- ⁵LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy; France.
- ⁶High Energy Physics Division, Argonne National Laboratory, Argonne IL; United States of America.
- ⁷Department of Physics, University of Arizona, Tucson AZ; United States of America.
- ⁸Department of Physics, University of Texas at Arlington, Arlington TX; United States of America.
- ⁹Physics Department, National and Kapodistrian University of Athens, Athens; Greece.
- ¹⁰Physics Department, National Technical University of Athens, Zografou; Greece.
- ¹¹Department of Physics, University of Texas at Austin, Austin TX; United States of America.
- ¹²(^a)Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; (^b)Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; (^c)Department of Physics, Bogazici University, Istanbul; (^d)Department of Physics Engineering, Gaziantep University, Gaziantep; Turkey.
- ¹³Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
- ¹⁴Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain.
- ¹⁵(^a)Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (^b)Physics Department, Tsinghua University, Beijing; (^c)Department of Physics, Nanjing University, Nanjing; (^d)University of Chinese Academy of Science (UCAS), Beijing; China.
- ¹⁶Institute of Physics, University of Belgrade, Belgrade; Serbia.
- ¹⁷Department for Physics and Technology, University of Bergen, Bergen; Norway.
- ¹⁸Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA; United States of America.
- ¹⁹Institut für Physik, Humboldt Universität zu Berlin, Berlin; Germany.
- ²⁰Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern; Switzerland.
- ²¹School of Physics and Astronomy, University of Birmingham, Birmingham; United Kingdom.
- ²²Centro de Investigaciones, Universidad Antonio Nariño, Bogota; Colombia.
- ²³(^a)Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna; (^b)INFN Sezione di Bologna; Italy.
- ²⁴Physikalisches Institut, Universität Bonn, Bonn; Germany.
- ²⁵Department of Physics, Boston University, Boston MA; United States of America.
- ²⁶Department of Physics, Brandeis University, Waltham MA; United States of America.
- ²⁷(^a)Transilvania University of Brasov, Brasov; (^b)Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; (^c)Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; (^d)National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; (^e)University Politehnica Bucharest, Bucharest; (^f)West University in Timisoara, Timisoara; Romania.
- ²⁸(^a)Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; (^b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice; Slovak Republic.
- ²⁹Physics Department, Brookhaven National Laboratory, Upton NY; United States of America.
- ³⁰Departamento de Física, Universidad de Buenos Aires, Buenos Aires; Argentina.
- ³¹Cavendish Laboratory, University of Cambridge, Cambridge; United Kingdom.
- ³²(^a)Department of Physics, University of Cape Town, Cape Town; (^b)Department of Mechanical

Engineering Science, University of Johannesburg, Johannesburg;^(c) School of Physics, University of the Witwatersrand, Johannesburg; South Africa.

³³Department of Physics, Carleton University, Ottawa ON; Canada.

³⁴(^a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (^b) Centre National de l'Énergie des Sciences Techniques Nucleaires (CNESTEN), Rabat; (^c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (^d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (^e) Faculté des sciences, Université Mohammed V, Rabat; Morocco.

³⁵CERN, Geneva; Switzerland.

³⁶Enrico Fermi Institute, University of Chicago, Chicago IL; United States of America.

³⁷LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand; France.

³⁸Nevis Laboratory, Columbia University, Irvington NY; United States of America.

³⁹Niels Bohr Institute, University of Copenhagen, Copenhagen; Denmark.

⁴⁰(^a) Dipartimento di Fisica, Università della Calabria, Rende; (^b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Italy.

⁴¹Physics Department, Southern Methodist University, Dallas TX; United States of America.

⁴²Physics Department, University of Texas at Dallas, Richardson TX; United States of America.

⁴³(^a) Department of Physics, Stockholm University; (^b) Oskar Klein Centre, Stockholm; Sweden.

⁴⁴Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen; Germany.

⁴⁵Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund; Germany.

⁴⁶Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden; Germany.

⁴⁷Department of Physics, Duke University, Durham NC; United States of America.

⁴⁸SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh; United Kingdom.

⁴⁹INFN e Laboratori Nazionali di Frascati, Frascati; Italy.

⁵⁰Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.

⁵¹II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.

⁵²Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.

⁵³(^a) Dipartimento di Fisica, Università di Genova, Genova; (^b) INFN Sezione di Genova; Italy.

⁵⁴II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen; Germany.

⁵⁵SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow; United Kingdom.

⁵⁶LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble; France.

⁵⁷Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA; United States of America.

⁵⁸(^a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; (^b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; (^c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; (^d) Tsung-Dao Lee Institute, Shanghai; China.

⁵⁹(^a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (^b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; Germany.

⁶⁰Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima; Japan.

⁶¹(^a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; (^b) Department of Physics, University of Hong Kong, Hong Kong; (^c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong; China.

⁶²Department of Physics, National Tsing Hua University, Hsinchu; Taiwan.

⁶³Department of Physics, Indiana University, Bloomington IN; United States of America.

- 64^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine;^(b) ICTP, Trieste;^(c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine; Italy.
- 65^(a) INFN Sezione di Lecce;^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce; Italy.
- 66^(a) INFN Sezione di Milano;^(b) Dipartimento di Fisica, Università di Milano, Milano; Italy.
- 67^(a) INFN Sezione di Napoli;^(b) Dipartimento di Fisica, Università di Napoli, Napoli; Italy.
- 68^(a) INFN Sezione di Pavia;^(b) Dipartimento di Fisica, Università di Pavia, Pavia; Italy.
- 69^(a) INFN Sezione di Pisa;^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.
- 70^(a) INFN Sezione di Roma;^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma; Italy.
- 71^(a) INFN Sezione di Roma Tor Vergata;^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma; Italy.
- 72^(a) INFN Sezione di Roma Tre;^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma; Italy.
- 73^(a) INFN-TIFPA;^(b) Università degli Studi di Trento, Trento; Italy.
- 74 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck; Austria.
- 75 University of Iowa, Iowa City IA; United States of America.
- 76 Department of Physics and Astronomy, Iowa State University, Ames IA; United States of America.
- 77 Joint Institute for Nuclear Research, Dubna; Russia.
- 78^(a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora;^(b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro;^(c) Universidade Federal de São João del Rei (UFSJ), São João del Rei;^(d) Instituto de Física, Universidade de São Paulo, São Paulo; Brazil.
- 79 KEK, High Energy Accelerator Research Organization, Tsukuba; Japan.
- 80 Graduate School of Science, Kobe University, Kobe; Japan.
- 81^(a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow;^(b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow; Poland.
- 82 Institute of Nuclear Physics Polish Academy of Sciences, Krakow; Poland.
- 83 Faculty of Science, Kyoto University, Kyoto; Japan.
- 84 Kyoto University of Education, Kyoto; Japan.
- 85 Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka ; Japan.
- 86 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata; Argentina.
- 87 Physics Department, Lancaster University, Lancaster; United Kingdom.
- 88 Oliver Lodge Laboratory, University of Liverpool, Liverpool; United Kingdom.
- 89 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana; Slovenia.
- 90 School of Physics and Astronomy, Queen Mary University of London, London; United Kingdom.
- 91 Department of Physics, Royal Holloway University of London, Egham; United Kingdom.
- 92 Department of Physics and Astronomy, University College London, London; United Kingdom.
- 93 Louisiana Tech University, Ruston LA; United States of America.
- 94 Fysiska institutionen, Lunds universitet, Lund; Sweden.
- 95 Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne; France.
- 96 Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid; Spain.
- 97 Institut für Physik, Universität Mainz, Mainz; Germany.
- 98 School of Physics and Astronomy, University of Manchester, Manchester; United Kingdom.
- 99 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
- 100 Department of Physics, University of Massachusetts, Amherst MA; United States of America.

- ¹⁰¹Department of Physics, McGill University, Montreal QC; Canada.
- ¹⁰²School of Physics, University of Melbourne, Victoria; Australia.
- ¹⁰³Department of Physics, University of Michigan, Ann Arbor MI; United States of America.
- ¹⁰⁴Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
- ¹⁰⁵B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk; Belarus.
- ¹⁰⁶Research Institute for Nuclear Problems of Byelorussian State University, Minsk; Belarus.
- ¹⁰⁷Group of Particle Physics, University of Montreal, Montreal QC; Canada.
- ¹⁰⁸P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow; Russia.
- ¹⁰⁹Institute for Theoretical and Experimental Physics (ITEP), Moscow; Russia.
- ¹¹⁰National Research Nuclear University MEPhI, Moscow; Russia.
- ¹¹¹D.V. Skobel'syn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow; Russia.
- ¹¹²Fakultät für Physik, Ludwig-Maximilians-Universität München, München; Germany.
- ¹¹³Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München; Germany.
- ¹¹⁴Nagasaki Institute of Applied Science, Nagasaki; Japan.
- ¹¹⁵Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya; Japan.
- ¹¹⁶Department of Physics and Astronomy, University of New Mexico, Albuquerque NM; United States of America.
- ¹¹⁷Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.
- ¹¹⁸Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam; Netherlands.
- ¹¹⁹Department of Physics, Northern Illinois University, DeKalb IL; United States of America.
- ¹²⁰(^a) Budker Institute of Nuclear Physics, SB RAS, Novosibirsk; (^b) Novosibirsk State University Novosibirsk; Russia.
- ¹²¹Department of Physics, New York University, New York NY; United States of America.
- ¹²²Ohio State University, Columbus OH; United States of America.
- ¹²³Faculty of Science, Okayama University, Okayama; Japan.
- ¹²⁴Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK; United States of America.
- ¹²⁵Department of Physics, Oklahoma State University, Stillwater OK; United States of America.
- ¹²⁶Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc; Czech Republic.
- ¹²⁷Center for High Energy Physics, University of Oregon, Eugene OR; United States of America.
- ¹²⁸LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.
- ¹²⁹Graduate School of Science, Osaka University, Osaka; Japan.
- ¹³⁰Department of Physics, University of Oslo, Oslo; Norway.
- ¹³¹Department of Physics, Oxford University, Oxford; United Kingdom.
- ¹³²LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris; France.
- ¹³³Department of Physics, University of Pennsylvania, Philadelphia PA; United States of America.
- ¹³⁴Konstantinov Nuclear Physics Institute of National Research Centre "Kurchatov Institute", PNPI, St. Petersburg; Russia.
- ¹³⁵Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA; United States of America.
- ¹³⁶(^a) Laboratório de Instrumentação e Física Experimental de Partículas - LIP; (^b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; (^c) Departamento de Física, Universidade de Coimbra, Coimbra; (^d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; (^e) Departamento

- de Física, Universidade do Minho, Braga;^(f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain);^(g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica; Portugal.
- ¹³⁷Institute of Physics, Academy of Sciences of the Czech Republic, Prague; Czech Republic.
- ¹³⁸Czech Technical University in Prague, Prague; Czech Republic.
- ¹³⁹Charles University, Faculty of Mathematics and Physics, Prague; Czech Republic.
- ¹⁴⁰State Research Center Institute for High Energy Physics, NRC KI, Protvino; Russia.
- ¹⁴¹Particle Physics Department, Rutherford Appleton Laboratory, Didcot; United Kingdom.
- ¹⁴²IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette; France.
- ¹⁴³Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA; United States of America.
- ¹⁴⁴^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; ^(b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso; Chile.
- ¹⁴⁵Department of Physics, University of Washington, Seattle WA; United States of America.
- ¹⁴⁶Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.
- ¹⁴⁷Department of Physics, Shinshu University, Nagano; Japan.
- ¹⁴⁸Department Physik, Universität Siegen, Siegen; Germany.
- ¹⁴⁹Department of Physics, Simon Fraser University, Burnaby BC; Canada.
- ¹⁵⁰SLAC National Accelerator Laboratory, Stanford CA; United States of America.
- ¹⁵¹Physics Department, Royal Institute of Technology, Stockholm; Sweden.
- ¹⁵²Departments of Physics and Astronomy, Stony Brook University, Stony Brook NY; United States of America.
- ¹⁵³Department of Physics and Astronomy, University of Sussex, Brighton; United Kingdom.
- ¹⁵⁴School of Physics, University of Sydney, Sydney; Australia.
- ¹⁵⁵Institute of Physics, Academia Sinica, Taipei; Taiwan.
- ¹⁵⁶Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei; Taiwan.
- ¹⁵⁷^(a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi; Georgia.
- ¹⁵⁸Department of Physics, Technion, Israel Institute of Technology, Haifa; Israel.
- ¹⁵⁹Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv; Israel.
- ¹⁶⁰Department of Physics, Aristotle University of Thessaloniki, Thessaloniki; Greece.
- ¹⁶¹International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo; Japan.
- ¹⁶²Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo; Japan.
- ¹⁶³Department of Physics, Tokyo Institute of Technology, Tokyo; Japan.
- ¹⁶⁴Tomsk State University, Tomsk; Russia.
- ¹⁶⁵Department of Physics, University of Toronto, Toronto ON; Canada.
- ¹⁶⁶^(a) TRIUMF, Vancouver BC; ^(b) Department of Physics and Astronomy, York University, Toronto ON; Canada.
- ¹⁶⁷Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba; Japan.
- ¹⁶⁸Department of Physics and Astronomy, Tufts University, Medford MA; United States of America.
- ¹⁶⁹Department of Physics and Astronomy, University of California Irvine, Irvine CA; United States of America.
- ¹⁷⁰Department of Physics and Astronomy, University of Uppsala, Uppsala; Sweden.
- ¹⁷¹Department of Physics, University of Illinois, Urbana IL; United States of America.
- ¹⁷²Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia; Spain.

- ¹⁷³Department of Physics, University of British Columbia, Vancouver BC; Canada.
- ¹⁷⁴Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada.
- ¹⁷⁵Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg; Germany.
- ¹⁷⁶Department of Physics, University of Warwick, Coventry; United Kingdom.
- ¹⁷⁷Waseda University, Tokyo; Japan.
- ¹⁷⁸Department of Particle Physics, Weizmann Institute of Science, Rehovot; Israel.
- ¹⁷⁹Department of Physics, University of Wisconsin, Madison WI; United States of America.
- ¹⁸⁰Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany.
- ¹⁸¹Department of Physics, Yale University, New Haven CT; United States of America.
- ¹⁸²Yerevan Physics Institute, Yerevan; Armenia.
- ^a Also at Department of Physics, University of Malaya, Kuala Lumpur; Malaysia.
- ^b Also at Borough of Manhattan Community College, City University of New York, NY; United States of America.
- ^c Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town; South Africa.
- ^d Also at CERN, Geneva; Switzerland.
- ^e Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
- ^f Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
- ^g Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
- ^h Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); Spain.
- ⁱ Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah; United Arab Emirates.
- ^j Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
- ^k Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.
- ^l Also at Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.
- ^m Also at Department of Physics, California State University, Fresno CA; United States of America.
- ⁿ Also at Department of Physics, California State University, Sacramento CA; United States of America.
- ^o Also at Department of Physics, King's College London, London; United Kingdom.
- ^p Also at Department of Physics, Nanjing University, Nanjing; China.
- ^q Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.
- ^r Also at Department of Physics, Stanford University; United States of America.
- ^s Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
- ^t Also at Department of Physics, University of Michigan, Ann Arbor MI; United States of America.
- ^u Also at Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.
- ^v Also at Giresun University, Faculty of Engineering, Giresun; Turkey.
- ^w Also at Graduate School of Science, Osaka University, Osaka; Japan.
- ^x Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; Romania.
- ^y Also at II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.
- ^z Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.
- ^{aa} Also at Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain.
- ^{ab} Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
- ^{ac} Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University

Nijmegen/Nikhef, Nijmegen; Netherlands.

ad Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.

ae Also at Institute of Particle Physics (IPP); Canada.

af Also at Institute of Physics, Academia Sinica, Taipei; Taiwan.

ag Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.

ah Also at Institute of Theoretical Physics, Ilia State University, Tbilisi; Georgia.

ai Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.

aj Also at Louisiana Tech University, Ruston LA; United States of America.

ak Also at Manhattan College, New York NY; United States of America.

al Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

am Also at National Research Nuclear University MEPhI, Moscow; Russia.

an Also at Near East University, Nicosia, North Cyprus, Mersin; Turkey.

ao Also at Novosibirsk State University, Novosibirsk; Russia.

ap Also at O Chadai Academic Production, Ochanomizu University, Tokyo; Japan.

aq Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.

ar Also at School of Physics, Sun Yat-sen University, Guangzhou; China.

as Also at The City College of New York, New York NY; United States of America.

at Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.

au Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

av Also at TRIUMF, Vancouver BC; Canada.

aw Also at Università di Napoli Parthenope, Napoli; Italy.

* Deceased