

Observation of charmless baryonic decays $B_{(s)}^0 \rightarrow p\bar{p}h^+h'^-$ R. Aaij *et al.**

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Decays of B^0 and B_s^0 mesons to the charmless baryonic final states $p\bar{p}h^+h'^-$, where h and h' each denote a kaon or a pion, are searched for using the LHCb detector. The analysis is based on a sample of proton-proton collision data collected at center-of-mass energies of 7 and 8 TeV, corresponding to an integrated luminosity of 3 fb^{-1} . Four-body charmless baryonic B_s^0 decays are observed for the first time. The decays $B_s^0 \rightarrow p\bar{p}K^+K^-$, $B_s^0 \rightarrow p\bar{p}K^\pm\pi^\mp$, $B^0 \rightarrow p\bar{p}K^\pm\pi^\mp$ and $B^0 \rightarrow p\bar{p}\pi^+\pi^-$ are observed with a significance greater than 5 standard deviations; evidence at 4.1 standard deviations is found for the $B^0 \rightarrow p\bar{p}K^+K^-$ decay and an upper limit is set on the branching fraction for $B_s^0 \rightarrow p\bar{p}\pi^+\pi^-$. Branching fractions in the kinematic region $m(p\bar{p}) < 2850 \text{ MeV}/c^2$ are measured relative to the $B^0 \rightarrow J/\psi(\rightarrow p\bar{p})K^*(892)^0$ channel.

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In recent years, studies by the LHCb Collaboration have greatly increased the knowledge of the decays of B mesons to final states containing baryons. The first observation of a baryonic B_c^+ decay was reported in 2014 [1], and LHCb recently reported the first observation of a baryonic B_s^0 decay [2], the last of the four B meson species for which a baryonic decay mode had yet to be observed.

Primary areas of interest in baryonic B decays include the hierarchy of branching fractions to the various decay modes, the presence of resonances and the existence of a threshold enhancement in the baryon-antibaryon mass spectrum [3,4]. The first evidence of CP violation in baryonic B decays has been reported from an analysis of $B^+ \rightarrow p\bar{p}K^+$ decays [5]. It is of great interest to search for further manifestations of CP violation in baryonic B decays, e.g. with so-called triple-product correlations (TPCs); see Ref. [6] and references therein. For certain decays, asymmetries of up to 20% are predicted [7]. Four-body decays are particularly suited for this approach since the definitions of the TPCs do not involve the spins of the final-state particles, unlike the TPCs in three-body decays [6,8].

This paper presents a search for the decays of B^0 and B_s^0 mesons to the four-body charmless baryonic final states $p\bar{p}h^+h'^-$, where h and h' each denote a kaon or a pion. The inclusion of charge-conjugate processes is implied, unless otherwise indicated. For simplicity, the charges of the $h^+h'^-$ combinations will be omitted unless necessary. The branching fractions of these baryonic decays are measured relative to the $B^0 \rightarrow J/\psi(\rightarrow p\bar{p})K^*(892)^0$ channel. So far only the

resonant decay $B^0 \rightarrow p\bar{p}K^*(892)^0$ has been seen by the BABAR [9] and Belle [10] collaborations, which measured its branching fraction to be $\mathcal{B}(B^0 \rightarrow p\bar{p}K^*(892)^0) = (1.24_{-0.25}^{+0.28}) \times 10^{-6}$ [11]. An upper limit $\mathcal{B}(B^0 \rightarrow p\bar{p}\pi^+\pi^-) < 2.5 \times 10^{-4}$ at 90% confidence level has been set by the CLEO Collaboration [12].

The data sample analyzed corresponds to an integrated luminosity of 1 fb^{-1} of proton-proton collision data at a center-of-mass energy of 7 TeV and 2 fb^{-1} at 8 TeV. The LHCb detector [13,14] is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing b or c quarks. The detector elements that are particularly relevant to this analysis are as follows: a silicon-strip vertex detector surrounding the proton-proton interaction region that allows c and b hadrons to be identified from their characteristically long flight distance; a tracking system that provides a measurement of momentum, p , of charged particles; two ring-imaging Cherenkov detectors that are able to discriminate between different species of charged hadrons; and calorimeter and muon systems for the measurement of photons and neutral hadrons, and the detection of penetrating charged particles. Simulated data samples, produced with software described in Refs. [15–20], are used to evaluate the response of the detector and to investigate possible sources of background.

Real-time event selection is performed by a trigger [21] that consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage, which performs a full event reconstruction. The hardware trigger stage requires events to have a muon with high transverse momentum, p_T , or a hadron, photon or electron with high transverse energy in the calorimeters. Signal candidates may come from events where the hardware trigger was caused either by signal particles or by other particles in the event. The software trigger requires a two-, three- or four-track secondary vertex with a

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significant displacement from any primary proton-proton interaction vertices (PVs). At least one charged particle must have $p_T > 1.6$ GeV/ c and be inconsistent with originating from a PV. A multivariate algorithm [22] is used for the identification of secondary vertices consistent with the decay of a b hadron.

The final selection of $B_{(s)}^0$ candidates, formed by combining four charged hadron candidates—a proton, an antiproton and an oppositely charged pair of light mesons—is carried out with a filtering stage, followed by requirements on the response of a boosted decision tree (BDT) classifier [23,24] and on particle identification (PID). The filtering stage includes requirements on the quality, p , p_T and χ_{IP}^2 of the tracks, loose PID requirements and an upper limit on the $p\bar{p}$ invariant mass; the χ_{IP}^2 is defined as the difference between the vertex-fit χ^2 of a PV reconstructed with and without the track in question. Each $B_{(s)}^0$ candidate must have a good-quality vertex that is displaced from the associated PV (that with which it forms the smallest χ_{IP}^2), must satisfy p and p_T requirements, and must have a reconstructed invariant mass close to that of a $B_{(s)}^0$ meson under the signal mass hypothesis. A requirement is also imposed on the angle ϑ_{dir} between the candidate momentum vector and the line between the associated PV and the candidate decay vertex.

There are 15 input quantities to the BDT classifier: p_T , η , χ_{IP}^2 , ϑ_{dir} and the flight distance of the $B_{(s)}^0$ candidate; the quality of the $B_{(s)}^0$ vertex fit; the p_T and χ_{IP}^2 of the tracks; and the largest distance of closest approach between any pair of tracks. The BDT is trained using simulated $B_{(s)}^0 \rightarrow p\bar{p}hh'$ signal candidates, generated with uniform distributions over phase space, and events in a high sideband of the $p\bar{p}K\pi$ invariant mass in data [$m(p\bar{p}K\pi)$ in the range 5450–5550 MeV/ c^2] to represent the background. Tight PID requirements are applied to all final-state particles to reduce the combinatorial background, suppress the cross-feed backgrounds between the different $p\bar{p}hh'$ final states—background from other signal decays where one particle is misidentified—and ensure that the data sets for the three $p\bar{p}hh'$ final states are mutually exclusive. For each final state individually, the requirements on the PID and BDT response are optimized for the signal significance using simulation samples for the signal. After all selection requirements are applied, approximately 3% of events with at least one candidate also contain a second candidate; a candidate is then selected at random. The efficiency of the full reconstruction and selection, including the acceptance and the trigger selection, is approximately 0.1%.

To reject contributions from intermediate charm states, candidates with hh' invariant mass consistent with a D^0 meson or phh' invariant mass consistent with a Λ_c^+ baryon are removed. The contribution from the charmonium region is removed by requiring the invariant mass of the $p\bar{p}$ pair to

be less than 2850 MeV/ c^2 , similar to the procedure in Refs. [5,25]. This last requirement is not applied to the normalization mode $B^0 \rightarrow J/\psi K^*(892)^0$, where the vector mesons are reconstructed in the $J/\psi \rightarrow p\bar{p}$ and $K^*(892)^0 \rightarrow K^+\pi^-$ decay modes. All the other steps of the selection for the signal and the normalization modes are shared in common.

The yields of the signal decays are obtained from a simultaneous unbinned extended maximum likelihood fit to the $B_{(s)}^0$ candidate invariant mass distributions in the three $p\bar{p}hh'$ final states in the range 5165–5525 MeV/ c^2 . This approach accounts for potential cross feed from one channel to another due to particle misidentification. Each signal component is modeled with a double-sided Crystal Ball (DSCB) function [26]. For each signal the tail parameters of the DSCB functions are determined from simulation. The peak position of the B^0 signals is common to the three final states, while the difference between the peak positions of the B^0 and B_s^0 signals is constrained to its known value [11]. The width of the B^0 signal is a free parameter in the $p\bar{p}K\pi$ final state and it is related to the width in the other two final states by scale factors determined from simulation. The same applies to the width of the B_s^0 signals, which is a free parameter only in the $p\bar{p}KK$ final state.

For each final state the dominant $B_{(s)}^0 \rightarrow p\bar{p}hh'$ cross-feed background is included: the $B^0 \rightarrow p\bar{p}K\pi$ mode in the $p\bar{p}KK$ and $p\bar{p}\pi\pi$ invariant mass distributions, and the $B^0 \rightarrow p\bar{p}\pi\pi$ mode in the $p\bar{p}K\pi$ spectrum. Each cross-feed background is modeled with a DSCB function with all the shape parameters fixed according to simulation; the yield is fixed relative to the yield in the correctly reconstructed final state taking into account the (mis)identification probabilities calibrated using data, as described below. In addition, a combinatorial background component modeled by an exponential function, with both parameters free to vary, is present for each final state.

The yield of the normalization decay is determined from a separate simultaneous fit to the $p\bar{p}K\pi$, $p\bar{p}$ and $K\pi$ invariant mass distributions in the ranges 5180–5380 MeV/ c^2 , 3047–3147 MeV/ c^2 and 642–1092 MeV/ c^2 , respectively. The $B^0 \rightarrow J/\psi K^*(892)^0$ component is parametrized in the $K\pi$ invariant mass distribution by a relativistic spin-1 Breit-Wigner function and in the $p\bar{p}K\pi$ and $p\bar{p}$ invariant mass distributions by DSCB functions with the tail parameters fixed from simulation. The $K\pi$ S-wave component is modeled in the $K\pi$ invariant mass distribution by the LASS parametrization [27,28] that describes nonresonant and $K_0^*(1430)^0$ S-wave contributions; this component is modeled in the $p\bar{p}K\pi$ and $p\bar{p}$ invariant mass distributions with the same shape as the $B^0 \rightarrow J/\psi K^*(892)^0$ component. A combinatorial background component modeled by a freely varying exponential function is also present in each spectrum.

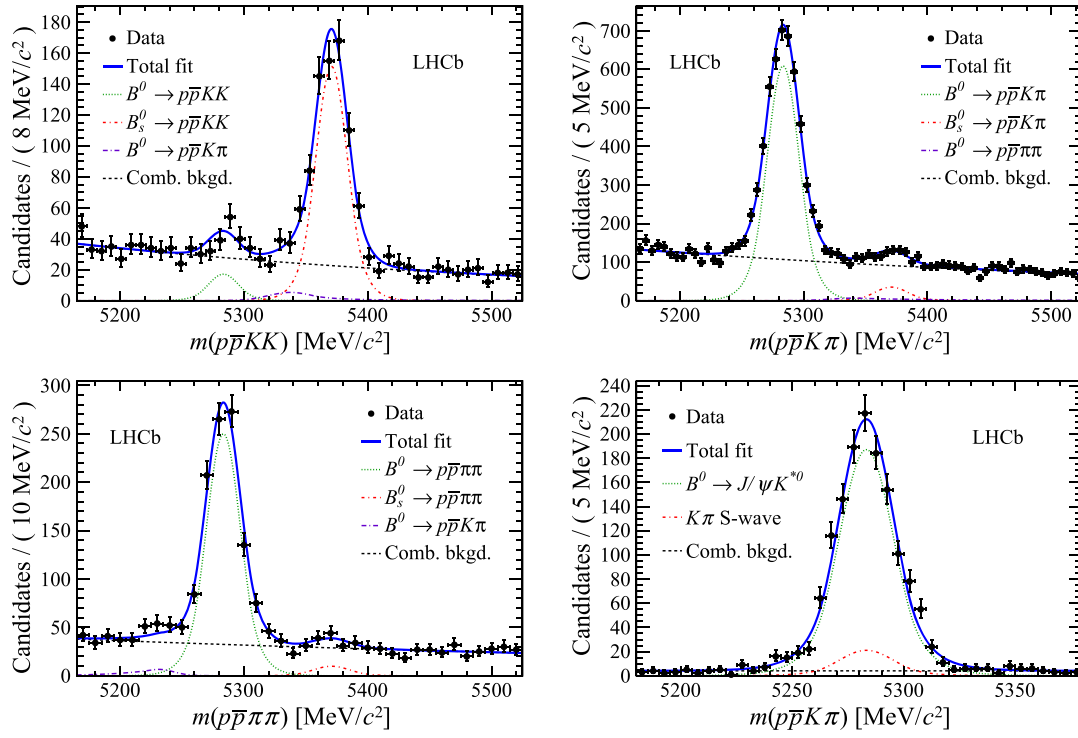


FIG. 1. Invariant mass distributions for $B_{(s)}^0$ candidates in the (top left) $p\bar{p}KK$, (top right) $p\bar{p}K\pi$, (bottom left) $p\bar{p}\pi\pi$ final state and (bottom right) invariant mass distribution of $B^0 \rightarrow J/\psi K^*(892)^0$ in the $p\bar{p}K\pi$ final state. The results of the fits are shown with blue solid lines. In the first three figures signals for B^0 and B_s^0 decays are shown, respectively, with green dotted and red dotted-dashed lines, combinatorial backgrounds are shown with black dashed lines and cross-feed backgrounds are shown with violet dotted-dashed lines. In the bottom right figure the normalization signal is shown with a green dotted line, the $K\pi$ S-wave component is displayed with a red dotted-dashed line and the combinatorial background with a black dashed line.

The $p\bar{p}hh'$ invariant mass distributions with the results of the fit overlaid are shown in Fig. 1 while the signal yields and the significances are collected in Table I. The significance of each of the signal modes is determined from the change in likelihood when the corresponding yield is fixed to zero, with systematic uncertainties taken into account [29]. The $B_s^0 \rightarrow p\bar{p}K\pi$, $B^0 \rightarrow p\bar{p}KK$ and $B_s^0 \rightarrow p\bar{p}\pi\pi$ modes are found to have significances of 6.5 standard deviations (σ), 4.1σ and 2.6σ , respectively, while the other signal modes have significances greater than 25σ .

The branching fractions of the $B_{(s)}^0 \rightarrow p\bar{p}hh'$ decays are determined relative to the visible branching fraction of the $B^0 \rightarrow J/\psi K^*(892)^0$ decay using

$$\frac{\mathcal{B}(B_{(s)}^0 \rightarrow p\bar{p}hh')}{\mathcal{B}_{\text{vis}}(B^0 \rightarrow J/\psi K^*(892)^0)} = \frac{\mathcal{N}^{\text{corr}}(B_{(s)}^0 \rightarrow p\bar{p}hh')}{\mathcal{N}^{\text{corr}}(B^0 \rightarrow J/\psi K^*(892)^0)} \left(\times \frac{f_d}{f_s} \right), \quad (1)$$

where $f_s/f_d = 0.259 \pm 0.015$ (included only for the B_s^0) is the ratio of b hadronization probabilities, f_q , to the hadron B_q [30], and $\mathcal{N}^{\text{corr}}$ denotes efficiency-corrected fitted signal yields. The yields are obtained from the mass fits, while simulation is used to evaluate the contribution to the

efficiency from each stage of the selection except for the effect of the PID criteria. The latter is determined from calibration data samples of kinematically identified pions, kaons and protons originating from the decays $D^{*+} \rightarrow D^0(\rightarrow K^-\pi^+)\pi^+$, $\Lambda \rightarrow p\pi^-$ and $\Lambda_c^+ \rightarrow pK^-\pi^+$ and weighted according to the kinematics of the signal particles [31,32]. For each final state the efficiencies are determined as a function of the position in phase space, and efficiency corrections for each candidate are applied using the method of Ref. [33] to take the variation over the phase space into account. Explicitly, $\mathcal{N}^{\text{corr}} = \sum_i \mathcal{W}_i / \varepsilon_i$, where the sum runs over the candidates in the fit, \mathcal{W}_i is the s Weight for candidate i determined with the s Plot method [34] and ε_i is the efficiency for the candidate i which depends only on its position in the five-dimensional phase space. The visible branching fraction of the normalization mode, defined as $\mathcal{B}(B^0 \rightarrow J/\psi K^*(892)^0) \times \mathcal{B}(J/\psi \rightarrow p\bar{p}) \times \mathcal{B}(K^*(892)^0 \rightarrow K^+\pi^-)$, is $\mathcal{B}_{\text{vis}}(B^0 \rightarrow J/\psi K^*(892)^0) = (1.68 \pm 0.12) \times 10^{-6}$, where the $B^0 \rightarrow J/\psi K^*(892)^0$ branching fraction is taken from Ref. [35] and the others from Ref. [11].

The branching fraction of each signal mode is reported in Table I. The significance for the $B_s^0 \rightarrow p\bar{p}\pi\pi$ mode is less than 3σ ; an upper limit on its branching fraction is found to be

TABLE I. Fitted yields, signal yield significances and branching fractions computed using Eq. (1). The uncertainties on the yields are statistical only. The first uncertainty on each branching fraction is statistical; the second is systematic; the third comes from the uncertainty on the branching fraction of the normalization mode; and the fourth, where present, is due to the uncertainty on f_d/f_s .

Decay channel	Yield \mathcal{N}	Significance (σ)	Branching fraction/ 10^{-6}
$B^0 \rightarrow p\bar{p}KK$	68 ± 17	4.1	$0.113 \pm 0.028 \pm 0.011 \pm 0.008$
$B^0 \rightarrow p\bar{p}K\pi$	4155 ± 83	>25	$5.9 \pm 0.3 \pm 0.3 \pm 0.4$
$B^0 \rightarrow p\bar{p}\pi\pi$	902 ± 35	>25	$2.7 \pm 0.1 \pm 0.1 \pm 0.2$
$B_s^0 \rightarrow p\bar{p}KK$	635 ± 32	>25	$4.2 \pm 0.3 \pm 0.2 \pm 0.3 \pm 0.2$
$B_s^0 \rightarrow p\bar{p}K\pi$	246 ± 39	6.5	$1.30 \pm 0.21 \pm 0.11 \pm 0.09 \pm 0.08$
$B_s^0 \rightarrow p\bar{p}\pi\pi$	39 ± 16	2.6	$0.41 \pm 0.17 \pm 0.04 \pm 0.03 \pm 0.02$
$B^0 \rightarrow J/\psi K^*(892)^0$	1216 ± 45	–	–

$$\mathcal{B}(B_s^0 \rightarrow p\bar{p}\pi\pi) < 6.6 \times 10^{-7} \text{ at } 90\% \text{ confidence level,}$$

by integrating the likelihood after multiplying by a prior probability distribution that is uniform in the region of a positive branching fraction. The values of the ratios of branching fractions between different $B_{(s)}^0 \rightarrow p\bar{p}hh'$ decay modes are reported in Table II.

The signal distributions in $m(hh')$ and $m(p\bar{p})$ are obtained by subtracting the background using the *sPlot* technique [34], with the $B_{(s)}^0$ candidate invariant mass as the discriminating variable. Per-candidate weights are applied to correct for the variation of the selection efficiency over the phase space. Figure 2 shows the hh' invariant mass distributions of the $B^0 \rightarrow p\bar{p}K\pi$, $B_s^0 \rightarrow p\bar{p}KK$ and $B^0 \rightarrow p\bar{p}\pi\pi$ decay modes. A peak from a vector meson is identifiable in each mass spectrum, corresponding to a $K^*(892)^0$, a $\phi(1020)$ and a $\rho(770)^0$ meson, respectively. The $p\bar{p}$ invariant mass distributions are also shown for the same decay modes. An enhancement near threshold, typical in baryonic B decays [3,4], is clearly visible in each case. Detailed amplitude analyses of the $B_{(s)}^0 \rightarrow p\bar{p}hh'$ decays will be of interest with larger samples.

The sources of systematic uncertainty on the absolute branching fractions and on the ratios of branching fractions arise from the fit model; the knowledge of the efficiencies; and, where appropriate, from the uncertainties on the branching fraction of the normalization mode and on the ratio of b -quark hadronization probabilities. Pseudoexperiments are used to estimate the effect of using alternative shapes for the fit components, or of including

TABLE II. Ratios of branching fractions among different $B_{(s)}^0 \rightarrow p\bar{p}hh'$ modes. The first uncertainty is statistical; the second is systematic; and the third, where present, comes from the uncertainty on f_d/f_s .

$\mathcal{B}(B^0 \rightarrow p\bar{p}KK)/\mathcal{B}(B^0 \rightarrow p\bar{p}K\pi)$	$0.019 \pm 0.005 \pm 0.002$
$\mathcal{B}(B^0 \rightarrow p\bar{p}\pi\pi)/\mathcal{B}(B^0 \rightarrow p\bar{p}K\pi)$	$0.46 \pm 0.02 \pm 0.02$
$\mathcal{B}(B_s^0 \rightarrow p\bar{p}K\pi)/\mathcal{B}(B^0 \rightarrow p\bar{p}K\pi)$	$0.22 \pm 0.04 \pm 0.02 \pm 0.01$
$\mathcal{B}(B_s^0 \rightarrow p\bar{p}K\pi)/\mathcal{B}(B_s^0 \rightarrow p\bar{p}KK)$	$0.31 \pm 0.05 \pm 0.02$

additional components in the fit. In particular, the effect of adding other cross-feed backgrounds, partially reconstructed backgrounds and components coming from Λ_b^0 decays have been investigated. These are the dominant sources of systematic uncertainty for the $B^0 \rightarrow p\bar{p}KK$ and $B_s^0 \rightarrow p\bar{p}\pi\pi$ modes. The effect of fixing the yields of the cross-feed backgrounds based on the (mis)identification probabilities is also assessed by varying these probabilities within their uncertainties. Intrinsic biases in the fitted yields are investigated with pseudoexperiments and are found to be negligible. Uncertainties on the efficiencies arise due to the limited size of the simulation samples, the uncertainty on their evaluated distributions across the phase space of the decays and from possible residual differences between data and simulation. The unknown decay kinematics are the principal source of systematic uncertainty for the $B_s^0 \rightarrow p\bar{p}K\pi$ mode, while for the $B_s^0 \rightarrow p\bar{p}KK$, $B^0 \rightarrow p\bar{p}K\pi$ and $B^0 \rightarrow p\bar{p}\pi\pi$ modes the dominant source of systematic uncertainty comes from the uncertainty on the efficiency of the hardware stage of the trigger. As the efficiencies depend on the signal decay-time distribution, the effect coming from the different lifetimes of the B_s^0 mass eigenstates has been evaluated. The systematic uncertainties due to the vetoes of charm hadrons are also included.

In summary, a search for the four-body charmless baryonic decays $B_{(s)}^0 \rightarrow p\bar{p}hh'$ has been carried out by the LHCb Collaboration with a sample of proton-proton collision data corresponding to an integrated luminosity of 3 fb^{-1} . First observations are obtained for the decays $B^0 \rightarrow p\bar{p}\pi\pi$, nonresonant $B^0 \rightarrow p\bar{p}K\pi$, $B_s^0 \rightarrow p\bar{p}KK$ and $B_s^0 \rightarrow p\bar{p}K\pi$, while first evidence is reported for the $B^0 \rightarrow p\bar{p}KK$ mode and an upper limit is set on the $B_s^0 \rightarrow p\bar{p}\pi\pi$ branching fraction. In particular, four-body baryonic B_s^0 decays are observed for the first time and a threshold enhancement in the baryon-antibaryon mass spectra is confirmed for baryonic B_s^0 decays [2].

The LHCb Collaboration has recently published studies of CP violation with four-body $\Lambda_b^0 \rightarrow ph^-h^+h^-$ decays studying triple-product correlations, and presented first evidence for CP violation in baryons [36]. The decays of B^0 and B_s^0 mesons to $p\bar{p}hh'$ final states reported in this

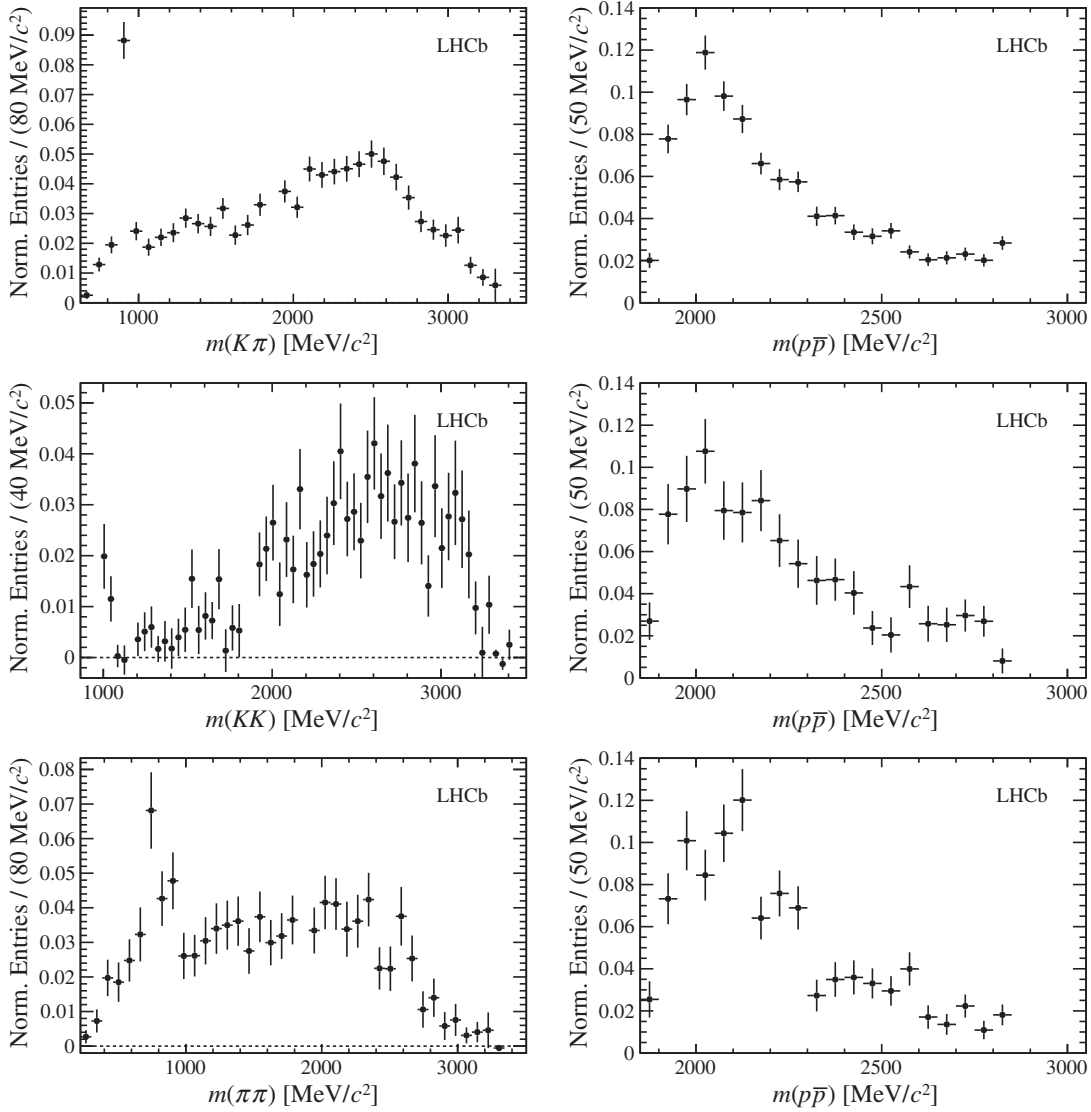


FIG. 2. Efficiency-corrected and background-subtracted (left) $m(hh')$ and (right) $m(p\bar{p})$ distributions from (top) $B^0 \rightarrow p\bar{p}K\pi$, (middle) $B_s^0 \rightarrow p\bar{p}KK$, and (bottom) $B^0 \rightarrow p\bar{p}\pi\pi$ candidates. Events with entries in the charmonium or D^0 mass regions have been removed from the samples. All distributions are normalized to unity.

paper may be used in the future for similar studies of CP violation in baryonic B decays.

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