π ⁻ + p \rightarrow A⁰ + Θ ⁰ ANGULAR DISTRIBUTION OF PRODUCTION $\rightarrow p + \pi^-$ **AND DECAY at** 1.12 **GeV/c** f

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Summary of maximum likelihood fit to the data

At 1.12 GeV/c, 236 A^0 -decays arising from $\pi^- + p \rightarrow A^0 + \theta^0$ were studied.

The data were fitted with the expression

$$
d\overline{N} = \varepsilon(\theta) [A_1 + A_2 \cos \theta + A_3 \cos^2 \theta + \xi \sin \theta (A_4 + A_5 \cos \theta)]
$$

$$
\times 2\pi d(\cos \theta) \frac{d\xi}{2}
$$
 (1)

- $\varepsilon(\theta)$ = probability that a Λ^0 produced at angle θ decays inside the chamber
- θ = centre of mass production angle of the A^0
- **f** $=$ $\left[p_{\pi \text{ inc}} \times p_{\pi} \cdot p_{\pi \text{ decay}}\right] / \text{(magnitude of same)}.$

This expression is what one obtains for production involving final *s*- and *p*-waves only ¹. (The centre of mass momentum is here 300 MeV/c.)

The likelihood function

$$
\mathfrak{L} = \prod_{ij} e^{-N_{ij}} (\bar{N})_{ij} \quad i = 1, 2 \dots 6 = \text{histogram interval} \n\text{in } \cos \theta
$$

$$
j = 1, 2 \dots 4 = \text{histogram interval}
$$

in ξ

$$
\overline{N}_{ij} = \text{counts from expression (1)}
$$

$$
N_{ij} = \text{observed counts}
$$

was formed and maximized by an iterative procedure.

The solution that maximises \mathfrak{L} is :

$$
\begin{array}{cccccc}\nA_1 & A_2 & A_3 & A_4 & A_5 \\
21.892 & -27.008 & 11.944 & 14.065 & -32.278\n\end{array}
$$

With this solution is associated a 5×5 error matrix $\left| \delta A_i \delta A_j \right|$

All five A 's are different from zero outside of experimental error.

The best fit to the data are shown in Fig. **1A,** and in Fig. **5** of Steinberger's report.

From the plot of

$$
aP(\theta) = \frac{\sin \theta \left(A_4 + A_5 \cos \theta \right)}{A_1 + A_2 \cos \theta + A_3 \cos^2 \theta} \tag{2}
$$

(Fig. **5)** a lower limit to *a* can be obtained as follows :

The maximum of aP is: $aP_{\text{max}} = 0.73 \pm 0.14$, where the error is obtained by differentiating (2) and using the error matrix given above.

Then, since the polarization cannot exceed unity, we have a lower limit on α :

$$
\left|a_{\rm min}\right| = 0.73 \pm 0.14
$$

Our data are consistent with any value of $|a|$ lying between **0.73** and unity.

t Appendix to Session 5. — Experimental.

Fig. 1a: $\pi^+ + p \to A^0 + \theta^0$ production angular distribution. Curve fitted to the experimental points using s- and p-waves.

Summary of s- and **p-wave** analysis of π ⁺+p $\rightarrow \Lambda$ +K[°] at 1.1 **GeV/c (236 A-decays)**

The paper of Lee et al. $¹$ gives the theoretical distribu-</sup> tion, $W(\theta, \xi)$, in ξ ("the up-down" direction cosine) and θ of the decay product of the *A⁰ .* The assumptions of this paper are *(a)* that parity is conserved in the production process, (b) but not conserved in the decay of the A^0 , and *(c)* that only *s*- and *p*-waves are present in the final Λ^0 , K^0 -state.

$$
W(\theta, \xi) d\xi d\Omega = \left\{ I(\theta) + a \xi I(\theta) P(\theta) \right\} \frac{d\xi}{2} d\Omega
$$

$$
I(\theta) = |\mathbf{a} + \mathbf{b} \cos \theta|^2 + |\mathbf{c}|^2 \sin^2 \theta
$$

$$
P(\theta) = \frac{2 \operatorname{Im} \mathbf{c}^*(\mathbf{a} + \mathbf{b} \cos \theta) \sin \theta}{I(\theta)}
$$

Now if *a* were known we could use the **5** least squares constants $A_1, A_2, \ldots A_5$ of

 $dN = {A_1 + A_2 \cos \theta + A_3 \cos^2 \theta + \xi \sin \theta (A_4 + A_5 \cos \theta)}$

 $\frac{d\xi}{2}$ 2*nd* cos θ

to determine a, b, c, φ and ψ , where

$$
\mathbf{a} = ae^{i\varphi},\n\mathbf{b} = be^{i\psi},\n\mathbf{c} = c.
$$

In order to obtain c , the spin-flip p -wave amplitude we must solve a cubic equation. Only two of the three solutions are real. These real solutions are called the $k = 0$ and $k = 2$ solutions. There is, for each of these two solutions, two alternative solutions for the phase angles, φ and ψ . The alternate solution is obtained by reflecting the " normal solution " about the imaginary axis in the complex plane.

a, b, and *c* have been normalized so that $\int I(\theta) d\Omega = 1$. φ and ψ are given in radians.

 $k = 0$ Solution (a = 0.95^(*)) Error Matrix

 $k = 2$ Solution ($\alpha = 0.95$) Error Matrix

^(*) From the plot of *a, b,* etc. vs *a* we see that these quantities are not rapidly varying and are almost independent of the value of *a.*

 $k = 0$ Solution

 $k=0$

Non-spin-flip cross-section

$$
\sigma_{\text{non-spin-flip}} = \sigma_{\text{tot}} \frac{1}{3} |b|^2 = \frac{(k=0)}{0.0273 \, , \, 0.0261} \, ,
$$

Spin-flip cross-section

$$
\sigma_{\text{spin-flip}} = \sigma_{\text{tot}} \frac{\frac{2}{3} |c|^2 = 0.0299, 0.0277}{(a^2 + \frac{1}{3}b^2 + \frac{2}{3}c^2)}
$$

 \mathcal{A}

 \cdot c

 \overline{c}

$$
\delta\left(\frac{\sigma_s}{\sigma_{\text{tot}}}\right) = 4\pi \, \delta(a^2) = 2 \times 4\pi \, a \, \delta a = 2 \cdot 4\pi (0.149) \, (0.044) \n= (0.082) \times 2 = 0.164
$$

$$
\delta \left(\frac{\sigma_{\text{NSF}}}{\sigma_{\text{tot}}} \right) = \frac{4\pi}{3} \delta(b)^2 = \frac{8\pi}{3} b \delta b = \frac{8\pi}{3} (0.286) (0.033) = 0.078
$$

$$
\delta \left(\frac{\sigma_{\rm SF}}{\sigma_{\rm tot}} \right) = \frac{8\pi}{3} \delta(c)^2 = \frac{16\pi}{3} c \delta c = \frac{16\pi}{3} (0.212)(0.033) = 0.116
$$

$$
k = 2
$$
\n
$$
\delta \left(\frac{\sigma_s}{\sigma_{\text{tot}}} \right) = 8\pi (0.161) (0.0628) = 0.254
$$
\n
$$
\delta \left(\frac{\sigma_{\text{NSF}}}{\sigma_{\text{tot}}} \right) = \frac{8\pi}{3} (0.280) (0.0417) = 0.0562
$$
\n
$$
\delta \left(\frac{\sigma_{\text{SF}}}{\sigma_{\text{tot}}} \right) = \frac{16\pi}{3} (0.204) (0.0530) = 0.181
$$

LIST OF REFERENCES

1. Lee, T. D., Steinberger, J., Feinberg, G., Kabir, P. K. and Yang, C. N., Phys. Rev., *106,* p. 1367, 1957.

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