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**Measurement of the inclusive differential jet cross section in pp collisions  
at  $\sqrt{s} = 2.76$  TeV**

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**Abstract**

The ALICE collaboration at the CERN Large Hadron Collider reports the first measurement of the inclusive differential jet cross section at mid-rapidity in pp collisions at  $\sqrt{s} = 2.76$  TeV, based on integrated luminosity of  $13.4 \text{ nb}^{-1}$ . Calculations based on Next-to-Leading-Order perturbative QCD are in good agreement with the measurements. The measured ratio of inclusive jet cross sections for jet cone radii  $R = 0.2$  and  $R = 0.4$  is studied as a function of the jet transverse momentum, and is also well reproduced by a Next-to-Leading Order perturbative QCD calculation.

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\*See Appendix A for the list of collaboration members



## 1 Introduction

A QCD jet is a collimated shower of particles arising from the hadronization of a highly virtual quark or gluon generated in a hard (high momentum transfer  $Q^2$ ) partonic scattering. Perturbative Quantum Chromodynamics (pQCD) calculations of inclusive jet cross sections agree with collider measurements over a wide kinematic range, for a variety of collision systems [1, 2, 3]. Jets are important tools for studying Standard Model and Beyond Standard Model physics, as well as hot and dense QCD matter that is created in high energy collisions of heavy nuclei. In heavy-ion collisions, large transverse momentum ( $p_T$ ) partons traverse the colored medium and lose energy via induced gluon radiation and elastic scattering, which modify jet longitudinal and transverse structure relative to jets generated in vacuum. These modifications (“jet quenching”) are observable experimentally, and can be calculated theoretically ([4] and references therein).

Measurements of the properties of the QCD medium generated in Pb–Pb collisions at the Large Hadron Collider (LHC) require reference data from more elementary collisions (pp and pPb), where such a medium is not expected to be created. In March 2011, the LHC undertook a three-day run with pp collisions at  $\sqrt{s} = 2.76$  TeV, the same center-of-mass energy as the initial Pb–Pb runs, to obtain initial measurements of such reference data. This paper reports the measurement of the inclusive differential jet cross section at mid-rapidity from that run, based on integrated luminosity of  $13.4 \text{ nb}^{-1}$ .

A QCD jet is not a uniquely defined physics observable. Jet measurements require specification of a jet reconstruction algorithm, for which there are various choices. A jet algorithm should be infrared and collinear safe, and should be applicable in a comparable way to both experimental measurements and theoretical calculations [5]. The anti- $k_T$  algorithm [5, 6] has the required properties and is used in this analysis. A key parameter of any jet reconstruction algorithm is the “cone radius”, or resolution parameter,  $R$ , which specifies the maximum distance in pseudorapidity  $\eta$  and azimuthal angle  $\phi$  over which the algorithm clusters constituent particles,  $\sqrt{\Delta\eta^2 + \Delta\phi^2} < R$ . The inclusive jet cross section as a function of  $R$  is sensitive to the transverse structure of jets. Inclusive jet cross sections and their ratio for different values of  $R$  can be calculated using pQCD at NLO [7, 8, 9, 10]. We compare our measurements at  $R = 0.2$  and  $R = 0.4$  to such calculations.

## 2 Data Set and Detector

The data was recorded by the ALICE detector for pp collisions at  $\sqrt{s} = 2.76$  TeV. ALICE detector [11] consists of two large-acceptance spectrometers: a central detector consisting of a high precision tracking system, particle identification detectors, and calorimetry, all located inside a large solenoidal magnet with field strength 0.5 T; and a forward muon spectrometer. Only the central detector is used for this analysis.

Several trigger detectors were utilized: the VZERO, consisting of segmented scintillator detectors covering the full azimuth over  $2.8 < \eta < 5.1$  (VZEROA) and  $-3.7 < \eta < -1.7$  (VZEROC); the SPD [12], a two-layer silicon pixel detector consisting of cylinders at radii 3.9 cm and 7.6 cm from the beam axis and covering the full azimuth over  $|\eta| < 2$  and  $|\eta| < 1.4$  respectively; and the EMCal [16, 17], an Electromagnetic Calorimeter covering 100 degrees in azimuth and  $|\eta| < 0.7$ . The EMCal consists of 10 supermodules with a total of 11520 individual towers, each covering an angular region  $\Delta\eta \times \Delta\phi = 0.014 \times 0.014$ . The EMCal “Single Shower (SSh)” trigger system generates the fast sum (800 ns) of energies in overlapping groups of  $4 \times 4$  adjacent EMCal towers, followed by comparison to a threshold energy. Event recording was initiated by two different trigger conditions: (i) the Minimum Bias (MB) trigger, requiring at least one hit in any of VZEROA, VZEROC, and SPD, in coincidence with the presence of a bunch crossing, and (ii) the EMCal trigger, requiring that the MB trigger condition was satisfied and that at least one SSh sum in the EMCal exceeded a nominal threshold energy of 3.0 GeV.

Events selected for offline analysis were required to have a primary vertex reconstructed within 10 cm from the center of the ALICE detector along the beam axis. After event selection cuts, the MB-triggered dataset corresponds to integrated luminosity of  $0.5 \text{ nb}^{-1}$ , while the EMCal-triggered dataset corresponds to  $12.9 \text{ nb}^{-1}$ .

For offline analysis, input to the jet reconstruction algorithm consists of charged particle tracks and EMCal clusters. Charged particle tracks are measured in the ALICE tracking system covering the full azimuth within  $|\eta| < 0.9$ . The tracking system consists of the ITS [12], a high precision, highly granular Inner Tracking System consisting of six silicon layers including the SPD, with inner radius 3.9 cm and outer radius 43.0 cm, and the TPC [13], a large Time Projection Chamber with inner radius 85 cm and outer radius 247 cm, that measures up to 159 independent space points per track.

Accepted tracks have a  $p_T$ -dependent minimum number of space points in the TPC, ranging from 70 at  $p_T = 0.15 \text{ GeV}/c$  to 100 for  $p_T > 20 \text{ GeV}/c$ . In order to achieve high and azimuthally uniform tracking efficiency required for jet reconstruction, charged track selection utilizes a hybrid approach that compensates local inefficiencies in the ITS. Two distinct track classes are accepted in the hybrid approach: (i) tracks containing at least three hits in the ITS, including at least one hit in the SPD, with momentum determined without including the primary vertex (90% of all accepted tracks), and (ii) tracks containing less than three hits in the ITS or no hit in the SPD, with the primary vertex included in the momentum determination (10% of all accepted tracks). Tracks with measured  $p_T > 0.15 \text{ GeV}/c$  are accepted. Tracking efficiency is  $\sim 70\%$  at  $p_T = 0.15 \text{ GeV}/c$ , increasing with  $p_T$  to  $\sim 85\%$  for  $3 < p_T < 40 \text{ GeV}/c$ . Charged track momentum resolution is estimated on a track-by-track basis using the covariance matrix of the Kalman track model fit [14] and by the invariant mass resolution of reconstructed  $\Lambda$  and  $K_s^0$  [15]. For the first class of tracks, the resolution is  $\delta p_T/p_T \sim 1\%$  at  $p_T = 1.0 \text{ GeV}/c$  and  $\sim 4\%$  at  $p_T = 40 \text{ GeV}/c$ , while for the second class of tracks the resolution is  $\sim 1\%$  at  $p_T = 1.0 \text{ GeV}/c$  and  $\sim 8\%$  at  $p_T = 40 \text{ GeV}/c$ . Fake jets due to incorrect track reconstruction are identified and removed using features of the jets themselves, as described below. Charged tracks with  $p_T > 40 \text{ GeV}/c$  make negligible contribution to the inclusive jet population considered in this analysis.

EMCal clusters are formed by a clustering algorithm that combines signals from adjacent EMCal towers, with cluster size limited by the requirement that each cluster contains only one local energy maximum. The systematic dependence on the clustering algorithm was explored with an alternative approach, in which clusters are strictly limited to  $3 \times 3 (\eta \times \varphi)$  adjacent EMCal towers. This comparison results in 4% systematic uncertainty in the inclusive jet cross section (Table 1). A noise suppression threshold of 0.05 GeV is imposed on individual tower energies, and the cluster transverse energy must exceed 0.3 GeV. Noisy towers, identified by their event-averaged characteristics and comprising about 1% of all EMCal towers, are removed from the analysis. Clusters with large apparent energy but anomalously small number of contributing towers are attributed to the interaction of slow neutrons or highly ionizing particles in the avalanche photodiode of the corresponding tower, and are removed from the analysis. Charged hadrons also deposit energy in the EMCal, most commonly via minimum ionization, but also via nuclear interactions, while electrons deposit their full energy in the EMCal via an electromagnetic shower. Both charged hadron and electron contributions to EMCal cluster energy are accounted for in order to avoid double counting a fraction of their energy in the measured jet energy. Charged tracks are propagated to a depth of  $10X_0$  in the EMCal [16] and matched to the closest cluster, within  $\Delta\eta = 0.015$  and  $\Delta\varphi = 0.03$ . Multiple charged hadrons can be matched to a single cluster, although the probability for this is less than 0.5% for pp collisions. Based on the measurement of negligible probability for an isolated charged track to generate an EMCal shower whose energy exceeds the track  $p_T$  [16], the energy of the matched cluster is corrected by removing 100% of the sum of all associated charged hadron momenta, and is set to zero if the remainder is negative. This procedure accurately corrects EMCal energy for clusters without contribution from photons or untracked charged particles (i.e. without ‘‘cluster pileup’’). The probability for cluster pileup in pp collisions is less than 1%. The accuracy of this correc-

tion algorithm is verified by simulation, and its systematic uncertainty is estimated by varying both the fraction of the momentum sum that is subtracted and the track-cluster matching criteria.

### 3 Jet Reconstruction and Trigger Bias

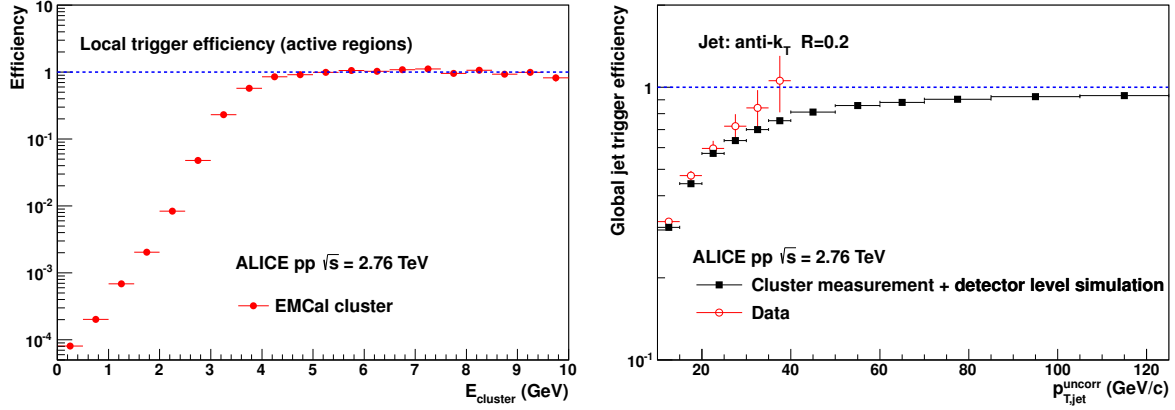
Jet reconstruction is carried out utilizing the FastJet anti- $k_T$  algorithm [5], with resolution parameters  $R = 0.2$  and  $0.4$ . A jet is accepted if its centroid lies within the calorimeter acceptance, with distance at least  $R$  to the calorimeter edge. The measured cross section is corrected to acceptance  $|\eta| < 0.5$  and  $0 < \varphi < 2\pi$ .

The charged particle tracking algorithm may misidentify low  $p_T$  decay daughters from secondary vertices as primary vertex tracks and assign them a much larger  $p_T$  value. Likewise, background processes in the EMCAL can generate false neutral clusters with large apparent  $p_T$  as described above. The cuts imposed at the track or cluster level to suppress such cases directly may not be fully efficient, leading to fake jets with large apparent  $p_{T,\text{jet}}$ . However, such false high  $p_T$  tracks or clusters will have little additional hadronic activity in their vicinity, since they are not a part of an energetic jet. These cases are identified by examining the distribution of  $z = p_{h,\text{proj}}/p_{\text{jet}}$ , the magnitude of the projection of the hadron 3-momentum on the jet axis, relative to the jet momentum. Jets whose  $p_T$  is carried almost entirely by a single hadron generate a peak near  $z \sim 1$  that is found to be discontinuous with the remainder of the distribution. The fake jet population due to single mis-measured tracks or clusters is therefore removed by requiring  $z_{\text{leading}} < 0.98$  independent of  $p_{T,\text{jet}}$ , where  $z_{\text{leading}}$  refers to the  $z$  value of the most energetic hadron in the jet. The effect of the  $z_{\text{leading}}$  cut on jet reconstruction efficiency is negligible for  $p_{T,\text{jet}} > 10$  GeV/c.

The jet bias imposed by the EMCAL SSh trigger must be determined, in order to extract the jet cross section from these data. Figure 1, left panel, shows the measured EMCAL cluster trigger turn-on curve from comparison of SSh-triggered clusters and clusters from MB-triggered data, for the regions of the EMCAL in which the trigger hardware was fully functional (about 90% of the acceptance). A plateau is observed for cluster energy  $E > 5$  GeV. The trigger efficiency in these regions is assumed to be 100% efficient at high  $p_T$ . Due to the limited statistics of the recorded MB data set, we make an independent estimate of the jet trigger efficiency using a data-driven approach that incorporates simulation. Jet events are simulated at the detector level using the PYTHIA6 (Perugia-2011) [18] and HERWIG (version 6.510) [19] event generators, followed by detailed GEANT3 [20] transport and detector response simulation. Each simulated EMCAL cluster in the event is accepted by the trigger with probability equal to the measured cluster trigger efficiency at that energy for the supermodule in which it is located. This procedure takes into account local variations in trigger efficiencies. A simulated event is accepted by the trigger if at least one EMCAL cluster in the event satisfies the trigger requirement. The global jet trigger efficiency is determined by comparing the inclusive jet spectrum for the triggered and MB populations in the simulation. Figure 1, right panel, shows the calculated trigger efficiency for  $R = 0.2$  jets (solid squares), compared to a measurement using the ratio of EMCAL-triggered and MB spectra from data (open circles). The distributions are consistent, within the precision obtained from the MB data set. The deficit at high  $p_T$  is due to the acceptance loss of the trigger system. The systematic uncertainty due to trigger efficiency arises from dependence on the hadronization model, which is assessed by comparing calculations incorporating the PYTHIA and HERWIG generators, as well as uncertainty in the online trigger threshold and in the relative scaling of SSh-triggered and MB cross sections. The resulting uncertainty increases significantly as jet  $p_T$  decreases.

The underlying event (UE) in the collision is distinct from jet fragments, and its contribution to jet transverse momentum is therefore removed for comparison to theoretical calculations. The UE energy density was estimated to be  $2.1 \pm 0.4$  GeV/c per unit area, using dijet measurements [21] over a limited kinematic range, supplemented by PYTHIA particle-level simulations. The assigned uncertainty is estimated by the difference between data and simulations. The corresponding systematic uncertainty of the measured

cross section is 5% at  $p_T = 20$  GeV/c and 1.5% at  $p_T = 100$  GeV/c for  $R = 0.4$  jets (see Table 1).



**Fig. 1:** Calculation of jet trigger efficiency. Left panel: EMCAL cluster trigger efficiency from data (see text). Right panel: Global jet trigger efficiency from data (red open circles) and from PYTHIA6 detector level simulations (black solid squares), using the measured cluster efficiency from the left panel.

#### 4 Correction to Particle Level

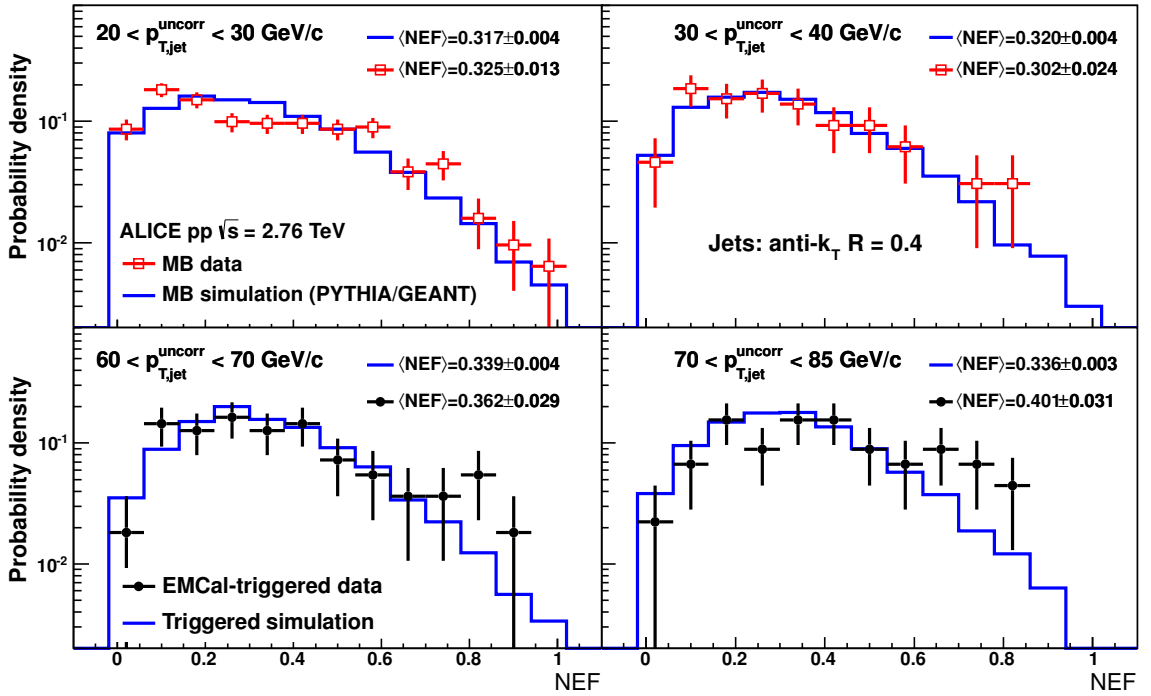
Jets are corrected to the particle level, without accounting for hadronization effects that may modify the energy in the jet cone at the particle level relative to the parton level [22]. This choice is made to facilitate future comparison to jet measurements in heavy-ion collisions, where correction to the parton level is difficult. A bin-by-bin technique [23], based on simulations utilizing PYTHIA6 and GEANT3, is used to correct the measured inclusive spectrum. The correction factor for a given bin is defined as

$$C_{MC} \left( p_T^{\text{low}}; p_T^{\text{high}} \right) = \frac{\int_{p_T^{\text{low}}}^{p_T^{\text{high}}} dp_T \frac{dF_{\text{meas}}^{\text{uncorr}}}{dp_T} \cdot \frac{d\sigma_{MC}^{\text{Particle}}/dp_T}{d\sigma_{MC}^{\text{Detector}}/dp_T}}{\int_{p_T^{\text{low}}}^{p_T^{\text{high}}} dp_T \frac{dF_{\text{meas}}^{\text{uncorr}}}{dp_T}}, \quad (1)$$

where  $d\sigma_{MC}^{\text{Particle}}/dp_T$  is the simulated inclusive jet spectrum at the particle level from PYTHIA6;  $d\sigma_{MC}^{\text{Detector}}/dp_T$  is the simulated inclusive jet spectrum at the detector level from PYTHIA6 and GEANT3, and includes all significant experimental effects;  $\frac{dF_{\text{meas}}^{\text{uncorr}}}{dp_T}$  is a parametrization of the measured, uncorrected inclusive jet distribution, which provides a weight function to minimize the dependence on the spectral shape of the simulation; and  $p_T^{\text{low}}$  and  $p_T^{\text{high}}$  are the bin limits.

The detector level simulation is validated by extensive comparisons to data. An example of this comparison is given in Fig. 2, which shows the distribution of jet neutral energy fraction (NEF) for data and detector level simulation for various intervals of  $p_{T,\text{jet}}$ . Good agreement between data and simulation is observed. Similar levels of agreement are achieved for other key comparisons of data and simulation, including the number of charged track and EMcal cluster constituents per jet, the  $z_{\text{leading}}$  distribution, and the mean  $p_T$  of clusters and tracks in jets.

The cumulative bin-by-bin correction factor  $C_{MC}$  arises from effects that either shift the jet energy scale (JES), or affect the shape or normalization of the inclusive spectrum. Table 1 presents the systematic uncertainty for each component of  $C_{MC}$ . Correction for unmeasured neutron and  $K_L^0$  energy is validated by comparison of the simulation to ALICE measurements of the spectra of protons and kaons in pp collisions at  $\sqrt{s} = 2.76$  TeV for  $p_T < 20$  GeV/c. The systematic uncertainty of  $C_{MC}$  due to this factor is estimated as the difference between PYTHIA6 and HERWIG calculations of the correction. The uncertainty in  $C_{MC}$  arising from the uncertainty in the particle-level spectrum shape is estimated by fitting the particle-level spectrum with a power law function  $\sim 1/p_T^n$  ( $n \sim 5$ ) and varying  $n$  by  $\pm 0.5$ ,



**Fig. 2:** Jet neutral energy fraction (NEF) distributions for minimum bias (MB) data (open squares), EMCal-triggered data (filled circles) and simulations (histograms), in four different  $p_{T,\text{jet}}$  intervals (as marked).

which covers the variation in  $n$  derived from different Monte Carlo models. The main sources of the JES uncertainty are from the corrections for unmeasured neutron and  $K_L^0$  energy and tracking efficiency. Direct contributions to the spectrum uncertainty impose smaller, but non-negligible, uncertainties to the cumulative correction factor. Uncertainties at different  $p_T$  are largely correlated. The components are added in quadrature to generate the cumulative uncertainty, which is labeled “Spectrum total systematic uncertainty” in Table 1, and “Systematic uncertainty” in Fig. 3 and 4.

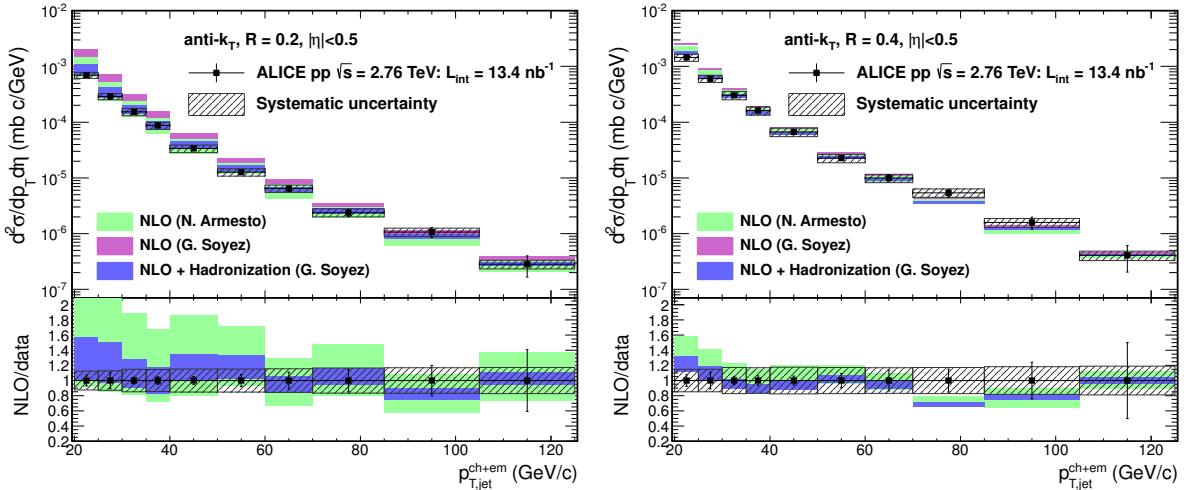
## 5 Results

Figure 3 shows the inclusive differential jet cross-sections at particle level for  $R = 0.2$  (left) and  $R = 0.4$  (right), together with the results of pQCD calculations at NLO. In order to limit sensitivity to the systematic uncertainty of the SSh trigger efficiency, MB data are used for  $p_T < 30$  GeV/c, whereas EMCal-triggered data are used for  $p_T > 30$  GeV/c. The calculation shown by the green band [9] is carried out at the parton level using parton distribution functions MSTW08 [24]. The magenta band [10] represents a parton level calculation utilizing CTEQ6.6 [25], while the blue band is the same calculation including hadronization corrections [10]. The bands indicate the theoretical uncertainty estimated by varying the renormalisation and factorisation scales from  $0.5 p_T$  to  $2.0 p_T$ . The lower panels of Fig. 3 show the ratio of the NLO pQCD calculations to data. The calculations for both  $R = 0.2$  and  $R = 0.4$  are seen to agree with data within uncertainties, when hadronization effects are included.

Figure 4 shows the ratio of the measured inclusive differential jet cross sections for  $R = 0.2$  and  $R = 0.4$ , which is sensitive to jet structure. The numerator and denominator were measured using disjoint subsets of the data, to ensure that they are statistically independent. The kinematic reach of this measurement is therefore less than that of the inclusive spectra themselves. The figure also shows parton-level pQCD calculations at Leading-Order (LO) (magenta band), NLO (orange band), and NLO with hadronization correction (blue band) [10]. This ratio allows a more stringent comparison of data and calculations

**Table 1:** Systematic uncertainty for each component of the corrections to the inclusive jet cross section measurement. Data at 25 GeV/c are from the MB dataset, whereas data at 100 GeV/c are from the SSh-triggered dataset. The upper section, labeled “JES uncertainty”, includes all corrections that can shift the Jet Energy Scale (JES), and the uncertainty is given as percentage of the JES. Due to the shape of the inclusive spectrum, a fractional variation in JES generates a fractional variation in cross section approximately five times larger. The lower section includes those corrections that affect the inclusive spectrum shape or normalization, with values referring to percent variation of cross section. “Spectrum total systematic uncertainty” is the cumulative uncertainty of the spectrum due to both sections of the Table, also as percent variation of cross section.

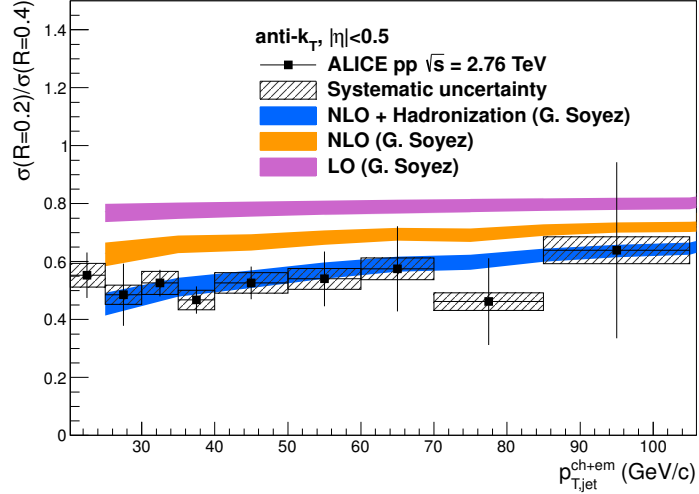
Systematic uncertainties	$R = 0.2$		$R = 0.4$	
	25 GeV/c	100 GeV/c	25 GeV/c	100 GeV/c
JES uncertainty				
Unmeasured neutron+ $K_L^0$	0.4%	0.7%	0.4%	0.5%
Tracking efficiency	1.4%	2.2%	1.8%	2.4%
Charged energy double-counting	0.7%	1.4%	0.6%	1.6%
Underlying event	0.2%	negligible	1.0%	0.3%
Energy scale of EMCal cluster	0.6%	0.6%	0.8%	0.8%
EMCal non-linearity	0.3%	negligible	0.6%	negligible
EMCal clustering algorithm	1.0%	1.0%	1.0%	1.0%
Momentum scale of charged tracks	negligible	negligible	negligible	negligible
Total JES uncertainty	2.0%	2.9%	2.6%	3.2%
Yield uncertainty				
Input PYTHIA spectrum shape	4%	6%	4%	7%
Dependence on fragmentation model	5%	5%	5%	5%
EMCal-SSh trigger efficiency	none	1.7%	none	1.8%
Momentum resolution of charged track	2%	2%	3%	3%
Energy resolution of EMCal cluster	1%	1%	1%	1%
Cross section normalization	1.9%	1.9%	1.9%	1.9%
Spectrum total systematic uncertainty	12.6%	17.1%	14.9%	18.4%



**Fig. 3:** Upper panels: inclusive differential jet cross section for  $R = 0.2$  (left) and  $R = 0.4$  (right). Vertical bars show the statistical error, while the shaded box shows the total systematic uncertainty (Table 1). The colored bands show the NLO pQCD calculations discussed in the text [9, 10]. Lower panels: ratio of NLO pQCD calculations to data. The ratio for the magenta band is not shown, for clarity.



than the individual inclusive cross sections, since systematic uncertainties that are common or highly correlated, most significantly trigger efficiency, tracking efficiency, and cross section normalization, will have smaller relative contribution to the uncertainty of the ratio. In addition, the pQCD calculation at NLO considers the ratio directly, rather than each distribution separately, and is effectively NNLO [10]. The NLO calculation with hadronization again agrees with the measurement within uncertainties.



**Fig. 4:** Ratio of inclusive differential jet cross sections for  $R = 0.2$  and  $R = 0.4$ , with pQCD calculations from [10].

## 6 Summary

In summary, we have presented the first measurement of the inclusive differential jet cross section at mid-rapidity in pp collisions at  $\sqrt{s} = 2.76$  TeV. These data provide an important reference for jet measurements in heavy-ion collisions at the same  $\sqrt{s_{NN}}$ , as well as a test of pQCD calculations at a previously unexamined energy. NLO pQCD calculations with hadronization agree well with both inclusive jet cross section measurements at  $R = 0.2$  and  $R = 0.4$ , and their ratio.

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## A The ALICE Collaboration

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