Dark Matter Searches at Colliders

Jalal Abdallah

On behalf of ATLAS and CMS Collaborations

11th Patras Workshop on Axions, WIMPs and WISPs

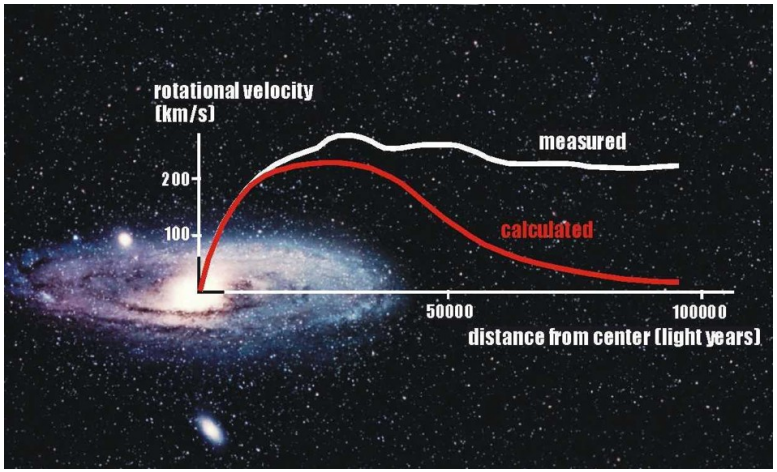
June 24th, 2015

- Motivation for Dark Matter at colliders
- EFT and simplified models
- ATLAS and CMS Dark Matter searches
- Prospective for Run-II
- Conclusions

All ATLAS and CMS public results:

https://twiki.cern.ch/twiki/bin/view/AtlasPublic
https://twiki.cern.ch/twiki/bin/view/CMSPublic

Motivation for dark matter

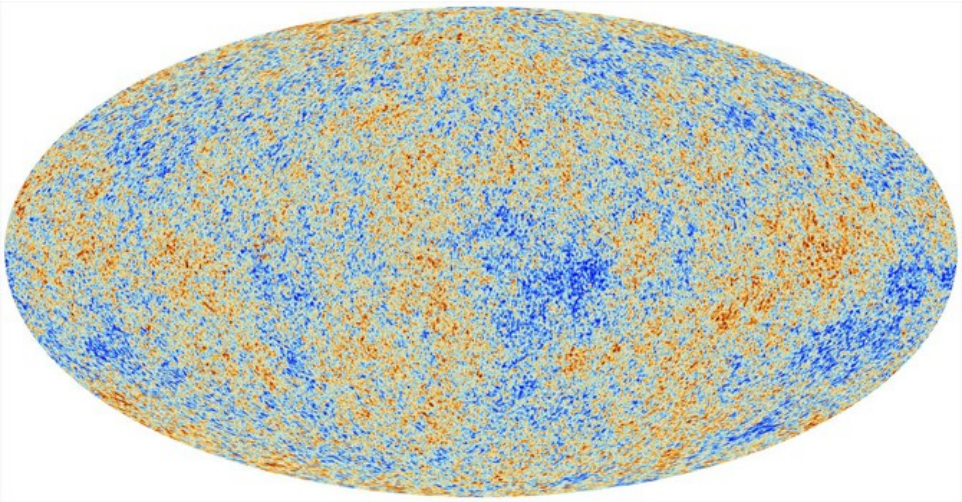


'Anomalous' galactic rotation curves

Gravitational lensing effect in deep field Hubble

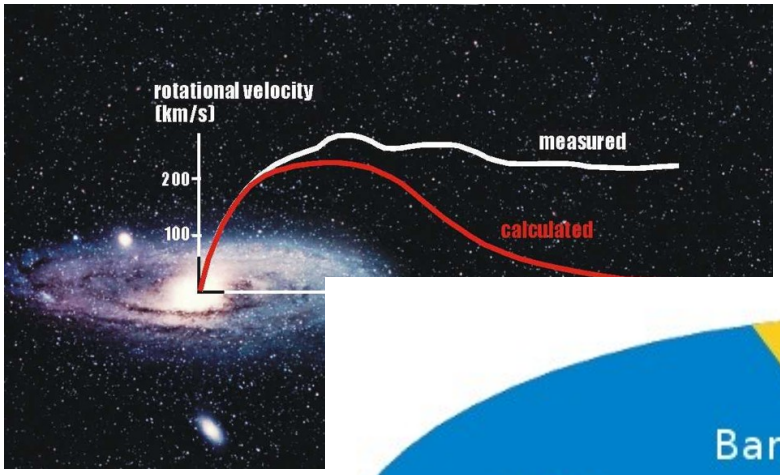
Should I really go through this slide?!

OK so dark matter exists and we all know it!

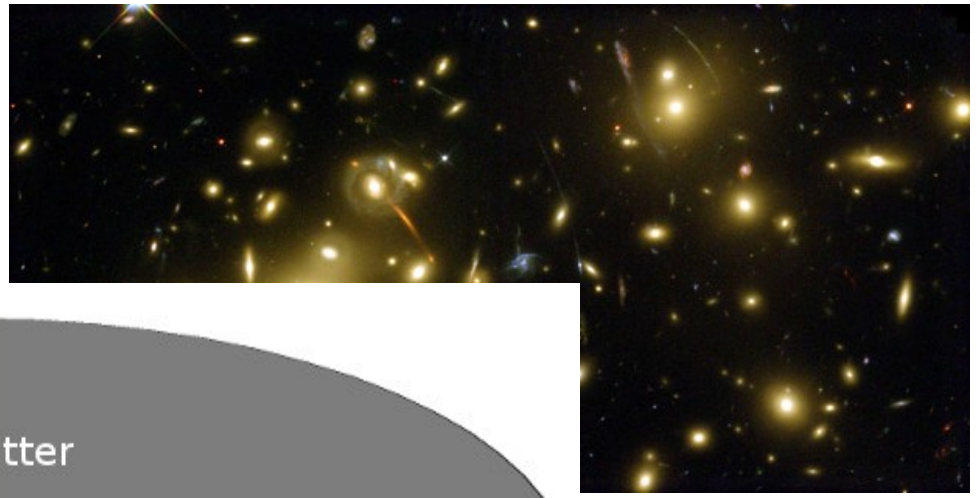


Planck 2013: Dark matter vs luminous matter from CMB anisotropy measurements

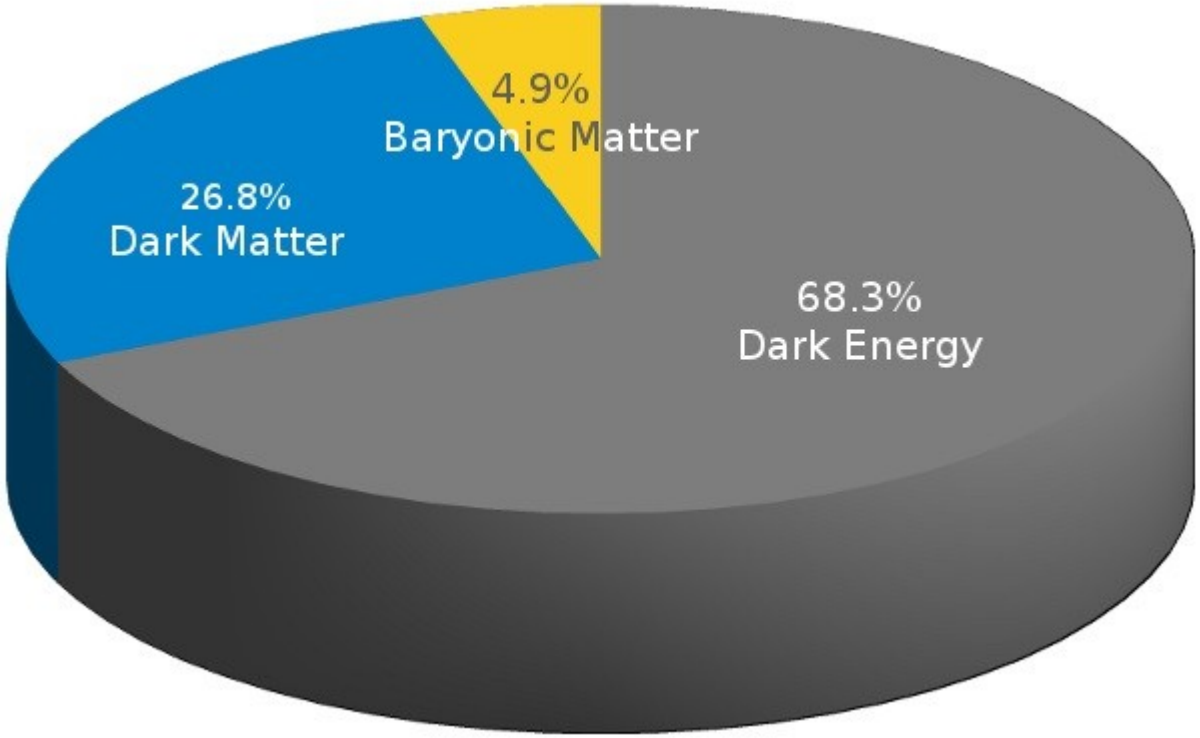
Motivation for dark matter



'Anomalous' ga



in deep field Hubble



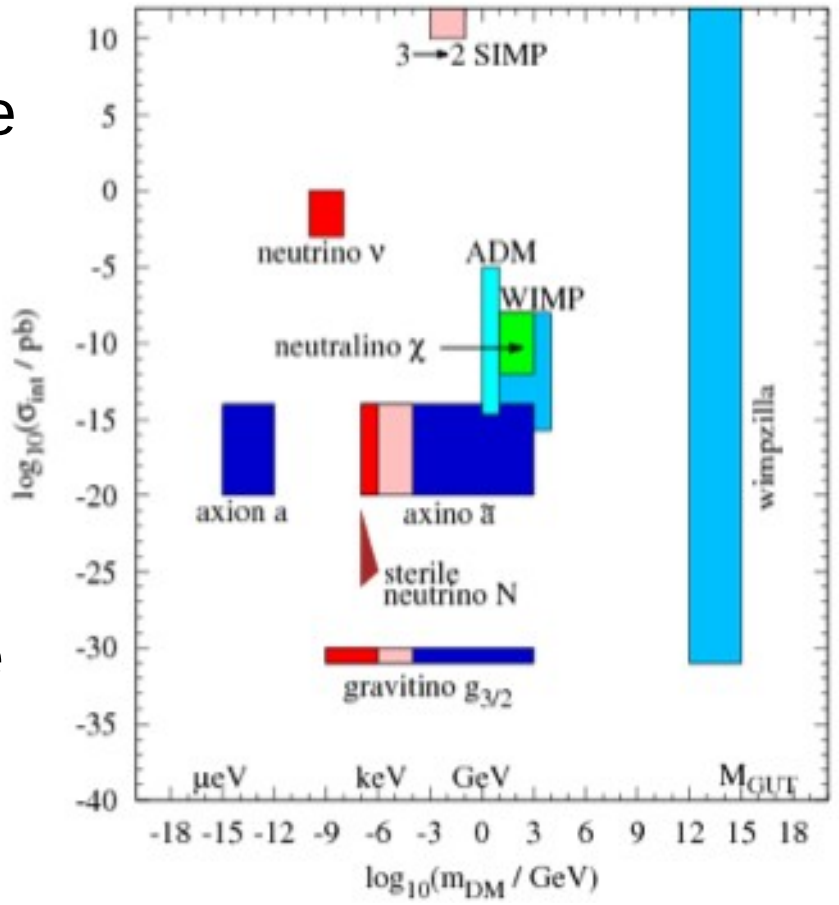
Should I real
slide?!

OK so dark matter exists and we all know it!

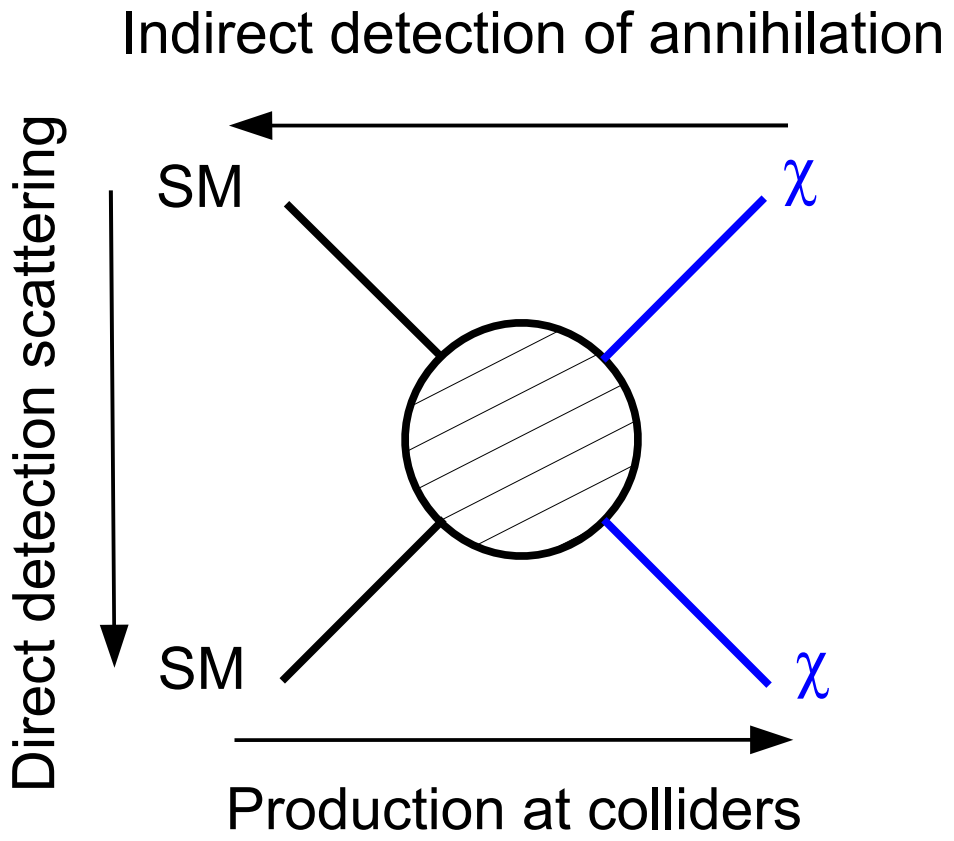
Planck 2013: Dark matter vs luminous matter from CMB anisotropy measurements

WIMP is “our” favourite candidate

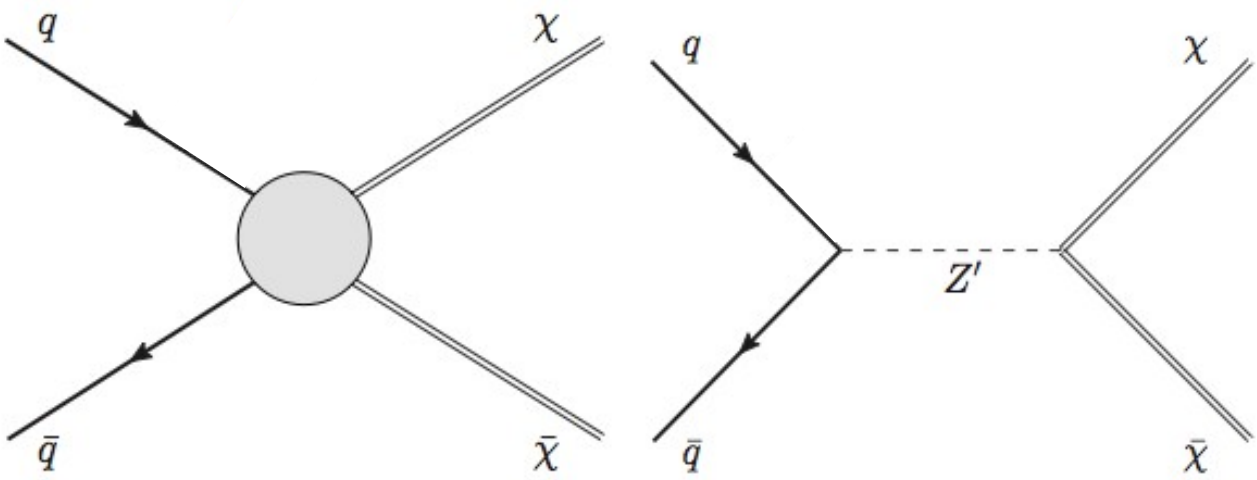
- WIMP is even capable of doing a miracle
  - A cold thermal relic with weak scale mass (few GeV to few TeV) and interactions ($<\sigma v> \approx 3 \times 10^{-26}$ cm^3/s)
  - Naturally accounts for the required relic abundance
  - The WIMP miracle!
- Other alternatives exist (axions, sterile neutrinos, etc) but they are not part of this talk. Also LHC SUSY searches involving missing transverse energy from LSP are not covered here.

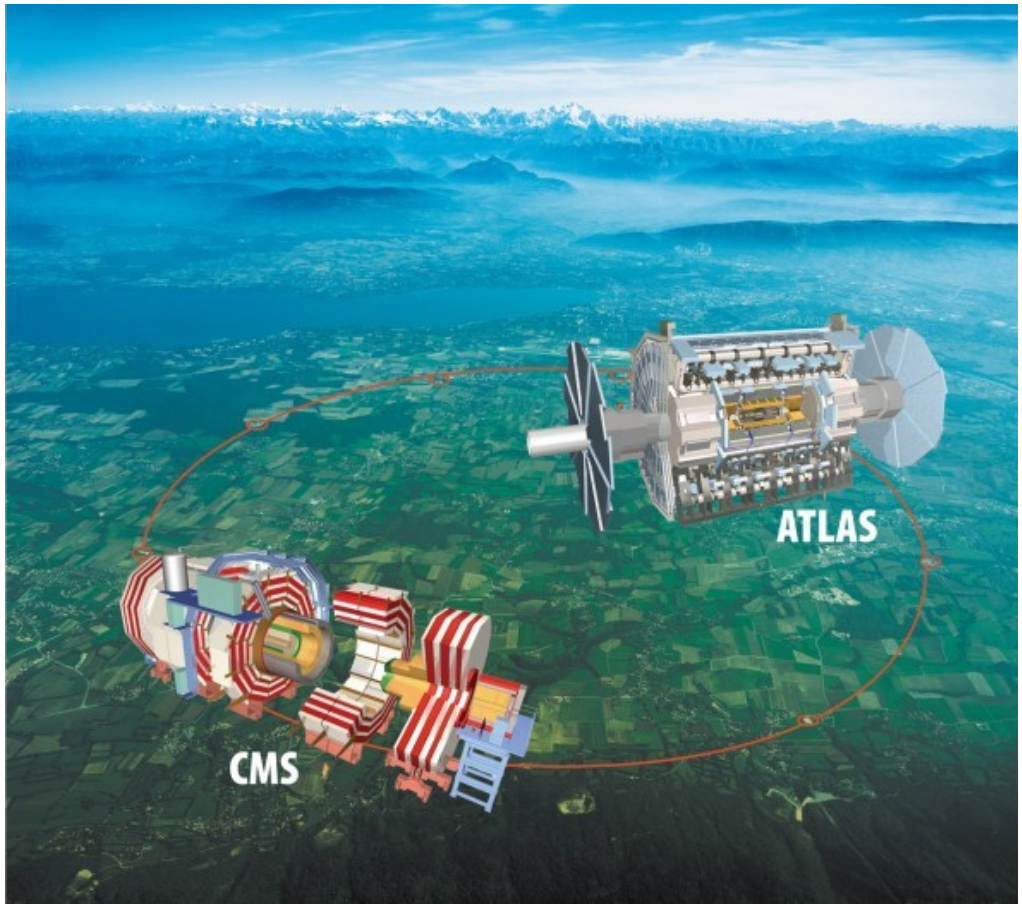


- Some kind of interaction is expected between WIMPs and SM particles
- Current experiments probe different scales:
  - Direct detection O(10 MeV)
  - Indirect detection O(100 GeV)
  - Colliders O(1TeV)
- They are complementary under some conditions



- EFT operators are convenient for representing DM - SM interactions with very few parameters: scale of the interaction M_* (or Λ) and DM mass, e.g. $L_{eff} = \frac{1}{M_\star^2}\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$
- These operators could arise from integrating out the heavy-mediator with mass $M_{\rm med} = \sqrt{g_{\rm q} g_\chi} M_\star$
- EFT approach is valid for a momentum transfer $Q_{\rm tr} < M_{\rm med}$
- Simplified models constitute a good alternative but involve additional parameters (mediator mass, couplings, and decay width)





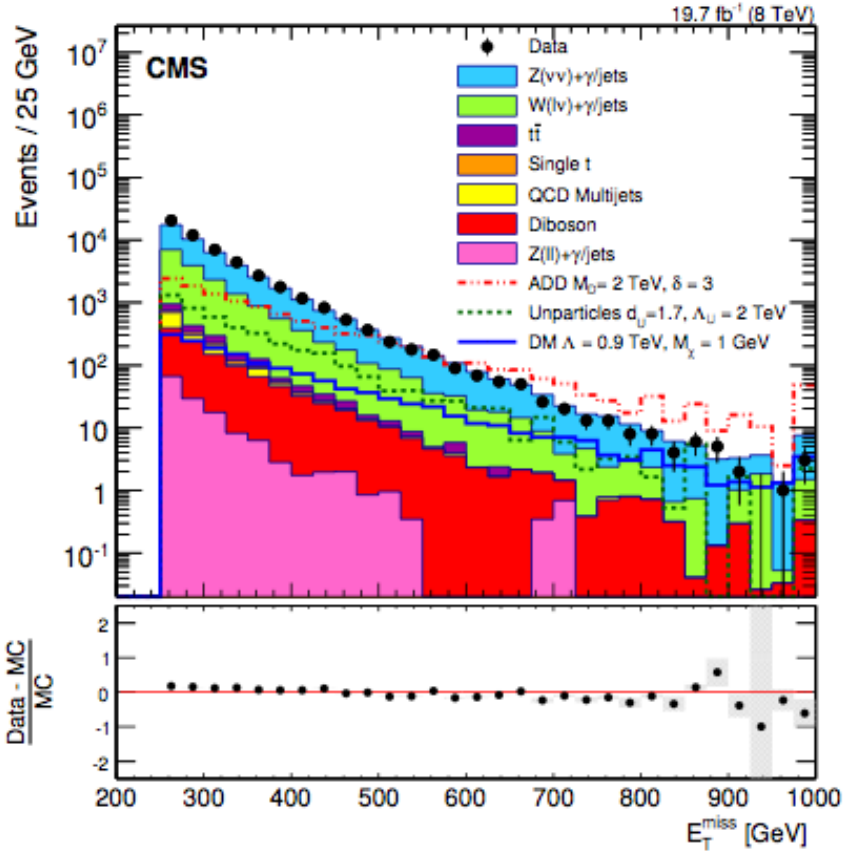
ATLAS and CMS are general-purpose experiments with nearly full solid angle coverage and excellent performance for tracking, calorimetry and muon spectrometer

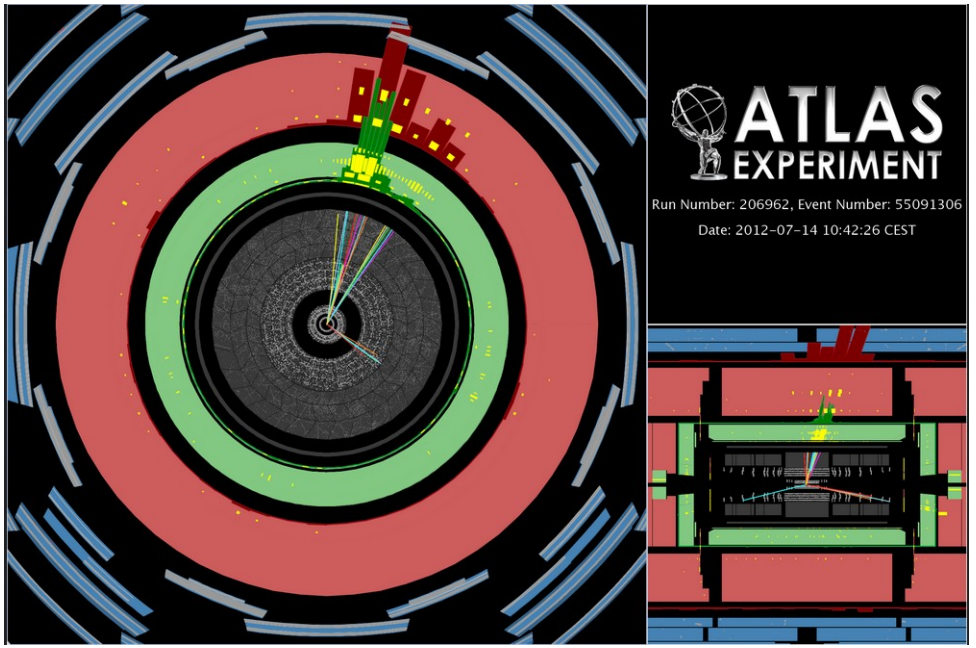
- Dark Matter searches at ATLAS and CMS could be categorized in three main lines (not fully independent though):
  - Mono-X searches (X=boson)
  - DM with heavy flavour
  - Higgs to invisible searches in Higgs-DM portal models
- Results are interpreted using EFT (with or without truncation $Q_{\mathrm{tr}} < M_{\mathrm{med}}$) and using simplified models

Monojet searches

ATLAS: arXiv:1502.01518

CMS: arXiv:1408.3583





Typical analysis strategy:

- Select events with high pT jet(s) and large transverse momentum
- Veto electrons/muons/(taus and tracks) to reject leptonic W/Z decays.
- Cut and count experiment at different pT and MET thresholds.

Monojet searches

ATLAS: arXiv:1502.01518

CMS: arXiv:1408.3583

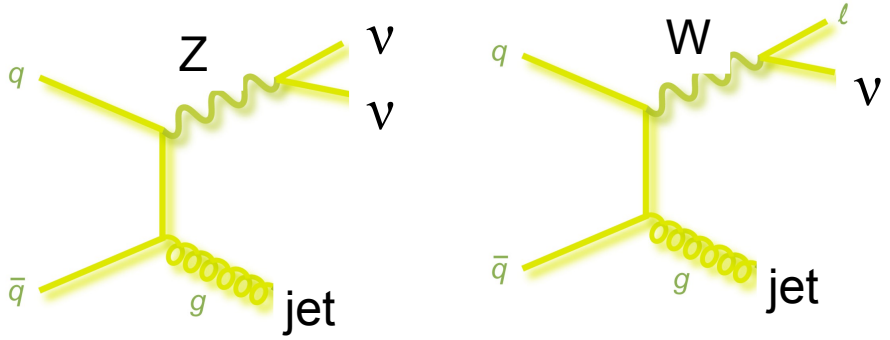
- Main backgrounds:

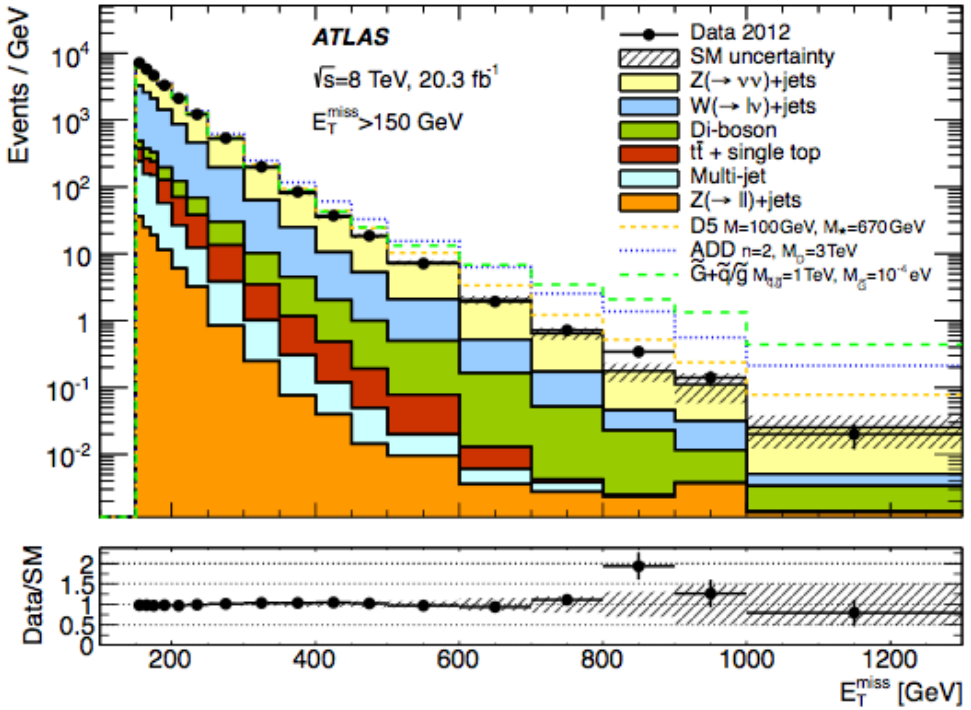
  Z(vv)+jets, W(lv)+jets with lost lepton
- Estimation of the EW background is based on leptonic W/Z control regions
- Transfer factor method:

MC TF

$$N_{SR}^{predicted} = (N_{CR}^{Data} - N_{Bkg}) \cdot C \cdot \frac{N_{SR}^{MC}}{N_{jet/E_T^{miss}}^{MC}}$$

Correction factor



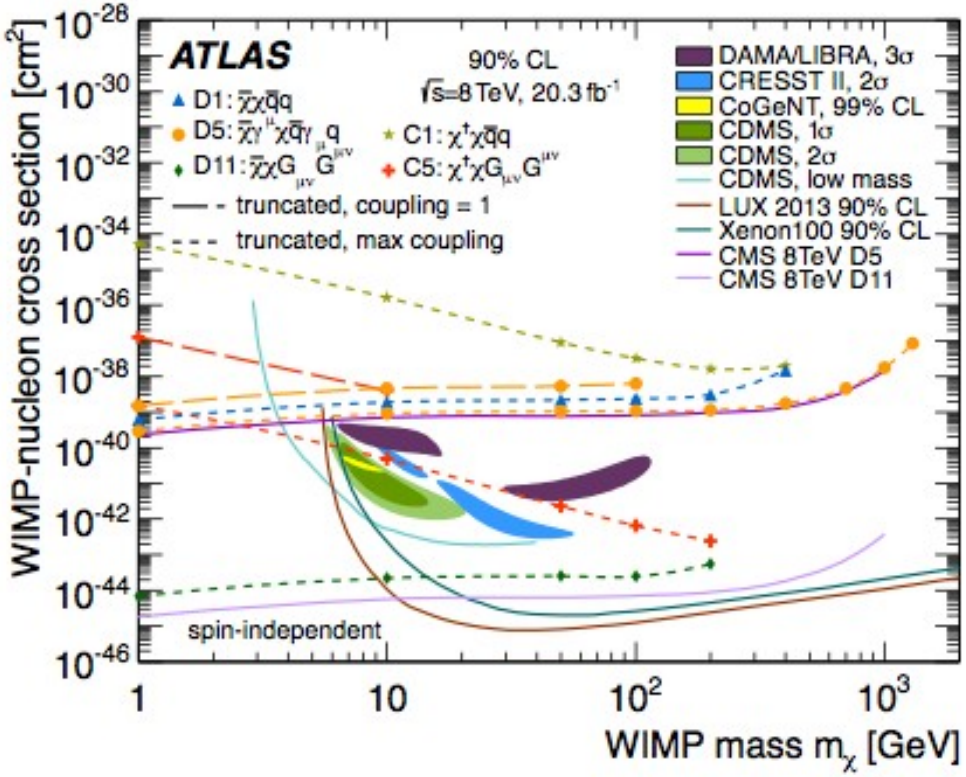


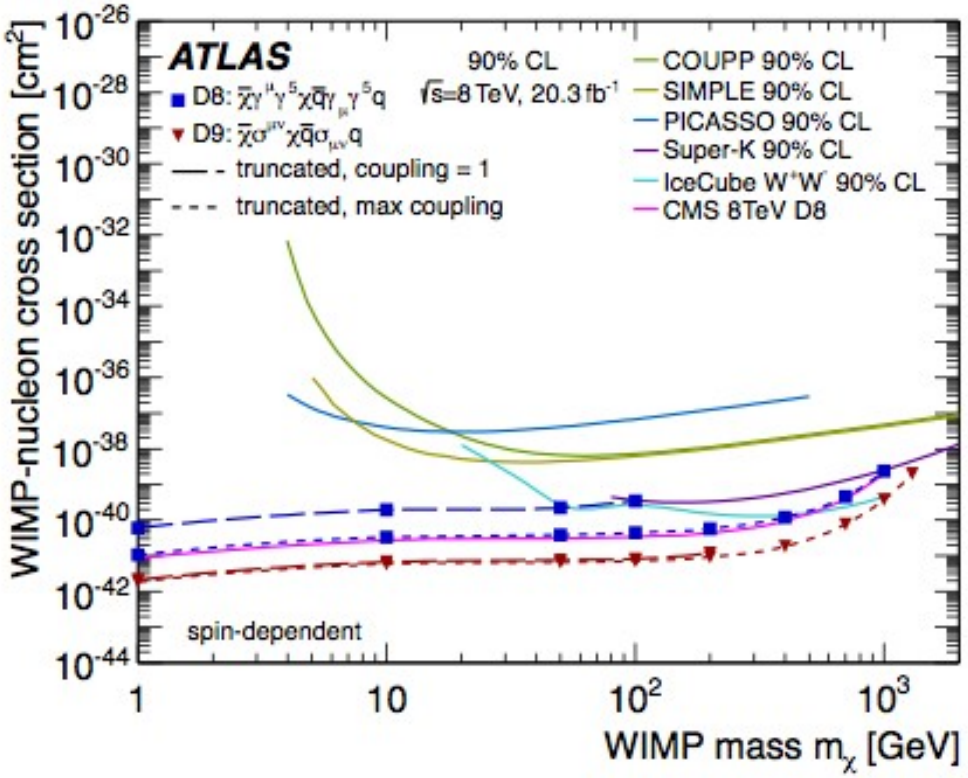
- Theoretical and experimental uncertainties largely cancel out in the TF
- Statistical uncertainties from CR usually take over the systematics (particularly at high MET)

Monojet EFT results

ATLAS: arXiv:1502.01518

CMS: arXiv:1408.3583



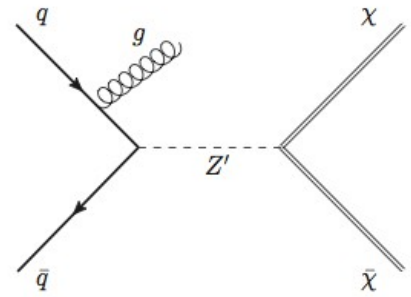


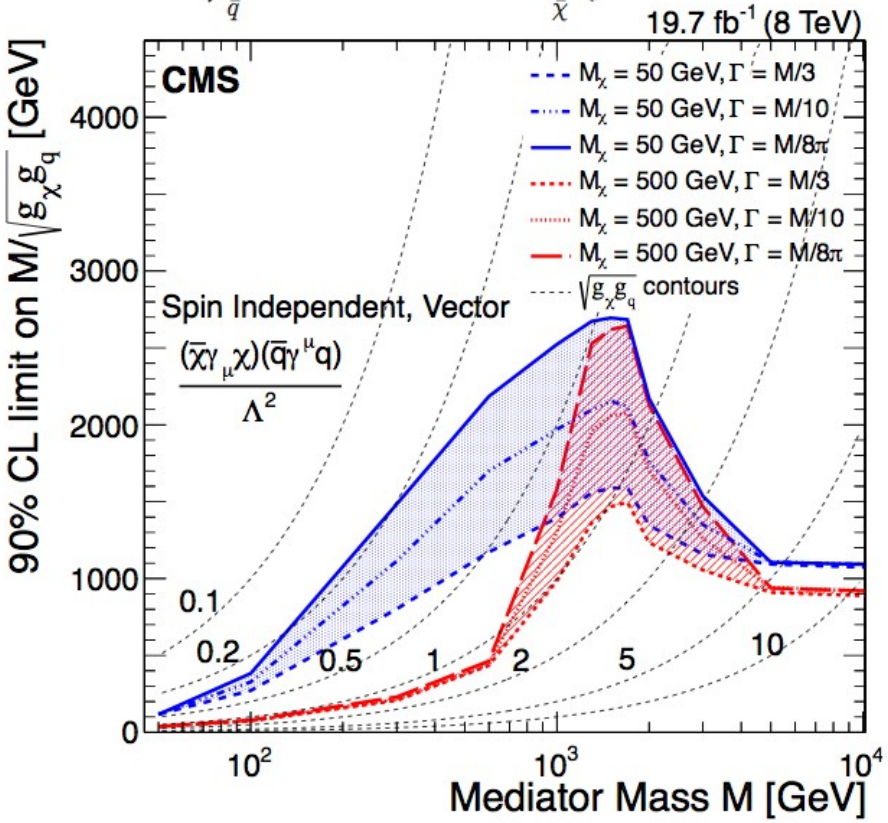
- Results are interpreted in terms of exclusion limits on WIMP-nucleon cross section
- ATLAS limits are done with the truncation $Q_{\mathrm{tr}} < M_{\mathrm{med}}$. Need agreement on common procedure.

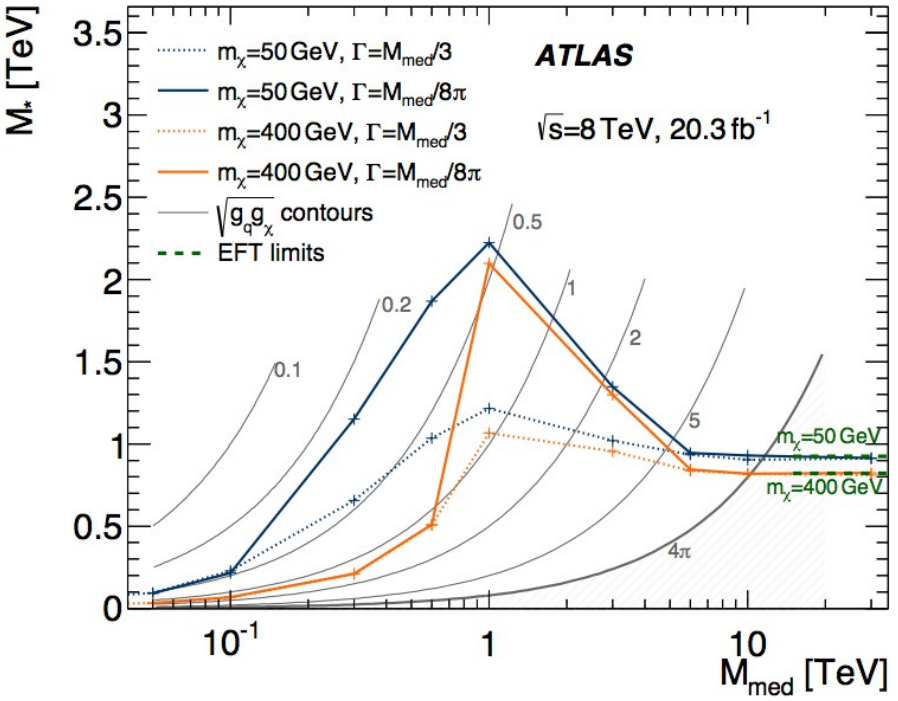
Monojet simplified model results

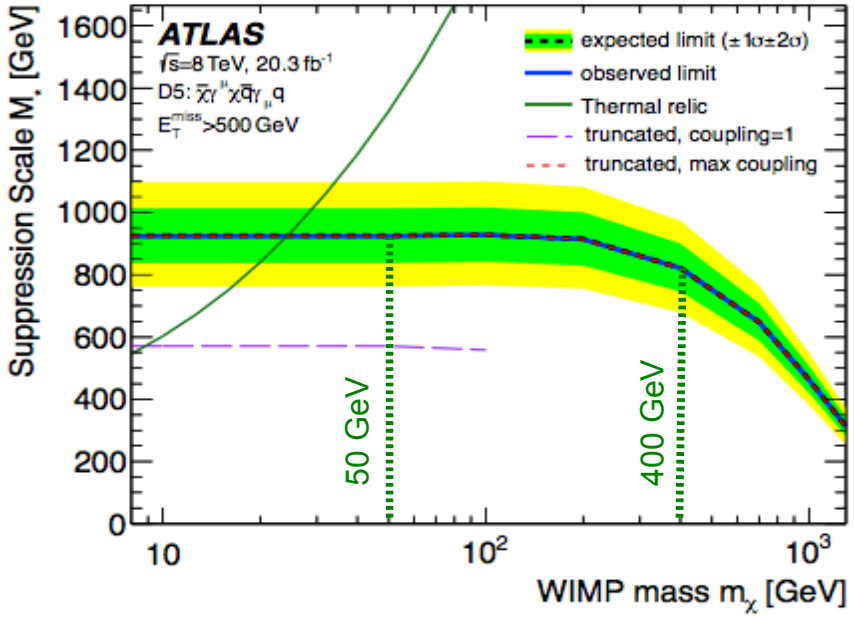
ATLAS: arXiv:1502.01518

CMS: arXiv:1408.3583





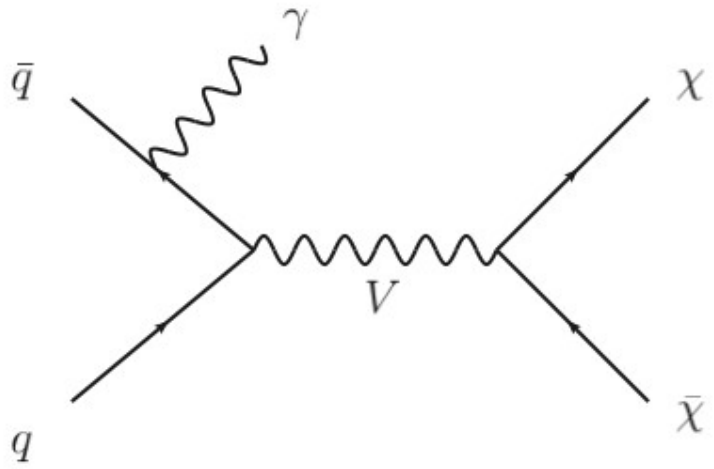


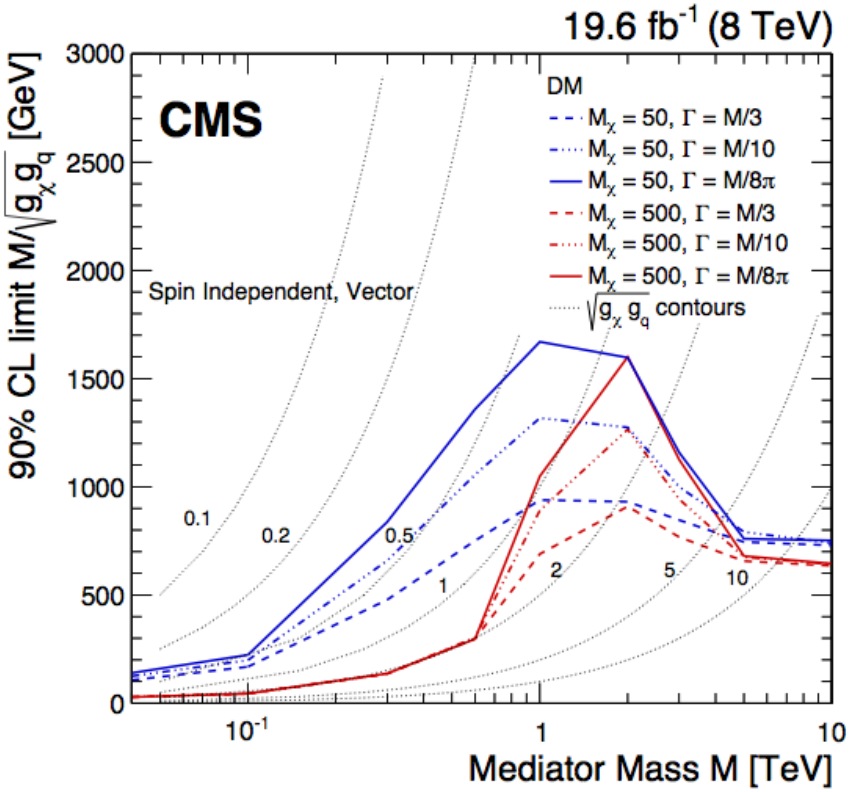


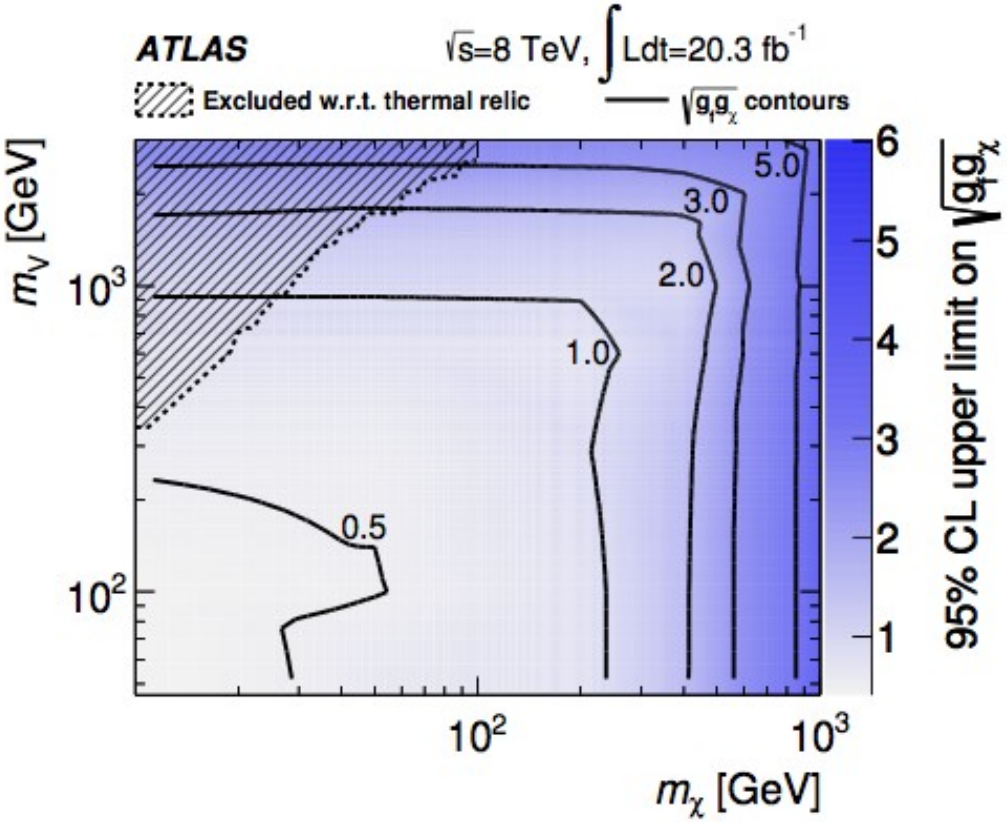
- Limits on M_* sensitive to mediator mass
- For $M_{med} > 5$ TeV, EFT limit of M_*~1 TeV is retrieved

Monophoton simplified model results

ATLAS: PRD 91, 012008 (2015)
CMS: arXiv:1410.8812

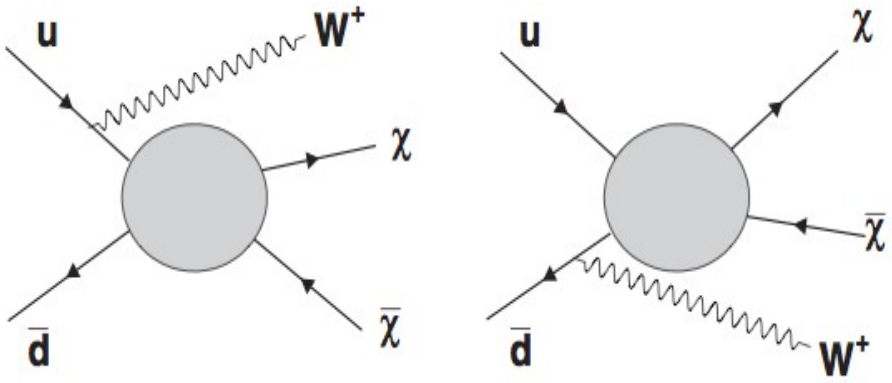




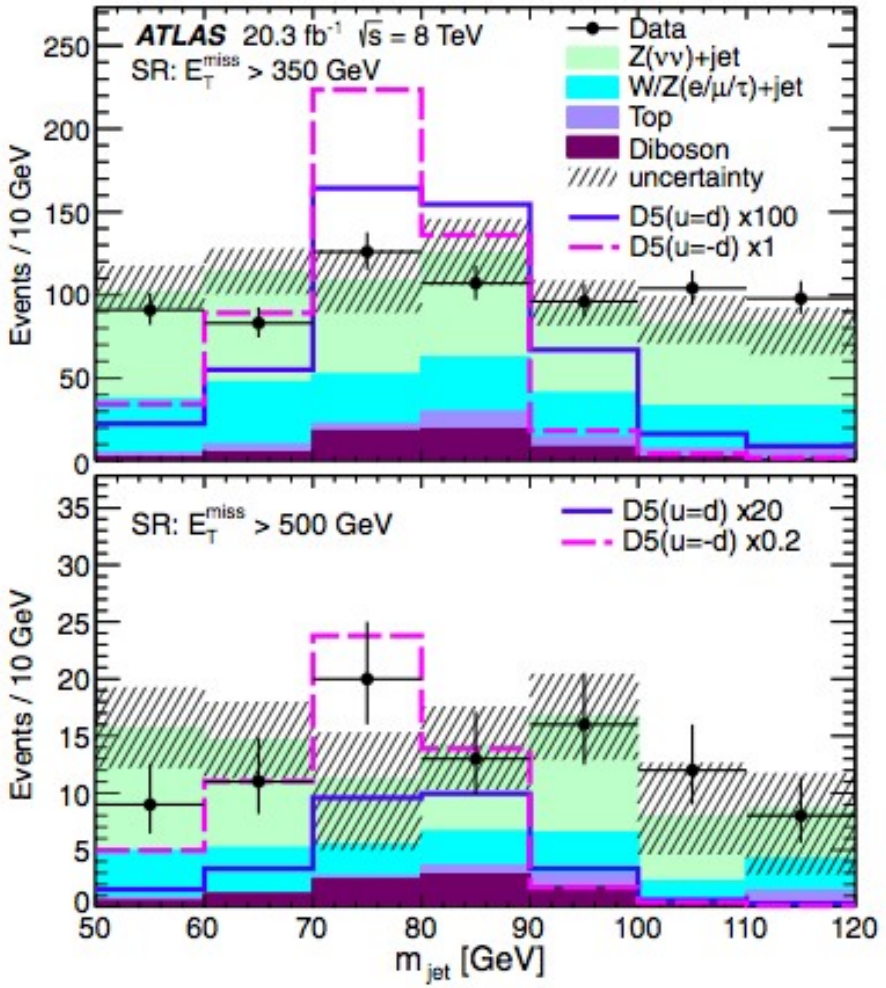


- Signal from photon emission in the initial state and large MET
- Background estimated from CR
- EFT results also available

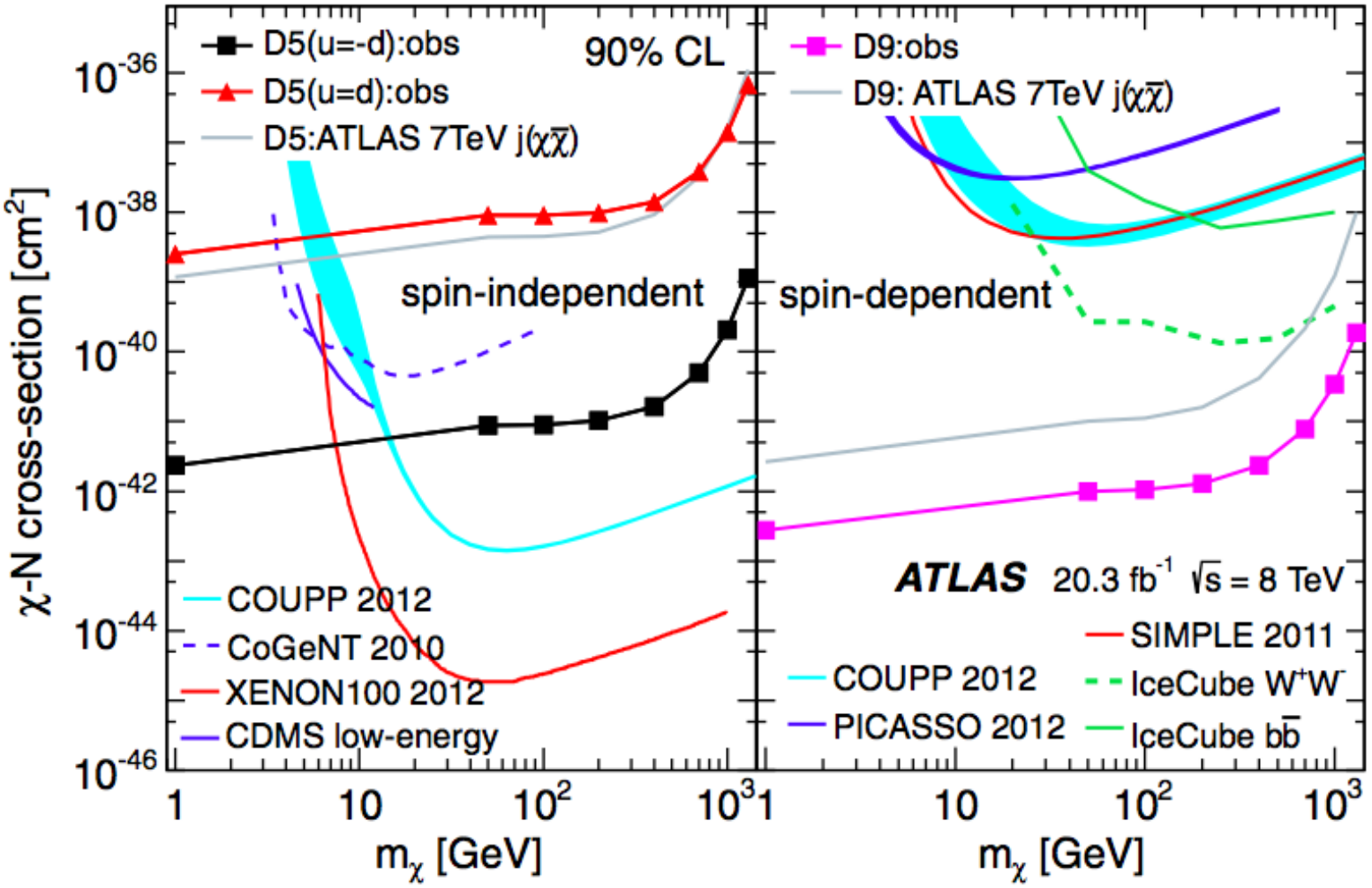
MET+W/Z → hadrons



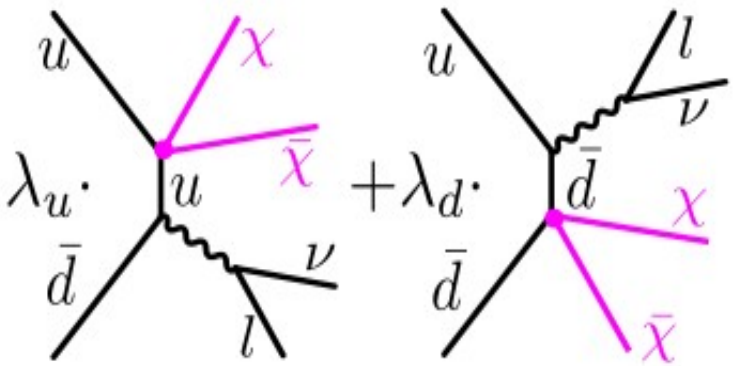
- High boson pT with high missing transverse energy
- Boosted regime: use single fat jet (ΔR=1.2) for W/Z → jj reconstruction
- Constructive and destructive interference in W diagrams play important role

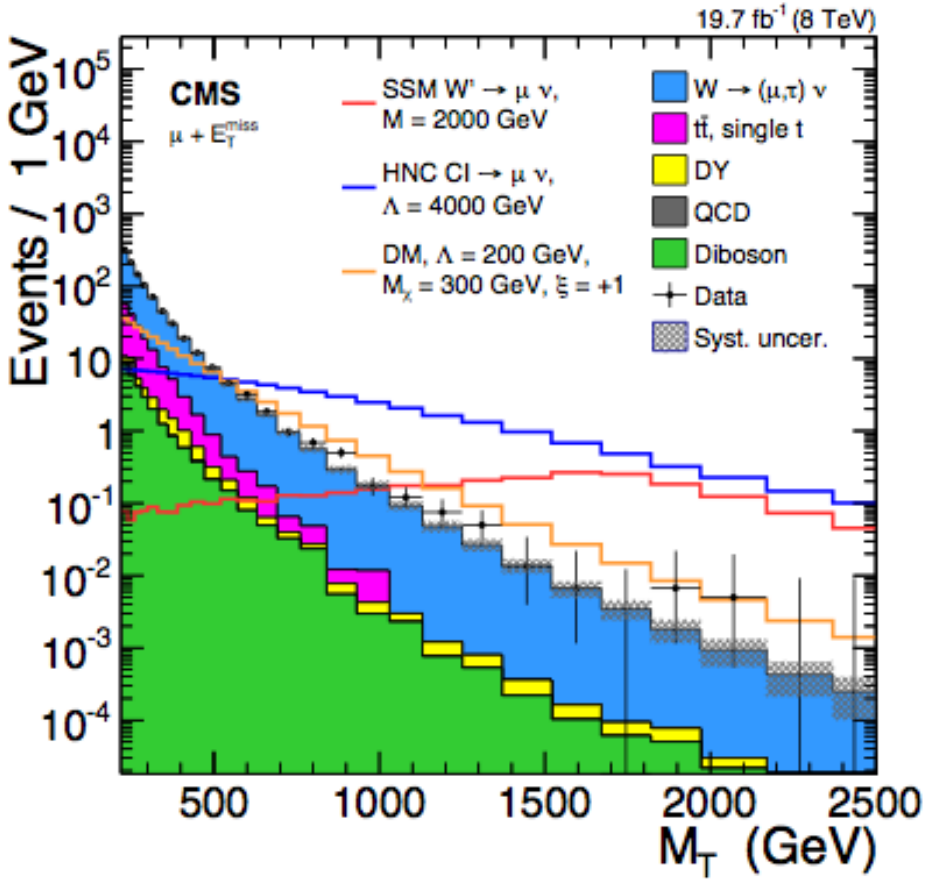


MET+W/Z → hadrons limits

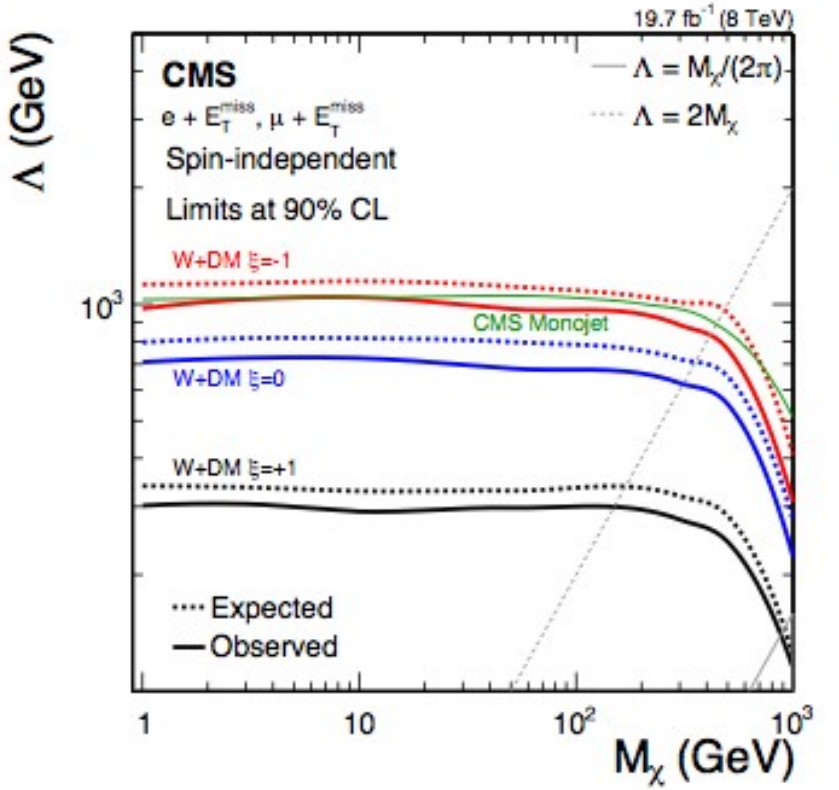


- Strong limits in the case of constructive interference.

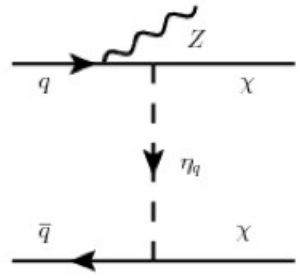


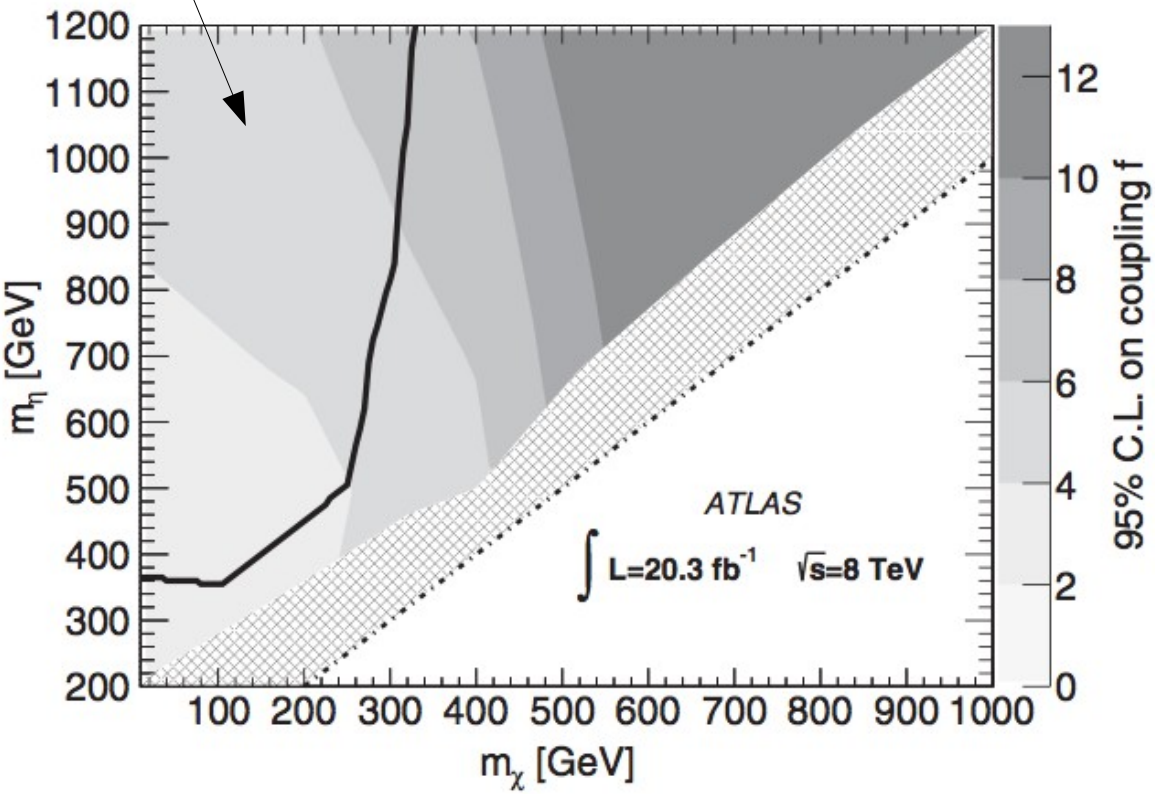


- Dark Matter production with W/Z-boson decaying to leptons
- Bosons recoil against the pair of dark matter → look for excess of events at high MET
- Use transverse mass M_T for W and invariant mass M_{ll} for Z as discriminant



Excluded by relic density measurement

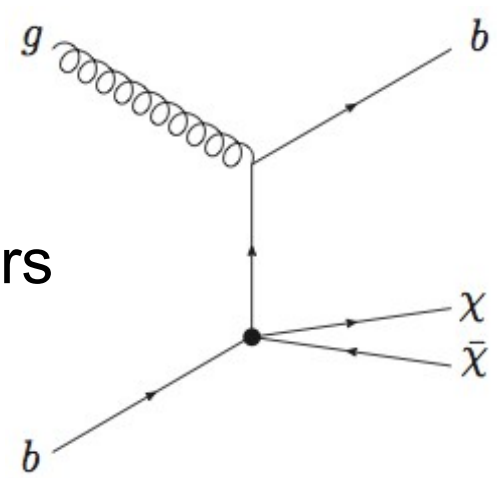


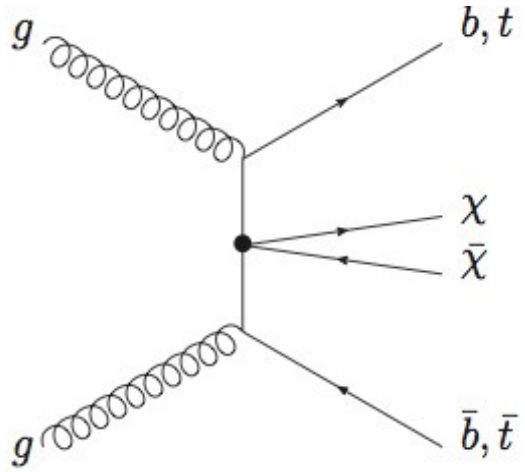


- Limits on the suppression scale as a function of dark matter mass
- Also simplified model limits in the scalar-mediator theory

ATLAS: EPJ 75:92

CMS: CMS-PAS-B2G-13-004, arxiv:1504.03198

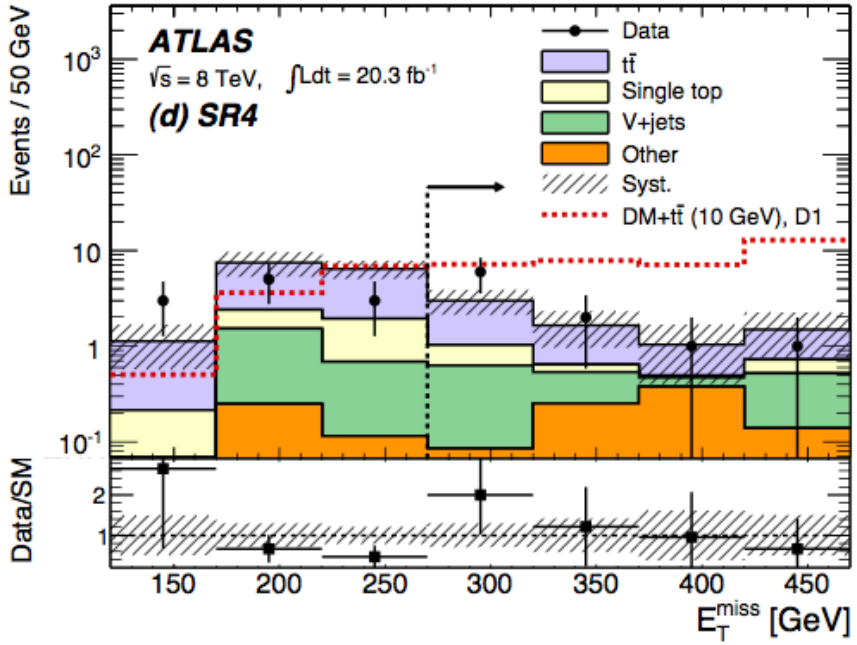




- Some EFT operators enhanced for heavy flavours

$$\mathcal{O}_{\text{scalar}} = \sum_q \frac{m_q}{M_*^N} \bar{q} q \bar{\chi} \chi,$$

- May arise from from integrating out a scalar mediator with couplings proportional to quark mass m_q.



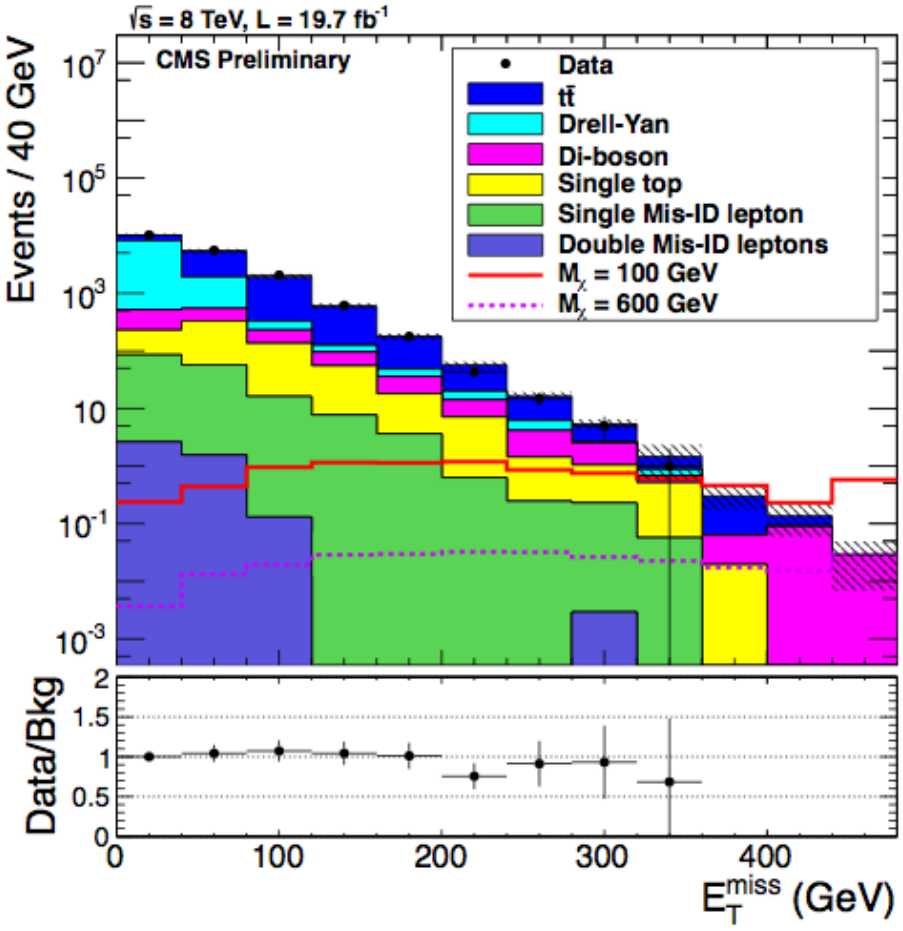
- Look for events with high MET and b-tagged jet(s)

- Signal regions normalized from control regions

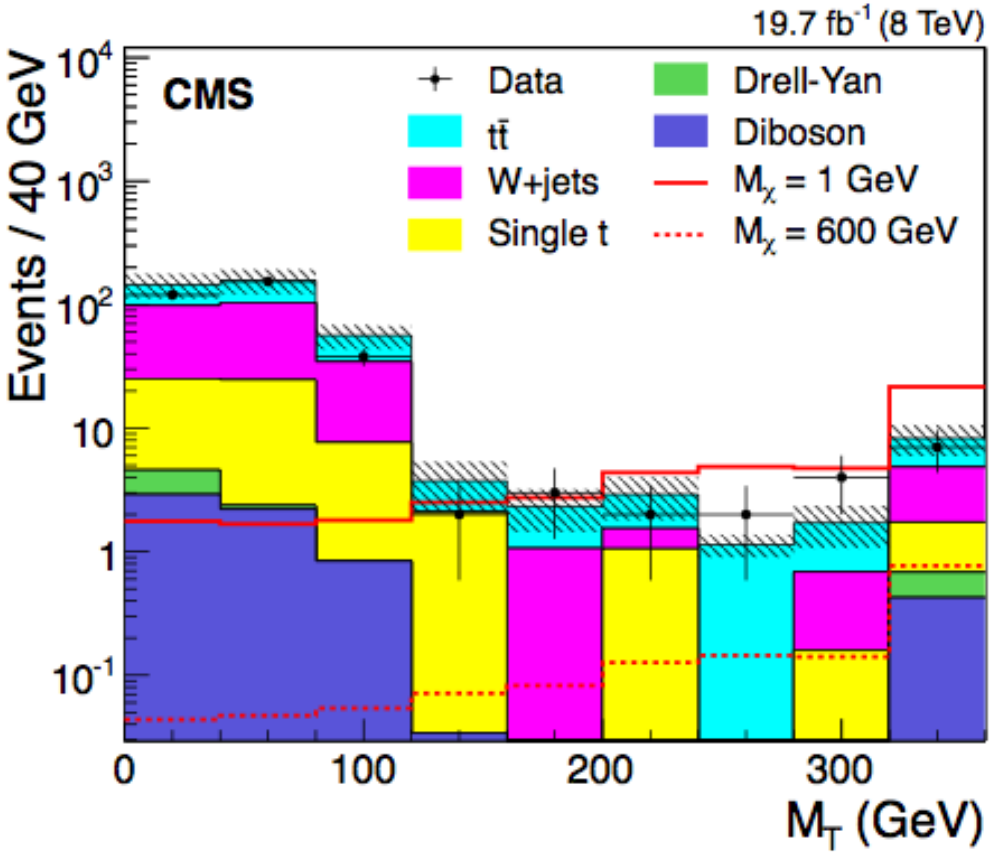
ATLAS: EPJ 75:92

CMS: CMS-PAS-B2G-13-004, arxiv:1504.03198

Full leptonic ttbar

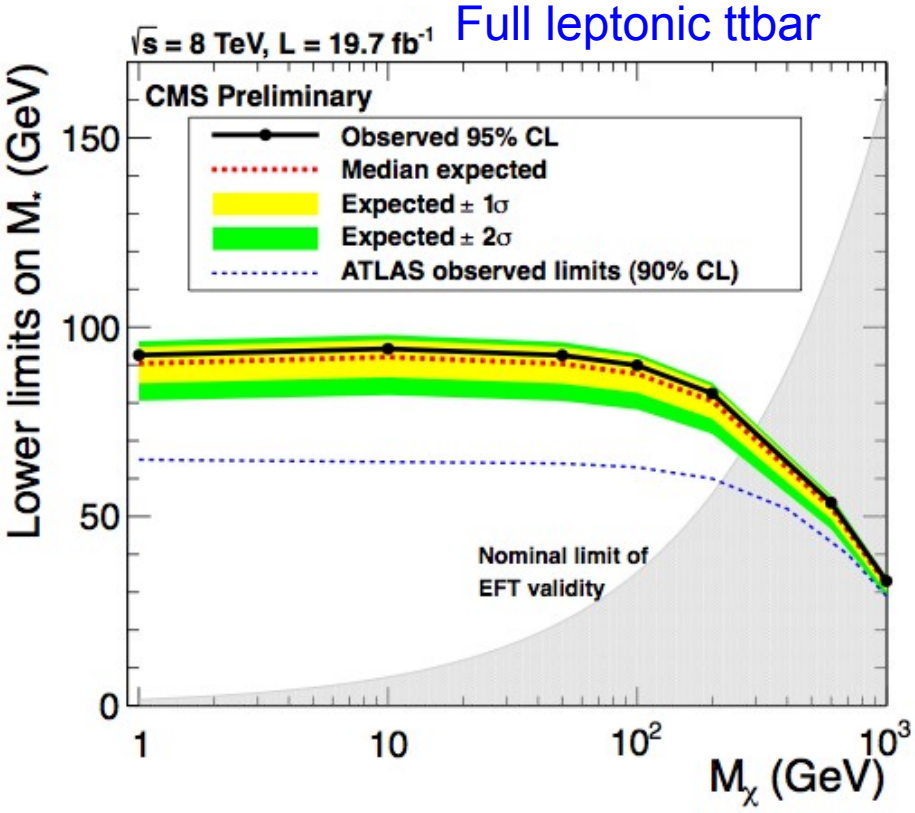


Semi-leptonic ttbar

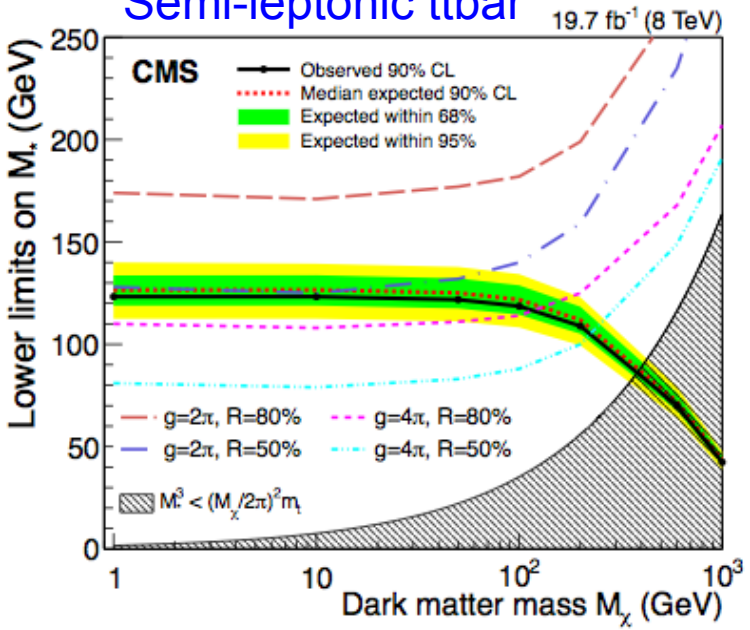


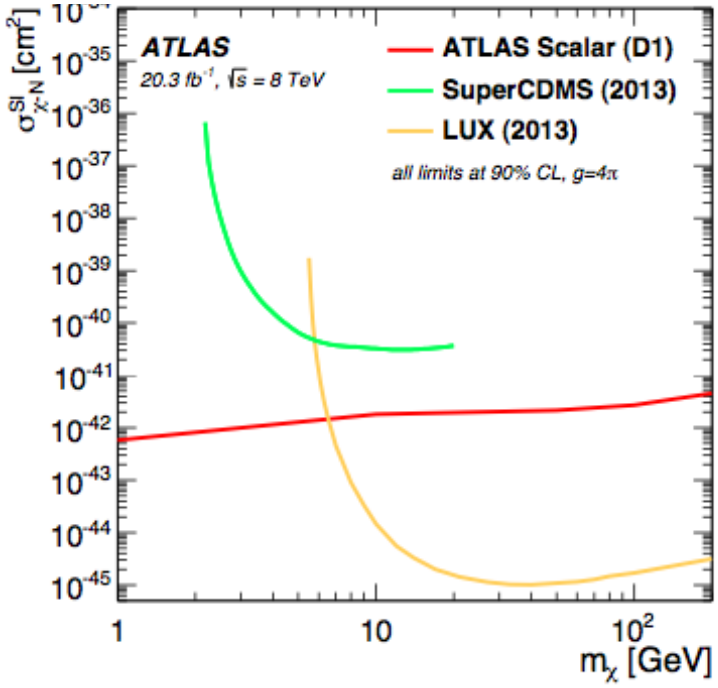
- Different analysis approaches for the different ttbar decay channels (all-hadronic, semi-leptonic and di-leptonic)

Full leptonic ttbar



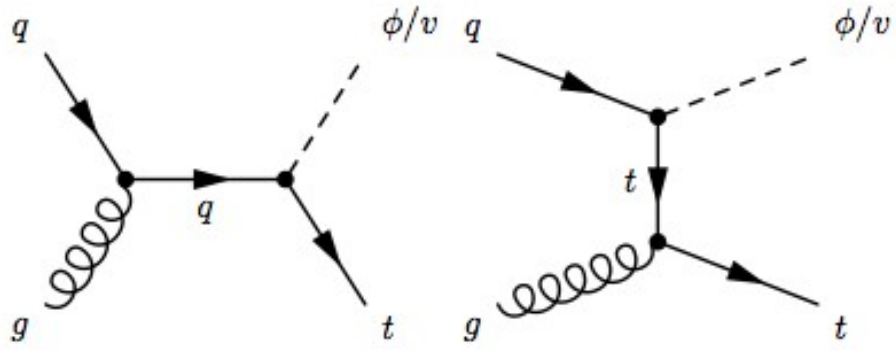
Semi-leptonic ttbar

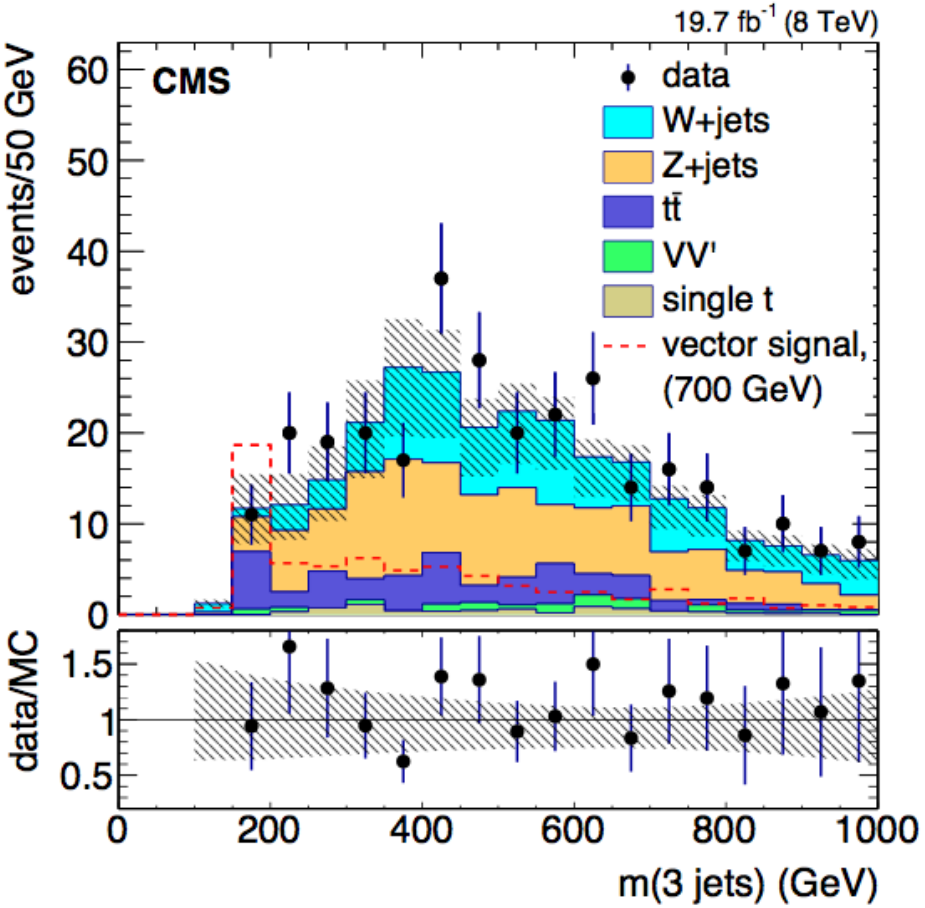


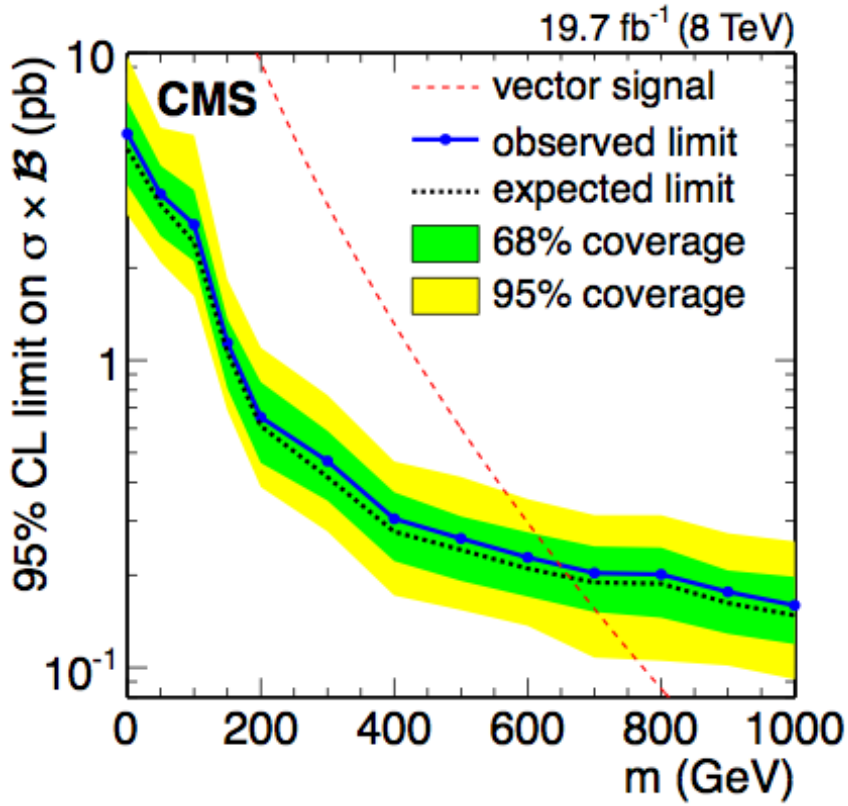


- Better limits from the semi-leptonic ttbar.
- Collider limits are competitive for DM mass below 10 GeV

- Single top+MET predicted by many BSM
- At least 3-jets with one b-tag

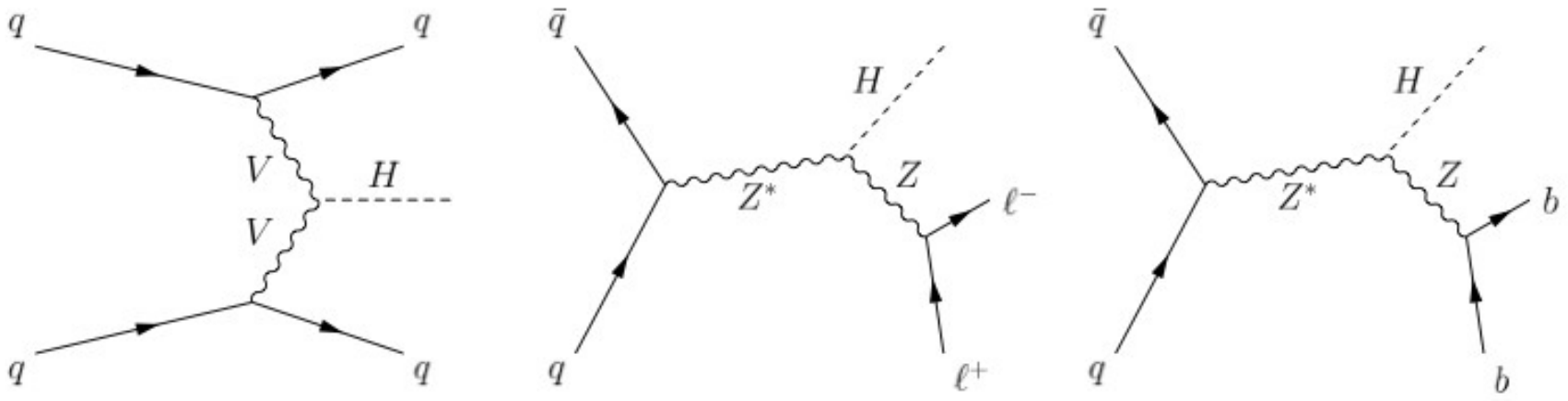






Invisible Higgs decay

ATLAS: ATLAS-CONF-2015-004, arXiv:1504.04324, PRL 112, 201802
CMS: EPJC 74 (2014) 2980, CMS-PAS-HIG-14-038

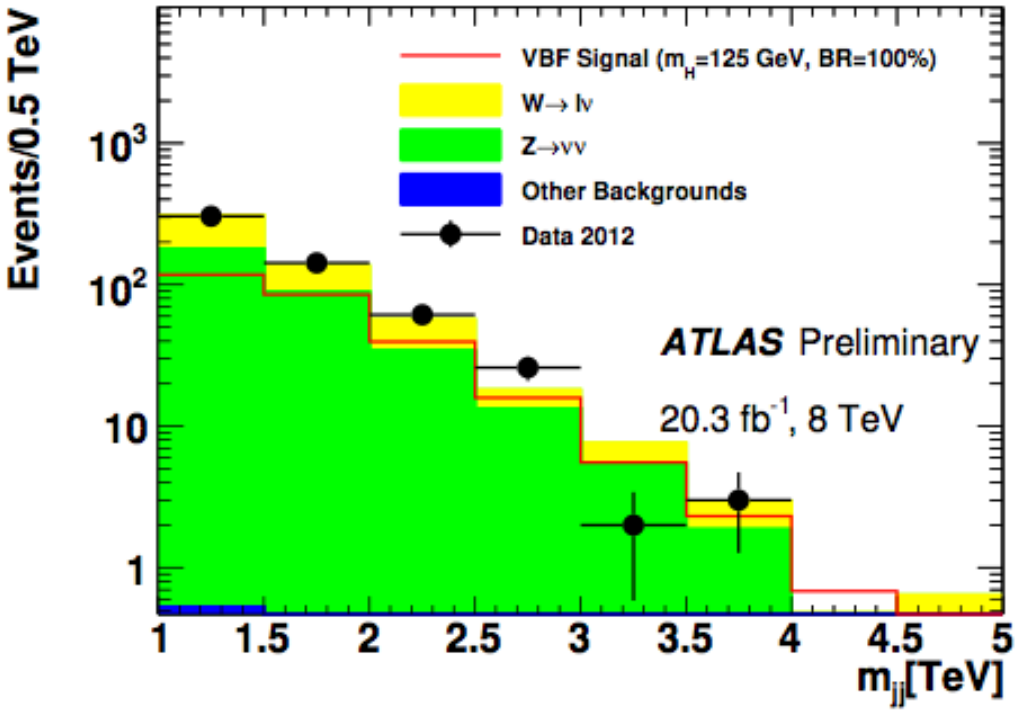


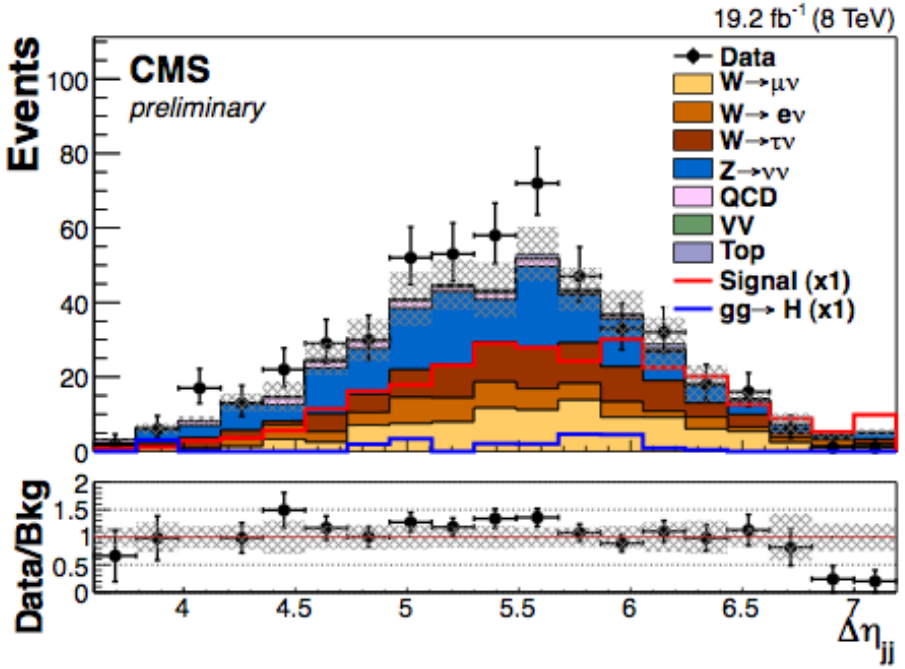
- The discovery of the Higgs at a mass of 125 GeV opens the door for realistic searches for invisible Higgs decays
- The SM H → ZZ → νννν has a BR of 1.2×10^{-3}, below the sensitivity of current analyses
- However DM could have substantially large Yukawa couplings to Higgs leading to BR(H → invisible) much larger than predicted in SM (if DM mass < $m_H/2$)
- Vector boson fusion (VBF) and associated vector boson production (VH) with different topologies have been explored

VBF H → invisible

ATLAS: ATLAS-CONF-2015-004

CMS: EPJC 74 (2014) 2980, CMS-PAS-HIG-14-038



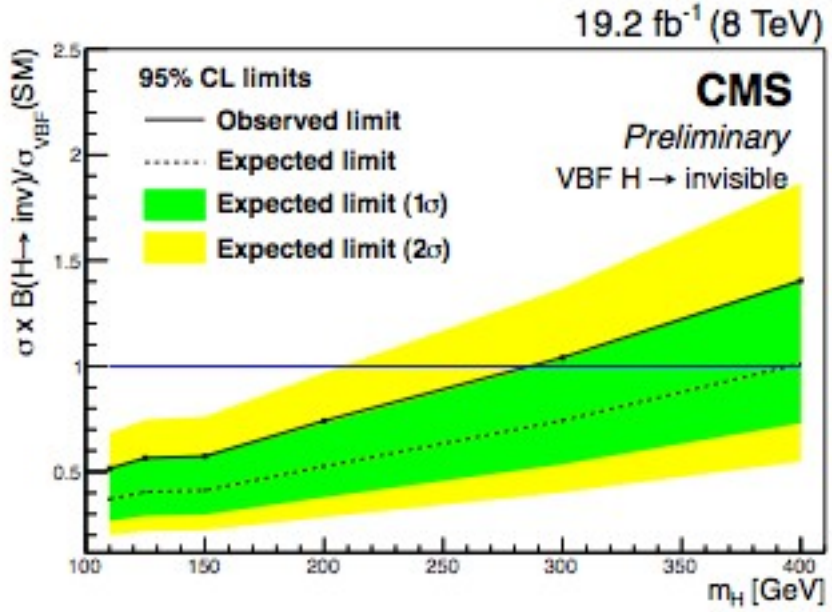


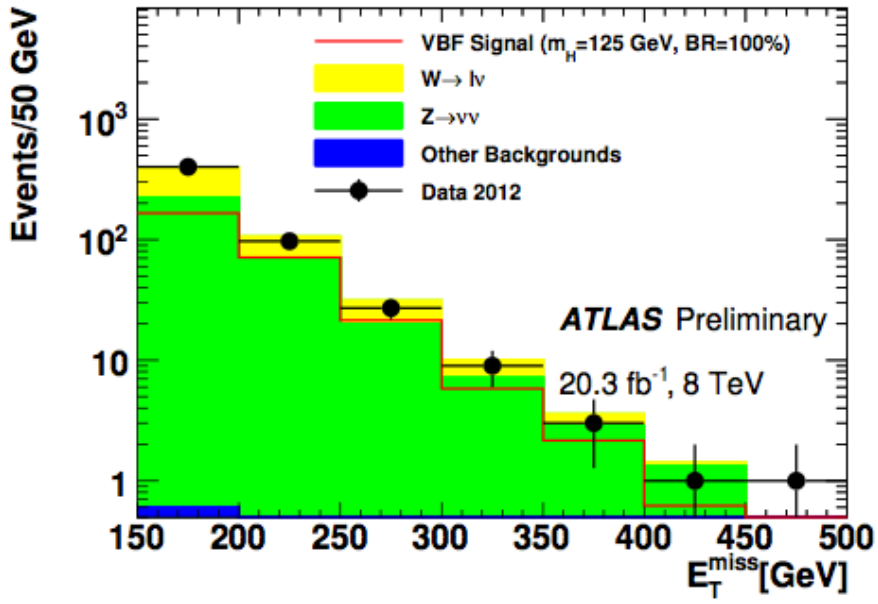
- Look for two forward jets with large invariant mass m_{jj} (above 1 TeV) and large η separation (typically $\Delta\eta_{jj}$ more than 3.5)
- Main background from V+jets and VBF EW production
- It is one of the most sensitive channels to look for invisible Higgs decays

VBF H → invisible

ATLAS: ATLAS-CONF-2015-004

CMS: EPJC 74 (2014) 2980, CMS-PAS-HIG-14-038





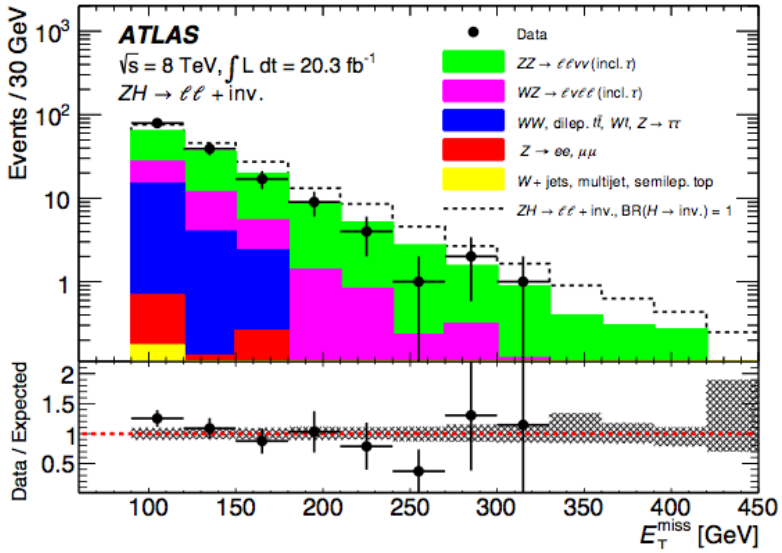
- Similar expected limits on BR(Higgs → invisible) between CMS (<35%) and ATLAS (< 29%) for m_H=125 GeV
- Observed: ATLAS <35%, CMS <47%

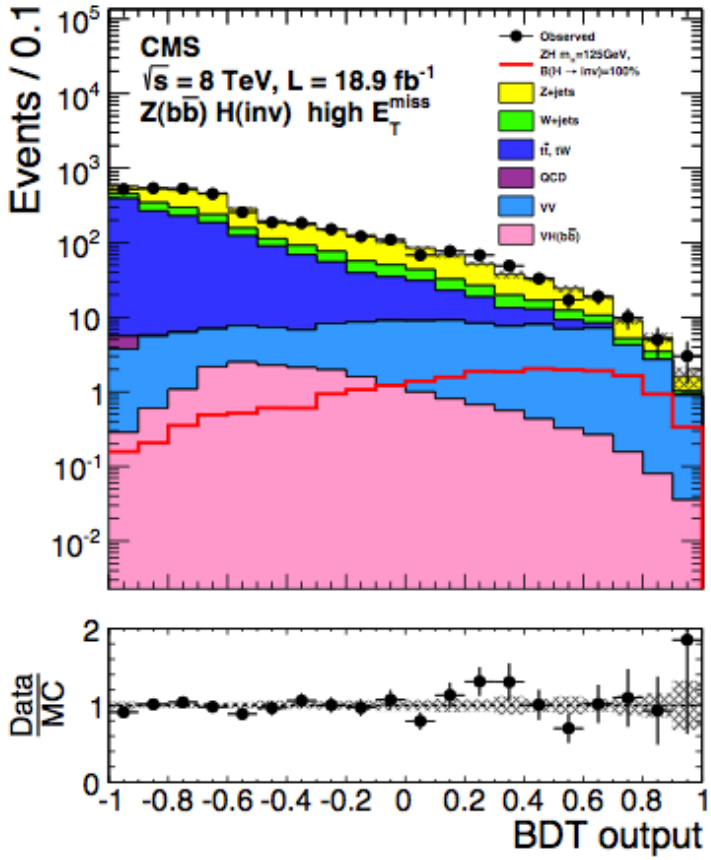
- Scan as a function of the Higgs mass assuming SM couplings.

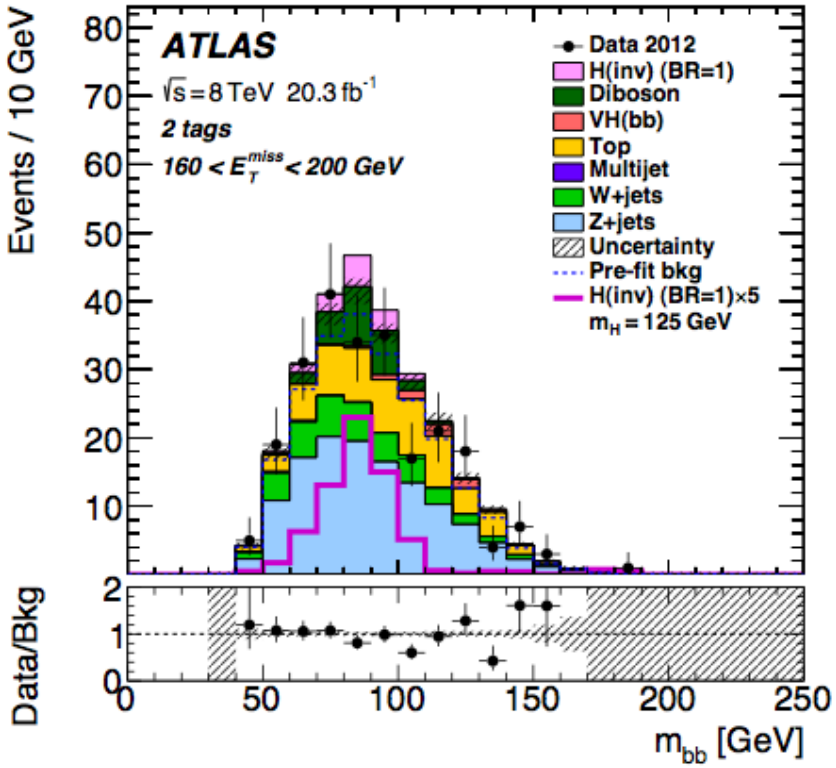
VH (H → invisible)

ATLAS: arXiv:1504.04324, PRL 112, 201802
CMS: EPJC 74 (2014) 2980

- Hadronic and leptonic boson decays + large MET
- Use b-tagging and/or BDT to enhance H(Z → bb) signal



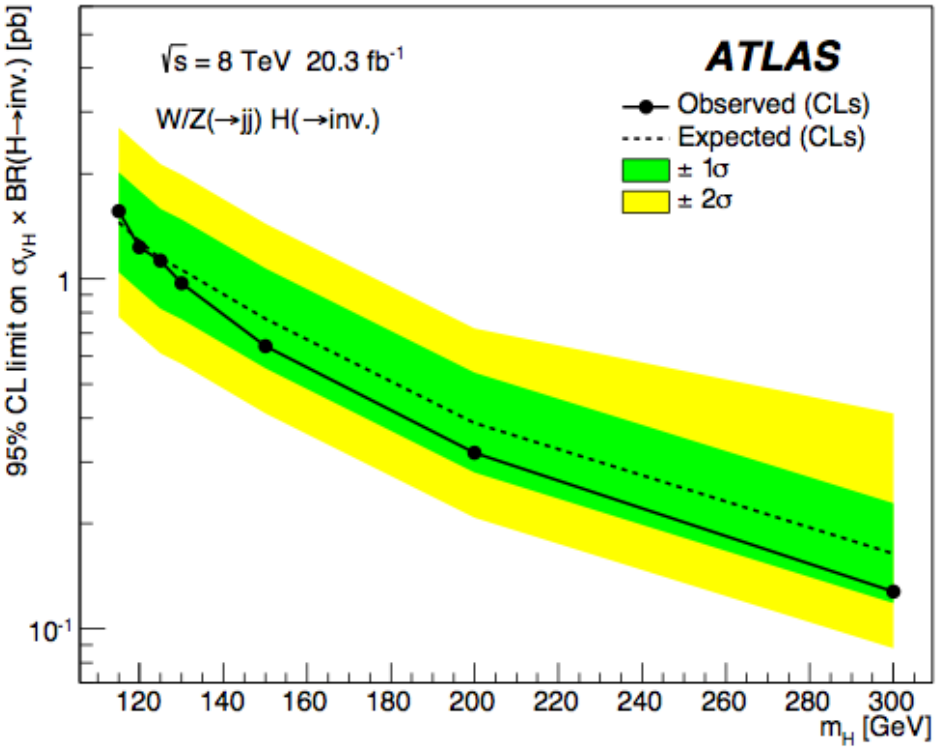


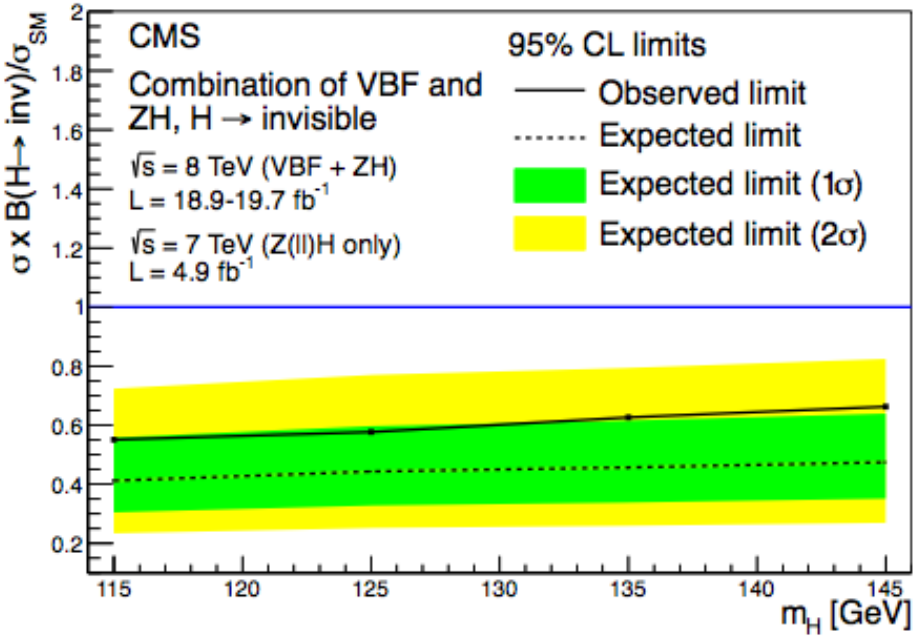


VH (H → invisible)

ATLAS: arXiv:1504.04324, PRL 112, 201802

CMS: EPJC 74 (2014) 2980





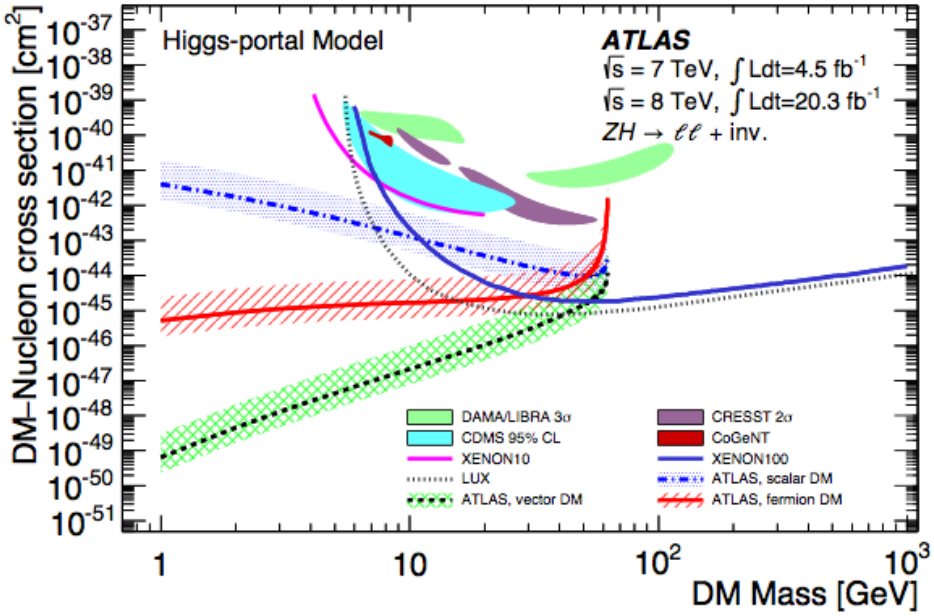
- Limit scan as a function of the Higgs mass assuming SM couplings.
- Limits on BR(Higgs → invisible) less stringent than VBF channel

DM-Higgs portal

ATLAS: PRL 112, 201802

CMS: EPJC 74 (2014) 2980

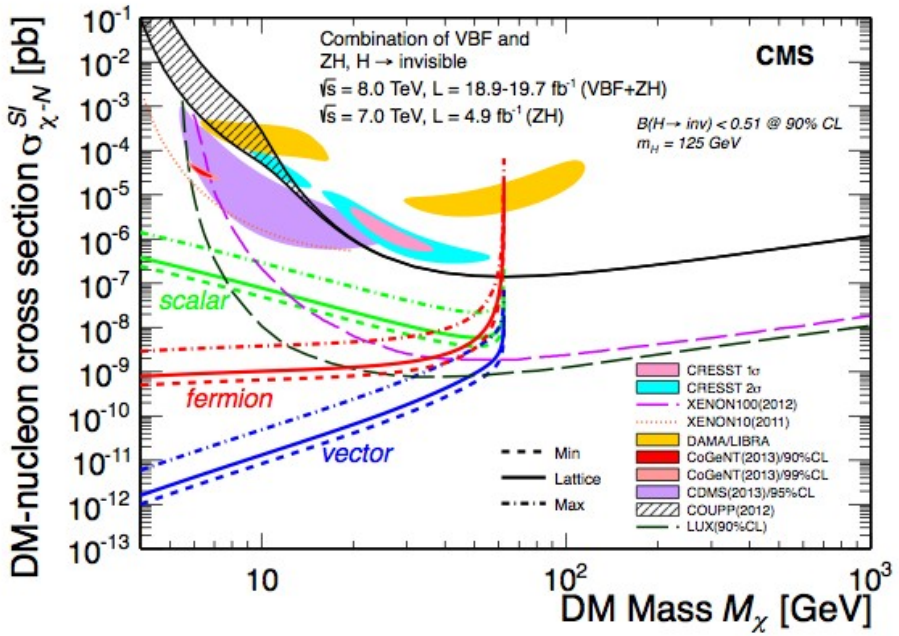
- If the DM particle has a mass below $m_H/2$ the invisible decay width of the Higgs can be directly translated to DM-nucleon spin independent cross section
- For DM masses below $m_H/2$ the collider limits are very competitive



$$\sigma^{SI}_{S-N} = \frac{4\Gamma_{inv}}{m_H^3 v^2 \beta} \frac{m_N^4 f_N^2}{(M_\chi + m_N)^2},$$

$$\sigma^{SI}_{V-N} = \frac{16\Gamma_{inv} M_\chi^4}{m_H^3 v^2 \beta (m_H^4 - 4M_\chi^2 m_H^2 + 12M_\chi^4)} \frac{m_N^4 f_N^2}{(M_\chi + m_N)^2},$$

$$\sigma^{SI}_{f-N} = \frac{8\Gamma_{inv} M_\chi^2}{m_H^5 v^2 \beta^3} \frac{m_N^4 f_N^2}{(M_\chi + m_N)^2}.$$

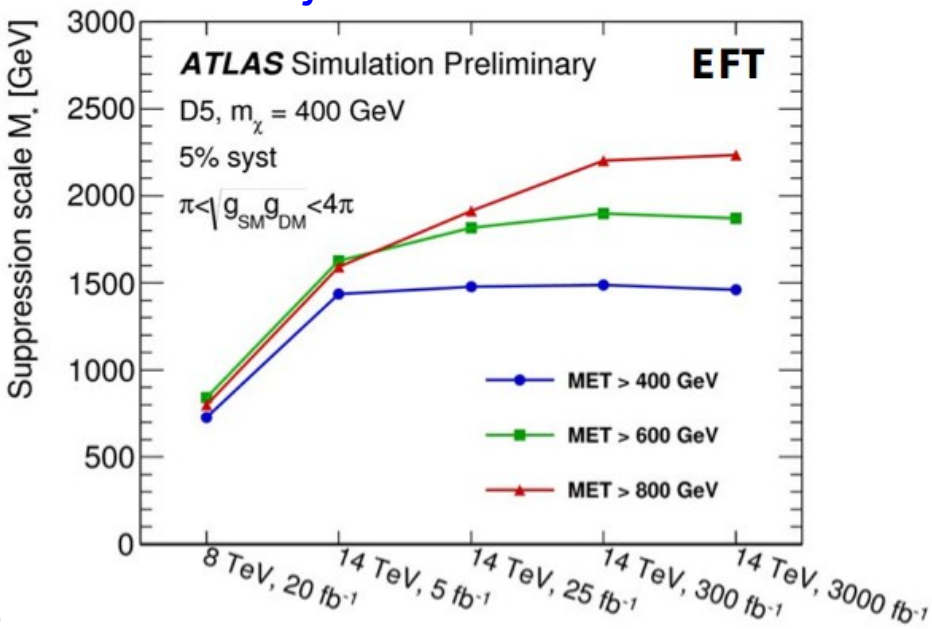


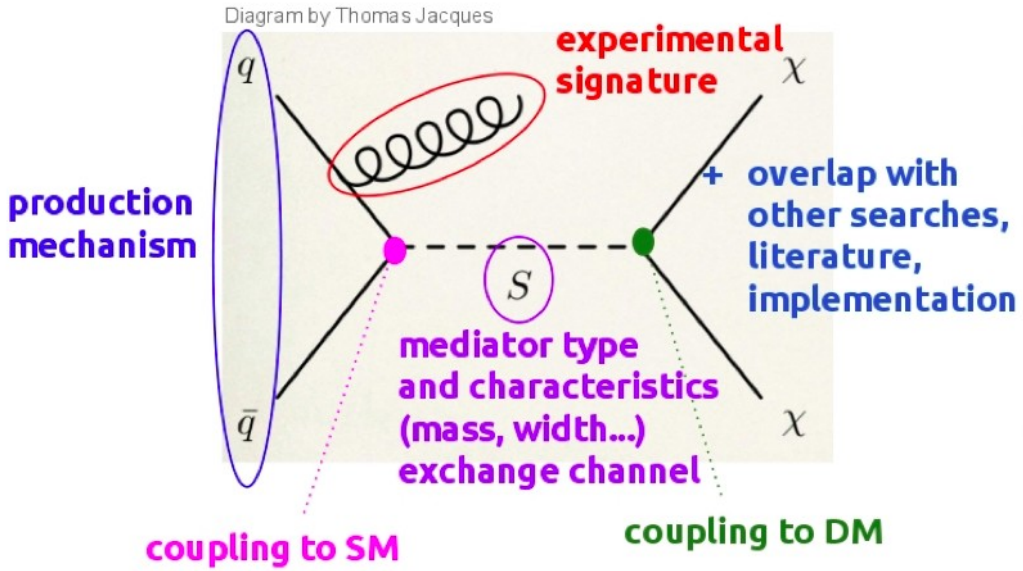
Prospects for Run-II dark matter searches

- Monojet analyses using similar strategies as in Run-I are expected to surpass previous limits with only 5/fb at 13 TeV
- Harmonization of DM searches ongoing within the ATLAS/CMS Dark Matter Forum

https://twiki.cern.ch/twiki/bin/view/LHCDMF/WebHome

Monojet selection



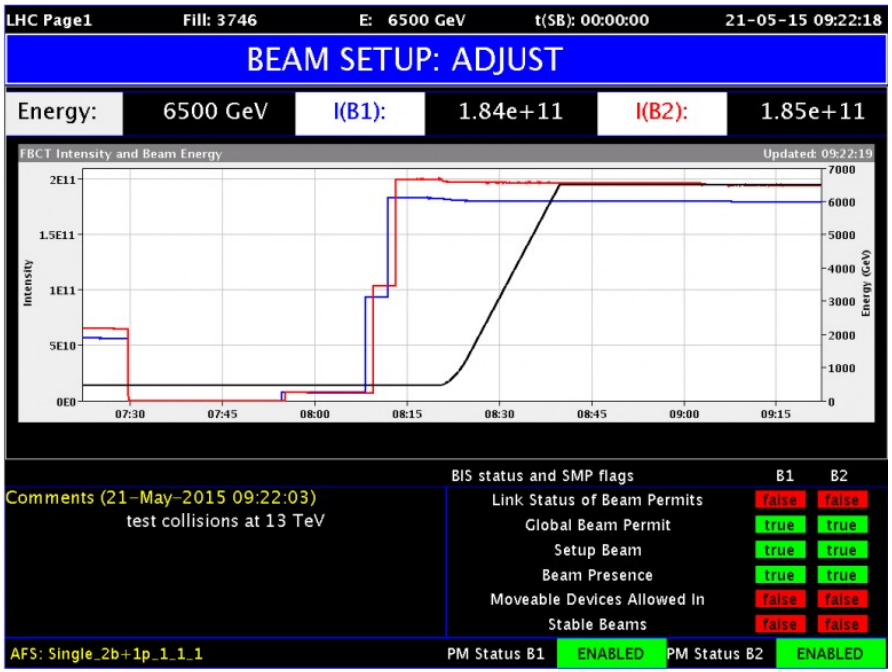


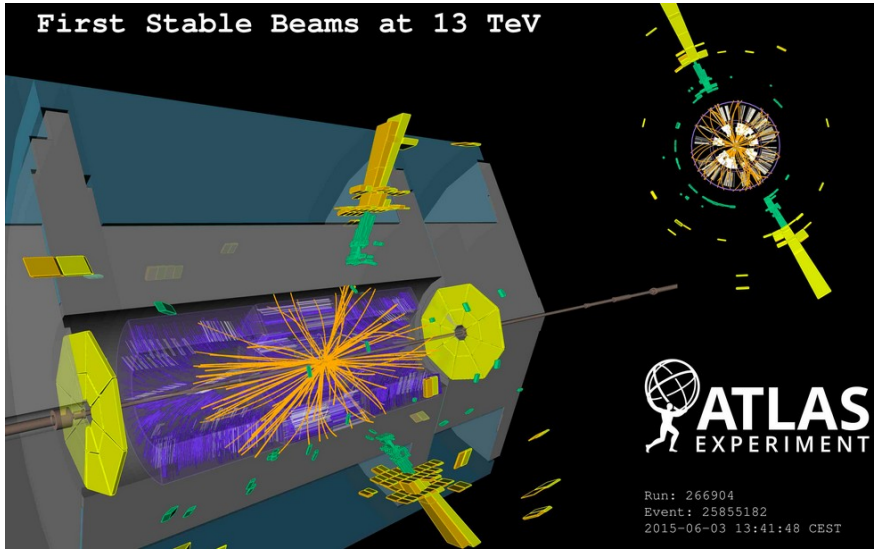
The Forum addresses:

- Set of prioritized simplified models
- Benchmark MC generation for common models (treatment of details, systematics, etc).
- Common procedure for EFT validity treatment

The LHC is up and running at 13 TeV!!

The CMS experiment team celebrated when the first collisions occurred





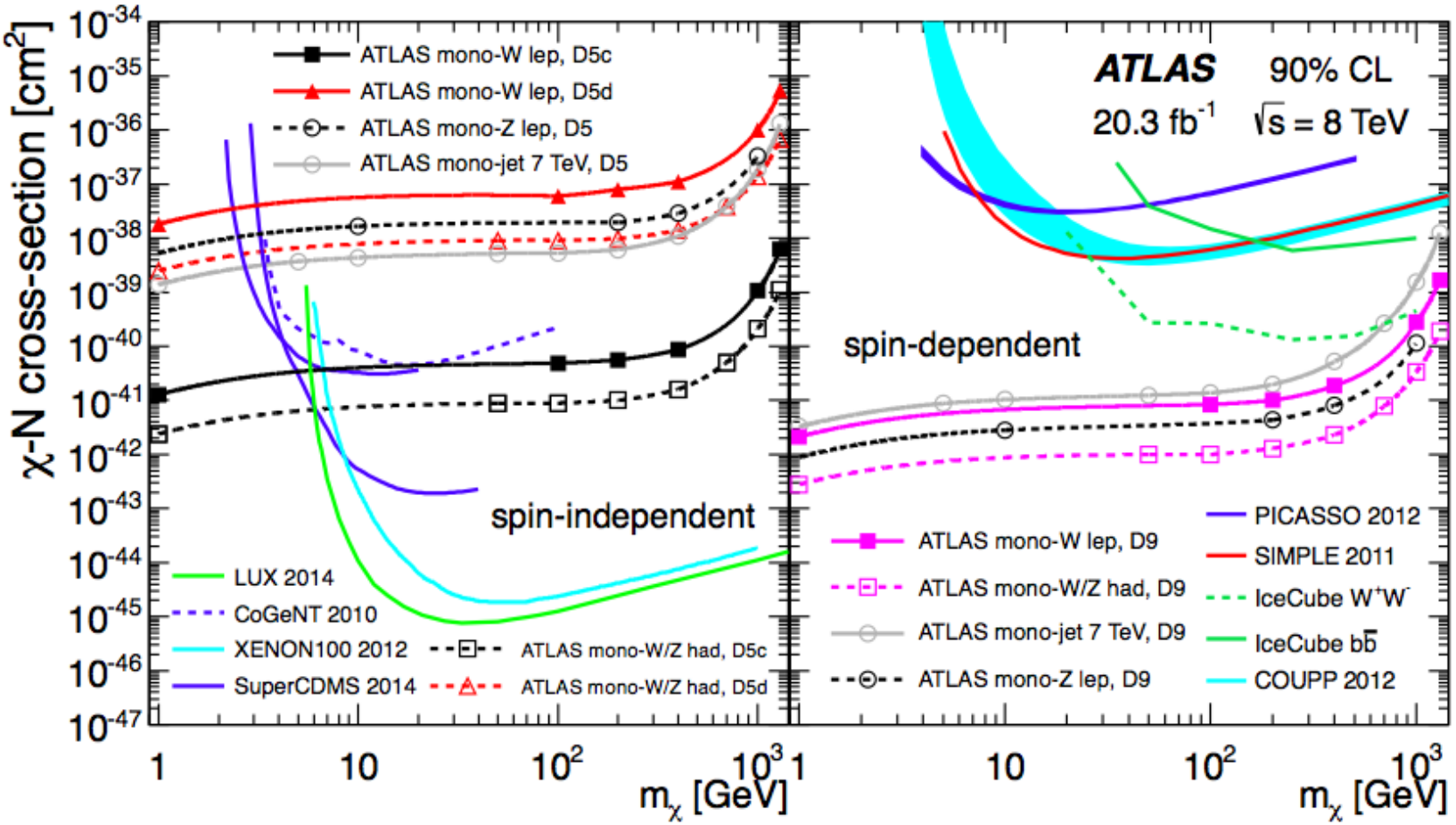
Conclusions

- Dedicated ATLAS and CMS dark matter searches using multiple final states.
- Interpretation of the results done in the context of EFT and simplified models.
- No signal discovered yet → complementary results to other Dark Matter experiments
- But run-II has just started... stay tuned!!

Additional material

ATLAS: PRD 90, 012004 (2014), JHEP 09 (2014) 037

CMS: arXiv:1408.2745



Name	Initial state	Type	Operator
C1	qq	scalar	$\frac{m_q}{M_\star^2}\chi^\dagger\chi\bar{q}q$
C5	gg	scalar	$\frac{1}{4M_\star^2}\chi^\dagger\chi\alpha_s(G^a_{\mu\nu})^2$
D1	qq	scalar	$\frac{m_q}{M_\star^3}\bar{\chi}\chi\bar{q}q$
D5	qq	vector	$\frac{1}{M_\star^2}\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$
D8	qq	axial-vector	$\frac{1}{M_\star^2}\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5 q$
D9	qq	tensor	$\frac{1}{M_\star^2}\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$
D11	gg	scalar	$\frac{1}{4M_\star^3}\bar{\chi}\chi\alpha_s(G^a_{\mu\nu})^2$