$\pi^{-} + p \rightarrow \Lambda^{0} + \Theta^{0} \qquad \text{ANGULAR DISTRIBUTION OF PRODUCTION}$ $\downarrow \rightarrow p + \pi^{-}$ AND DECAY at 1.12 GeV/c †

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Summary of maximum likelihood fit to the data

At 1.12 GeV/c, 236 Λ^0 -decays arising from $\pi^- + p \rightarrow \Lambda^0 + \theta^0$ were studied.

The data were fitted with the expression

$$d\overline{N} = \varepsilon(\theta) [A_1 + A_2 \cos \theta + A_3 \cos^2 \theta + \xi \sin \theta (A_4 + A_5 \cos \theta)] \\ \times 2\pi d(\cos \theta) \frac{d\xi}{2} \qquad (1)$$

- $\varepsilon(\theta) = \text{probability that a } \Lambda^0 \text{ produced at angle } \theta \text{ decays}$ inside the chamber
- θ = centre of mass production angle of the Λ^0
- $\xi = [\mathbf{p}_{\pi \text{ inc}} \times \mathbf{p}_A \cdot \mathbf{p}_{\pi \text{ decay}}] / \text{ (magnitude of same).}$

This expression is what one obtains for production involving final *s*- and *p*-waves only ¹⁾. (The centre of mass momentum is here 300 MeV/c.)

The likelihood function

$$\mathfrak{L} = \prod_{ij} e^{-\overline{N}_{ij}} (\overline{N})_{ij} \quad i = 1, 2 \dots 6 = \text{histogram interval}$$

in cos θ

$$j = 1, 2 \dots 4$$
 = histogram interval in ξ

$$\overline{N}_{ij}$$
 = counts from expression (1)
 N_{ij} = observed counts

was formed and maximized by an iterative procedure.

The solution that maximises \mathfrak{L} is :

$$A_1$$
 A_2 A_3 A_4 A_5
21.892 - 27.008 11.944 14.065 - 32.278

With this solution is associated a 5 \times 5 error matrix $|\delta A_i \delta A_j|$

	1	2	3	4	5
1	5.6480	-0.2248	-9.2172	+2.2051	-1.8025
2		8.8182	-7.1432	-2.0162	+1.7779
3			30.2028	-2.0248	-5.4913
4				11.5237	-9.5883
5					50.3981

All five A's are different from zero outside of experimental error.

The best fit to the data are shown in Fig. 1A, and in Fig. 5 of Steinberger's report.

From the plot of

$$aP(\theta) = \frac{\sin\theta \left(A_4 + A_5 \cos\theta\right)}{A_1 + A_2 \cos\theta + A_3 \cos^2\theta}$$
(2)

(Fig. 5) a lower limit to α can be obtained as follows :

The maximum of αP is: $\alpha P_{\text{max}} = 0.73 \pm 0.14$, where the error is obtained by differentiating (2) and using the error matrix given above.

Then, since the polarization cannot exceed unity, we have a lower limit on α :

$$\left|a_{\min}\right| = 0.73 \pm 0.14$$

Our data are consistent with any value of $|\alpha|$ lying between 0.73 and unity.

[†] Appendix to Session 5. - Experimental.



Fig. 1a: $\pi^- + p \rightarrow \Lambda^0 + \theta^0$ production angular distribution. Curve fitted to the experimental points using s- and p-waves.

Summary of s- and p-wave analysis of $\pi^- + p \rightarrow \Lambda + K^0$ at 1.1 GeV/c (236 Λ -decays)

The paper of Lee et al.¹⁾ gives the theoretical distribution, $W(\theta, \xi)$, in ξ ("the up-down" direction cosine) and θ of the decay product of the Λ^0 . The assumptions of this paper are (a) that parity is conserved in the production process, (b) but not conserved in the decay of the Λ^0 , and (c) that only s- and p-waves are present in the final Λ^0 , K^0 -state.

$$W(\theta,\xi) d\xi d\Omega = \left\{ I(\theta) + a\xi I(\theta) P(\theta) \right\} \frac{d\xi}{2} d\Omega$$
$$I(\theta) = |\mathbf{a} + \mathbf{b} \cos \theta|^2 + |\mathbf{c}|^2 \sin^2 \theta$$
$$2 \operatorname{Im} \mathbf{c}^*(\mathbf{a} + \mathbf{b} \cos \theta) \sin \theta$$

$$P(\theta) = \frac{2 \operatorname{Im} \mathbf{c}^* (\mathbf{a} + \mathbf{b} \cos \theta) \sin \theta}{I(\theta)}$$

Now if a were known we could use the 5 least squares constants $A_1, A_2, \ldots A_5$ of

 $dN = \{A_1 + A_2 \cos \theta + A_3 \cos^2 \theta + \xi \sin \theta (A_4 + A_5 \cos \theta)\}$

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to determine a, b, c, φ and ψ , where

In order to obtain c, the spin-flip p-wave amplitude we must solve a cubic equation. Only two of the three solutions are real. These real solutions are called the k = 0 and k = 2 solutions. There is, for each of these two solutions, two alternative solutions for the phase angles, φ and ψ . The alternate solution is obtained by reflecting the "normal solution" about the imaginary axis in the complex plane.

a, b, and c have been normalized so that $\int I(\theta) d\Omega = 1$. φ and ψ are given in radians.

$a = 0.149 \pm 0.044$	$\frac{\overline{\delta a^2}}{0.001902}$	$\overline{\delta a \delta b}$ $-$ 0.001078	$\frac{\overline{\delta a \delta c}}{-0.001214}$	$\overline{\delta a \delta \varphi}$ 0.008778	$\overline{\delta a \delta \psi}$ — 0.02174
$b = 0.286 \pm 0.033$		$\frac{\delta b^2}{0.001098}$	<i>δbδc</i> 0.0005559	$\delta b \delta arphi \ -$ 0.005680	δ <i>bδψ</i> 0.01167
$c = 0.212 \pm 0.033$			$\frac{\overline{\delta c^2}}{0.001067}$	$\delta c \delta arphi \ -$ 0.005429	$\overline{\delta c \delta \psi}$ 0.01520
$arphi = 2.34 \pm 0.31$ (0.70 \pm 0.31 alt)				$\overline{\delta arphi^2}$ 0.09767	$\overline{\delta arphi \delta \psi} = 0.06420$
$\psi = 5.24 \pm 0.59$ (4.18 \pm 0.59 alt)					$\frac{\overline{\delta\psi^2}}{0.34911}$

k = 0 Solution (a = 0.95 (*))

Error Matrix

k = 2 Solution (a = 0.95)

Error Matrix

$a = 0.161 \pm 0.063$	$\frac{\overline{\delta a^2}}{0.003945}$	δ <i>aδb</i> — 0.002193	$\overline{\delta a \delta c}$ $-$ 0.003125	$\overline{\delta a \delta arphi} = 0.007259$	$\overline{\delta a \delta \psi} \ - 0.03127$
$b = 0.280 \pm 0.042$		$\frac{\overline{\delta b^2}}{0.001743}$	$\overline{\delta b \delta c}$ 0.001584	$\overline{\delta b \delta \varphi}$ 0.004829	$\overline{\delta b \delta \psi}$ 0.01608
$c = 0.204 \pm 0.053$			$\frac{\overline{\delta c^2}}{0.002810}$	$\overline{\delta c \delta \varphi}$ 0.005371	$\overline{\frac{\delta c \delta \psi}{0.02623}}$
$arphi = 2.37 \pm 0.25 \ (0.77 \pm 0.25 \text{ alt})$				$\overline{\delta arphi^2}$ 0.06358	$\overline{\delta arphi \delta \psi} \ 0.01102$
$arphi \ = 5.09 \pm 0.58 \ (4.30 \pm 0.58 \ ext{alt})$					$\overline{\delta\psi^2}$ 0.3314

^(*) From the plot of a, b, etc. vs a we see that these quantities are not rapidly varying and are almost independent of the value of a.

K = 0 SOLUTION



s-wave cross-section $|\mathbf{a}|^2 = \overset{(k=0)}{0.0222}, \overset{(k=2)}{0.0259}$ $\sigma_{s-wave} = \sigma_{tot}$ $|a|^{2}+\frac{1}{3}|b|^{2}+\frac{2}{3}|c|^{2}$ $\overline{4\pi}$ k = 0k = 2 σ_{s-wave} $\textbf{0.28} \pm \textbf{0.16}$ 0.33 ± 0.25 ____ σtot

Non-spin-flip cross-section

$$\sigma_{\text{non-spin-flip}} = \sigma_{\text{tot}} \frac{\frac{1}{3}|\mathbf{b}|^2 = 0.0273, \ 0.0261}{(\mathbf{a}^2 + \frac{1}{3}\mathbf{b}^2 + \frac{2}{3}\mathbf{c}^2)}$$

(k - 0)

(k-2)

	$\kappa = 0$	$\kappa = 2$
$\frac{\sigma_{\text{non-spin-flip}}}{\sigma_{\text{tot}}}$	0.34 ± 0.08	0.33 ± 0.06

Spin-flip cross-section

$$\sigma_{\text{spin-flip}} = \sigma_{\text{tot}} \frac{\frac{2}{3} |\mathbf{c}|^2 = 0.0299}{(\mathbf{a}^2 + \frac{1}{3} \mathbf{b}^2 + \frac{2}{3} \mathbf{c}^2)}$$

	k = 0	k = 2
$\frac{\sigma_{\text{spin-flip}}}{\sigma_{\text{tot}}}$	0.38 ± 0.12	0.35 ± 0.18

k = 0 $\delta\left(\frac{\sigma_s}{\sigma_{\text{tot}}}\right) = 4\pi \,\delta(a^2) = 2 \times 4\pi \, a \,\delta a = 2 \cdot 4\pi (0.149) \,(0.044)$ $= (0.082) \times 2 = 0.164$

c

С

(ALTERNATE)

$$\delta\left(\frac{\sigma_{\rm NSF}}{\sigma_{\rm tot}}\right) = \frac{4\pi}{3}\,\delta(b)^2 = \frac{8\pi}{3}\,b\,\delta\,b = \frac{8\pi}{3}\,(0.286)\,(0.033) = 0.078$$

$$\delta\left(\frac{\sigma_{\rm SF}}{\sigma_{\rm tot}}\right) = \frac{8\pi}{3}\,\delta(c)^2 = \frac{16\pi}{3}\,c\,\delta\,c = \frac{16\pi}{3}(0.212)\,(0.033) = 0.116$$

$$k = 2$$

$$\delta \left(\frac{\sigma_s}{\sigma_{\text{tot}}} \right) = 8\pi (0.161) (0.0628) = 0.254$$

$$\delta \left(\frac{\sigma_{\text{NSF}}}{\sigma_{\text{tot}}} \right) = \frac{8\pi}{3} (0.280) (0.0417) = 0.0562$$

$$\delta \left(\frac{\sigma_{\text{SF}}}{\sigma_{\text{tot}}} \right) = \frac{16\pi}{3} (0.204) (0.0530) = 0.181$$

LIST OF REFERENCES

1. Lee, T. D., Steinberger, J., Feinberg, G., Kabir, P. K. and Yang, C. N., Phys. Rev., 106, p. 1367, 1957.

Appendix I

