

## ANOMALOUS SINGLE PRODUCTION OF THE FOURTH GENERATION QUARKS AT FUTURE $ep$ AND $\gamma p$ COLLIDERS

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**Abstract.** - Possible anomalous single productions of the fourth standard model generation up and down type quarks at LHC based  $ep$  and  $\gamma p$  colliders are studied. Some decay channels are considered. Signatures for signals and corresponding standard model backgrounds are discussed. Discovery limits for quark mass and achievable values of anomalous coupling strength are determined.

Colliders with TeV enegy are needed for discovery of the fourth SM generation fermions if they exist (see [1] and references therein). The fourth generation quarks will be produced in pairs copiously at the Large Hadron Collider (LHC). Lepton colliders are the best place for pair production of the fourth generation charged lepton and neutrino. The discovery capacity of lepton collider could be enlarged if the anomalous interactions of the fourth generation fermions with the first three ones exist. Such anomalous interactions seems to be quite natural due to large masses of the fourth generation fermions. These anomalous interactions could provide also single production of the fourth generation fermions at future lepton-hadron [2, 3] and gamma-hadron colliders. Lepton-hadron colliders with  $\sqrt{s} = 1.3-1.4$  TeV are named QCD Explorer or LHeC depending on electrons provided by linac or ring, respectively [4]. Reference [5] has given estimations of the luminosity values for various configuration and parameter sets of LHC and electron accelerator. For the choice of a cw superconducting energy recovery e-linac with a upgrated LHC, the integrated luminosity for a year has been estimated  $10fb^{-1}$ . The integrated luminosity for a year has been estimated  $4.1fb^{-1}$ , for the choice of pulsed superconducting e-linac with a

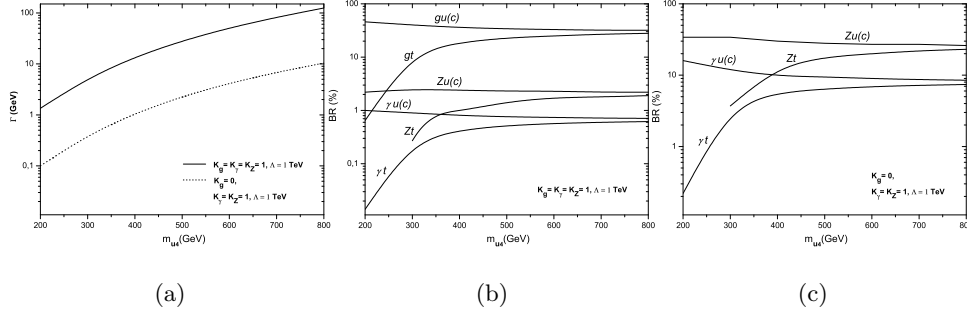


Figure 1: (a) The total and (b), (c) the partial decay widths of the fourth SM generation up type quarks as a function of the quark mass

upgraded LHC. While the first choice has better luminosity, it is not possible to found a gamma-p collider on base of it. The real gamma production through Compton back scattering technique on the base of the ep collider is possible on the second choice. In this case, the integrated luminosity of  $\gamma p$  collider reaches the maximum value of  $2.665 fb^{-1}$ . In this study the possible anomalous single productions of the fourth generation up and down type quarks at LHC based  $ep$  and  $\gamma p$  colliders are considered.

The effective Lagrangian for the flavor changing neutral current (FCNC) interactions of  $u_4$  and  $d_4$  quarks can be rewritten from [6, 7] with minor modifications as:

$$\mathcal{L} = \left( \frac{\kappa_\gamma^{q_i}}{\Lambda} \right) e_q g_e \bar{q}_4 \sigma_{\mu\nu} q_i F^{\mu\nu} + \left( \frac{\kappa_Z^{q_i}}{2\Lambda} \right) g_Z \bar{q}_4 \sigma_{\mu\nu} q_i Z^{\mu\nu} + \left( \frac{\kappa_g^{q_i}}{\Lambda} \right) g_s \bar{q}_4 \sigma_{\mu\nu} T^a q_i G_a^{\mu\nu} + H.c. \quad (1)$$

where  $i = 1, 2, 3$  denotes the generation index.  $\kappa_\gamma^{q_i}$ ,  $\kappa_Z^{q_i}$  and  $\kappa_g^{q_i}$  are anomalous couplings for the electromagnetic, the weak (neutral current) and the strong interactions, respectively.  $\Lambda$  is the cutoff scale for the new physics and  $e_q$  is the quark charge.  $g_e$ ,  $g_Z$  and  $g_s$  are the electroweak and the strong coupling constants. In the above equa-

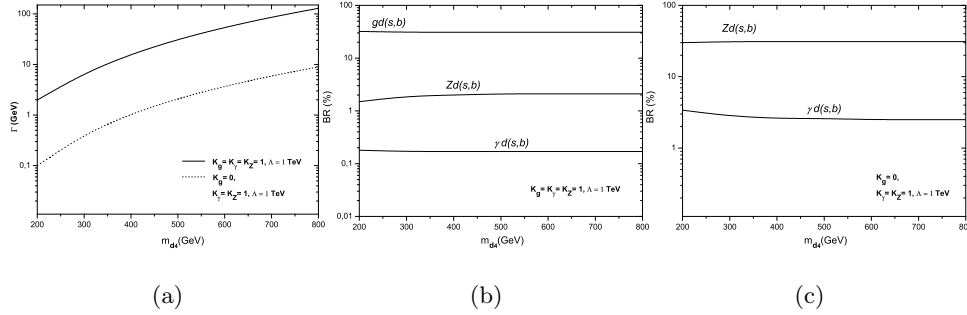


Figure 2: (a) The total and (b), (c) the partial decay widths of the fourth SM generation down type quarks as a function of the quark mass

Table 1: Signal and SM background cross sections for  $ep \rightarrow u_4(d_4)X \rightarrow q\ell^+\ell^-X$  and  $\gamma p \rightarrow u_4(d_4)X \rightarrow q\gamma X$  processes

$m_4$ (GeV)	Quark Type	Signal $\sigma$ (fb) for $ep \rightarrow u_4(d_4)X \rightarrow q\ell^+\ell^-X$		Signal $\sigma$ (fb) for $\gamma p \rightarrow u_4(d_4)X \rightarrow q\gamma X$	
		$\kappa/\Lambda = 1 \text{ TeV}^{-1}$	$\kappa_g/\Lambda = 0$	$\kappa/\Lambda = 1 \text{ TeV}^{-1}$	$\kappa_g/\Lambda = 0$
No Cut					
300	$u_4$	8.68	163	564	7340
	$d_4$	2.44	40	33.9	540
600	$u_4$	1.45	19	136	1820
	$d_4$	0.32	5	5.13	81
SM Bck. $\sigma$ (fb)		97.7		$3.19 \cdot 10^8$	
With Cut					
300	$u_4$	—	3.37	380	4940
	$d_4$	—	1.75	—	390
600	$u_4$	—	0.94	61	800
	$d_4$	—	0.35	—	38
SM Bck. $\sigma$ (fb)		0.64		1220	

tion,  $\sigma_{\mu\nu} = i(\gamma_\mu\gamma_\nu - \gamma_\nu\gamma_\mu)/2$ .  $F^{\mu\nu}$ ,  $Z^{\mu\nu}$  and  $G_a^{\mu\nu}$  are field strength tensors of the photon, the Z boson and gluons, respectively.  $T_a$  is the Gell-Mann matrices.

We have calculated the anomalous single production cross sections of the fourth SM generation quarks at the linac-LHC and  $\gamma p$  colliders based on it using CompHEP with CTEQ6L1 [8]. The total decay width  $\Gamma$  of the fourth generation up (down) type quarks and the relative branching ratios are plotted with assumption of  $\kappa_\gamma^{q_i} = \kappa_Z^{q_i} = \kappa_g^{q_i}$  and  $\kappa_\gamma^{q_i} = \kappa_Z^{q_i} = 1$ ,  $\kappa_g^{q_i} = 0$  in Fig. 1 (Fig. 2).  $\Lambda = 1\text{TeV}$  is selected at calculations. Single anomalous production cross sections of the fourth generation quarks are given for ep collider at Fig. 3a and  $\gamma p$  collider at Fig. 3b.

In this study,  $ep \rightarrow u_4(d_4)X \rightarrow qZ X \rightarrow q\ell^+\ell^-X$  and  $\gamma p \rightarrow u_4(d_4)X \rightarrow q\gamma X$  processes (and their H.c.) are considered as a signature of anomalous interactions of

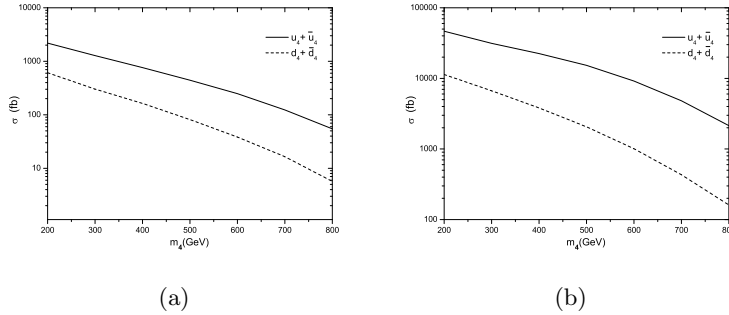


Figure 3: The total production cross sections of the fourth SM generation up and down type quarks at (a)  $ep$  and (b)  $\gamma p$  colliders

Table 2: Statistical significances (SS) for for  $ep \rightarrow u_4(d_4)X \rightarrow q\ell^+\ell^-X$  and  $\gamma p \rightarrow u_4(d_4)X \rightarrow q\gamma X$  processes by using  $\kappa_\gamma^{q_i} = \kappa_Z^{q_i} = 1$ ,  $\kappa_g^{q_i} = 0$

$m_4$ (GeV)	Quark Type	SS for $ep \rightarrow d_4(u_4)X \rightarrow q\ell^+\ell^-X$		SS for $\gamma p \rightarrow d_4(u_4)X \rightarrow q\gamma X$
		$L = 10 \text{ fb}^{-1}$	$L = 4.1 \text{ fb}^{-1}$	$L = 2.665 \text{ fb}^{-1}$
300	$u_4$	13.3	8.5	231
	$d_4$	6.9	4.4	18.2
600	$u_4$	3.7	2.4	37.4
	$d_4$	1.4	0.9	1.8

the fourth generation up and down type quarks ( $q$  is  $u$  or  $c$  for  $u_4$  and  $d, s$  or  $b$  for  $d_4$  and  $\ell$  is  $e$  or  $\mu$ ). The SM background for this processes is potentially much larger than the signal. However, after applying some kinematic cuts, it is possible to decrease background to the reasonable levels. We choose the following cuts for the first process:  $P_T > 10$  GeV for the scattered electron,  $P_T > 80$  GeV for leptons coming from Z boson,  $P_T > 20$  GeV for jet;  $|\eta_{j,l}| < 2.5$ ,  $\Delta R > 0.4$  between the leptons and jet. For the second process similar cuts are applied. Only leptons on the cuts related to leptons coming from Z boson are replaced by photons. The calculated signal and SM background cross sections are given in Table 1. The statistical significance (SS) values, evaluated from  $SS = (\sigma_S/\sqrt{\sigma_B})\sqrt{L_{int}}$ , where  $L_{int}$  is the integrated luminosity of the collider, are presented in Table 2 for both processes after applying above cuts.

As a result of this study it is shown that when anomalous coupling for strong interactions is close to one,  $ep$  and  $\gamma p$  colliders are almost blind to anomalous interactions. These colliders give possibility to investigate effects of both anomalous couplings of electromagnetic and weak interactions for  $\kappa_g^{q_i} = 0$ . Observation limits (at  $3\sigma$ ) as low as 0.33 (0.52)  $\text{TeV}^{-1}$  are reachable for the  $(\kappa_\gamma/\Lambda)$  at  $m_4 = 300$  (600) GeV for the up and down type quarks (combined) at  $\gamma p$  collider. Meanwhile, 0.38 (0.76)  $\text{TeV}^{-1}$  for the  $(\kappa_Z/\Lambda)$  is reachable at  $m_4 = 300$  (600) GeV for the ep collider with  $L = 10 \text{ fb}^{-1}$ .

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