

## Monojetlike Searches for Top Squarks with a $b$ Tag

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The LHC searches for light compressed top squarks have resulted in considerable bounds in the case where the top squark decays to a neutralino and a charm quark. However, in the case where the top squark decays to a neutralino, a bottom quark, and two fermions via an off-shell  $W$  boson, there is currently a significant unconstrained region in the top-squark–neutralino mass plane, still allowing for top squark masses in the range 90–140 GeV. In this Letter we propose a new monojetlike search for light top squarks, optimized for the four-body decay mode, in which at least one  $b$ -tagged jet is required. We show that, by using the existing 8 TeV LHC data set, such a search would cover the entire unconstrained region. Moreover, in the process of validating our tools against an ATLAS monojet search, we show that the existing limit can be extended to exclude also top squark masses below 100 GeV.

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*Introduction.*—The top quark gives rise to the leading quantum correction that destabilizes the electroweak scale in the standard model (SM). One way to solve this so-called hierarchy problem is to extend the ordinary spacetime symmetries by supersymmetry (SUSY) and introduce new physics at a low scale in the form of superpartners of the SM particles. In order to cancel the leading quantum correction, the superpartner of the top quark should have a mass of the order of the electroweak scale and hence would be observable at the Large Hadron Collider (LHC). Light top squarks have been subject to intense recent studies both in theory [1–36] and in the experimental community [37–50].

Taking a simplified model approach, a key strategy to test  $R$ -parity conserving SUSY is to consider only the lightest top squark mass eigenstate  $\tilde{t}_1$ , decaying to the lightest superpartner, the neutralino  $\tilde{\chi}_1^0$ , and taking all other superpartners to be sufficiently heavy and effectively decoupled. In such a simplified model, where the top squark mass  $m_{\tilde{t}_1}$  and the neutralino mass  $m_{\tilde{\chi}_1^0}$  are the only

two parameters, the only relevant SUSY production mode is top squark pair production, for which the cross section is determined by  $m_{\tilde{t}_1}$ .

In the case where the mass difference between the top squark and the neutralino is larger than the top mass,  $\Delta m = m_{\tilde{t}_1} - m_{\tilde{\chi}_1^0} > m_t$ , each of the pair produced top squarks decays in two-body form via an on-shell top quark,  $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ . In this case, for neutralino masses below around 250 GeV, the current LHC limits exclude top squark masses below 600–750 GeV [41,50], while for larger neutralino masses there are no bounds.

For mass splittings in the range  $m_W + m_b < \Delta m < m_t$ , the top squark decays in three-body form via an off-shell top quark,  $\tilde{t}_1 \rightarrow bW\tilde{\chi}_1^0$ , and the top squark mass limits reach up to around 200–300 GeV [38,41,44]. It should be noted that the limits in the different mass splitting regions are not continuously connected to each other and, close to the mass thresholds at  $\Delta m \sim m_t$  and  $\Delta m \sim m_W + m_b$ , the bounds become weak or disappear.

For even smaller mass splittings,  $\Delta m < m_W + m_b$ , the top squark can have two different decay modes, either the four-body decay, via an off-shell  $W$ ,  $\tilde{t}_1 \rightarrow bff'\tilde{\chi}_1^0$ , or the two-body decay  $\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ . While the former decay mode is phase space suppressed, the latter is one-loop and Cabibbo-Kobayashi-Maskawa suppressed; see, e.g., Refs. [20,33,51,52] for discussions. The branching ratios

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(BRs) for these two competing decay modes depend strongly on the flavor structure of the squark soft masses, such as the off-diagonal top-squark–charm squark mixing mass term, as well as the masses of the superpartners that enter the loop. In order to be able to restrict ourselves to the minimal set of parameters, it is customary to consider these two decay modes separately, and we assume a 100% branching ratio in each case.

In the small mass splitting case, several LHC searches have been optimized for the two-body charm decay mode. Under the assumption  $\text{BR}(\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0) = 1$ , top squark masses below 250 to 300 GeV have been excluded [42,45]. In contrast, no existing LHC analysis has been optimized for the four-body top squark decay mode. Currently, even though the four-body case is partly covered by the ATLAS searches [41,42], there exists a significant unconstrained region in the top-squark–neutralino mass plane, still allowing for top squark masses in the range 90–140 GeV (for neutralino masses below 60 GeV). Since very light top squarks could be hiding there, it is important to find a means for probing it. In this Letter we show that, by augmenting an existing monojet search with a  $b$ -tag requirement, this unconstrained region could already be covered completely by the existing 8 TeV LHC data set.

*Existing top squark four-body searches.*—We start by reviewing the existing searches relevant for top squarks in the mass range  $m_b < \Delta m < m_W + m_b$  for which  $\text{BR}(\tilde{t}_1 \rightarrow bff'\tilde{\chi}_1^0) = 1$ . In contrast to the case where the top squark decays via the two-body charm decay mode, which has been probed by Tevatron [53], CMS [45], and ATLAS [42], neither Tevatron nor CMS have performed any search that places a bound in the four-body decay case (Searches for the top squark decay  $\tilde{t}_1 \rightarrow b\ell\tilde{\nu}$  might have some sensitivity to the four-body decay mode that we study. However, since the results are not presented in terms of the four-body top squark decay, we do not include them in our

summary of the existing bounds). Therefore we will focus our discussion on two searches performed by ATLAS.

The first search that places a limit in the top squark four-body decay case is a monojet search in which events are required to contain at least one hard jet and a large amount of missing transverse energy ( $\cancel{E}_T$ ) [42]. The exclusion curve arising from this ATLAS search is indicated by the blue dashed curve in Fig. 1 (left panel). As can be seen in the figure, this search is most sensitive to the case where the mass splitting  $\Delta m$  is small (For  $\Delta m$  smaller than about 20 GeV, the partial width for the top squark four-body decay decreases to the point where the top squark either decays via a displaced vertex or, if other decay channels are present, the four-body branching ratio is strongly suppressed. However, in the spirit of simplified models, we follow the same strategy as ATLAS and present our results for a 100% four-body decay branching ratio, assuming that the top squark will always decay promptly) and the top squark four-body decay products are soft. The required hard jet arises from initial state radiation (ISR), against which the pair-produced top squarks recoil. The ISR jet boosts the two top squarks, which are no longer produced back to back, thereby increasing the  $\cancel{E}_T$  in the event.

The abrupt end of the blue curve at  $m_{\tilde{t}_1} = 100$  GeV in Fig. 1 (left panel) is simply due to the fact that ATLAS does not provide the limit for smaller top squark masses. Given the requirements in this search, one would expect that the exclusion curve should continuously extend diagonally down to the left, reaching the LEP limit [54], which is indicated by the black dashed curve. In the section results below we discuss this issue further and provide the expected limits for top squark masses below 100 GeV.

The second ATLAS search that places a bound in the top squark case under consideration is a search in the final state with one lepton, jets, and  $\cancel{E}_T$  [41]. Since this search relies on the presence of a lepton, arising from the top squark

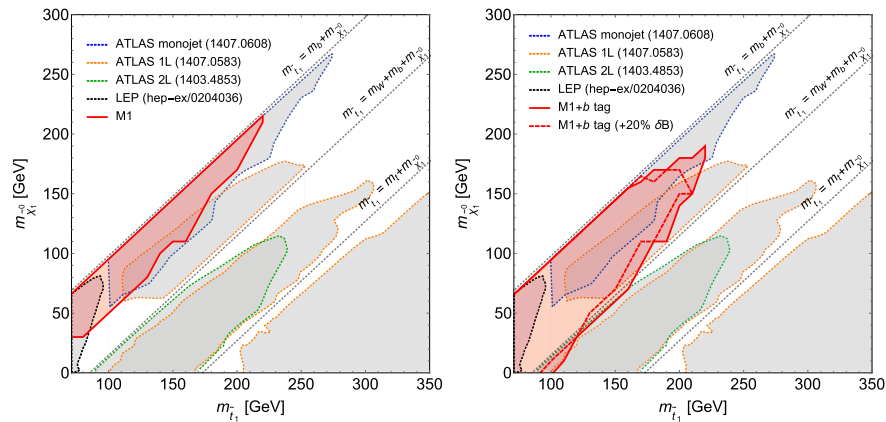


FIG. 1 (color online). (Left panel) Summary of the existing limits in the top-squark–neutralino mass plane, as well as our exclusion curve for the M1 signal region of the ATLAS search [42], in which we also extend the range of excluded top squark masses below 100 GeV. (Right panel) Exclusion arising from our proposed search M1 +  $b$  tag, indicated by the red solid curve. The red dashed curve denotes the change caused by a 20% increase of the total background error.

four-body decay, it is most sensitive to the case where the mass splitting  $\Delta m$  is at least sufficiently large to allow for a reconstruction of the lepton. The exclusion curve arising from this search is indicated by the orange curve in Fig. 1 (left panel). Note that this curve ends at neutralino and top squark masses of around 60 and 110 GeV, respectively. Unlike the monojet search, in which the abrupt ending of the exclusion curve was due to the lack of signal samples, the reason for this ending has a physical origin. As one moves diagonally down to the left, i.e., keeping fixed  $\Delta m$ , the  $\bar{E}_T$  spectrum becomes softer and the search loses sensitivity [55]. Hence, while the cross section increases, as a consequence of the decreasing top squark mass, the acceptance times efficiency decreases faster, as a consequence of the decreasing neutralino mass.

Figure 1 (left panel) summarizes the current experimental status concerning searches for light top squarks that dominantly decay in a four-body final state. We see that there is a triangle shaped unconstrained region for top squark masses in the range 90–140 GeV, with boundaries given by the exclusion curves from LEP [54], the ATLAS monojet search [42], and the ATLAS one-lepton search [41].

*Proposed top squark four-body search.*—The search we propose is a simple extension of the ATLAS inclusive monojet search [42] where, in addition, we require the presence of at least one  $b$ -tagged jet. This  $b$ -tag requirement is motivated by how efficiently it reduces the leading background processes in Ref. [42], with respect to the top squark four-body decay signal process, thus enhancing the sensitivity to light compressed top squarks.

In order to display the potential gain in sensitivity that can be achieved, we consider the addition of a  $b$  tag to the definition of the signal region “M1” of the ATLAS monojet search [42]. The M1 selections require the presence of at most three jets (including  $b$  jets) with  $p_T > 30$  GeV and  $|\eta| < 2.8$ . The leading jet  $p_T$  must be greater than 280 GeV and the  $\bar{E}_T$  in the event must be above 220 GeV. There is a veto on muons with  $p_T > 10$  GeV and  $|\eta| < 2.4$  and on electrons with  $p_T > 20$  GeV and  $|\eta| < 2.47$ . Moreover, there is a condition requiring that the azimuthal angle  $\Delta\phi$  between the  $\bar{E}_T$  vector and each of the jet  $p_T$  vectors should be larger than 0.4. The signal region that selects events according to these cuts is denoted by M1, in accordance with the ATLAS notation.

We remark that the selections that define the signal region M1 might not be optimal over the entire range of mass spectra that we consider. For instance, one could

apply the same argument of adding  $b$ -tagging information to the other signal regions defined in Ref. [42], in which harder  $\bar{E}_T$  cuts are employed. Moreover, as the mass splitting between top squark and neutralino varies, one could optimize further the number of jets that defines the signal region. This is done, for instance, in the dark matter search in Ref. [56], where two signal regions with at least one  $b$  jet and large  $\bar{E}_T$  are employed. However, we find the M1 selections sufficient to make the case for the addition of a  $b$  tag that we propose and hence, in this Letter, we choose to concentrate on this signal region. The benefit of M1 is that it is the signal region in Ref. [42] with the lowest  $\bar{E}_T$  cut and therefore it is the most promising for low mass top squarks. Furthermore this signal region is defined more inclusively on the number of jets than the signal regions of Ref. [56], which allows us to make more reliable predictions for the signal and background, and also to better reproduce the ATLAS results.

The background and signal simulations were performed using MADGRAPH5 [57], PYTHIA6 [58], FASTJET3 [59,60], and DELPHES3 [61] with the ATLAS standard detector specification. Jets (including  $b$  jets) are reconstructed using the anti- $k_r$  clustering algorithm [62] with a jet radius parameter of 0.4. We have used the parton distribution functions (PDF) CTEQ6L1 PDF sets [63] and MLM jet matching [64] throughout the analysis.

The leading SM backgrounds for our analysis are the  $t\bar{t}$ ,  $Z(\rightarrow \nu\nu) + \text{jets}$ ,  $W(\rightarrow \ell\nu) + \text{jets}$  (where  $\ell = e, \mu, \tau$ ), and diboson processes. For the normalization of the cross sections of these processes at 8 TeV LHC, we have used the same theoretical predictions as ATLAS in Ref. [42], obtained from Refs. [65–71]. We have validated our simulations against the M1 signal region of the ATLAS monojet search. Comparing our estimates for the background rates to the corresponding ATLAS estimates in Table VIII of Ref. [42], we find that the central values from our simulations are within 20% of the corresponding ATLAS numbers. This agreement is quite satisfactory and makes us confident that reliable conclusions can be drawn using our tools. To further improve our estimates, we normalize our predictions for all of the leading background processes to exactly match the ATLAS background estimates, which, for the  $Z + \text{jets}$  and  $W + \text{jets}$  samples, have the extra advantage of incorporating data-driven reweighting of the simulated events. The expected backgrounds are summarized in Table I, written as  $B \pm \delta B$ , where  $B$  is the central value and  $\delta B$  the  $1\sigma$  error. For all of

TABLE I. Estimated numbers of background events with  $20.3 \text{ fb}^{-1}$  of 8 TeV LHC data.

Background	$t\bar{t}$	$Z(\rightarrow \nu\nu)$	$W(\rightarrow \ell\nu)$	Dibosons	Others	Total
M1 (ATLAS [42])	$780 \pm 73$	$17400 \pm 720$	$14100 \pm 337$	$650 \pm 99$	$565 \pm 301$	$33450 \pm 960$
M1 + $b$ tag	$307 \pm 57$	$261 \pm 22$	$144 \pm 7$	$55 \pm 17$	...	$767 \pm 64$

the backgrounds shown in Table I, we quote the  $1\sigma$  errors from ATLAS.

The crucial ingredient in our analysis is that we extend the M1 set of selections by adding the requirement that the events passing the M1 cuts must contain at least one  $b$ -tagged jet with  $p_T$  in the range 30–300 GeV, and with  $|\eta| < 2.5$ . The  $b$  tagging is parametrized through DELPHES3 [61] using a “mild” working point characterized by a light flavor jet rejection of 1/1000 and an efficiency for actual  $b$  quarks of around 0.4 for central jets. We emphasize that for the  $p_T$  range that we have chosen for the  $b$ -tagged jet, the calibration of the  $b$ -tagging algorithms in the ATLAS experiment is data driven, minimizing the systematic uncertainties on  $b$  tagging coming from Monte Carlo simulations [72]. We denote this signal region by “M1 +  $b$  tag.”

In the M1 +  $b$  tag row in Table I, we show the number of background events we get in the proposed signal region. The comparison of the events expected for the M1 and M1 +  $b$ -tag signal regions gives us an idea of the power of the  $b$ -tag requirement to enhance the sensitivity to the signal. In fact, the top squark signal is expected to behave similarly to the  $t\bar{t}$  background—hence to be only mildly reduced by the  $b$ -tag requirement. Typical signal efficiencies to the  $b$ -tag requirement are around 20%. In contrast, we observe a great suppression of the  $W$  and  $Z$  boson backgrounds, which are the leading backgrounds in the M1 signal region. We stress that, in our work, the efficiency of signal and backgrounds to the  $b$ -tag requirement is given by the parametrization implemented in DELPHES3 [61], which is expected to be reliable. In order to draw conclusions on safe grounds, we take the relative error of the M1 +  $b$ -tag background prediction to be *twice* the relative error quoted by ATLAS for the M1 region. This gives background relative errors slightly smaller than those quoted by ATLAS [42] in the signal regions involving  $c$  tags. Given that  $b$  tags are under better experimental control than  $c$  tags, we expect our estimate of the uncertainties to be fair.

The signal process has been simulated for a grid of points with  $m_{\tilde{t}_1}$  varying from 70 to 250 GeV, and  $m_{\tilde{\chi}_1^0}$  varying from 0 to 200 GeV, in steps of 10 GeV. The only relevant SUSY production mode is top squark pair production, with a cross section given in Table II. For top squark masses in the range 100–250 GeV we have used the next-to-leading-order (NLO) plus next-to-leading-logarithm top squark cross sections used by ATLAS and given by the LHC SUSY Cross Section Working Group [73], to which an uncertainty of around 16% is assigned. For the top squark mass points below 100 GeV, which are not given by ATLAS, we

computed the NLO cross sections using PROSPINO [74] and normalized them with the available ATLAS cross sections for top squark masses above 100 GeV. We compared the expected number of signal events we obtained to the corresponding numbers reported by ATLAS [42] in several points and we found a systematic overshoot of around 20%. Therefore we normalized the expected number of signal events by decreasing it by 20% to match the ATLAS central values.

The efficiencies in the analysis are rather small, requiring the generation of a large number of Monte Carlo events. For the backgrounds in the M1 case we are able to generate a sufficient amount of fully jet-matched events to keep the statistical error below 10%. However, for the signal, given the amount of points in the grid we aim to cover, we do not have the computer resources to generate fully matched events at a similar level of statistical uncertainty. We solve this problem by performing the analysis in two steps. Since all events that pass the cuts of our signal regions contain a hard jet, for each point of the grid, we generated both the exclusive zero-jet leading-order top squark pair production process  $pp \rightarrow \tilde{t}_1\tilde{t}_1$  and the one-jet process  $pp \rightarrow \tilde{t}_1\tilde{t}_1j$ , with jet  $p_T > 200$  GeV. The ratio of these two cross sections is used to calculate the efficiency of the  $p_T > 200$  GeV cut. The unmatched one-jet sample was then used for the grid in the analysis. We have checked the validity of this procedure in a few points by generating fully matched inclusive samples, with sufficient statistics, and we found the two procedures to be in agreement within one statistical standard deviation.

*Results.*—In order to estimate whether a point in the top-squark–neutralino mass plane is excluded, we compute the number of expected events for that point in a given signal region. We expect that the experiments can put a 95% C.L. exclusion for those mass points that yield a number of signal events greater than  $N_{95} = 1.96\delta B$ , where  $\delta B$  is the total error in Table I.

Following this limit-setting procedure, we start by calculating the exclusion we obtain using the M1 signal region. The resulting exclusion curve is given by the red solid curve in the left panel of Fig. 1. We remark that our exclusion curve follows quite closely the one given by ATLAS, which further validates our procedure. It should be noted that our curve extends down to the LEP bound for  $m_{\tilde{\chi}_1^0} > m_{\tilde{t}_1} - 40$  GeV, thereby covering part of the unconstrained region between the blue ATLAS curve and the black LEP curve. Hence, by considering top squark masses below 100 GeV, the existing ATLAS bound arising from the M1 signal region can be extended.

TABLE II. Inclusive top squark pair production cross sections used in the analysis.

$m_{\tilde{t}_1}$ (GeV)	70	80	90	100	110	120	130	140	150	160	170	180	190	200	210	220	230	240	250
$\sigma(\tilde{t}\tilde{t})$ (pb)	2797	1550	912	560	362	240	163	113	80.3	58.0	42.6	31.9	24.2	18.5	14.3	11.2	8.78	6.97	5.58



Following the same limit-setting procedure for the signal region  $M1 + b$  tag, we obtain the red solid exclusion curve in the right panel of Fig. 1. This is the main result of this Letter. The dashed red curve corresponds to a 20% increase of the total error on the background. The comparison of the two red curves gives an idea of the sensitivity of our result to (a) the uncertainty associated with the signal cross sections, (b) possible contributions from subleading backgrounds not evaluated for  $M1 + b$  tag. We see that our proposed search  $M1 + b$  tag covers the entire unconstrained region. Moreover, it slightly extends the existing LHC limits for a top squark mass of around 200 GeV.

It is worth mentioning that our simulations suggest that further sensitivity is gained by removing the  $\Delta\phi$  condition, but keeping the same  $b$ -tag requirement as in  $M1 + b$  tag. In the  $M1$  signal region, the  $\Delta\phi$  condition is introduced to reduce the pure-QCD multijet background, for which the  $\cancel{E}_T$  originates from jet mismeasurements. However, the  $b$ -tag requirement can be seen as an alternative to the  $\Delta\phi$  cut since it is expected, already by itself, to dramatically reduce the multijet background. With our simulation tools, the estimation of the multijet background would not be reliable; therefore, we do not attempt to estimate the gain in sensitivity to light top squarks that could be achieved by employing a looser  $\Delta\phi$  cut. Instead we content ourselves with simply encouraging the experimental collaborations to also consider a signal region with a looser  $\Delta\phi$  cut.

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[1] T. Plehn, M. Spannowsky, M. Takeuchi, and D. Zerwas, *J. High Energy Phys.* **10** (2010) 078.  
 [2] S. Bornhauser, M. Drees, S. Grab, and J. S. Kim, *Phys. Rev. D* **83**, 035008 (2011).  
 [3] C. Brust, A. Katz, S. Lawrence, and R. Sundrum, *J. High Energy Phys.* **03** (2012) 103.  
 [4] Y. Kats and D. Shih, *J. High Energy Phys.* **08** (2011) 049.  
 [5] X.-J. Bi, Q.-S. Yan, and P.-F. Yin, *Phys. Rev. D* **85**, 035005 (2012).

[6] B. He, T. Li, and Q. Shafi, *J. High Energy Phys.* **05** (2012) 148.  
 [7] M. Drees, M. Hanussek, and J. S. Kim, *Phys. Rev. D* **86**, 035024 (2012).  
 [8] Y. Bai, H.-C. Cheng, J. Gallicchio, and J. Gu, *J. High Energy Phys.* **07** (2012) 110.  
 [9] T. Plehn, M. Spannowsky, and M. Takeuchi, *J. High Energy Phys.* **08** (2012) 091.  
 [10] D. S. M. Alves, M. R. Buckley, P. J. Fox, J. D. Lykken, and C.-T. Yu, *Phys. Rev. D* **87**, 035016 (2013).  
 [11] Z. Han, A. Katz, D. Krohn, and M. Reece, *J. High Energy Phys.* **08** (2012) 083.  
 [12] D. E. Kaplan, K. Rehermann, and D. Stolarski, *J. High Energy Phys.* **07** (2012) 119.  
 [13] C. Brust, A. Katz, and R. Sundrum, *J. High Energy Phys.* **08** (2012) 059.  
 [14] K. Ghosh, K. Huitu, J. Laamanen, and L. Leinonen, *Phys. Rev. Lett.* **110**, 141801 (2013).  
 [15] J. A. Evans and Y. Kats, *J. High Energy Phys.* **04** (2013) 028.  
 [16] C. Kilic and B. Tweedie, *J. High Energy Phys.* **04** (2013) 110.  
 [17] M. L. Graesser and J. Shelton, *Phys. Rev. Lett.* **111**, 121802 (2013).  
 [18] K. Krizka, A. Kumar, and D. E. Morrissey, *Phys. Rev. D* **87**, 095016 (2013).  
 [19] R. Franceschini and R. Torre, *Eur. Phys. J. C* **73**, 2422 (2013).  
 [20] A. Delgado, G. F. Giudice, G. Isidori, M. Pierini, and A. Strumia, *Eur. Phys. J. C* **73**, 2370 (2013).  
 [21] B. Dutta, T. Kamon, N. Kolev, K. Sinha, K. Wang, and S. Wu, *Phys. Rev. D* **87**, 095007 (2013).  
 [22] M. R. Buckley, T. Plehn, and M. Takeuchi, *J. High Energy Phys.* **08** (2013) 086.  
 [23] A. Chakraborty, D. K. Ghosh, D. Ghosh, and D. Sengupta, *J. High Energy Phys.* **10** (2013) 122.  
 [24] I. Low, *Phys. Rev. D* **88**, 095018 (2013).  
 [25] Y. Bai, H.-C. Cheng, J. Gallicchio, and J. Gu, *J. High Energy Phys.* **08** (2013) 085.  
 [26] G. Belanger, D. Ghosh, R. Godbole, M. Guchait, and D. Sengupta, *Phys. Rev. D* **89**, 015003 (2014).  
 [27] R. Boughezal and M. Schulze, *Phys. Rev. D* **88**, 114002 (2013).  
 [28] Z. Han and A. Katz, [arXiv:1310.0356](https://arxiv.org/abs/1310.0356).  
 [29] B. Dutta, W. Flanagan, A. Gurrola, W. Johns, T. Kamon, P. Sheldon, K. Sinha, K. Wang, and S. Wu, *Phys. Rev. D* **90**, 095022 (2014).  
 [30] M. Papucci, K. Sakurai, A. Weiler, and L. Zeune, *Eur. Phys. J. C* **74**, 3163 (2014).  
 [31] M. R. Buckley, T. Plehn, and M. J. Ramsey-Musolf, *Phys. Rev. D* **90**, 014046 (2014).  
 [32] M. Czakon, A. Mitov, M. Papucci, J. T. Ruderman, and A. Weiler, *Phys. Rev. Lett.* **113**, 201803 (2014).  
 [33] R. Grober, M. Mühlleitner, E. Popenza, and A. Wlotzka, [arXiv:1408.4662](https://arxiv.org/abs/1408.4662).  
 [34] T. Eifert and B. Nachman, *Phys. Lett. B* **743**, 218 (2015).  
 [35] W. S. Cho, J. S. Gainer, D. Kim, K. T. Matchev, F. Moortgat, L. Pape, and M. Park, [arXiv:1411.0664](https://arxiv.org/abs/1411.0664).  
 [36] A. Choudhury and A. Datta, *Mod. Phys. Lett. A* **27**, 1250188 (2012).

- [37] G. Aad *et al.* (ATLAS Collaboration), *J. High Energy Phys.* **10** (2013) 189.
- [38] G. Aad *et al.* (ATLAS Collaboration), *J. High Energy Phys.* **06** (2014) 124.
- [39] G. Aad *et al.* (ATLAS Collaboration), *Eur. Phys. J. C* **74**, 2883 (2014).
- [40] G. Aad *et al.* (ATLAS Collaboration), *J. High Energy Phys.* **09** (2014) 015.
- [41] G. Aad *et al.* (ATLAS Collaboration), *J. High Energy Phys.* **11** (2014) 118.
- [42] G. Aad *et al.* (ATLAS Collaboration), *Phys. Rev. D* **90**, 052008 (2014).
- [43] S. Chatrchyan *et al.* (CMS Collaboration), *Phys. Rev. Lett.* **111**, 221801 (2013).
- [44] S. Chatrchyan *et al.* (CMS Collaboration), *Eur. Phys. J. C* **73**, 2677 (2013).
- [45] S. Chatrchyan *et al.* (CMS Collaboration), CMS Report No. CMS-PAS-SUS-13-009, 2013.
- [46] S. Chatrchyan *et al.* (CMS Collaboration), *Phys. Rev. Lett.* **112**, 161802 (2014).
- [47] V. Khachatryan *et al.* (CMS Collaboration), *Phys. Lett. B* **736**, 371 (2014).
- [48] S. Chatrchyan *et al.* (CMS Collaboration), *Phys. Lett. B* **743**, 503 (2015).
- [49] CMS Collaboration, [arXiv:1412.7706](https://arxiv.org/abs/1412.7706).
- [50] S. Chatrchyan *et al.* (CMS Collaboration), CMS Report No. CMS-PAS-SUS-14-011, 2014.
- [51] C. Boehm, A. Djouadi, and Y. Mambrini, *Phys. Rev. D* **61**, 095006 (2000).
- [52] S. P. Das, A. Datta, and M. Guchait, *Phys. Rev. D* **65**, 095006 (2002).
- [53] T. Aaltonen *et al.* (CDF Collaboration), *J. High Energy Phys.* **10** (2012) 158.
- [54] A. Heister *et al.* (ALEPH Collaboration), *Phys. Lett. B* **537**, 5 (2002).
- [55] S. Strandberg (private communication).
- [56] G. Aad *et al.* (ATLAS Collaboration), *Eur. Phys. J. C* **75**, 92 (2015).
- [57] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, H.-S. Shao, and T. Stelzer *et al.*, *J. High Energy Phys.* **07** (2014) 079.
- [58] T. Sjöstrand, S. Mrenna, and P. Z. Skands, *J. High Energy Phys.* **05** (2006) 026.
- [59] M. Cacciari, G. P. Salam, and G. Soyez, *Eur. Phys. J. C* **72**, 1896 (2012).
- [60] M. Cacciari and G. P. Salam, *Phys. Lett. B* **641**, 57 (2006).
- [61] J. de Favereau, C. Delaere, P. Demin, A. Giammanco, V. Lemaitre, A. Mertens, and M. Selvaggi, *J. High Energy Phys.* **02** (2014) 057.
- [62] M. Cacciari, G. P. Salam, and G. Soyez, *J. High Energy Phys.* **04** (2008) 063.
- [63] J. Pumplin, D. R. Stump, J. Huston, H. L. Lai, P. M. Nadolsky, and W. K. Tung, *J. High Energy Phys.* **07** (2002) 012.
- [64] M. L. Mangano, M. Moretti, F. Piccinini, and M. Treccani, *J. High Energy Phys.* **01** (2007) 013.
- [65] M. Czakon (ATLAS and CMS Collaborations), [https://twiki.cern.ch/twiki/bin/view/LHCPhysics/TbarNNLO#Top\\_quark\\_pair\\_cross\\_sections\\_at](https://twiki.cern.ch/twiki/bin/view/LHCPhysics/TbarNNLO#Top_quark_pair_cross_sections_at).
- [66] M. Czakon and A. Mitov, *Comput. Phys. Commun.* **185**, 2930 (2014).
- [67] M. Czakon, P. Fiedler, and A. Mitov, *Phys. Rev. Lett.* **110**, 252004 (2013).
- [68] J. M. Campbell, R. K. Ellis, and C. Williams, MCFM—Monte Carlo for FeMtobarn processes, <http://mcfm.fnal.gov/>.
- [69] J. M. Campbell, R. K. Ellis, and C. Williams, *J. High Energy Phys.* **07** (2011) 018.
- [70] S. Catani and M. Grazzini, *Phys. Rev. Lett.* **98**, 222002 (2007).
- [71] S. Catani, L. Cieri, G. Ferrera, D. de Florian, and M. Grazzini, *Phys. Rev. Lett.* **103**, 082001 (2009).
- [72] A. Coccaro (private communication).
- [73] S. Padhi (LHC SUSY Cross Section Working Group), <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections8TeVstoppingbottom>.
- [74] W. Beenakker, M. Krämer, T. Plehn, M. Spira, and P. M. Zerwas, *Nucl. Phys.* **B515**, 3 (1998).