

MULTIPLICITY DEPENDENCE OF TRANSVERSE MOMENTUM SPECTRA

IN  $pp$ ,  $\bar{p}p$ ,  $dd$  AND  $\alpha\alpha$  COLLISIONS AT ISR ENERGIES

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ABSTRACT

We analyze the variation of the average transverse momentum,  $\langle p_T \rangle$ , with the multiplicity of charged particles produced in  $pp$ ,  $\bar{p}p$ ,  $dd$  and  $\alpha\alpha$  collisions at ISR energies. An increase of  $\langle p_T \rangle$  with increasing particle density  $\rho = \Delta n/\Delta y$  for charged particles produced in the central region is observed. The energy dependence of this effect and its dependence on the type of colliding particles are discussed.

## 1. INTRODUCTION

The experimental results obtained in the last decade at the CERN ISR and at lower energies showed that the transverse momenta,  $p_T$ , of particles produced in hadron-hadron collisions are concentrated at low values, and that, at low  $p_T$ , the  $p_T$  distributions depend only weakly on energy and on the type of incident particles. Recent experiments at the CERN proton-antiproton collider at a centre of mass energy  $\sqrt{s} = 540$  GeV [1] and at the highest ISR energy,  $\sqrt{s} = 62$  GeV [2] have shown that there is an increase of the mean transverse momentum  $\langle p_T \rangle$  with the number of charged particles produced per unit rapidity,  $\rho = \Delta n / \Delta y$ , in the central region of rapidity and that the increase becomes larger as  $\sqrt{s}$  increases. A number of possible explanations for this observation have been proposed [3-5] in the context of geometrical models, of thermodynamical models, of mini-jets (semi hard scattering), etc. In thermodynamical models the average  $p_T$  is a measure of the temperature, while the central particle density is a measure of the entropy density. The study of  $\langle p_T \rangle$  versus  $\rho$  may yield information on possible phase transitions. In geometrical models high multiplicity events occur preferentially at small impact parameters and correspond to higher temperatures and thus have larger average transverse momentum. Due to lack of phase space, the increase of  $\langle p_T \rangle$  should slow down at a certain value of  $\rho$ , dependent on energy. The high  $p_T$  region is essentially thought to arise from constituent scattering, which should yield particles in a characteristic jet structure and higher multiplicities. Thus at higher  $p_T$  we should have higher multiplicities. As the energy increases the proportion of hard scattered events increases and hence the slope of the dependence of  $\langle p_T \rangle$  on  $\rho$  should increase. The full comprehension of the phenomenon requires precise data at different energies, comparison of different types of collisions as well as analyses of related phenomena.

In this paper we present the results of an analysis of charged particle production in proton-proton, antiproton-proton, deuteron-deuteron and alpha-alpha interactions at nucleon-nucleon centre of mass energies of  $\sqrt{s_{NN}} = 62$  and 31 GeV. Preliminary results on this subject were already presented [6].

## 2. EXPERIMENTAL

The experiment was performed at the CERN Intersecting Storage Rings (ISR) using the Split Field Magnet (SFM) detector. The magnetic volume of the detector (with a maximum field strength of 1 Tesla) was filled with Multiwire Proportional Chambers. The momenta of charged particles were measured in nearly the full solid angle. The performance of the detector is described in previous publications [7]. Table 1 gives the details of the available data. Note that the statistics for pp data are about an order of magnitude larger than for the other data.

This experiment used a "minimum bias" trigger, which essentially required the presence of at least one candidate charged track in the detector. This trigger accepted  $\sim 95\%$  of the inelastic cross section. The events were reconstructed with the standard SFM chain of computer programs (an older version was used for the  $\bar{p}p$  and some pp data). A good vertex fit, defined by a minimum of two charged tracks, was required for each event. This reduced the acceptance to about 90% of the inelastic cross section.

A number of selection criteria were applied to the data:

- (a) The interaction vertex was required to be in the overlap region of the two colliding beams. This condition is particularly important for  $\bar{p}p$  interactions, where backgrounds from beam-gas and beam-wall interactions were high. Only 15,000  $\bar{p}p$  interactions were obtained with the minimum bias trigger discussed above; these events came from about one million triggers; the other 68,000 interactions were obtained with a quasi-minimum bias trigger, which required at least one forward and one backward track. It was checked that the two triggers gave the same results after the requirement of a good vertex fit.
- (b) Tracks were accepted if the relative precision in the measurement of their momenta was  $\Delta p/p < 0.8$ .
- (c) In order to reduce further some background from mismeasured tracks, only particles with  $p_T < 2.5$  GeV/c were used.

Furthermore, events with less than four tracks are not considered for this analysis, in order to reduce the contribution of diffractive processes.

### 3. ANALYSIS

Fig. 1 shows the uncorrected transverse momentum distributions for events with low (4-6) and high (> 12) multiplicities of charged particles produced in the central rapidity region ( $|y| < 1.5$ ) for  $\alpha\alpha$  and  $dd$  interactions at a centre of mass energy  $\sqrt{s_{NN}} = 31$  GeV. The data follow an approximate exponential behaviour for  $0.2 < p_T < 0.8$  GeV/c, while at larger values of  $p_T$  the distributions flatten out; there are substantial losses for  $p_T < 0.15$  GeV/c [8]. It is apparent that the curves at lower multiplicities are steeper than the ones at higher multiplicities. Moreover, the largest differences involve medium and high  $p_T$  values.

The average transverse momentum,  $\langle p_T \rangle$ , was evaluated for the  $\Delta n$  charged tracks remaining after the previously mentioned cuts ( $p_T < 2.5$  GeV/c and  $\Delta p/p < 0.8$ ), in the rapidity interval  $-1.5 < y < 1.5$ . Furthermore, we applied a low  $p_T$  cut ( $p_T > 0.15$  GeV/c) in order to reduce the effects of acceptance losses. The quantity  $\langle p_T \rangle$  may be computed using two methods.

In the first method we perform a fit of the measured  $p_T$  distributions in the range  $0.3 < p_T < 0.8$  GeV/c to an exponential form

$$\frac{dn}{dp_T^2} = A e^{-bp_T} \quad . \quad (1)$$

The average transverse momentum is then given by

$$\langle p_T \rangle = 2/b \quad . \quad (2)$$

In the second method one computes algebraically the mean values of  $p_T$ . The value of  $\langle p_T \rangle$  corresponding to a given  $\Delta n$  ( $\Delta n$  is the number of central particles having  $p_T < 2.5$  GeV/c and  $\Delta p/p < 0.8$ ) is then calculated using only those particles that, among the  $\Delta n$ , have  $p_T > 0.15$  GeV/c. This gives a lower statistical error, but yields values of  $\langle p_T \rangle$  slightly larger than (2) because of losses of tracks at small  $p_T$  (fig. 1). This second method is nevertheless useful for the analysis of the variations of  $\langle p_T \rangle$  versus

particle density  $\rho = \Delta n/\Delta y$  and for the comparison of the average transverse momenta for particles produced in different initial states. The two methods are not exactly equivalent, since the second takes into account contributions from the high  $p_T$  tail.

#### 4. RESULTS

Fig. 2 shows the average transverse momenta of charged particles produced in the central region of rapidity ( $|y| < 1.5$ ) versus the observed charged multiplicity density in the same central region. The average transverse momentum was computed algebraically in the range  $0.15 < p_T < 2.5$  GeV/c. Therefore  $\langle p_T \rangle$  is overestimated and the absolute  $\langle p_T \rangle$  scale is affected by a systematic shift of about 20% for all multiplicities. In fact, we estimate that the average  $p_T$  should be:

- (a) reduced by about 18% because of the cut at  $p_T = 0.15$  GeV/c,
- (b) increased by less than 1% because of the cut at 2.5 GeV/c and
- (c) decreased by 2% because of some background in the interval  $2 < p_T < 2.5$  GeV/c. We have also estimated that acceptance losses in the central region reduce the average multiplicity by about 30%.

All the data in fig. 2 exhibit an increase of the average transverse momentum with increasing charged particle density  $\rho$ . The increase is, within errors, equal for pp and  $\bar{p}p$  (the slight difference between the absolute values is due to the different versions of the reconstruction program); the pp data at  $\sqrt{s} = 31$  GeV have a smaller rate of increase than at 62 GeV; at 31 GeV the  $\alpha\alpha$  data [9] have a larger increase than dd and pp. The increase could be non-linear for pp and  $\bar{p}p$ , for which the dependence seems to weaken for  $\rho > 4-5$  [10].

The  $p_T$  distributions in fig. 1 indicate that the dependence of the average transverse momentum on charged multiplicity density should be enhanced if one uses data with  $p_T > 0.3$  GeV/c. This is indeed so, as shown in fig. 3.

The average transverse momenta computed on the basis of exponential fits to the  $p_T$  distributions in the region  $0.3 < p_T < 0.8$  GeV/c with  $|y| < 2$  are shown in fig. 4 for pp and  $\bar{p}p$  interactions at  $\sqrt{s} = 62$  GeV.

The pp sample (240,000 events) was reconstructed with exactly the same version of the SFM chain of programs as that used for  $\bar{p}p$ . The data of fig. 4 show that  $\bar{p}p$  and pp at  $\sqrt{s} = 62$  GeV are equal within errors. Notice also that the absolute value has a smaller systematic error compared to that shown in fig. 2.

In conclusion, we have observed an increase of the average transverse momentum with charged multiplicity density in the central rapidity region at ISR energies for pp,  $\bar{p}p$ , dd and  $\alpha\alpha$  collisions. The increase is the same for pp and  $\bar{p}p$  collisions. At  $\sqrt{s} = 31$  GeV the rate of increase for  $\alpha\alpha$  collisions is larger than that for dd, which is comparable with that for pp. The amount of increase for pp collisions at ISR energies rises with c.m. energy. The magnitude of the effect at the ISR is considerably smaller than at the CERN  $\bar{p}p$  collider, where there is still a sizeable energy dependence [11].

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TABLE 1

Details of the event samples used for the present analysis. The centre of mass energy,  $\sqrt{s_{NN}}$ , is the nucleon-nucleon c.m. energy. The last column gives the nominal magnetic field strength in the SFM detector.

Interaction	$\sqrt{s_{NN}}$ (GeV)	Number of events ( $\times 10^3$ )	Magnetic field (Tesla)
pp	62	816	1
$p\bar{p}$	62	83	1
pp	31	1400	0.5
$\alpha\alpha$	31	74	1
dd	31	117	1

FIGURE CAPTIONS

- Fig. 1 Transverse momentum distributions for low (4-6) and high (> 12) multiplicities of charged particles produced in the central region of rapidity,  $|y| < 1.5$ , of  $\alpha\alpha$  and  $dd$  collisions at  $\sqrt{s_{NN}} = 31$  GeV.
- Fig. 2 The average transverse momentum of charged particles (with  $0.15 < p_T < 2.5$  GeV/c) produced in the central region of rapidity ( $|y| < 1.5$ ) plotted versus charged particle multiplicity density. The graph shows data for  $pp$  and  $p\bar{p}$  interactions at  $\sqrt{s} = 62$  GeV and for  $pp$ ,  $dd$  and  $\alpha\alpha$  interactions at  $\sqrt{s} = 31$  GeV per nucleon. The cuts applied to the data are described in the text.
- Fig. 3 Data as in fig. 2, but using the  $p_T$  range  $0.3 < p_T < 2.5$  GeV/c.
- Fig. 4 The average transverse momentum for charged particles produced in the central region ( $|y| < 2$ ) for  $pp$  and  $p\bar{p}$  interactions at  $\sqrt{s} = 62$  GeV computed using a parametrization of the form of eq. (1).

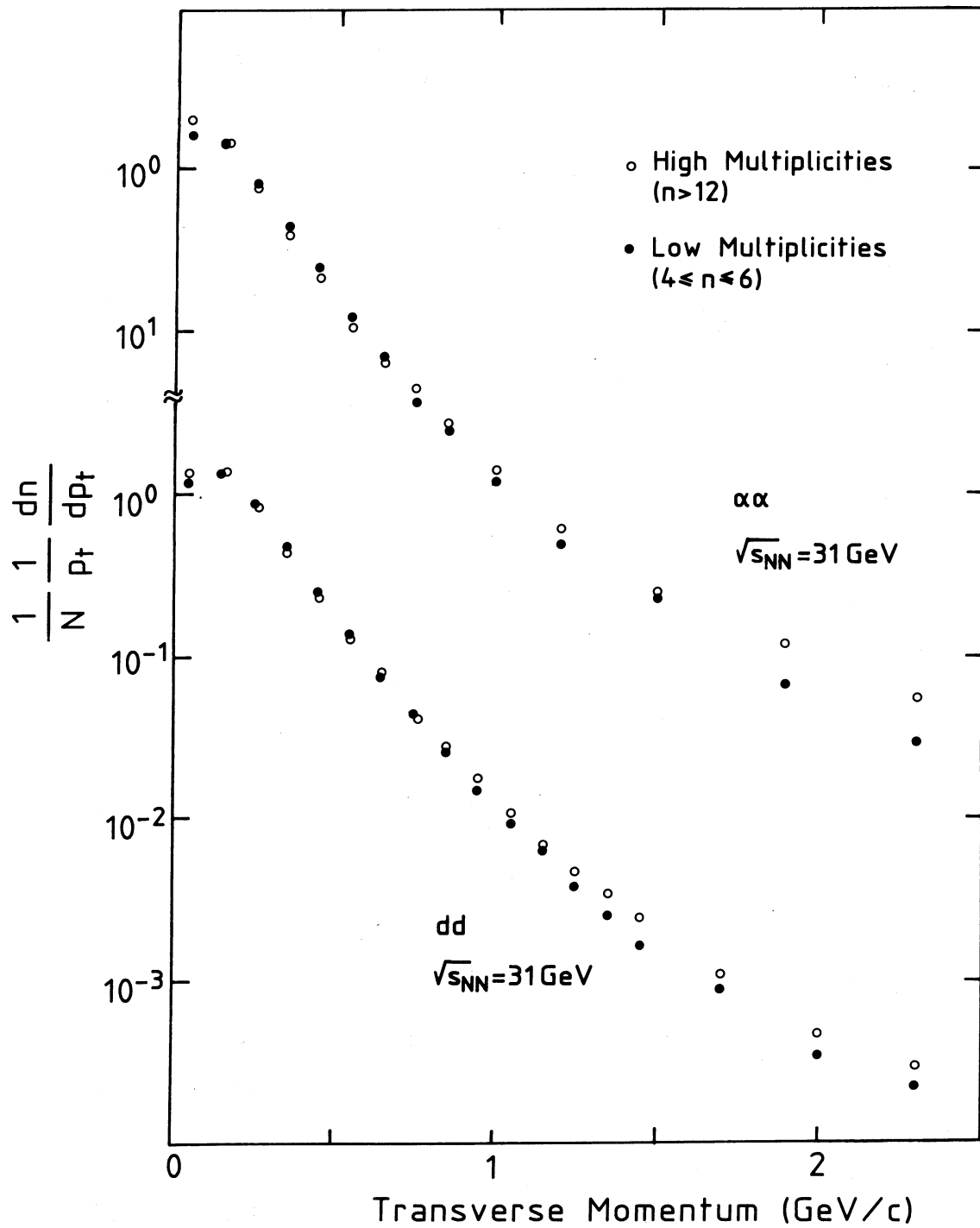


Fig. 1

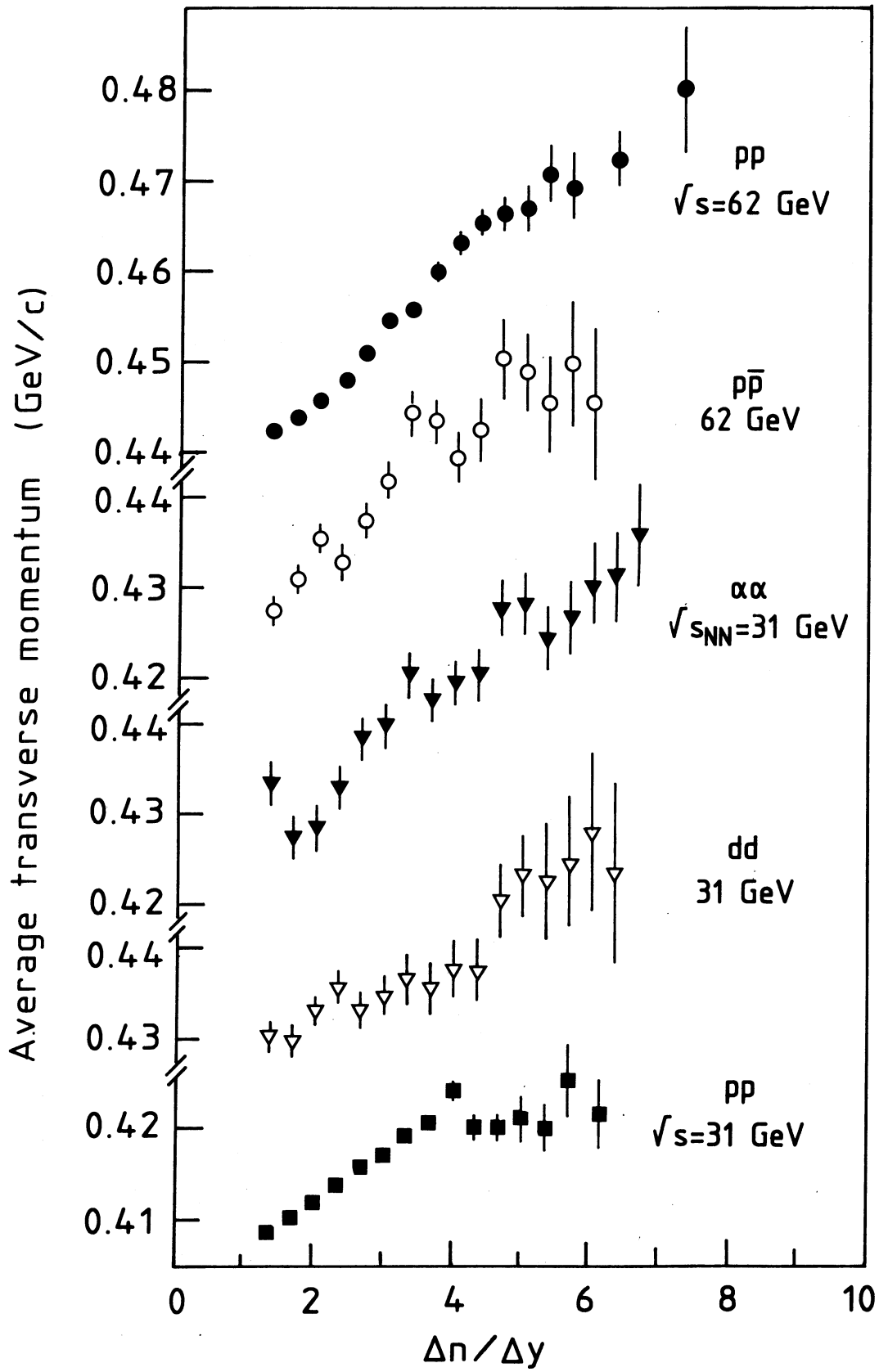


Fig. 2

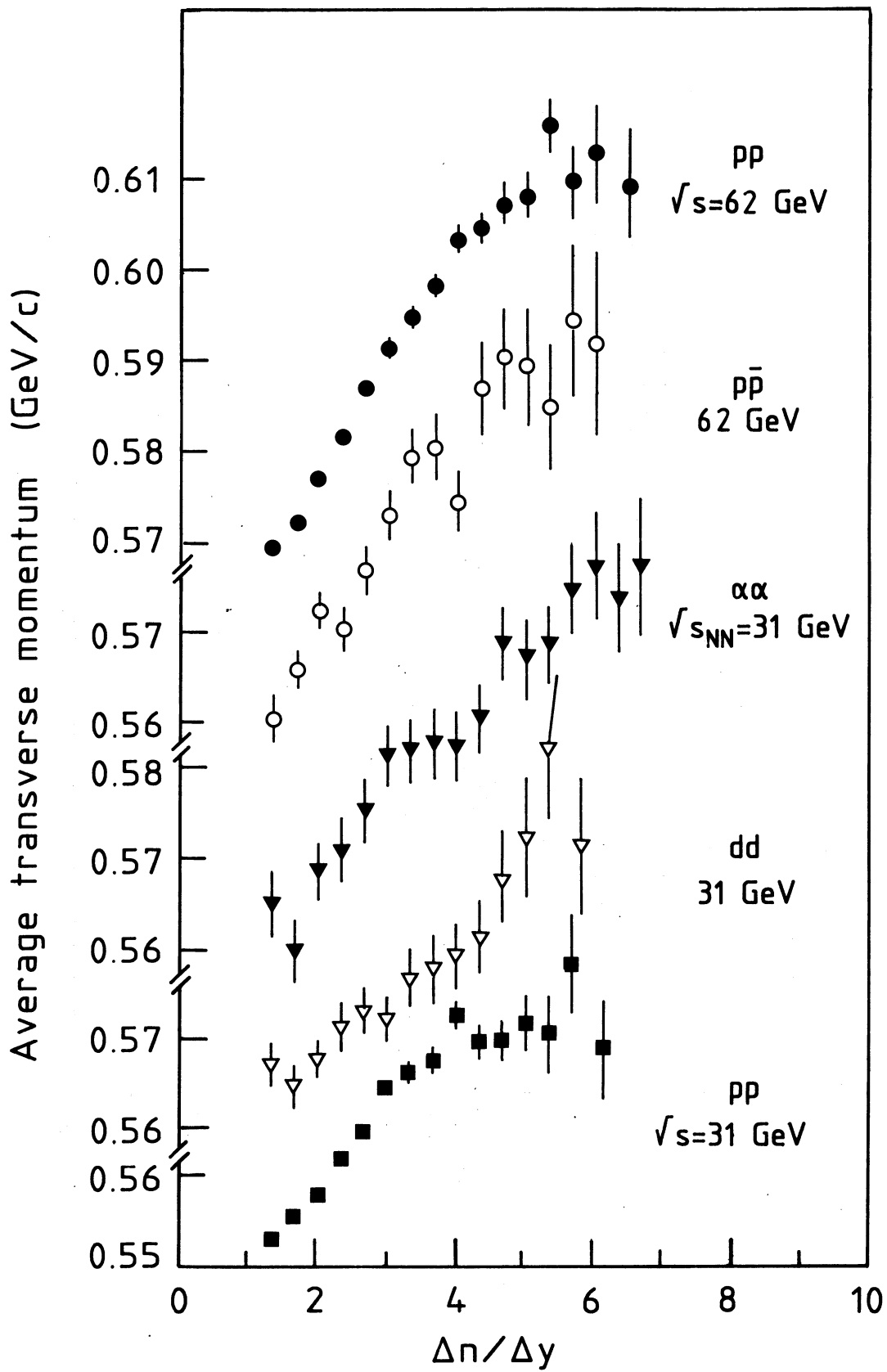


Fig. 3

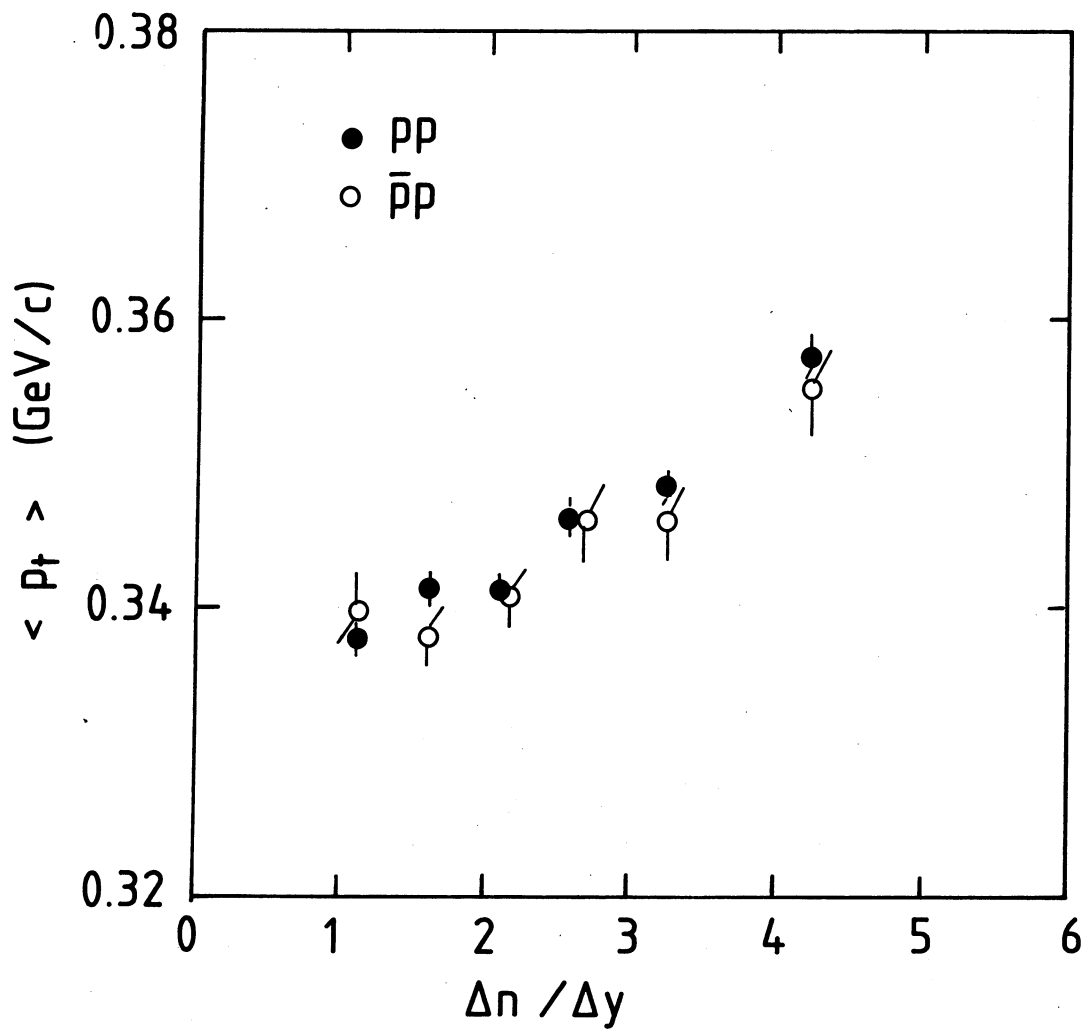


Fig. 4