

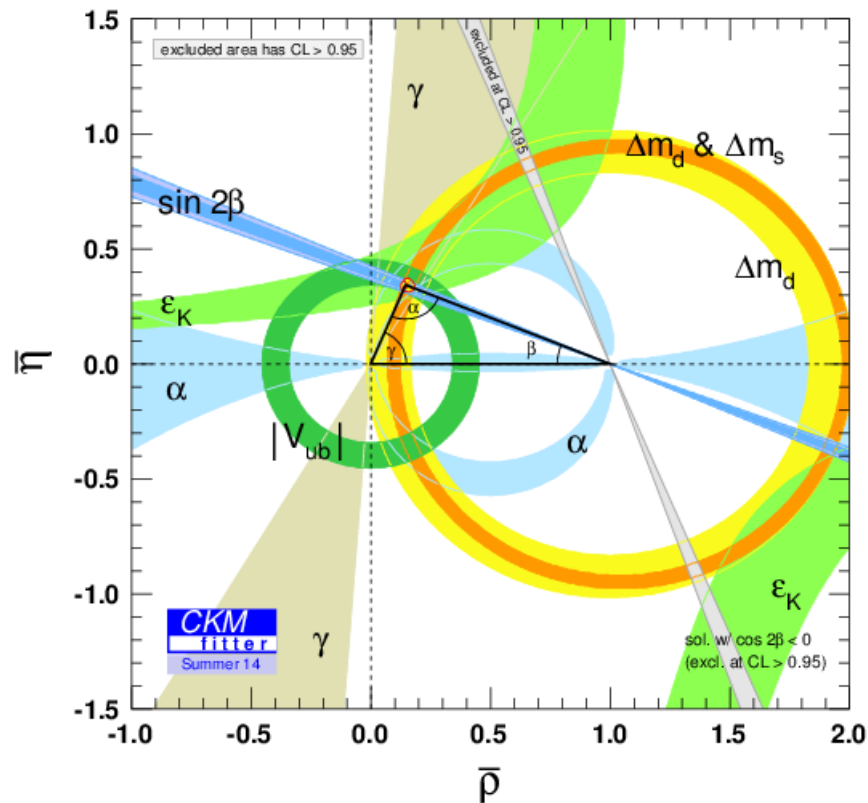
Measurements of φ_3/γ at LHCb

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On behalf of the LHCb collaboration

B2TiP meeting Krakow 27th -29th April 2015

The CKM angle γ



CKM matrix parametrises quark couplings

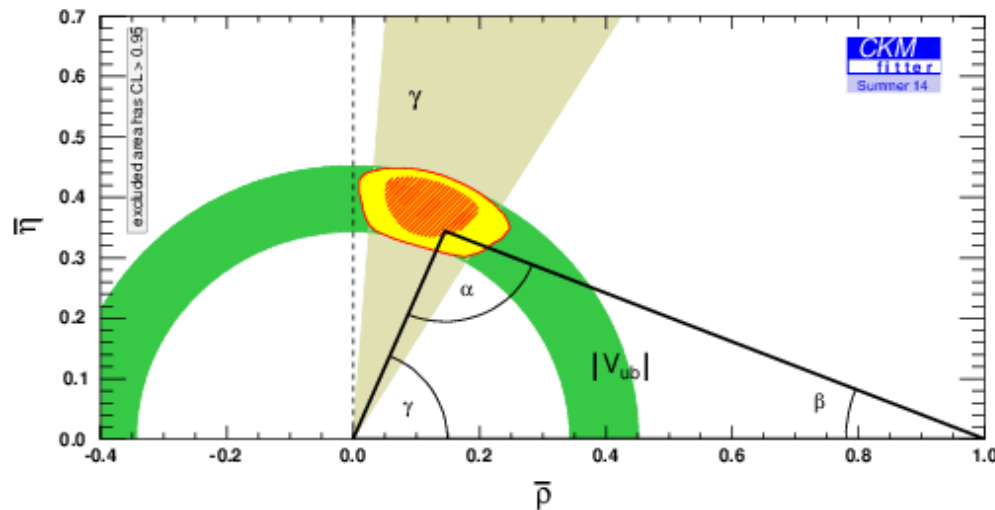
The matrix has one complex phase that results in CPV

Unitarity triangle is a representation of this CPV

$$\gamma = -\arg\left(\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}\right)$$

γ is the least well known angle

Precision measurement → New Physics test bed



γ is the only angle directly accessible in tree decays

Measurement of the “Standard Model”. Theoretical uncertainty: 10^{-7}

Direct measurement: $\gamma = 73.2^{+6.3}_{-7.0}$

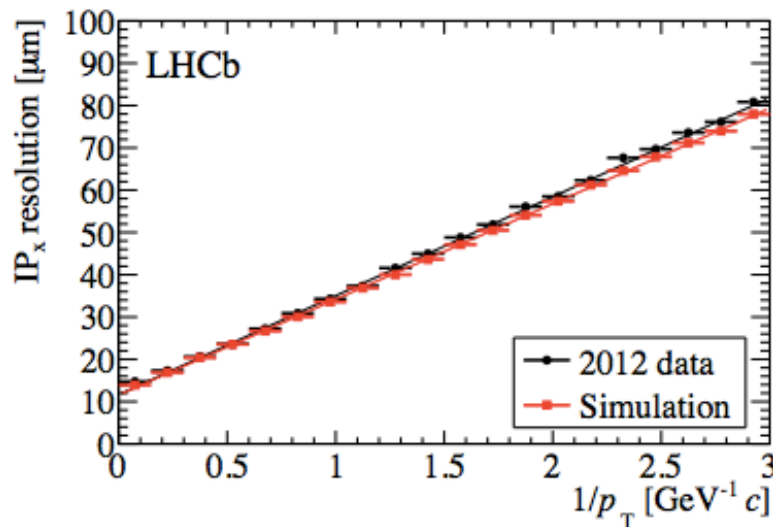
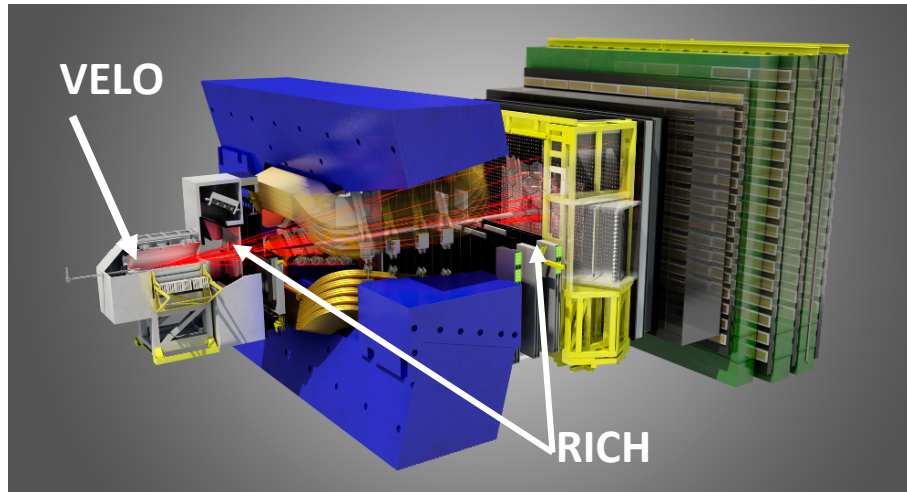
Indirect precision from global CKM fit – includes loop based measurements

$$\gamma = 66.9^{+1.0}_{-3.7}$$

Despite recent progress in the direct measurements better precision still required to test for New Physics.

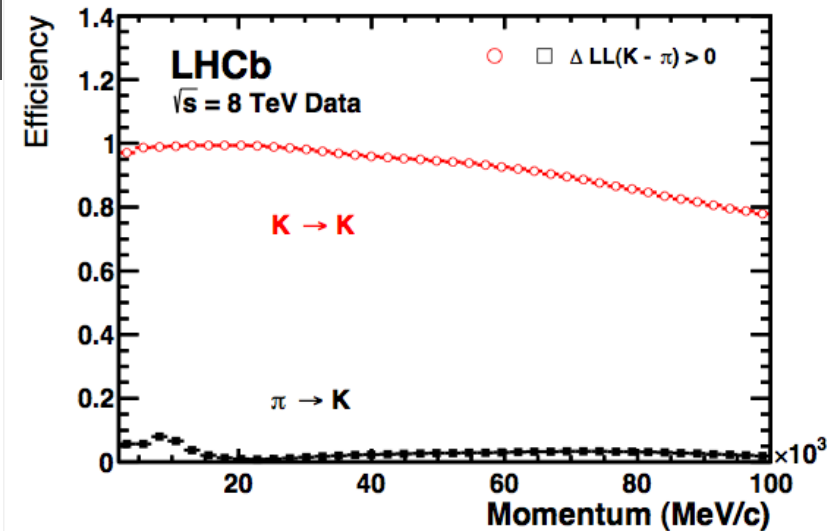
Goal : Improve the direct precision

The LHCb detector



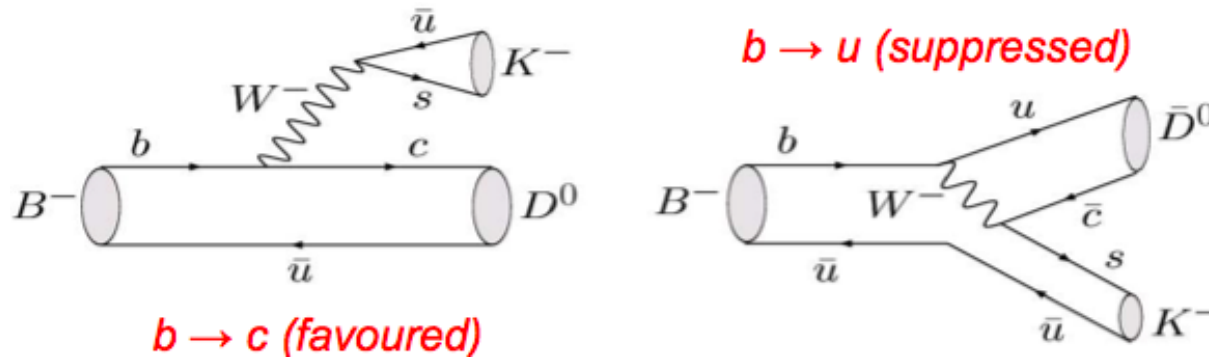
All measurements profit from the VELO for displaced vertices and the RICH for particle identification

$B_s \rightarrow D_s K$ benefits in particular from the excellent time resolution.



Int. J. Mod. Phys. A 30 (2015) 1530022

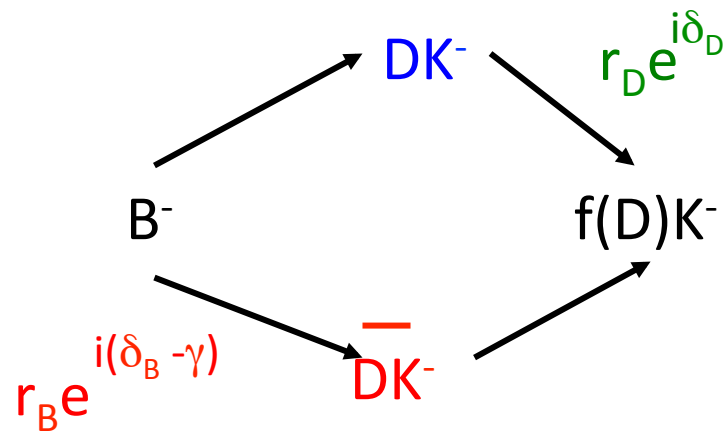
B → DK



Sensitivity to γ from $b \rightarrow c$ and $b \rightarrow u$ interference

Require D^0 and \bar{D}^0 to decay to same final state

$$\gamma = -\arg\left(\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*}\right)$$

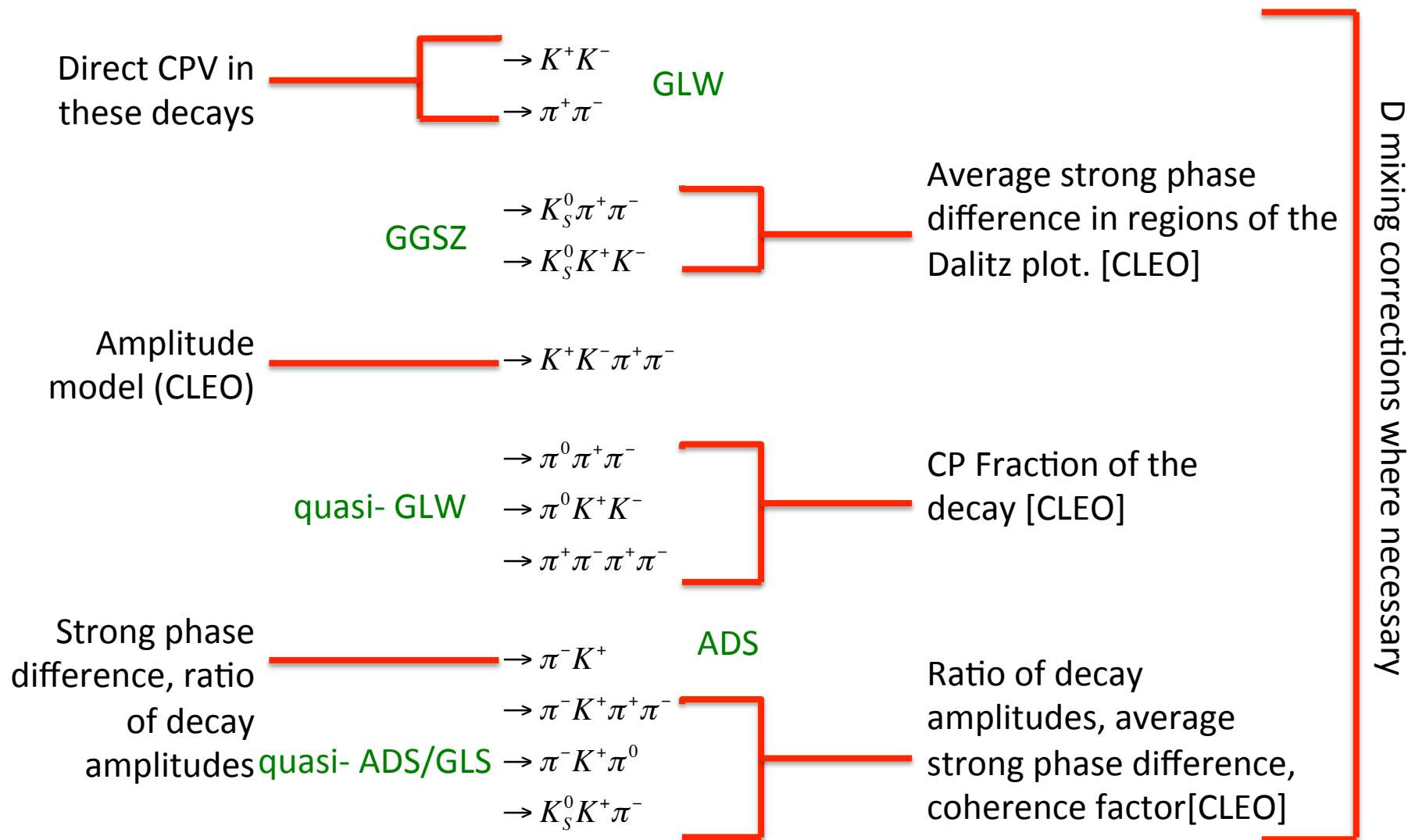


$r_B \sim 0.1$

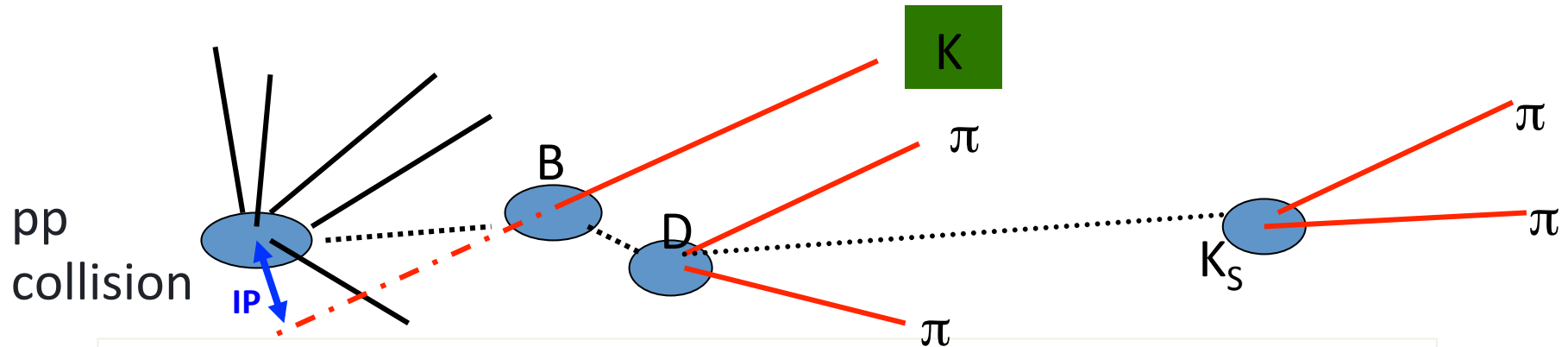
$B \rightarrow D\pi$ also possible $r_B \sim 0.01$ but higher statistics.

Measured observables are asymmetries and ratio of yields.

Pick your D channel



Selection



Separate the topology of interest from random combinations

Useful variables include:

Impact parameters

Flight distances. (B travels a \sim cm, K_S many cm)

Vertex quality

Particle ID

Specific vetos against particular backgrounds

π^0 reconstructed in decay to 2 photons

All analyses shown here employ similar strategies

Using multibody final states – $hh^{(\prime)}\pi^0$

Intermediate resonances dilute the impact of the observables.

$$\begin{aligned}
 R_{\text{ADS}(K)}^{K\pi\pi^0} &\approx (r_B)^2 + (r_D^{K\pi\pi^0})^2 + 2\kappa_D^{K\pi\pi^0} r_B r_D^{K\pi\pi^0} \cos(\delta_B + \delta_D^{K\pi\pi^0}) \cos \gamma \\
 A_{\text{ADS}(K)}^{K\pi\pi^0} &\approx \left[2\kappa_D^{K\pi\pi^0} r_B r_D^{K\pi\pi^0} \sin(\delta_B + \delta_D^{K\pi\pi^0}) \sin \gamma \right] / R_{\text{ADS}(K)}^{K\pi\pi^0} \\
 R_{\text{qGLW}}^{h'h'\pi^0} &= 1 + (r_B)^2 + (2F_+^{h'h'\pi^0} - 1) \cdot 2r_B \cos \delta_B \cos \gamma \\
 A_{\text{qGLW}(K)}^{h'h'\pi^0} &= (2F_+^{h'h'\pi^0} - 1) \cdot 2r_B \sin \delta_B \sin \gamma / R_{\text{qGLW}}^{h'h'\pi^0}
 \end{aligned}$$

Dilution factors
Parameters of interest

Dilution factors all measured

They are large – all three decay modes pursued

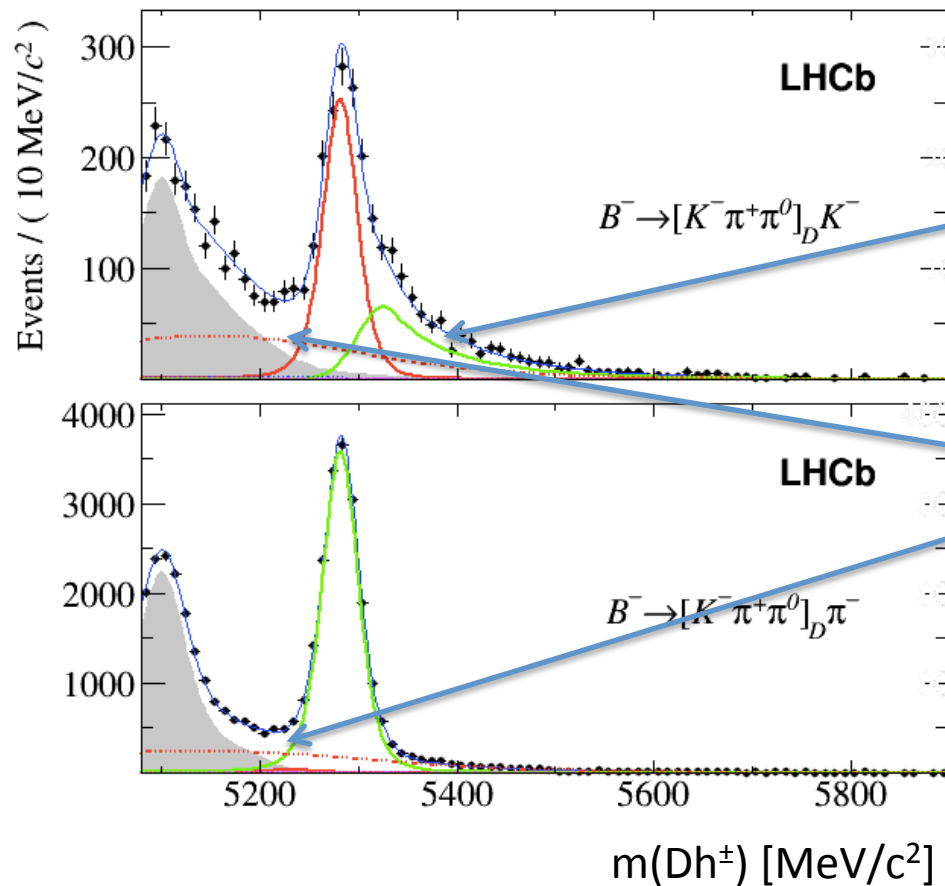
$$\kappa^{K\pi\pi^0} = 0.82 \pm 0.07$$

$$F^{\pi\pi\pi^0} = 0.973 \pm 0.017$$

$$F^{KK\pi^0} = 0.732 \pm 0.055$$

PLB 731 (2014) 197, arXiv :1504.05878

Favoured modes



Mass fit shapes driven by the favoured control channel

PID efficiencies constrain the fraction of $B \rightarrow D\pi$ in the $B \rightarrow DK$ channel

Presence of the π^0 reduces purity in comparison to all charged final states.

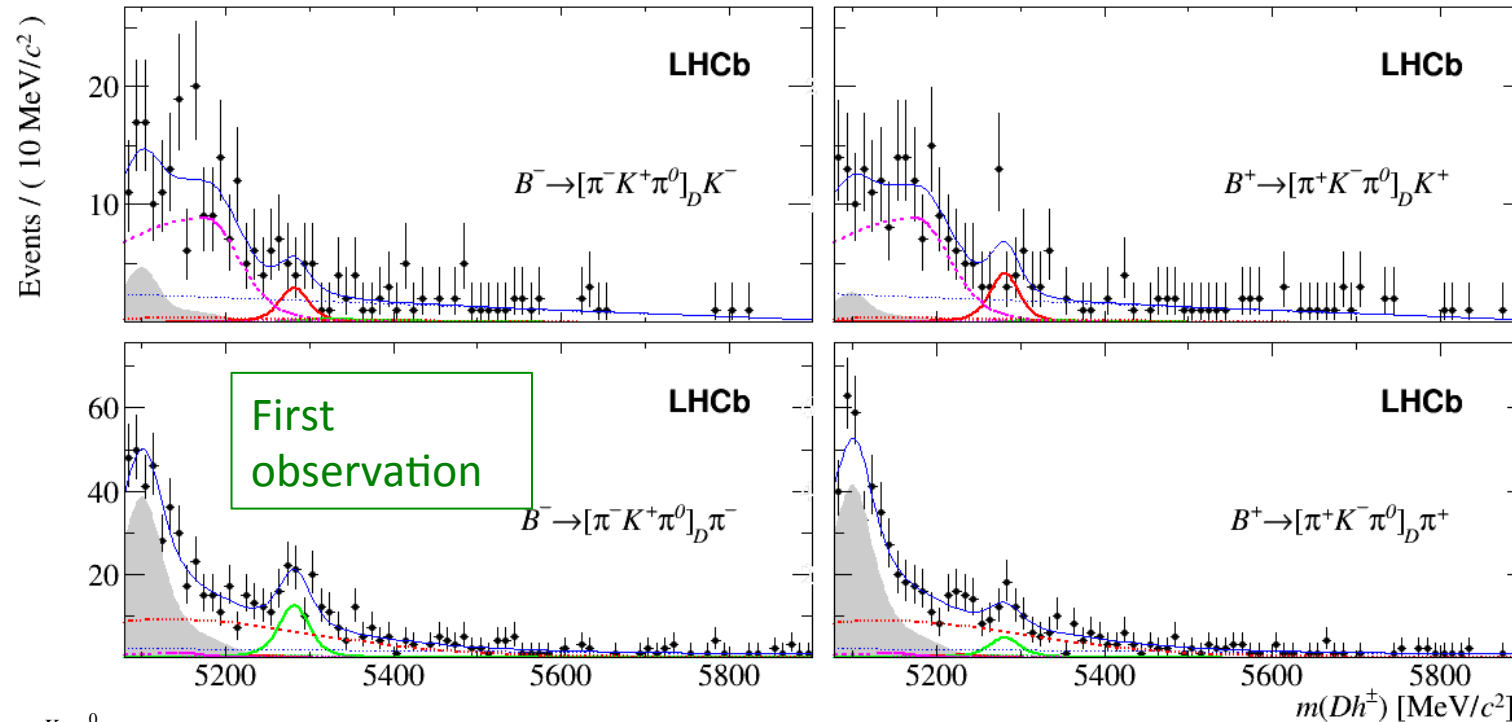
-Wrongly reconstructed D background

Detection and detector asymmetries input – fit observables corrected for these.

Production asymmetries is determined from the favoured decay.

arXiv:1504.05442

ADS-like modes



$$R_{ADS(K)}^{K\pi\pi^0} = 0.0140 \pm 0.0047 \pm 0.0021$$

$$R_{ADS(\pi)}^{K\pi\pi^0} = 0.00235 \pm 0.00049 \pm 0.00006$$

$$A_{ADS(K)}^{K\pi\pi^0} = -0.20 \pm 0.27 \pm 0.03$$

$$A_{ADS(\pi)}^{K\pi\pi^0} = 0.438 \pm 0.190 \pm 0.009$$

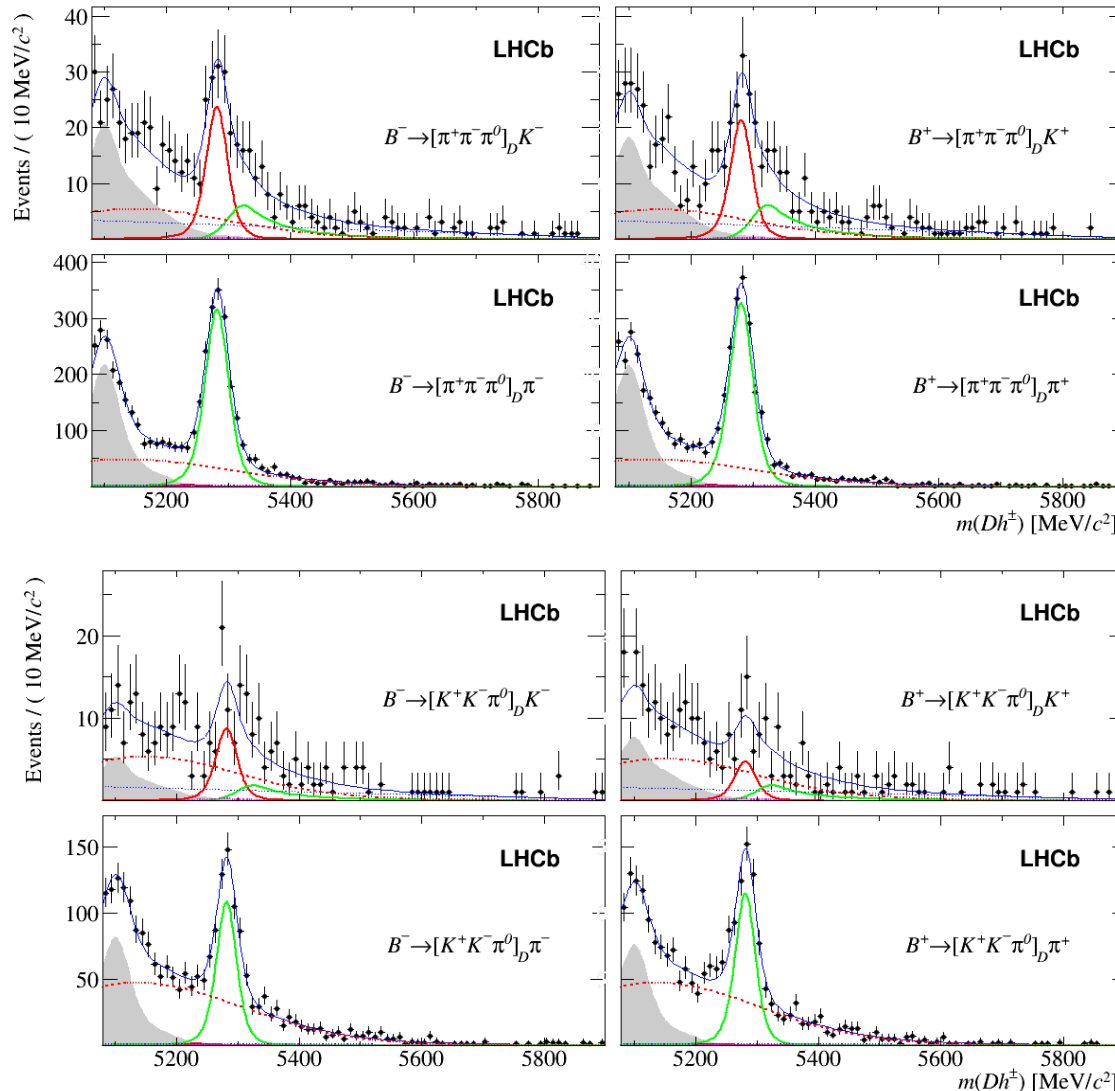
Uncertainties smaller than an analysis from Belle

Significance of $B \rightarrow DK$ signal is 2.8σ

Significance of $B \rightarrow D\pi$ signal is 5.3σ

arXiv:1504.05442

GLW-like modes



Statistics not sufficient for significant asymmetries

First observation of $B \rightarrow [KK\pi^0] \pi$. ($>10 \sigma$)

First evidence of $B \rightarrow [KK\pi^0] K$ (4.5σ)

$$A_{qGLW(K)}^{\pi\pi\pi^0} = 0.054 \pm 0.091 \pm 0.011$$

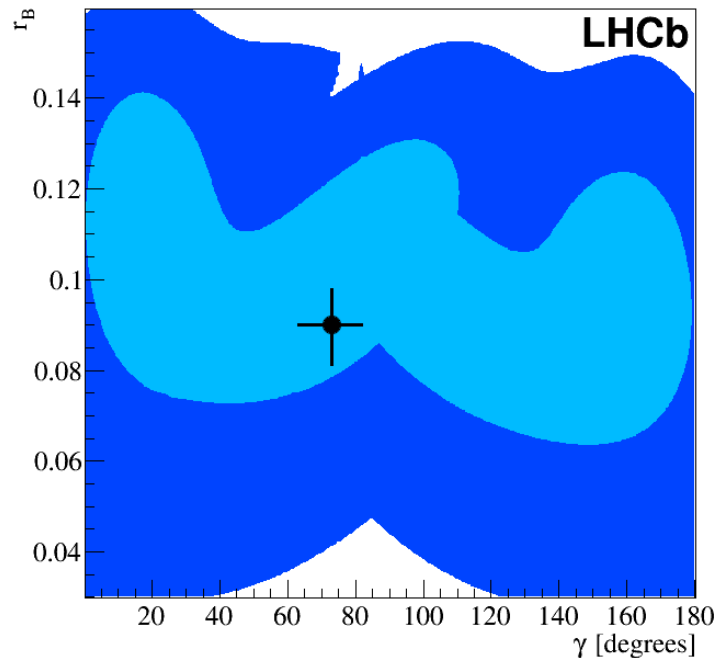
$$A_{qGLW(\pi)}^{\pi\pi\pi^0} = -0.016 \pm 0.020 \pm 0.004$$

$$A_{qGLW(K)}^{KK\pi^0} = 0.030 \pm 0.020 \pm 0.02$$

$$A_{qGLW(\pi)}^{KK\pi^0} = -0.030 \pm 0.040 \pm 0.005$$

arXiv:1504.05442

Interpretation of these results



2D likelihood scans of γ - r_B

1, 2 σ contours shown

Cross marks the best of combination of other LHCb results

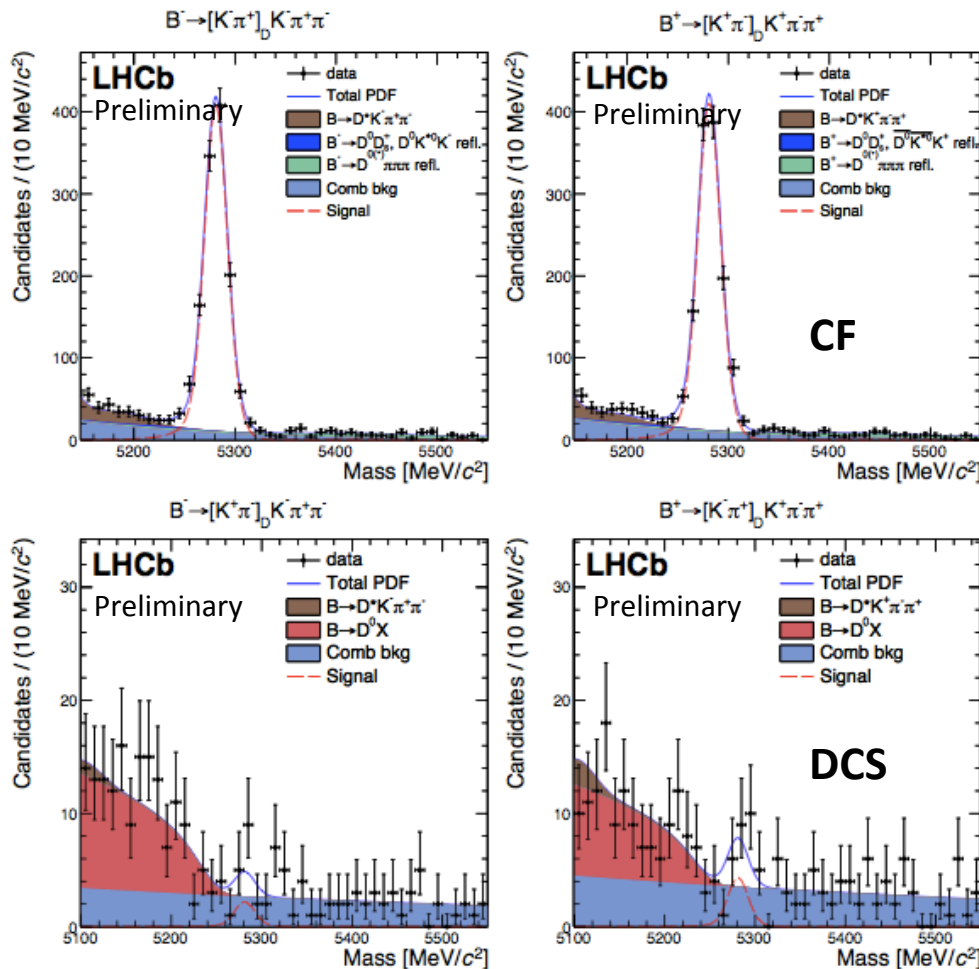
Although with this set of measurements γ is unconstrained $r_B = 0.11 \pm 0.03$

Will help further constrain γ , r_B , δ_B when entered into the combination

arXiv:1504.05442



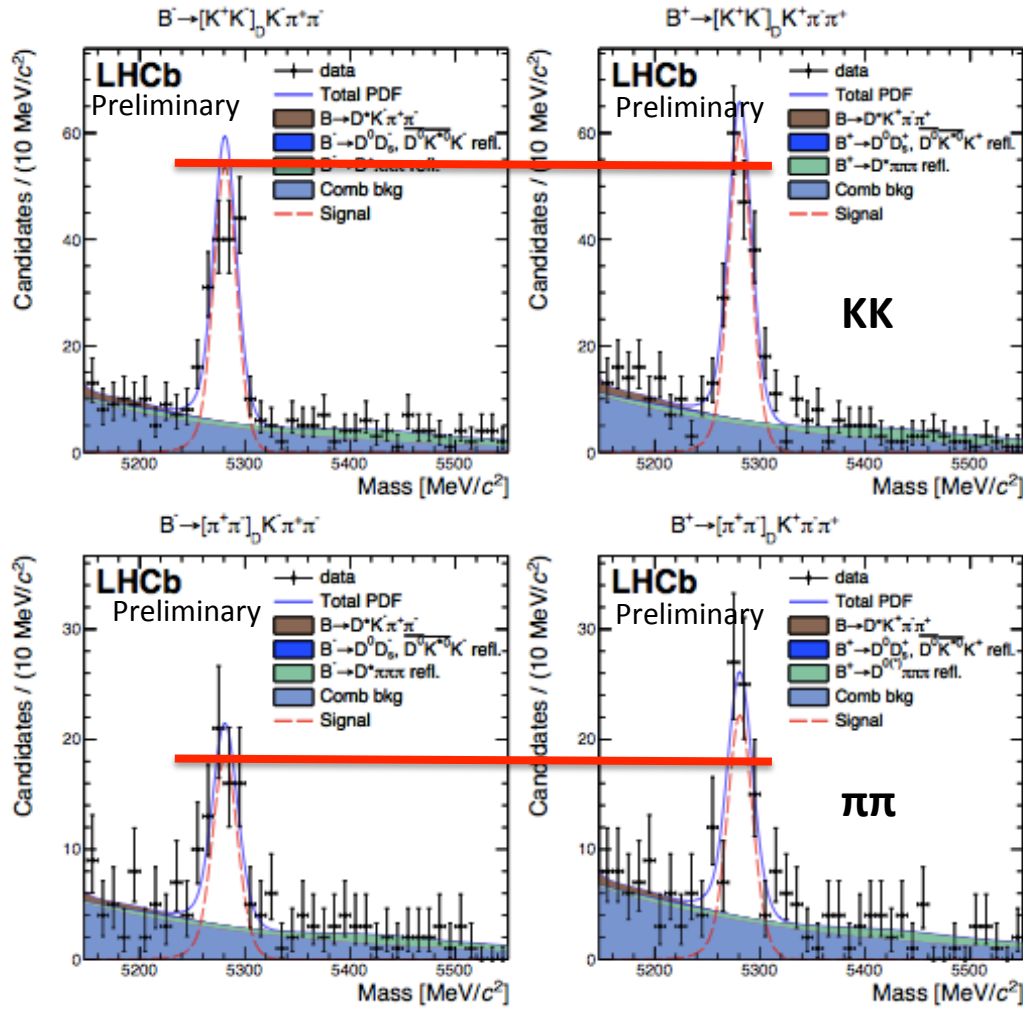
Other B decays: $B \rightarrow DK\pi\pi$ (1)



- Analyse KK, $\pi\pi$, πK decay modes of the D
- r_B, δ_B are different to $B \rightarrow DK, \gamma$ is common
- Additional coherence factor due to different resonances in $B \rightarrow DK\pi\pi$ and $B \rightarrow \overline{DK}\pi\pi$
- Determined alongside the other physics parameters of interest
- $K\pi\pi$ system phasespace is restricted to enhance coherence
 - $M(K\pi\pi) < 2 \text{ GeV}/c^2$
 - $M(K\pi)$ within 100 MeV/c² of K^*
- $B \rightarrow D\pi\pi\pi$ also studied for CPV



Other B decays: $B \rightarrow Dh\pi\pi$ (2)



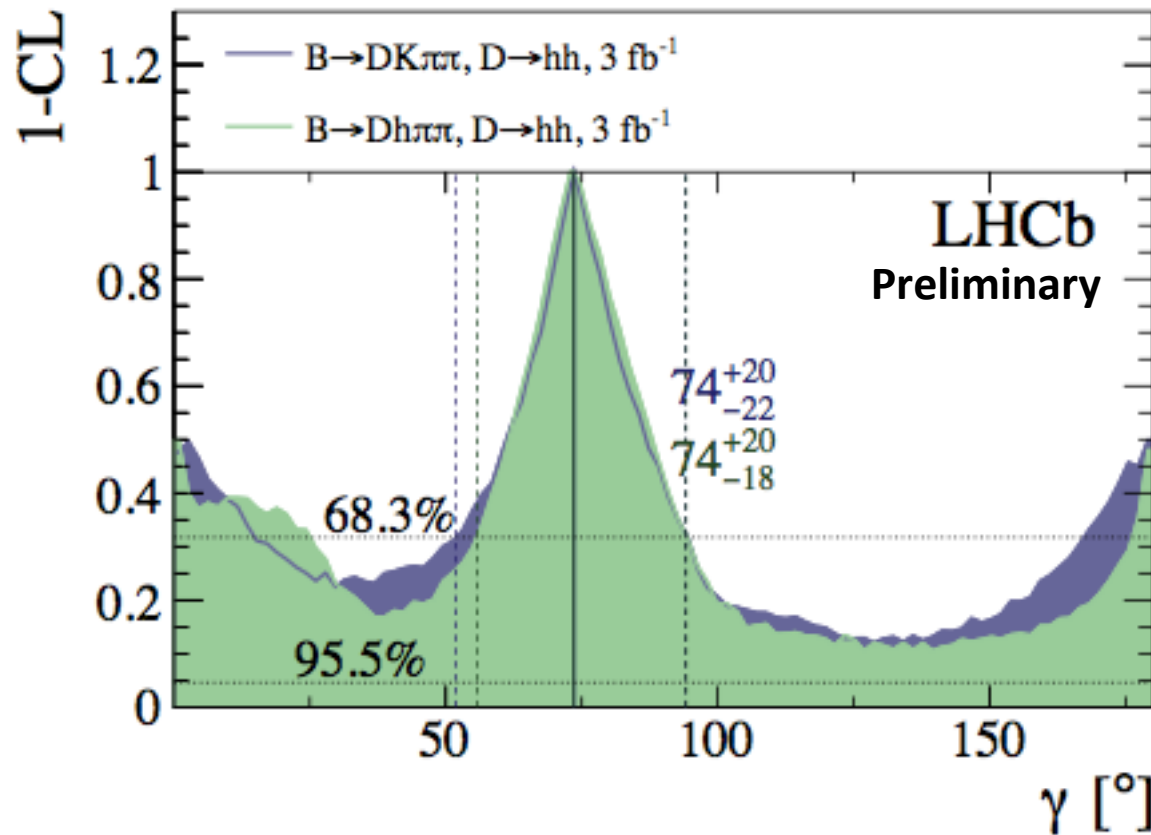
Different selection used for the GLW and ADS modes

Looser selection possible in the GLW modes due to higher signal yields.

First observation of all these decays



Constraints on γ



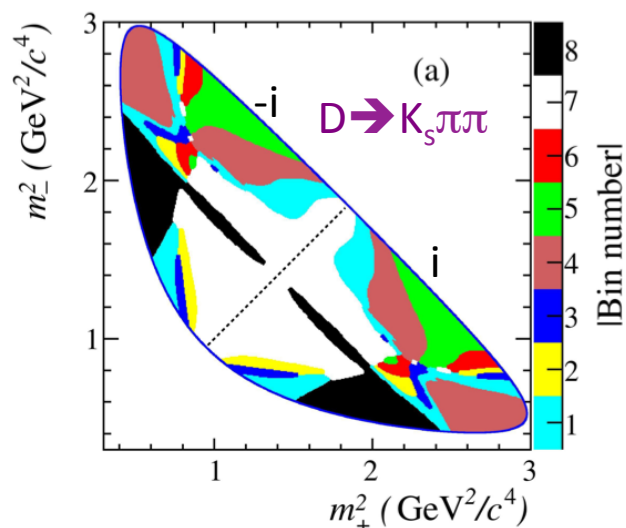
Strong constraints set on γ at the 68.3% CL.

This decay channel has potential

Additional information from $B \rightarrow D\pi\pi\pi$ adds a little sensitivity

Coherence factor close to 1 though no constraints set in the physical limit [0,1]

The golden modes – $K_S\pi\pi$, $K_S KK$



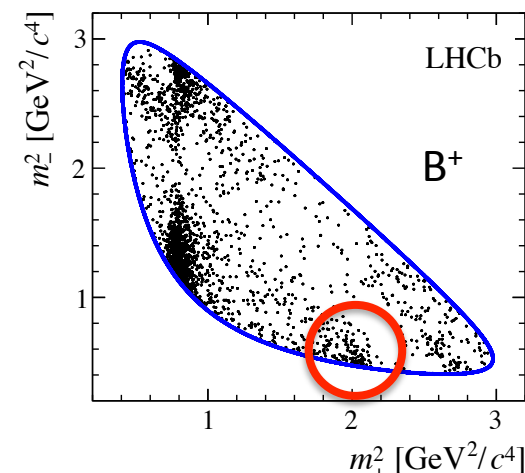
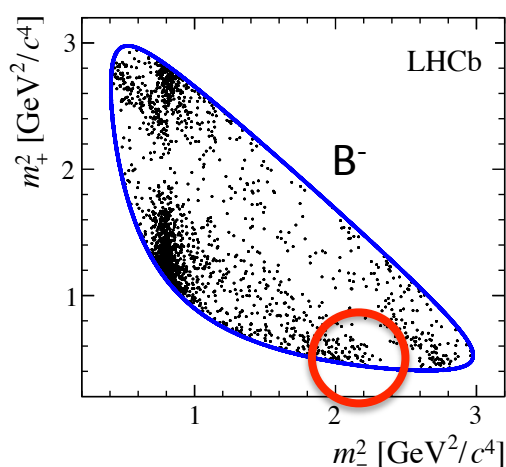
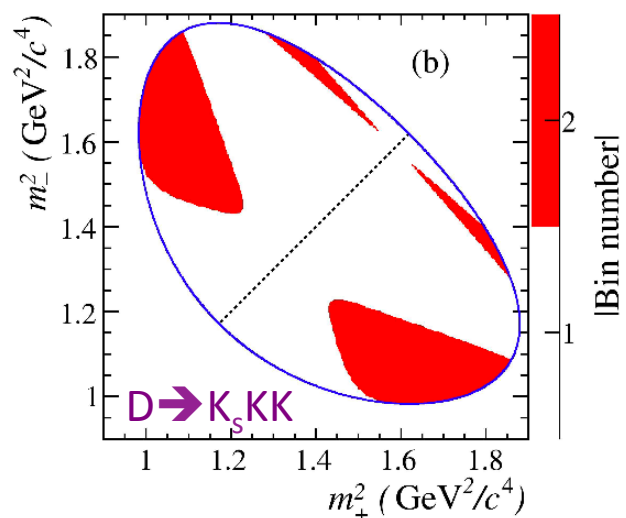
These decays analysed in regions of the Dalitz plot – not much sensitivity if analysed inclusively.

Binning shapes optimised for statistical sensitivity.

Symmetry of Dalitz plot defines positive and negative bins.

Reduces the analysis to a counting experiment in bins of the Dalitz plot - model independent

Data from $D \rightarrow K_S KK$ easily added as two additional bins.



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Analysis strategy

To determine γ : Count the number of observed events in a region of the Dalitz plot.

$$N_i^\pm = h \left(K_{\pm i} + r_B^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} \left[x_\pm c_i \pm y_\pm s_i \right] \right)$$

D from B^\pm events in bin i of Dalitz plot

Fraction of events in bin for pure D^0 sample with the efficiency profile of signal

$x_\pm = r_B \cos(\delta_B \pm \gamma)$
 $y_\pm = r_B \sin(\delta_B \pm \gamma)$

K_i are inputs from other LHCb decays - use $B^0 \rightarrow D^{*+} \mu^- \nu$, $D^{*+} \rightarrow D^0 \pi^+$

Charge of the π tags the D^0 flavour.

c_i and s_i are inputs from CLEO - they are measurements of the cosine and sine of the average strong phase difference

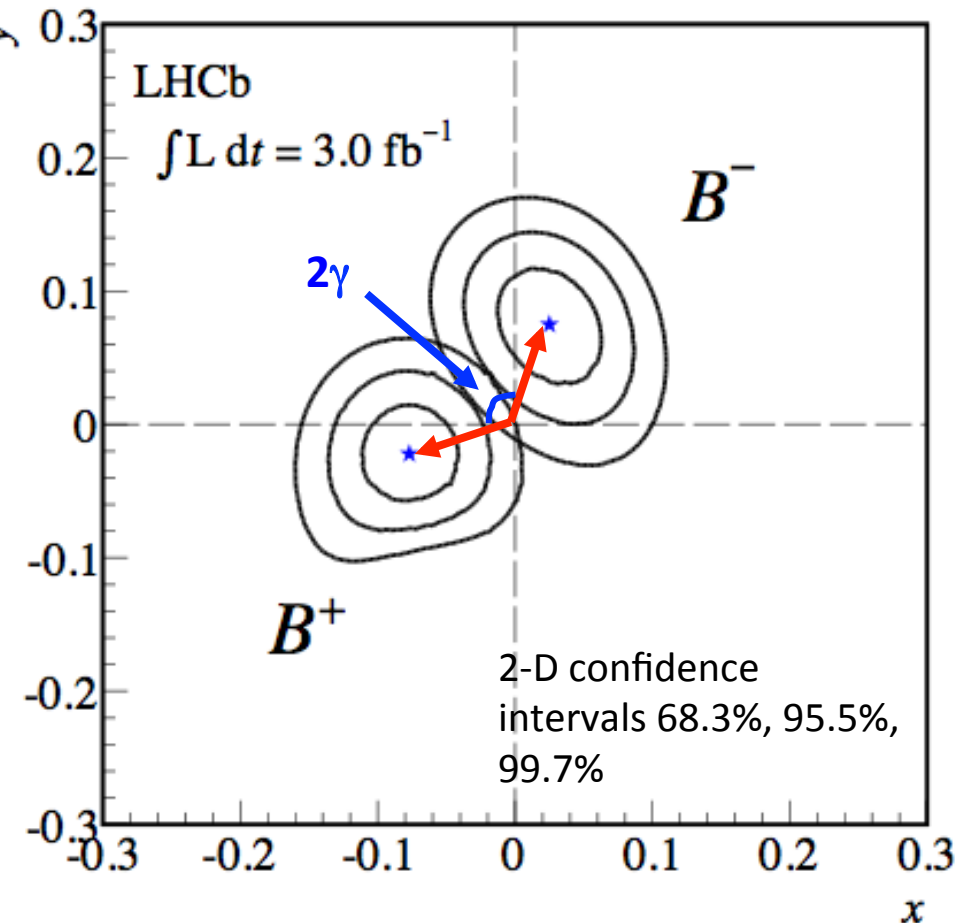
Simultaneous mass fit to candidates in all bins to extract best x, y

Combined measurement of 3 fb^{-1} data

Results

$$\begin{aligned}x_+ &= (-7.7 \pm 2.4 \pm 1.0 \pm 0.4) \times 10^{-2}, \\x_- &= (2.5 \pm 2.5 \pm 1.0 \pm 0.5) \times 10^{-2}, \\y_+ &= (-2.2 \pm 2.5 \pm 0.4 \pm 1.0) \times 10^{-2}, \\y_- &= (7.5 \pm 2.9 \pm 0.5 \pm 1.4) \times 10^{-2},\end{aligned}$$

Corrections for D mixing, K_s CPV
ignored - negligible effect



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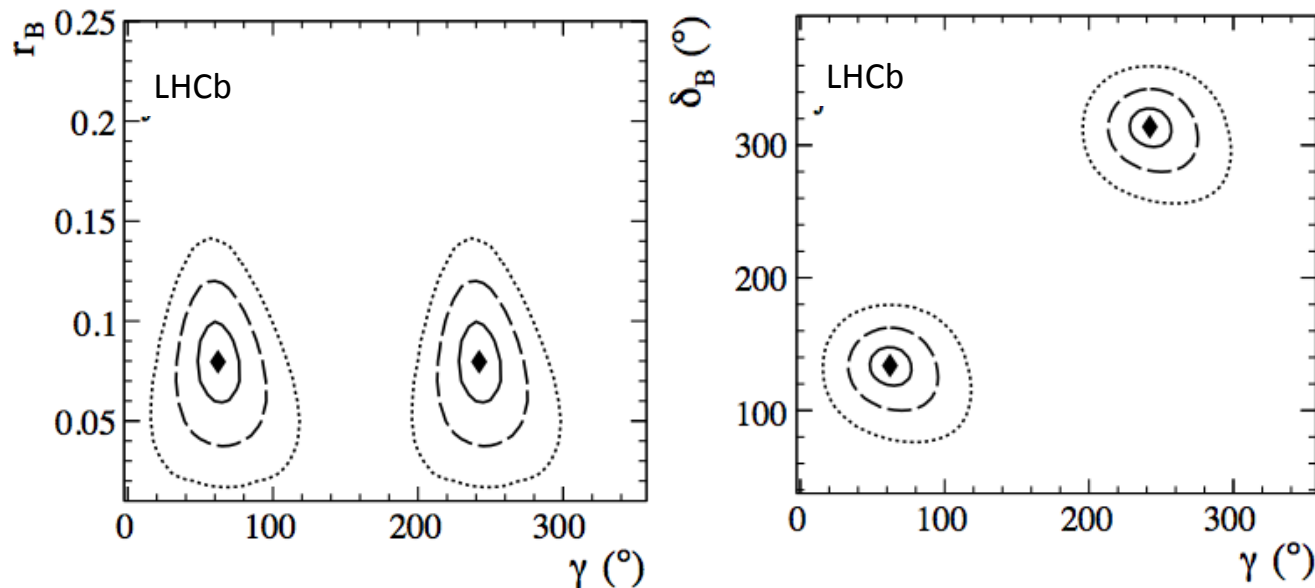
Constraints on γ

$$\gamma = (62_{-14}^{+15})^\circ$$

$$r_B = (8.0_{-2.1}^{+1.9}) \times 10^{-2}$$

$$\delta_B = (134_{-15}^{+14})^\circ$$

Precision matches that of either B factory γ combinations

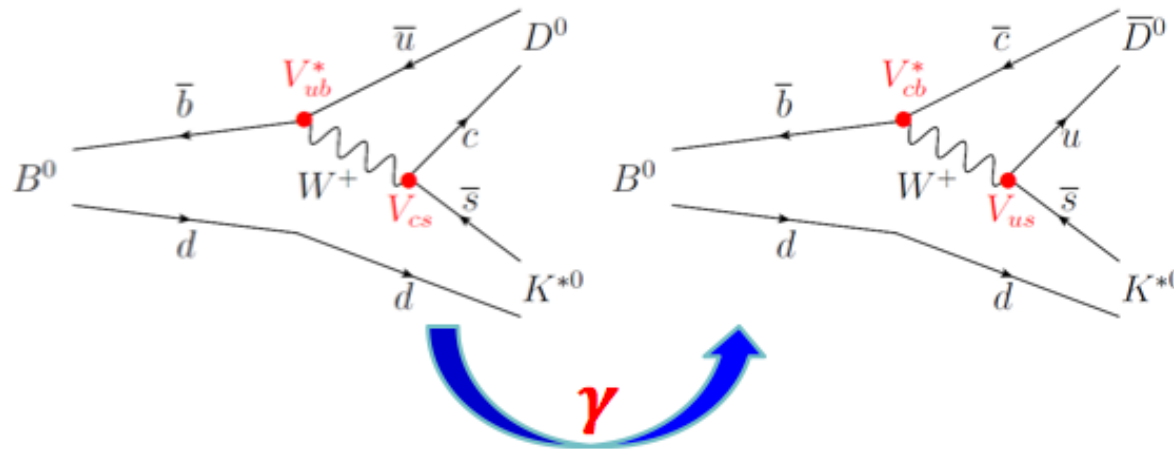


Projection of contours onto 1-D gives 68.3, 95.5, 99.7 % CL

Two-fold ambiguity from trigonometric relations

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$B^0 \rightarrow DK^*$



Both diagrams colour suppressed \rightarrow large r_B (large interference)

Initial B^0 flavour tagged by the charge of the kaon in the $K^{0*} \rightarrow K\pi$

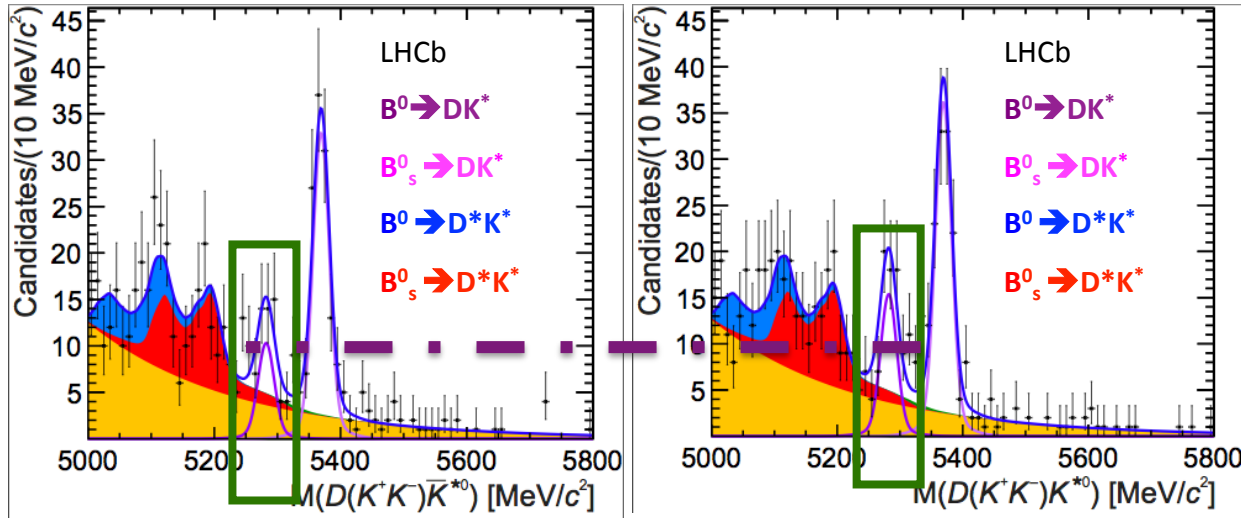
(no need for time-dependent analysis)

Interference requires same D final state

So far two-body D decay modes considered with 3 fb^{-1}

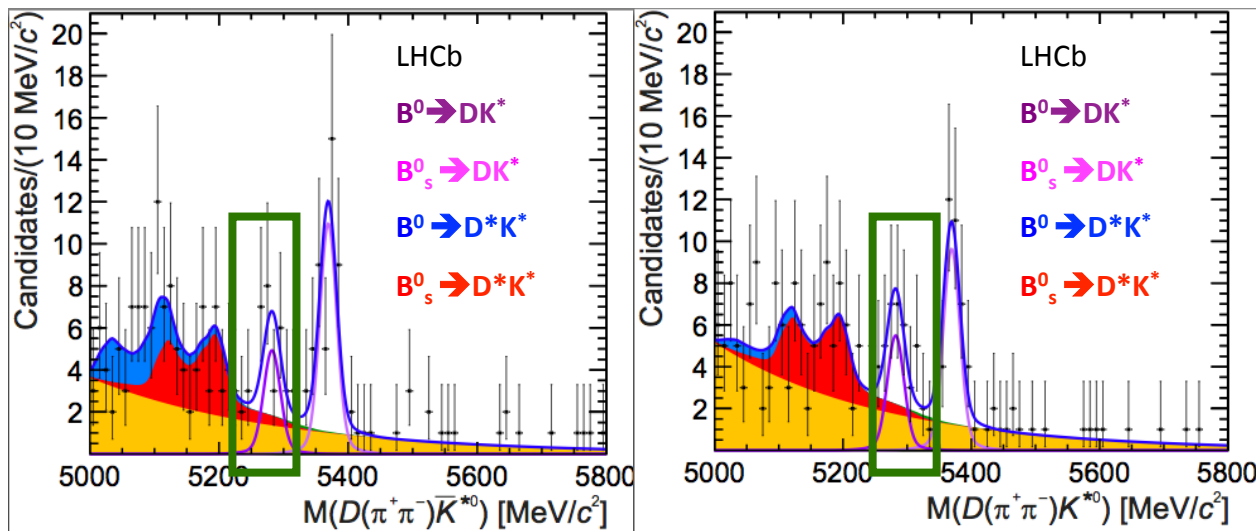
γ common to all analyses; here different B decay means different r_B and δ_B

KK, $\pi\pi$ in $B^0 \rightarrow DK^{*0}$



KK - significance of signal: 8.6σ

$\pi\pi$ - significance of signal: 5.8σ

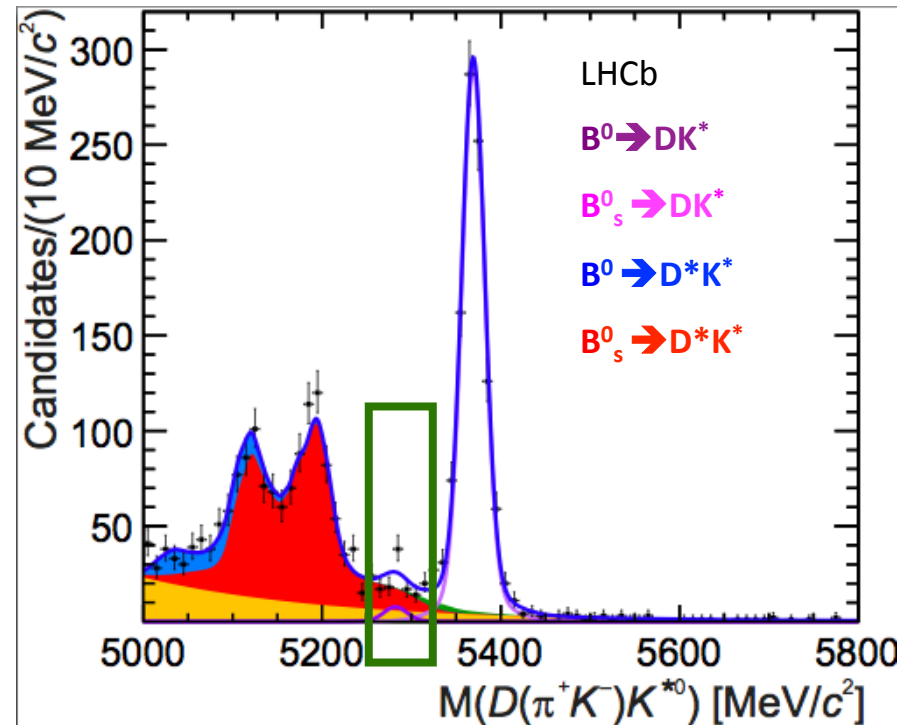
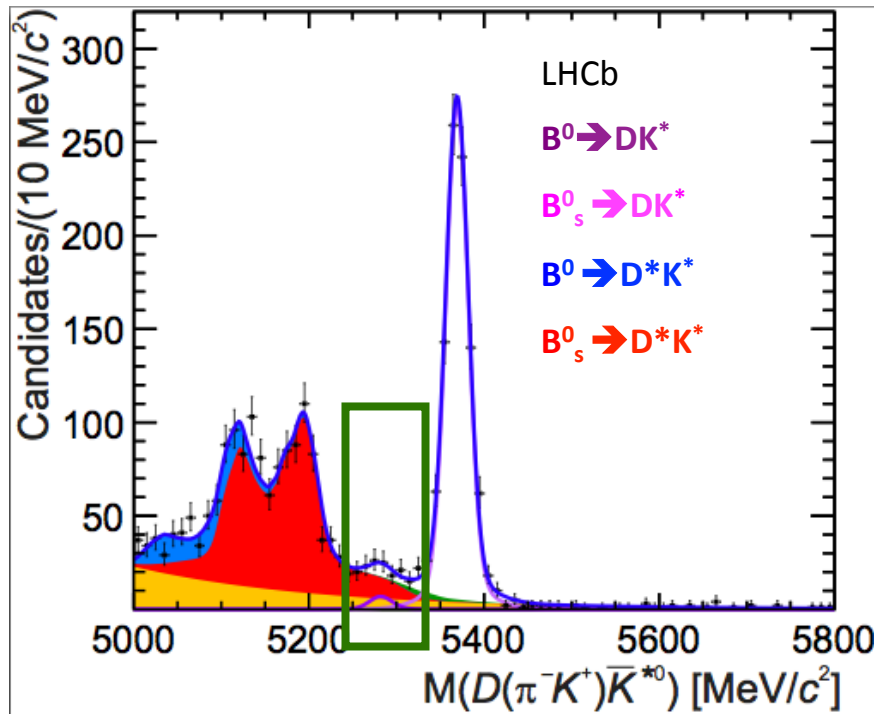


$$A_d^{KK} = -0.198^{+0.144+0.019}_{-0.145-0.020}$$

$$A_d^{\pi\pi} = -0.092^{+0.217+0.019}_{-0.217-0.019}$$

PRD 90 (2014) 112002

πK in $B^0 \rightarrow DK^{*0}$



Combined signal significance is 2.9σ .

$$R_d^+ = 0.057^{+0.029+0.009}_{-0.027-0.012}$$

$$R_d^- = 0.056^{+0.032+0.009}_{-0.030-0.012}$$

PRD 90 (2014) 112002

Interpretation

Coherence factor κ determined from simulation of a realistic model of the resonance content of $B^0 \rightarrow DK\pi$

$$\kappa = 0.95 \pm 0.03$$

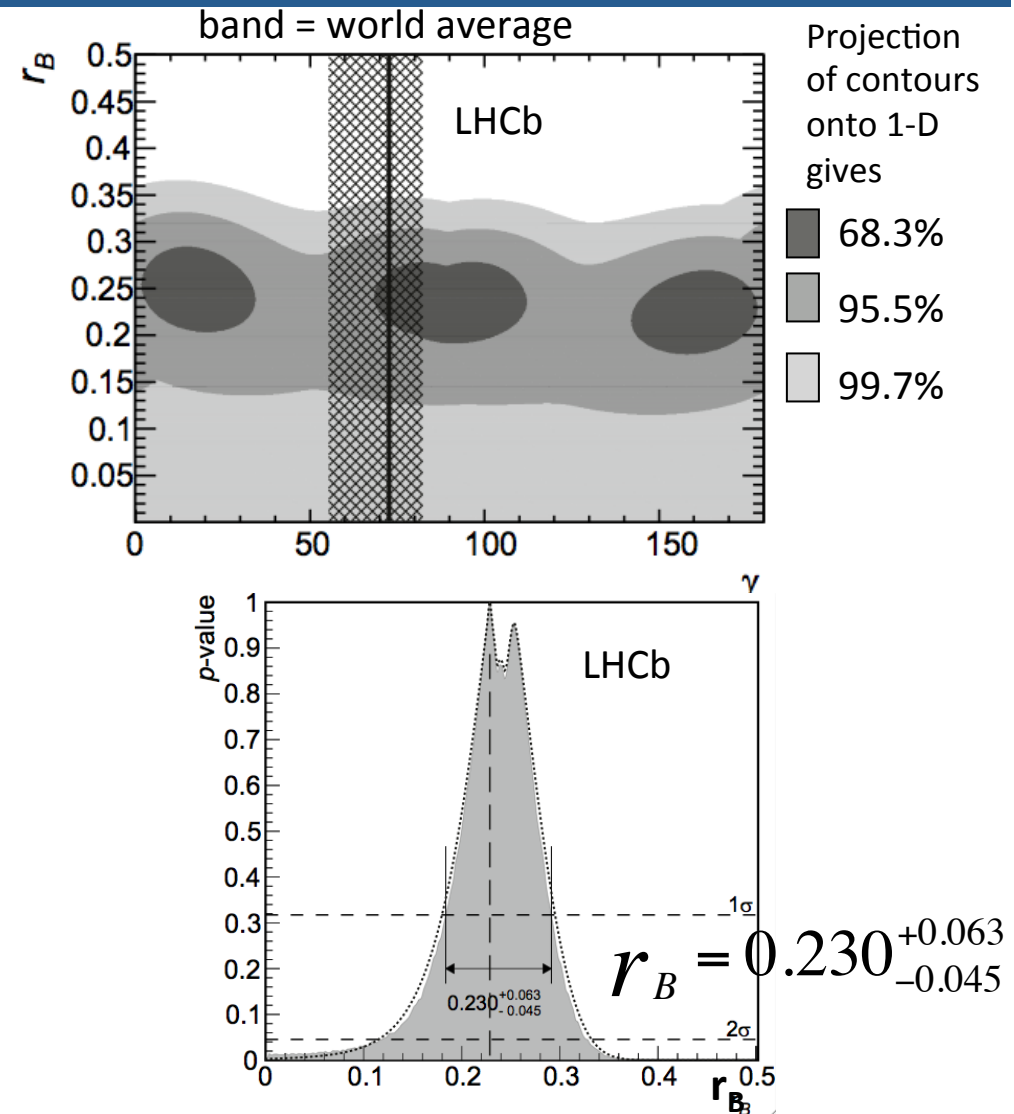
Measurements of all observables combined to determine r_B, δ_B, γ

Ambiguities from the trigonometric relations

Some constraints can be set at the 68.3% CL

r_B value larger than that for $B \rightarrow DK$

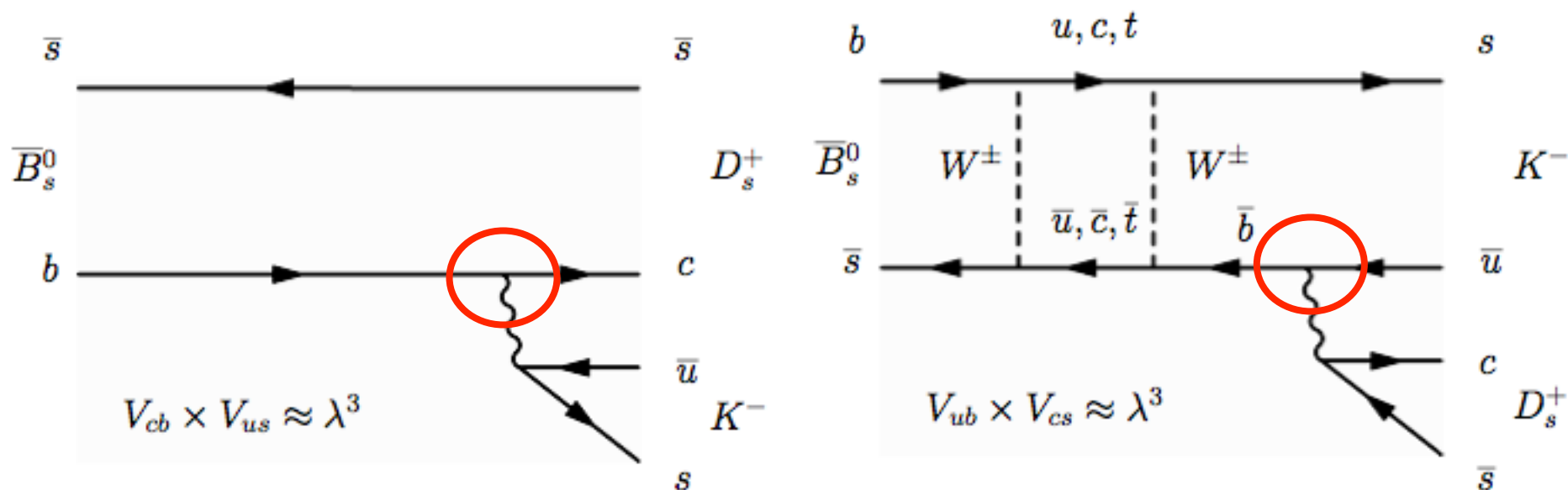
Promising decay to study further



PRD 90 (2014) 112002

Time dependent methods

Measure CP violation in the interference of mixing and decay



Both decay amplitudes $\sim \lambda^3 \rightarrow$ Large interference

Tree level process like other analyses shown

Time-dependence increases the complexity of the analysis

Decay rate relations

$$\frac{d\Gamma_{B_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2}|A_f|^2(1 + |\lambda_f|^2)e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t) \right],$$

$$\frac{d\Gamma_{\bar{B}_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2}|A_f|^2 \left| \frac{p}{q} \right|^2 (1 + |\lambda_f|^2)e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) - C_f \cos(\Delta m_s t) + S_f \sin(\Delta m_s t) \right],$$

$$\lambda_f \equiv \frac{q}{p} \left(\frac{\bar{A}_f}{A_f} \right)$$

A_f is the decay amplitude for B_s to decay to final state f

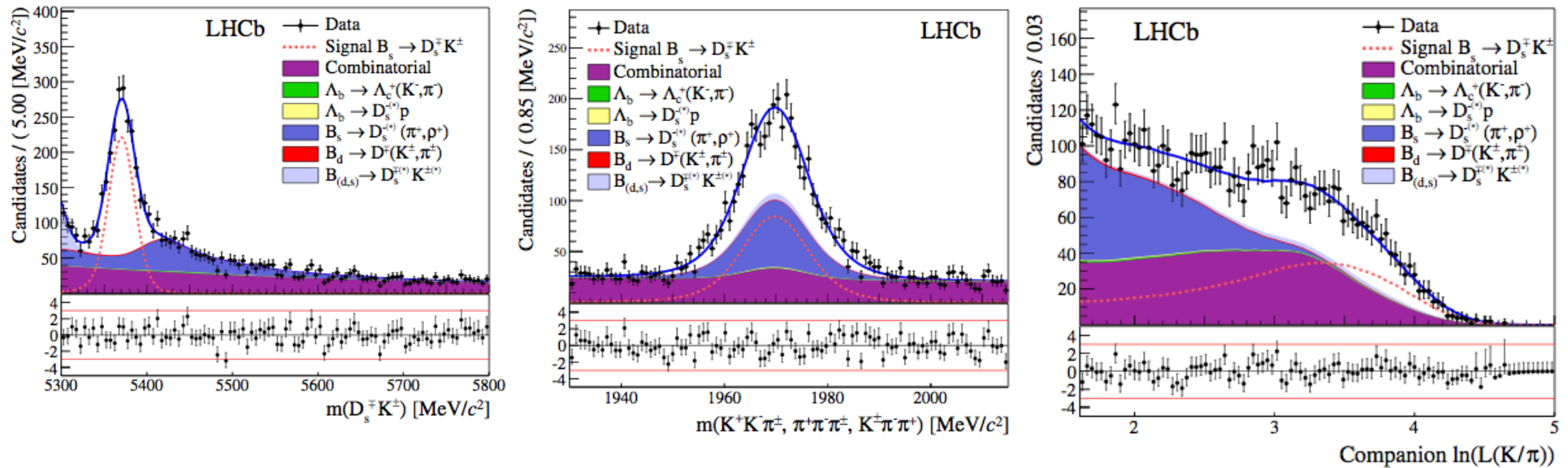
$$C_f = \frac{1 - r_{D_s K}^2}{1 + r_{D_s K}^2},$$

$$A_f^{\Delta\Gamma} = \frac{-2r_{D_s K} \cos(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad \bar{A}_f^{\Delta\Gamma} = \frac{-2r_{D_s K} \cos(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2},$$

$$S_f = \frac{2r_{D_s K} \sin(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad \bar{S}_f = \frac{-2r_{D_s K} \sin(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}.$$

β_s - mixing phase

Parametrising signal



Three D_s^- decays considered: $K^- K^+ \pi^-$, $\pi^- \pi^+ \pi^-$, $K^- \pi^+ \pi^-$: Plots show all D_s^- states combined
 Simultaneous fit in 3 variables: $M(B_s)$, $M(D_s^-)$ and PID variable on the Kaon from the B
 Allows for signal/background discrimination, and for determination of signal weights

Tagging and decay time performance

Flavour tagging:

Combination of SS and OS taggers

Efficiency of tagging an event = 67.5%

Effective tagging power = 5.07%

Decay time acceptance:

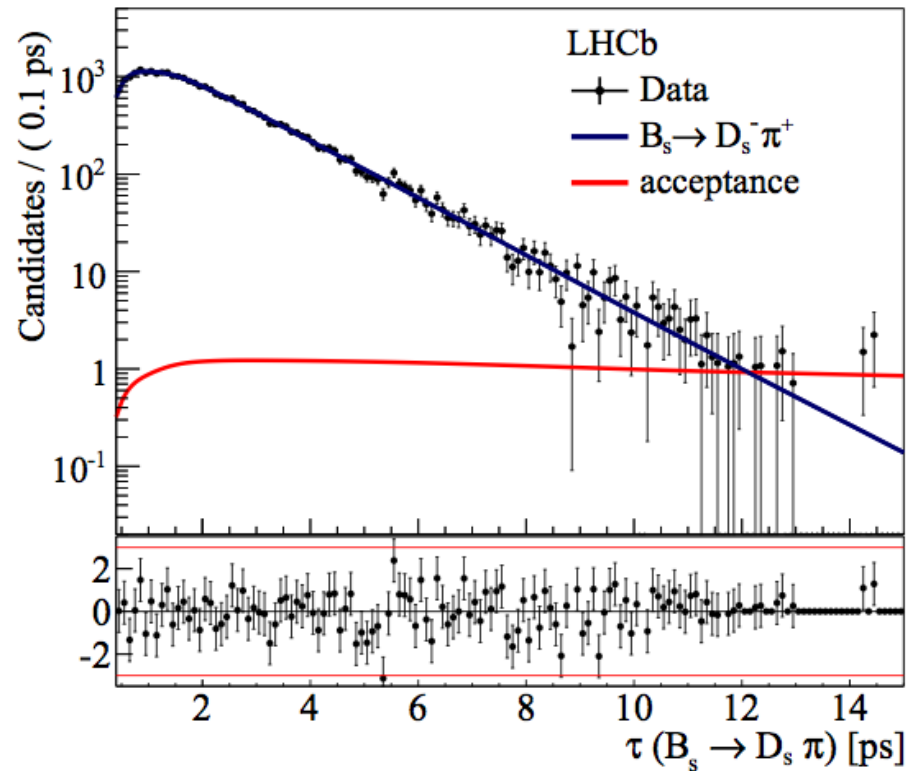
$B_s \rightarrow D_s \pi$ with additional corrections from simulation

Decay time resolution:

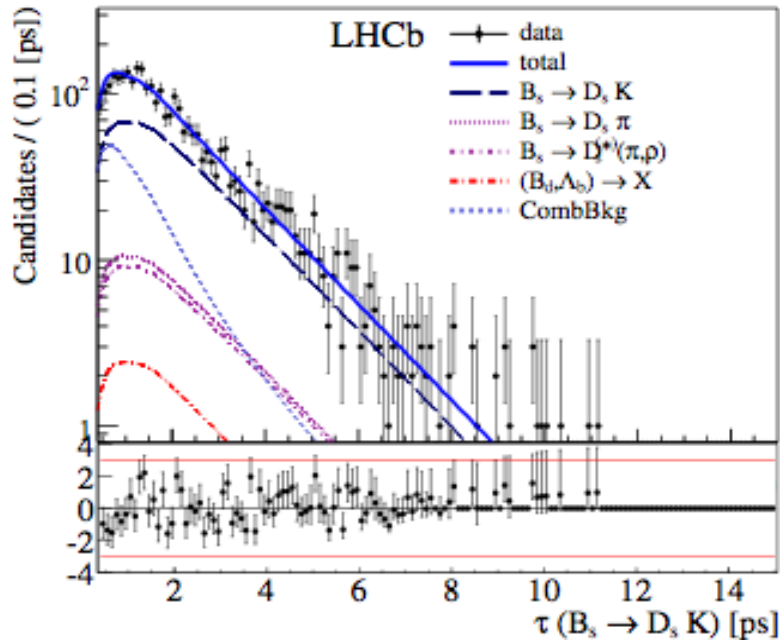
Use the per-event error. Average resolution is 47 fs

External Inputs:

$\Gamma_s, \Delta\Gamma_s, \Gamma_d, \Gamma_{\Delta b}, \Delta m_s$ all fixed from other measurements.



Results



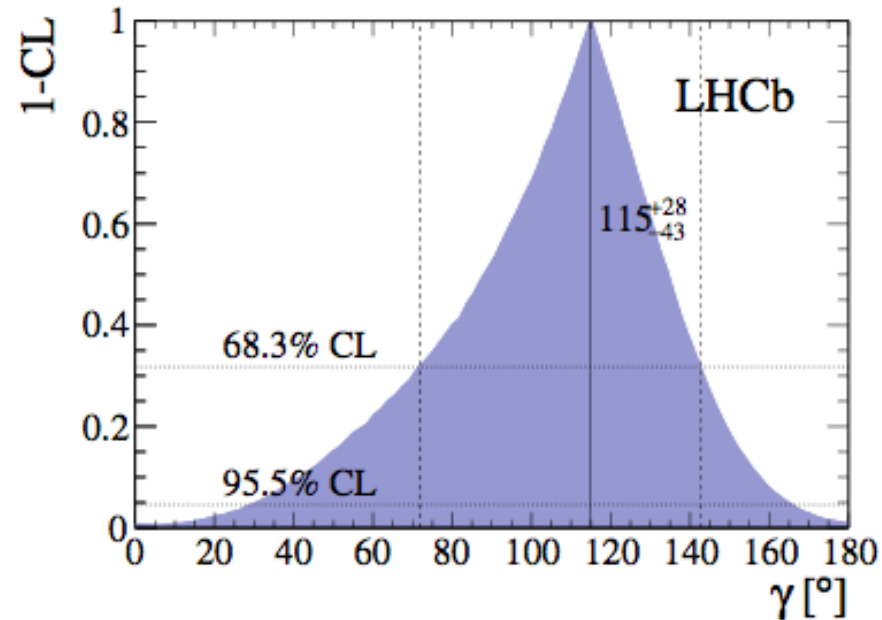
Parameter	<i>cFit</i> fitted value
C_f	$0.53 \pm 0.25 \pm 0.04$
$A_f^{\Delta\Gamma}$	$0.37 \pm 0.42 \pm 0.20$
$A_{\bar{f}}^{\Delta\Gamma}$	$0.20 \pm 0.41 \pm 0.20$
S_f	$-1.09 \pm 0.33 \pm 0.08$
$S_{\bar{f}}$	$-0.36 \pm 0.34 \pm 0.08$

Observables converted to parameters of interest.

β_s external input.

First measurement from $B_s \rightarrow D_s K$

Only 1fb^{-1} - more available



$$\gamma = (115^{+28}_{-43})^\circ$$

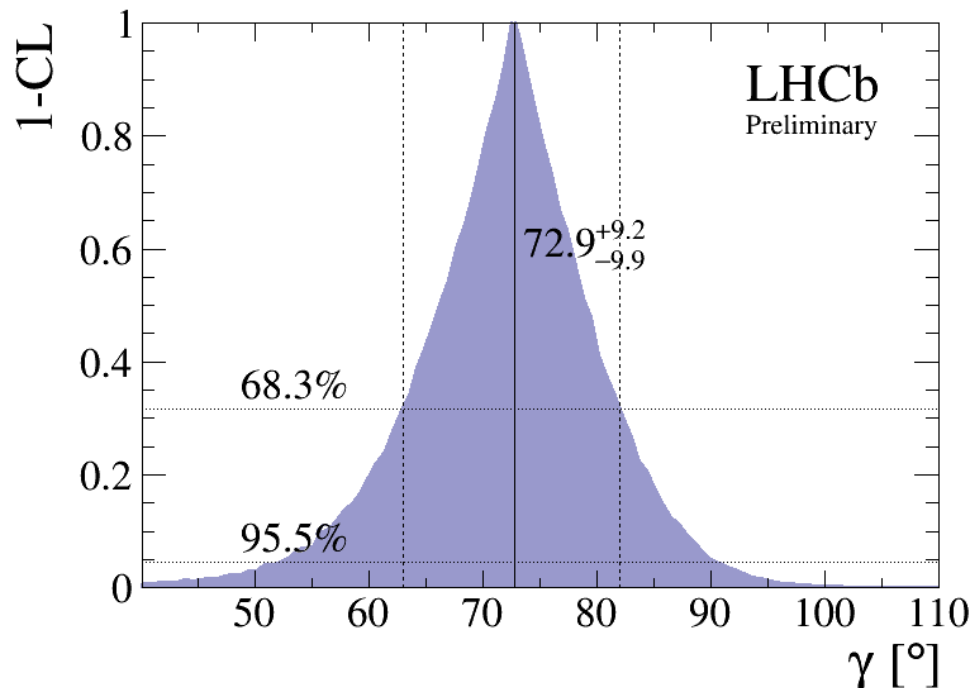
$$r_{D_s K} = (0.53^{+0.17}_{-0.16})$$

$$\delta_{D_s K} = (3^{+19}_{-20})^\circ$$

JHEP 11 (2014) 060

The sum is greater than the parts

$$\gamma = 72.9^{+9.2}_{-9.9}$$



Most observables treated Gaussian
Assume Gaussian systematic fluctuations

Non-physical regions excluded.

External parameters and D mixing taken into account

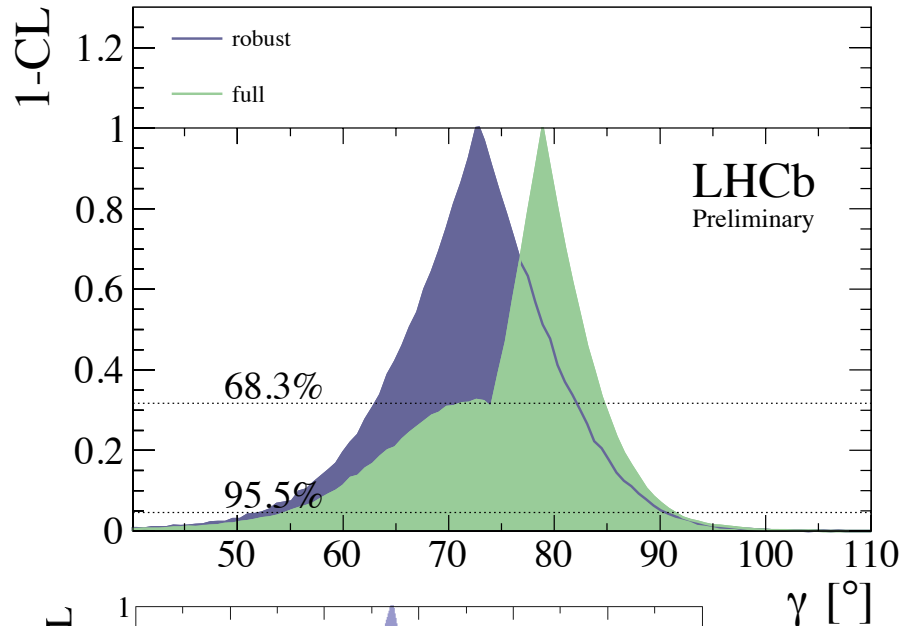
Does not yet include the $D \rightarrow hh^{(\prime)}\pi^0$,
 $B \rightarrow Dh\pi\pi$ results shown earlier.

Plugin method for nominal results –
exclude $B \rightarrow D\pi$ measurements

Bayesian interpretation in good agreement.

LHCb-CONF-2014-004

Including $B \rightarrow D\pi$

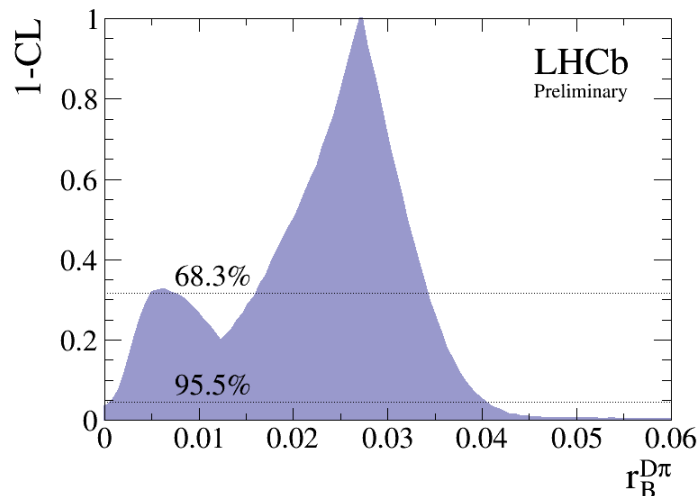


Central value shifts up and peak sharpens when $B \rightarrow D\pi$ is added – ADS+GLW from 1fb^{-1}

Second solution result understandable by looking at the r_B^π values.

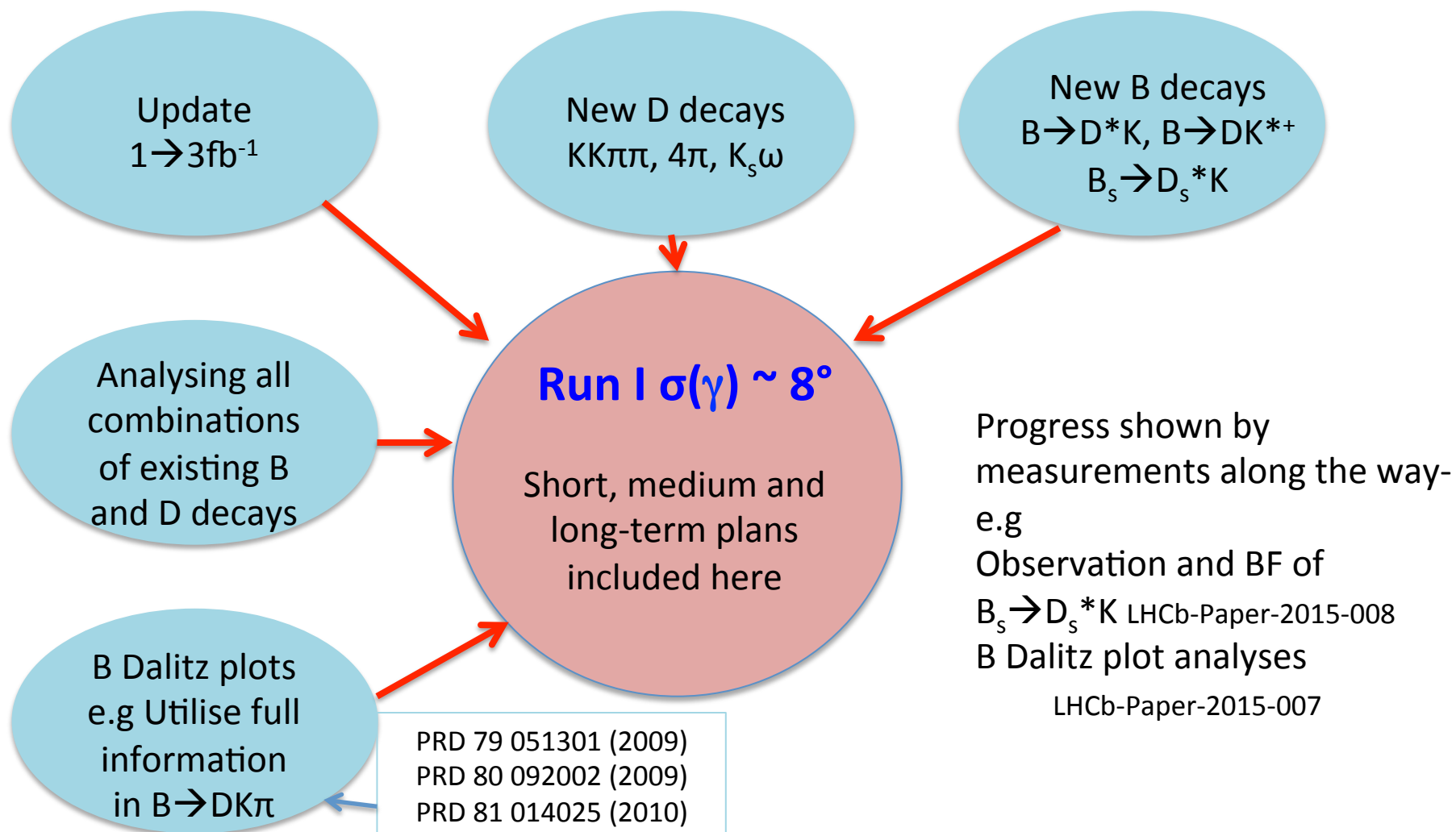
Fit favours an unexpectedly high r_B^π value

More $B \rightarrow D\pi$ results soon – will help resolve this feature.



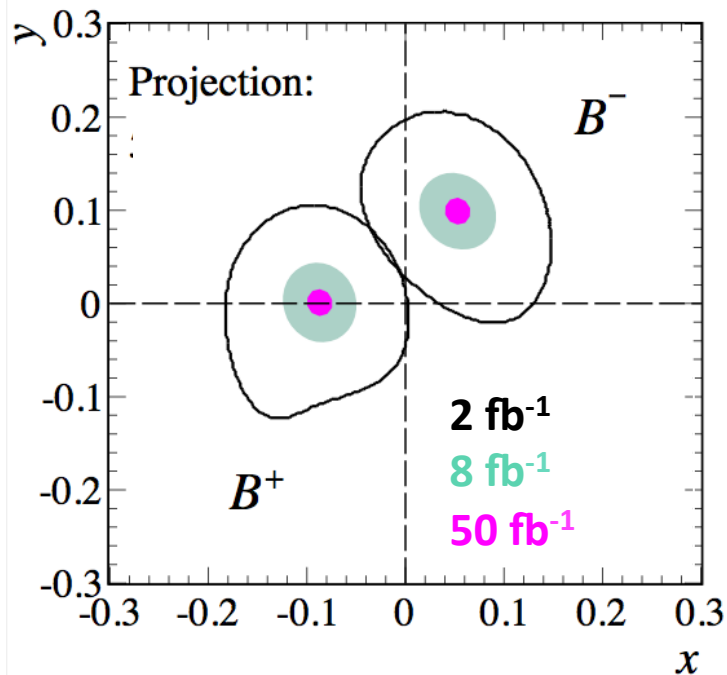
LHCb-CONF-2014-004

Run 1 final say



Prospects beyond Run I

Projections for the GGSZ contours as more data is gathered



Expect to quadruple our dataset in Run II (2015 -2017)

Current γ sensitivity at end of Run I $\sim 8^\circ$

Expected Run II sensitivity 4°

LHCb upgrade will take considerably more data from 2018

Final combined sensitivity $\sim 1^\circ$

Currently statistically limited

Controlling systematic uncertainties is the key

Systematic uncertainties

External inputs

Updated charm inputs from BES-III required

- Large suite of results from CLEO-c all could use improvement in precision

c_i, s_i - Limiting uncertainty of 2°

- Current update when finished would bring uncertainty down to $\sim 1^\circ$

- There are smarter things to be done

- Optimal binning based on experiment background model
- Take experimental efficiency model into account
- Finer binning
- These would **improve the statistical reach** of a given dataset, and remove other smaller systematic uncertainties

Higher order physics effects that can be controlled

D mixing, K_s^0 mixing, CPV and regeneration

Understanding the CPV in backgrounds that start to become significant.

Experimental uncertainties

- How much to trust MC simulation for e.g Dalitz plot efficiencies.
- Controlling the level of detection and production asymmetries.

Summary

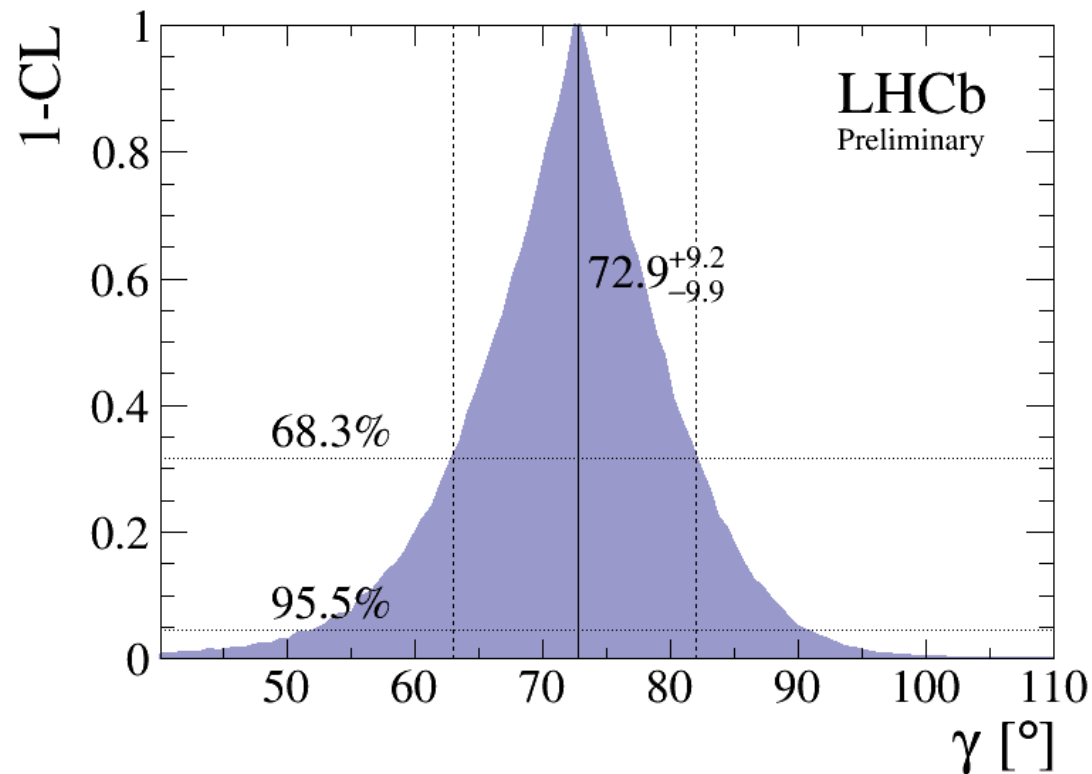
LHCb benefits from high statistics with high purity. Exploring new B decays

Success of analysis with π^0

LHCb can also target unique B_s decay modes

Expected sensitivity by year end $\gamma \sim 8^\circ$

On target to meet Run II projected sensitivity



LHCB-CONF-2014-004