

Studies on the benefit of extended capabilities of the driver accelerator for EURISOL

K.-H. Schmidt^a, A. Kelić^a, S. Lukić^a, M. V. Ricciardi^a, M. Veselsky^b

^a*Gesellschaft für Schwerionenforschung m.b.H., Planckstr. 1, 64291 Darmstadt, Germany*

^b*Institute of Physics, Slovak Academy of Sciences, Bratislava, Slovakia*

(Received August 25, 2006)

Possibilities are studied for the optimization of EURISOL rare nuclide yields in specific regions of the nuclear chart by building the driver accelerator in a way that enables accelerating several additional beam species, to specific energies, besides the baseline 1 GeV proton beam. Nuclide production rates with these driver beams are compared to the production rates expected with the 1 GeV proton beam in the direct-production and the high-power-converter scenarios. Arguments are presented to show that several additional driver-beam scenarios could provide substantial benefit for the production of nuclides in specific regions of the nuclear chart. The quantitative values in this report are preliminary in the sense that they depend on assumptions on the values of some key parameters which are subject to technical development, e.g. maximum beam intensities or limits on the target heat load. The different scenarios are compared from the aspect of nuclide yields. The arguments presented here, when complemented by the corresponding technological and financial considerations, are intended to serve as a part of the basis for decisions on the design of the driver accelerator of EURISOL.

PACS: 29.25.Rm, 25.70.Mn, 25.85.-w, 24.10.-i,

I. INTRODUCTION

Projects for future secondary-beam facilities aim at exploiting the very technological limits of the production of rare nuclide beams. Such a goal requires careful work on the optimization of the production techniques. That is an interdisciplinary task, as it involves the fields of accelerator technology, nuclear reactions, extraction techniques, target handling and others.

EURISOL [1] is the European project for constructing a secondary-beam facility based on the ISOL approach, which should provide the highest beam intensities and which will give access to the most exotic nuclides within the technological limits. The present work is dedicated to provide arguments for optimising the technical parameters of the driver accelerator, which is one of the most urgent problems to be solved. It has been performed within task 11, “Beam-intensity calculations”, of the EURISOL design study.

Guided by the experience with the long-term operation of the ISOLDE facility [2,3], a 1 GeV proton beam has been chosen as the baseline option for the driver accelerator. Using different target material, this option allows for producing a large number of isotopes of many elements. However, this option might not be the optimum solution for all nuclides of interest. Extended capabilities of the driver accelerator in order to provide additional beam species and more complex technical approaches might result in a benefit for producing nuclides in certain regions of the chart of the nuclides. The present report investigates this question in a systematic way. For this purpose, the benefit of extended capabilities of the driver accelerator is considered in connection with a quantitative discussion of nuclear-reaction aspects and the technical limitations of the ISOL method.

II. TECHNICAL CONSTRAINTS FOR THE DRIVER ACCELERATOR

Several concepts for the heavy-ion driver have been mentioned in the EURISOL Report [4] resulting from the Feasibility Study RTD. These concepts define the technical constraints on eventual extensions of the driver accelerator:

1. The three HI driver-beam options described in the report of the Driver-Accelerator group (App. B, Sect. 4.2 of the EURISOL RTD report [4])
 - a. $A/q = 2$ at 43 A MeV and $A/q = 3$ at 28 A MeV
 - b. $A/q = 2$ at 500 A MeV
 - c. $A/q = 3$ at 100 A MeV and $A/q = 2$ at 150 A MeV

2. The HI driver beam option described in the Target-and-Ion-Source group (App. C, Sect. 2.2.3 of the EURISOL RTD report [4])

$A/q = 6$ at 166.5 A MeV and $A/q = 3$ at 333 A MeV

A ${}^3\text{He}^{2+}$ beam, accelerated to 2 GeV in the main linac with minor cavity modifications, will also be considered.

There is an important distinction between options 1 and 2 on one side and the last option on the other side: The nuclides of interest are projectile-like residues in options 1 and 2, but target-like residues in the last option, like in the baseline proton-beam concept.

In addition, the option of a deuteron beam of $E \leq 200$ MeV and a current of ≥ 0.5 mA for a neutron-converter option, as originally proposed by J. Nolen for the ANL secondary-beam facility [5,6,7] will be investigated.

All these options will be compared with the baseline 1 GeV proton-driver case. This baseline option comprises two scenarios:

1. 200 μA beam current for the direct-target option.
2. 3-4 mA beam current for the converter-target option.

III. CHARACTERISTICS OF THE EURISOL BASELINE OPTIONS

As an ISOL-based facility, EURISOL has to face a few technical conditions like a limited choice on the target material, element-dependent overall efficiencies, and decay losses for short-lived species. As a starting point of the present investigation, it is necessary to have a view on these characteristics. Together with the features of the in-target nuclide production by a 1-GeV proton beam, this will allow analyzing, whether the secondary-beam production at the EURISOL facility can be enhanced in specific regions of the nuclear chart if the baseline driver-beam options (1 GeV proton beam of 3-4 MW on a converter target and 200 kW on a direct target) are complemented by additional driver beams.

A. Target material

The two baseline proton-driver options, direct target and high-power fission target, rely on the use of a thick target, where the radioactive species are produced by the primary protons and by the secondary neutrons, respectively. Thus, it is the same material in which the radioactive species are produced and from which they have to be extracted. For this reason, at present only a limited number of elements is well suited as a component of an ISOL target, e.g. due to the required thermal stability. The following list of target materials used at ISOLDE and some other ISOL facilities has been compiled and discussed by Ulli Köster in 2001 [8]:

- **Molten metal targets**
Typical elements used for ISOL molten-metal targets: Ge, Sn, La, Pb, Bi, U.
- **Metal-foil targets**
Typical elements used for ISOL metal-foil targets are: Ti, Nb, Ta, W.
- **Carbide targets¹**
Typical carbide targets are: C, Al₄C₃, SiC, VC, LaC₂, ThC₂, UC₂.
Other possible carbide targets are: ScC₂, CaC₂, TiC, ZrC, GdC_x.
- **Oxide targets**
Typical oxide targets are: MgO, Al₂O₃, CaO, TiO₂, SrO, ZrO₂, BaO, La₂O₃, CeO₂, ThO₂
- **Other targets**
Some other target material was used, at least for tests: BN, AlN, CaB₆, Zr₅Si₃, Zr₅Ge₃, BaB₆, BaZrO₃, CeB₆, CeS, Hf₃Si₂, Hf₅Ge₃, Hf₅Sn₃, Ta₃Si₃, Na-zeolite, Ca-zeolite, Ti-zeolite.

The restrictions are more severe than it may appear from this list because there are further restrictions on the target material for specific elements to be extracted as secondary beams, e.g. due to conditions on chemical reactivity or vapour pressure of the target material. These limitations are explicitly discussed in ref. [8].

Figure 1 shows a graphical presentation, which illustrates that important gaps exist between the different elements used as ISOL targets. As a general rule one may conclude that these gaps usually extend over about 10 elements. If specific target types are required for optimum extraction conditions of certain elements, one may even encounter gaps about 20 elements wide. This implies some restrictions on the optimization of the target material in view of nuclear-reaction aspects governing the production of specific secondary beams.

It has to be stressed that these limitations can be overcome in some cases by dedicated research and development. However, the general impression given by figure 1 will remain valid to a certain extent also for the operation of EURISOL.

¹ Often the stoichiometric carbide is diluted in excess graphite to provide more favourable release characteristics.

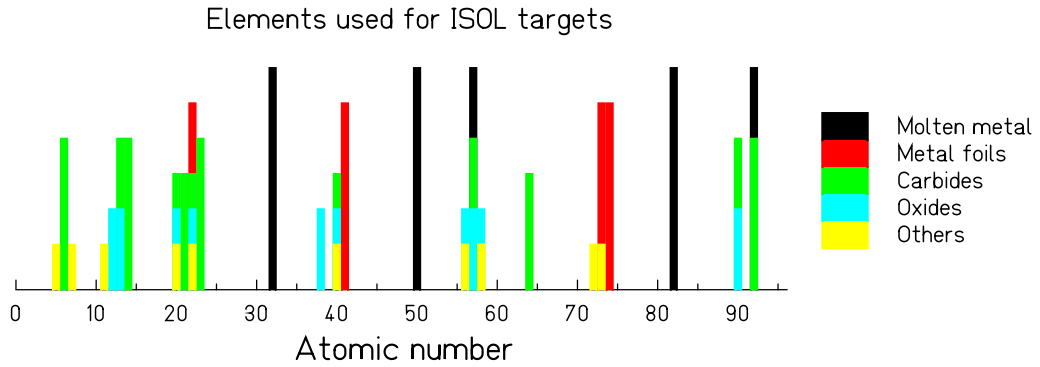


Figure 1: Graphical presentation of the elements used for ISOL targets (data from ref. [8]). (The columns are drawn with different heights for graphical reasons only.)

B. ISOL efficiencies, available elements

Not all elements can be extracted from ISOL targets with high overall efficiencies. Even for sufficiently long-lived isotopes, which are not affected by decay losses during the diffusion and effusion processes occurring in the target – ion-source system, losses may occur in the different stages of the extraction process. Some elements, e.g. the refractory elements, have even so low volatility that they are extracted with very low efficiencies or not observed at all. A general impression of the strong variation of the ISOL efficiency on the atomic number of the nuclear species can be obtained from figure 2, where the result of careful systematic measurements of the overall efficiencies of the OSIRIS target [9] is shown. New target – ion-source developments have led to higher efficiencies in the mean time. We refer to the intense study of overall efficiencies and intensities of beams for the seven NuPECC reference elements beryllium, argon, nickel, gallium, krypton, tin, and francium reported in Appendix C of the EURISOL RTD report [4]. However, a strong dependence on the chemical properties of the different elements as globally demonstrated in figure 2 remains important.

The elements, which were available as ISOLDE beams up to 1994, are shown in figure 3. It gives a good impression on the possibilities and limitations of the ISOL method.

It is particularly penalizing that the elements between ytterbium ($Z = 39$) and palladium ($Z = 46$) and elements around phosphorus ($Z = 15$) and chromium ($Z = 24$), for which very neutron-rich isotopes can be produced by fission of actinides, have very low overall-efficiencies.

The ISOL method relies on complex processes, with characteristics depending strongly on the element considered. Developing methods with optimized efficiency for specific elements requires intensive research, and the overall efficiencies reached are not always high, in particular if the element has low volatility and if the half-life of the isotope of interest is short.

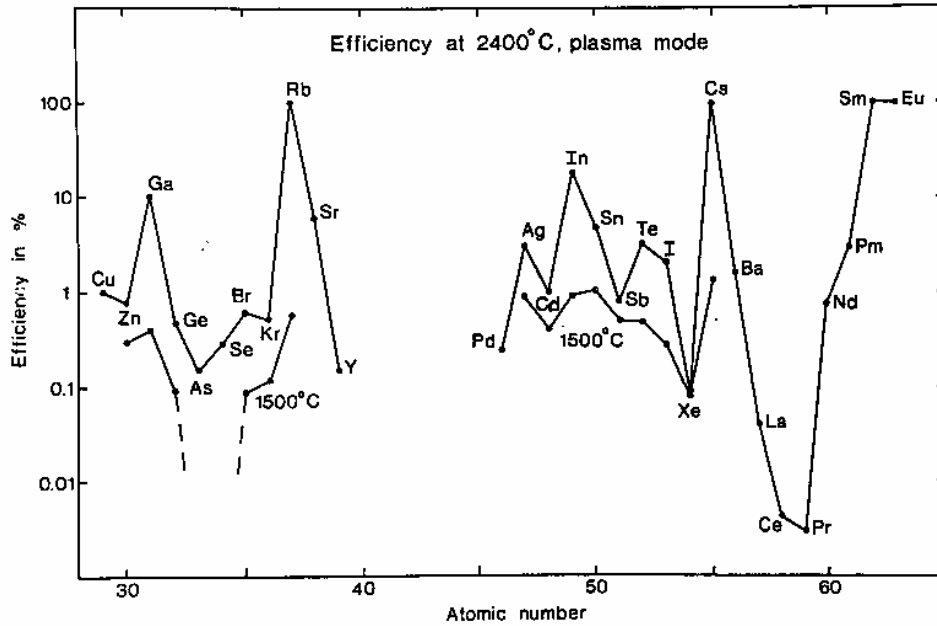


Figure 2: Overall efficiency of the OSIRIS target and ion source as a function of atomic number of fission-product nuclei. The upper curve is measured at a target temperature of 2400 °C. For comparison, the lower curve shows the efficiencies of the previously used system at 1500 °C. The figure is taken from ref. [9].

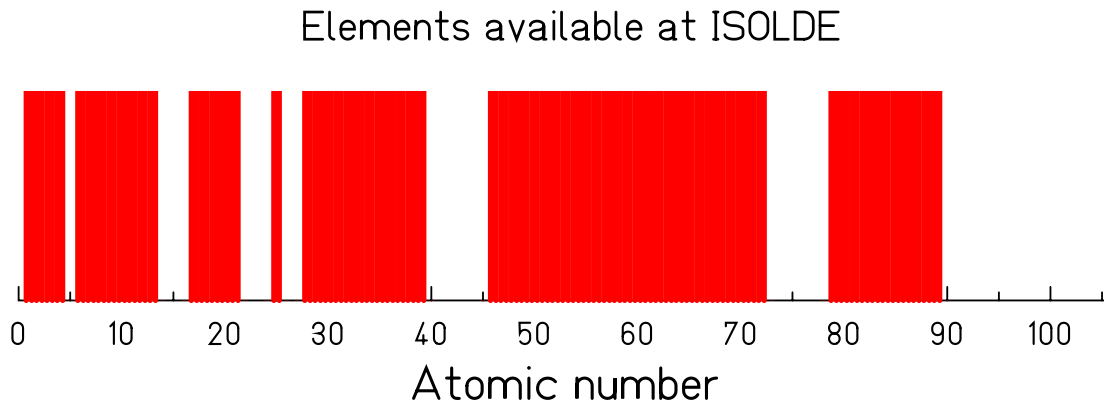


Figure 3: Elements available as ISOLDE beams up to 1994. (Data are taken from ref. [3].) It should be noted that the efficiencies can be rather different for the different elements.

In order to fill some of these “gaps”, a novel approach involving fragmentation and in-flight separation of a re-accelerated ISOL beam [10] is intensively investigated. However, since this scenario is not related to the driver accelerator, it is beyond the scope of the present work.

C. Extraction times

High primary reaction rates are reached by using rather large target volumes. Unfortunately, this leads almost inevitably to long transport times. Figure 4 illustrates the restrictions on the length of isotopic chains to be prepared as secondary beams, which are imposed by the limitations on half lives longer than 100 s, 10 s and 1 s, respectively. Although the figure is incomplete, because measurements of shorter half lives are scarce, the figure gives an impression on the step decrease of the nuclear life times with increasing distance from the valley of beta stability.

Figure 5 shows a compilation of the elements for which efficiencies of the order of 1% or more are reached for half lives above 30 ms, 1 s and 1 min, respectively, from measurements at the GSI mass separator. One should take into account that the targets used at the GSI mass separator were rather thin, around 1 mg/cm², because they were adapted to the range of heavy ions at UNILAC energies (≈ 5 to 10 *A* MeV). Extraction times from targets with larger volume as used at EURISOL can be appreciably longer.

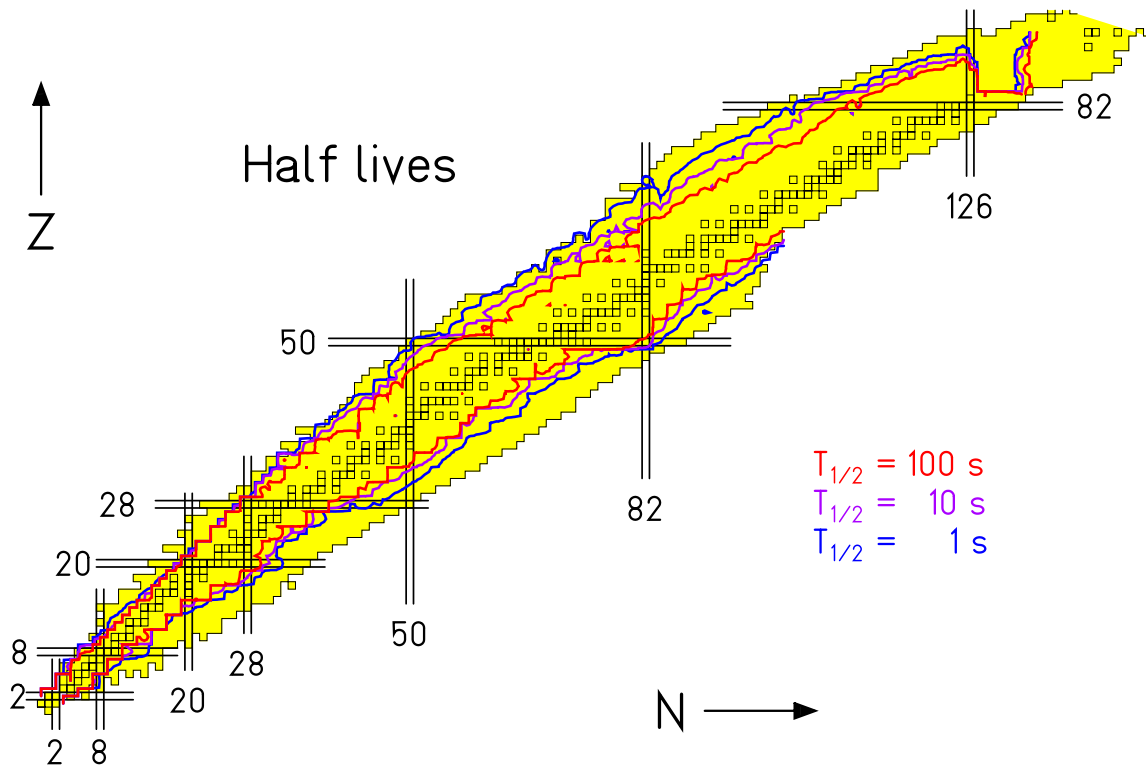


Figure 4: Chart of the nuclides. The yellow area covers the nuclides observed up to now. Black open squares denote stable nuclides. The contour lines mark regions of equal measured nuclear life times: 1, 10 and 100 seconds. On the neutron-rich side, the contour lines are missing above $Z = 65$, since the experimental data for that region do not reach short half-lives.

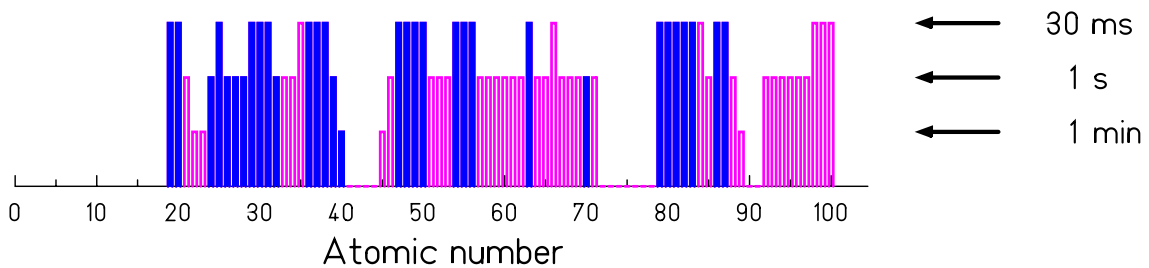


Figure 5: Systematic overview on release-time properties of the different elements with $Z > 18$. The heights of the columns denote the elements for which efficiencies of the order of 1% or more are reached for half lives above 30 ms, 1 s and 1 min, respectively. Blue columns denote those elements for which Z selection is possible without laser ionization. (Compilation by R. Kirchner from measurements at the GSI mass separator, 2001.)

According to a systematic analysis of ISOLDE-SC yields performed in ref. [11], the overall efficiency ε can be represented by a universal analytical function:

$$\varepsilon = \frac{\varepsilon_s}{1 + \left(\frac{t_{1/2}}{\tau}\right)^{-\alpha}} \quad (1)$$

Thus, the efficiency scales roughly with $(t_{1/2}/\tau)^\alpha$ for half lives $t_{1/2}$, which are shorter than the corresponding time parameter τ . The exponent α shows some variation around the mean value of about 1.5. The value of the parameter τ depends on both the element to be extracted and the material properties of the ISOL target. Typical values found for the operation of ISOLDE reach from 1 s to several 100 s and more.

Based on these results, one expects the overall efficiency to decrease roughly *by a factor of 4 from one isotope to the next* in direction towards the drip lines, once the half lives $t_{1/2}$ become smaller than the time parameter τ . This estimate is valid in the region around $A = 100$, while the overall efficiencies vary more strongly for lighter nuclei.

Obviously, fast extraction is very important for keeping decay losses of short-lived nuclides small, since these are in most cases the most exotic and thus the most interesting species, for which already the in-target production rates decrease steeply when approaching the drip lines. But also alternative, technically more complex solutions, which avoid the decay losses, might be envisaged. In the design of the RIA facility, projectile-fragmentation reactions induced by stable nuclear beams and in-flight separation as well as the use of a gas catcher have been considered to overcome this problem, see e.g. ref. [12]. As already mentioned, for the EURISOL project, the fragmentation and in-flight separation of a re-accelerated neutron-rich ISOL beam [10] is being investigated, which should allow producing extremely neutron-rich nuclides. However, a detailed discussion of this option is beyond the scope of the present paper.

D. Nuclear-reaction aspects

Different types of nuclear reactions are exploited for the nuclide production in the baseline design of the EURISOL project, driven by a 1-GeV proton beam. In the direct-target option, spallation-evaporation and spallation-fission residues are produced in mostly direct interactions of the primary protons with different target material. In the high-power fission-target option, secondary neutrons from a converter target induce fission in a uranium or thorium target at excitation energies mostly up to a few MeV above the fission barrier.

From recent experiments performed at GSI in inverse kinematics, there is a good empirical knowledge on the general characteristics of these reactions. In addition, the development of realistic nuclear-reaction codes has strongly advanced on the basis of these new data.

Figure 6, figure 7, and figure 8 present the measured nuclide distributions for 1 GeV proton-induced reactions on ^{56}Fe , ^{136}Xe and ^{238}U , respectively, measured in inverse kinematics [13, 14, 15, 16, 17, 18, 19]. In addition, the nuclide distribution produced in fission of ^{239}U at the typical excitation energy of 2 MeV above the fission barrier (corresponding to the fission of ^{238}U induced by 3.5 MeV neutrons) is shown in figure 9 as calculated with the ABRABLA code [20, 21, 22, 23].

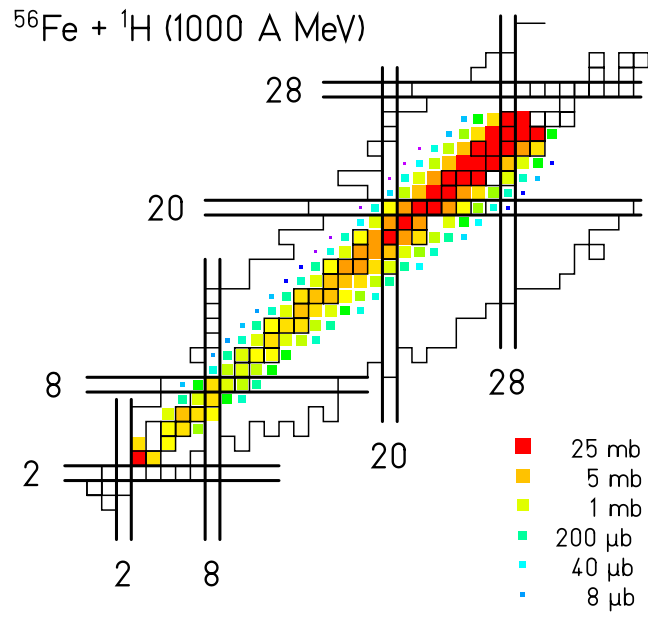


Figure 6: Measured nuclide distribution from the reaction ^{56}Fe (1 A GeV) + ^1H [13, 14].

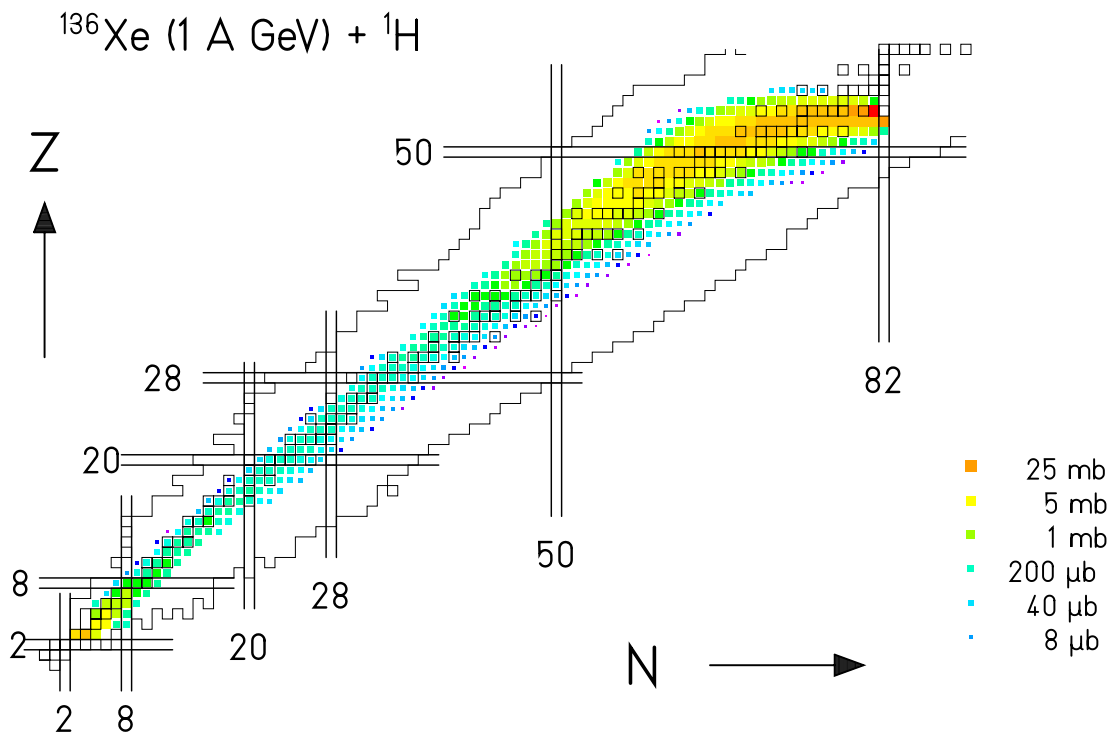


Figure 7: Measured nuclide distribution from the reaction ^{136}Xe (1 A GeV) + ^1H [15].

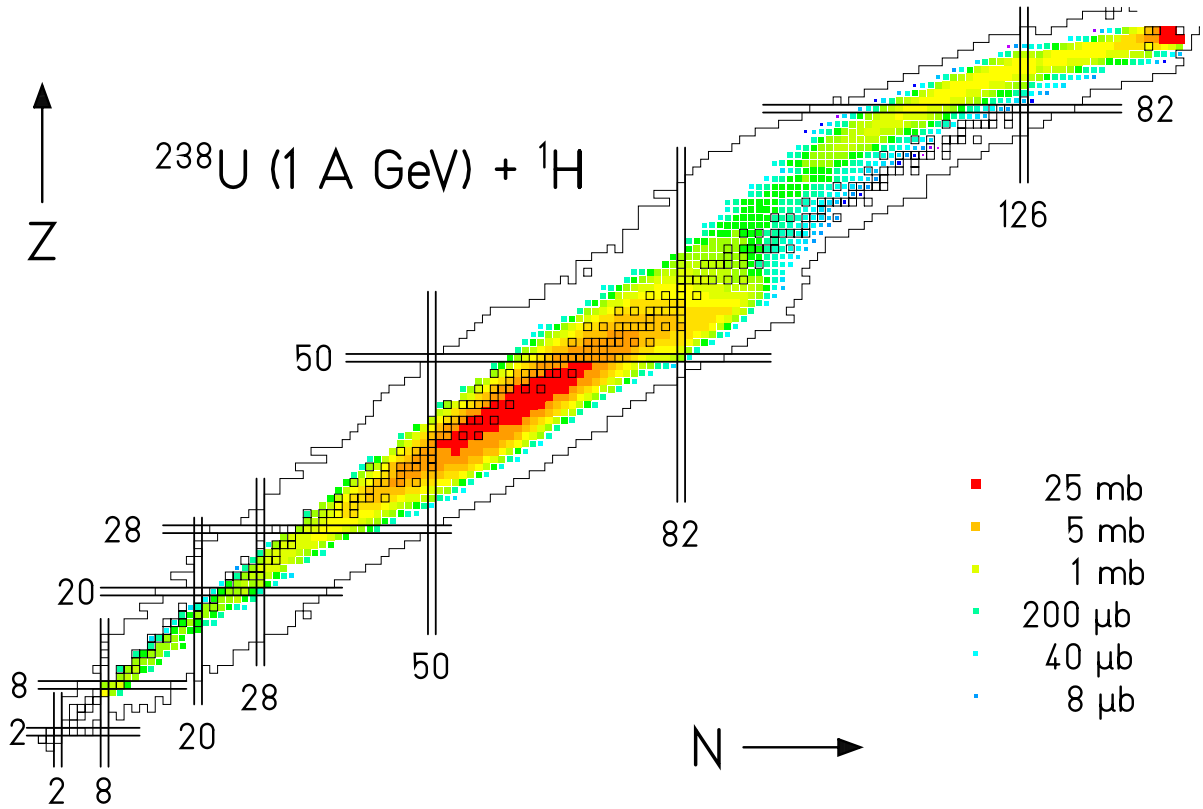


Figure 8: Measured nuclide distribution from the reaction $^{238}\text{U} (1 \text{ A GeV}) + ^1\text{H}$ [16, 17, 18, 19].

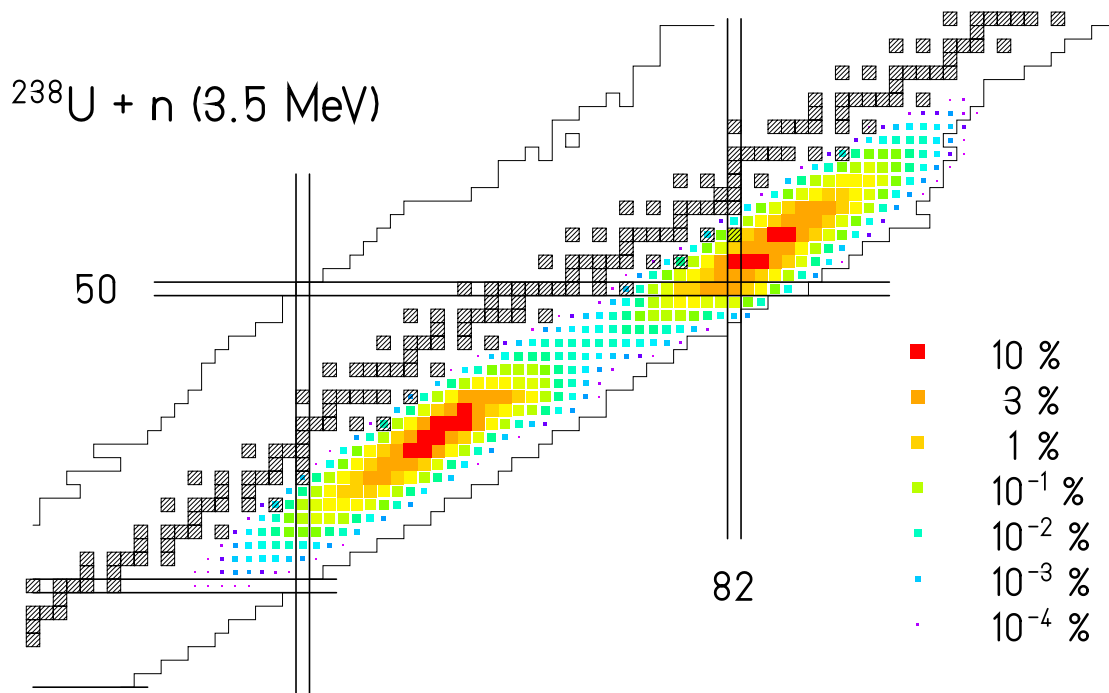


Figure 9: Nuclide distribution of fission products from the neutron-induced fission of ^{238}U for fixed neutron energy of $E_n = 3.5 \text{ MeV}$, calculated with the nuclear-reaction code ABRABLA. The colour code gives the yields in percent.

This overview reveals the following general characteristics:

- Spallation-evaporation residues populate neutron-rich isotopes of a few elements below the target²⁾ with strongly decreasing cross sections by cold-fragmentation reactions [24]. This is an option for the production of neutron-rich isotopes above the fission region. However, the limited choice of target elements imposes a severe restriction on this option.
- The main nuclide production by spallation-evaporation reactions gradually tends towards the proton-rich side and extends to roughly 10 to 15 elements below the target element. Saturation in the production of the most neutron-deficient isotopes is reached around 5 elements below the target, where the cross sections already decrease. The limited choice of target elements imposes again some restriction on this option to produce isotopes at the proton drip line. Lighter residues are situated closer to the beta-stability line.
- Spallation-fission reactions with actinide targets (uranium and thorium are technical options) produce neutron-rich isotopes in a wide element range (extending from about $Z = 10$ to $Z = 60$ and further) with high cross sections, see figure 8.
- Spallation-fission reactions with lighter targets (e.g. lead, see figure 11) are not interesting for the production of radioactive beams, because they produce nuclides rather close to the beta-stability line. Moreover, the cross sections are rather low.
- The nuclide production by fission of a ^{238}U target induced by neutrons of a few MeV is focussed on about 8 very neutron-rich isotopes of a limited number of elements between $Z \approx 30$ to $Z \approx 62$ with a dip around symmetric fission ($Z \approx 46$), see figure 9.

These characteristics were already qualitatively known before and were exploited e.g. for the production of radioactive species at the ISOLDE facility. However, the new experimental results and the improved codes provide a better quantitative and systematic knowledge, which is important for the prediction of yields.

The high-power target option deserves some additional remarks: Due to the deposition of the primary-beam energy in the converter target and the concentration of the production to a very limited number of nuclides, one obtains very high in-target production rates. Best conditions for very high secondary-beam intensities by the ISOL method are provided for nuclides around ^{132}Sn . The double shell closure is responsible for several peculiarities:

3. High production yields in the heavy fission-fragment group go in line with an enhancement of the N/Z ratio due to strong charge polarization, see figure 9.
4. Long half lives extend to rather large values of the neutron excess, see figure 4.

Luckily, the mostly produced elements of the heavy fission-fragment group (tin to lanthanum) can be extracted from the production target with rather high efficiencies by the ISOL method. In addition, their long half lives allow for using rather large target volumes, which are connected with long extraction times, without introducing excessive decay losses.

The nuclides in the light fission-fragment group are less favoured: Firstly, they are less neutron-rich; secondly, they have shorter half lives. However, also here the elements from zinc ($Z = 30$) to strontium ($Z = 38$) can be extracted from the production target with rather high efficiencies by the ISOL method. However, moderate target volumes should be used in order to limit the decay losses.

²⁾ Although the experiment was performed in inverse kinematics, we refer here to normal kinematics, which is applied in the EURISOL scenario. Therefore, the heavy reaction partner is denoted as the target nucleus.

IV. SOME POSSIBILITIES TO ENHANCE THE PRODUCTION YIELDS FOR SPECIFIC NUCLIDES

In the course of the intense discussion on the design of future secondary-beam facilities, different scenarios for the production of secondary beams have been introduced. In particular, numerous ideas have been developed, which exploit specific aspects of nuclear reactions and technical solutions, e.g. with respect to limitations on heat load, ISOL efficiencies and background intensities. The following sections present, in a *qualitative* way, the basic ideas of several of these propositions, which seem to be promising for the EURISOL facility. The assessment of the *quantitative* benefit in comparison with the standard 1 GeV proton-driver option of the EURISOL project will be given in chapter 5.

A. Extending the mass range of spallation-evaporation residues by a 2 GeV primary beam

It has been observed that the intensities of several nuclides increase strongly if the energy of the primary beam grows over a certain range. Figure 10 demonstrates this behaviour.

A better understanding of the nuclear-reaction characteristics, which are responsible for this phenomenon, can be obtained on the basis of systematic measurements on the nuclide production induced by 500 MeV and 1 GeV protons as well as 2 GeV deuterons on ^{208}Pb [25, 26, 27, 28], recently performed in inverse-kinematics experiments at GSI, Darmstadt. In all three cases, which are shown in figure 11, the nuclide distributions consist of two components. Spallation followed by evaporation populates the heavy component, which extends from the projectile to lighter neutron-deficient nuclei. Spallation followed by a fission process (in addition to evaporation before and after fission) populates another region close to beta-stability, centred somewhat below one half the mass of the projectile.

We can observe that the spallation-evaporation distribution extends to lighter masses, when the centre-of-mass energy of the system is increased, while the spallation-fission distribution is less sensitive to the energy. For a centre-of-mass energy of about 500 MeV, the spallation-evaporation distribution extends over essentially 5 elements, for about 1 GeV, essentially 10 elements are formed, while for about 2 GeV more than 15 elements are formed. Thus, the increase in the cross sections shown in figure 10 originates from the extension of the spallation-evaporation distribution to lighter masses with higher energy.

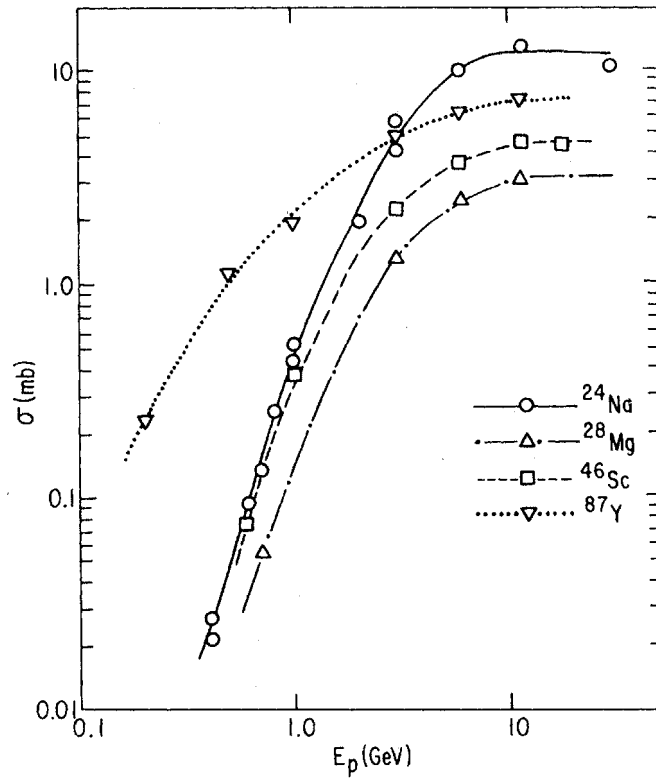


Figure 10: Measured excitation function for the fragmentation of gold by protons leading to some typical light products (taken from ref. [29]).

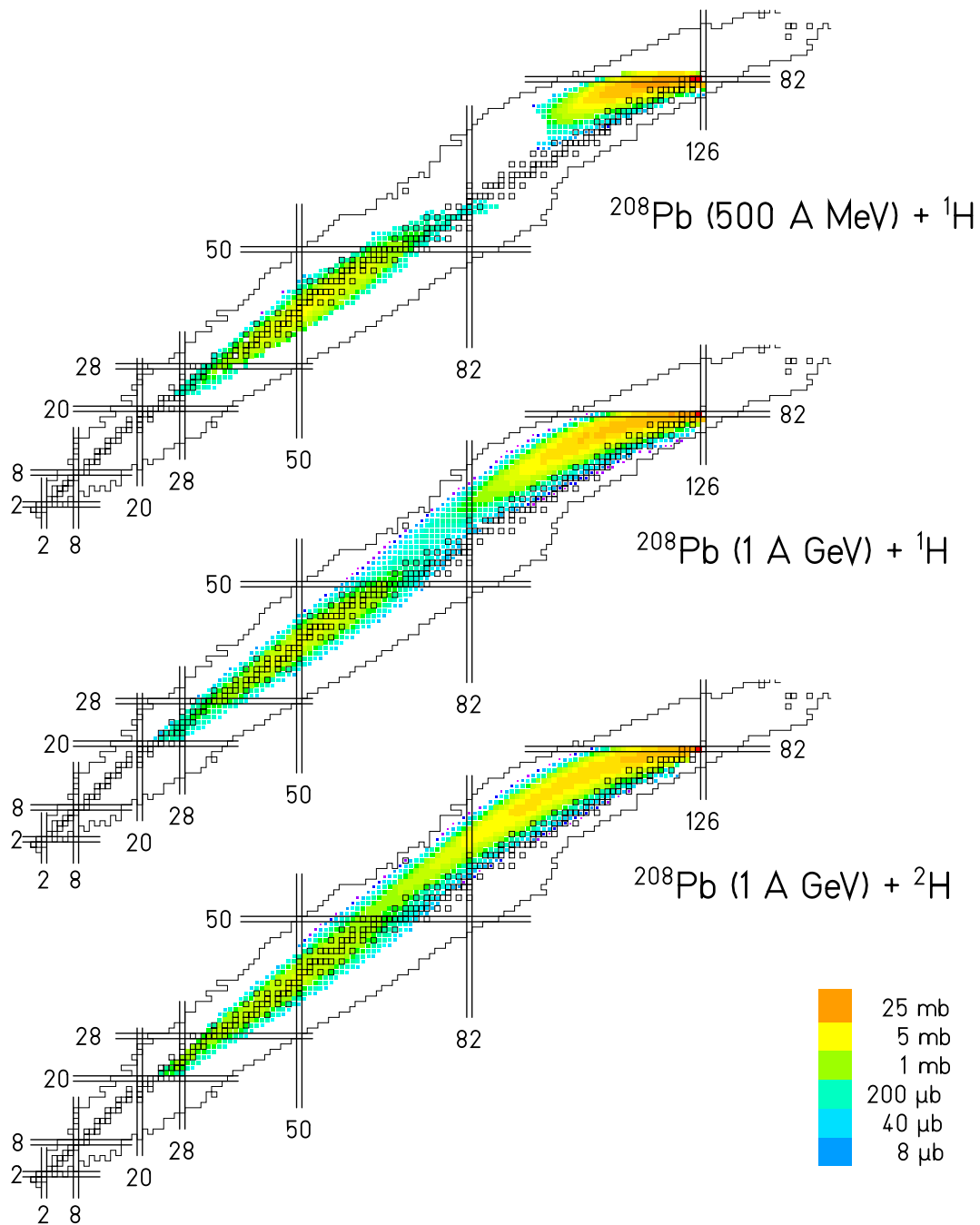


Figure 11: Nuclide distributions of residues produced in the spallation of ^{208}Pb by 500 MeV protons, 1-GeV protons and 2-GeV deuterons, measured at GSI in inverse kinematics [25, 26, 27, 28]. Elements below $Z \approx 20$ were not covered in the experiments.

From this observation, one may conclude that the most neutron-deficient isotopes of all elements that are available by ISOL techniques can be produced with high intensities, if the beam energy is chosen sufficiently high to bridge the gaps between the different target materials. A quantitative discussion of this aspect will follow in the next chapter.

B. Enhancing IMF emission to produce light neutron-rich nuclides by a 2 GeV primary beam in a fissile target

While the total fission cross section stays almost the same when the beam energy is increased from 1 GeV to 2 GeV, the production cross section of light binary products below $Z = 30$ is rather sensitive to the energy deposited in the uranium nucleus. This is clearly seen in some measurements performed at ISOLDE. Figure 12 represents production cross sections of sodium isotopes in uranium at 1.0 and 1.4 GeV proton energy in a uranium-carbide target, deduced from detailed measurements of release curves at ISOLDE [30, 31]. Production cross sections for $^{21,24,25}\text{Na}$ by 600 MeV protons in uranium are also shown. These have been deduced from the ISOLDE SC final yield data [2,32] and not from the detailed release curves, so that it was impossible to correct for the decay losses. However, $^{21,24,25}\text{Na}$ are sufficiently long lived that their decay losses can be neglected (typically 1-5% for $^{21,25}\text{Na}$, and even less for ^{24}Na).

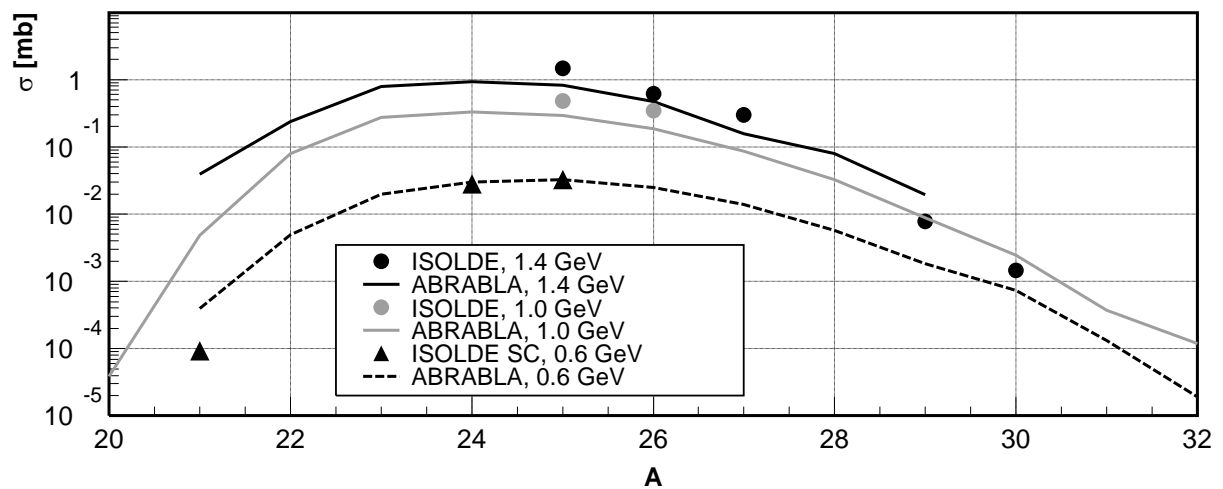


Figure 12: Na isotope production cross sections in uranium deduced from the release measurements at 1.0 and 1.4 GeV proton energies. For comparison, ABRABLA results for the same reaction are also included. Production cross sections for $^{21,24,25,26}\text{Na}$ by 600 MeV protons in uranium deduced from the ISOLDE SC yields are shown, too.

For comparison, production cross sections of sodium isotopes obtained by the ABRABLA code are also shown. The energy dependence of the deduced cross sections is clearly seen, which confirms the ABRABLA predictions. Thus, an important gain in cross section for nuclei in this mass range in a fissile target is to be expected, if the energy of the primary beam is increased.

C. Extended fission-fragment distribution in neutron-induced fission with $E_n = 20$ to 80 MeV by the deuteron-converter option

For the EURISOL facility, a high-power converter target is being designed. It consists of several essentially cylindrical parts, see figure 13. The inner cylinder is the converter, which acts as a spallation neutron source, converting the 1 GeV proton beam into a neutron flux. The converter is surrounded by the production target of a fissile material, where the neutrons induce fission reactions that produce the desired neutron-rich nuclides.

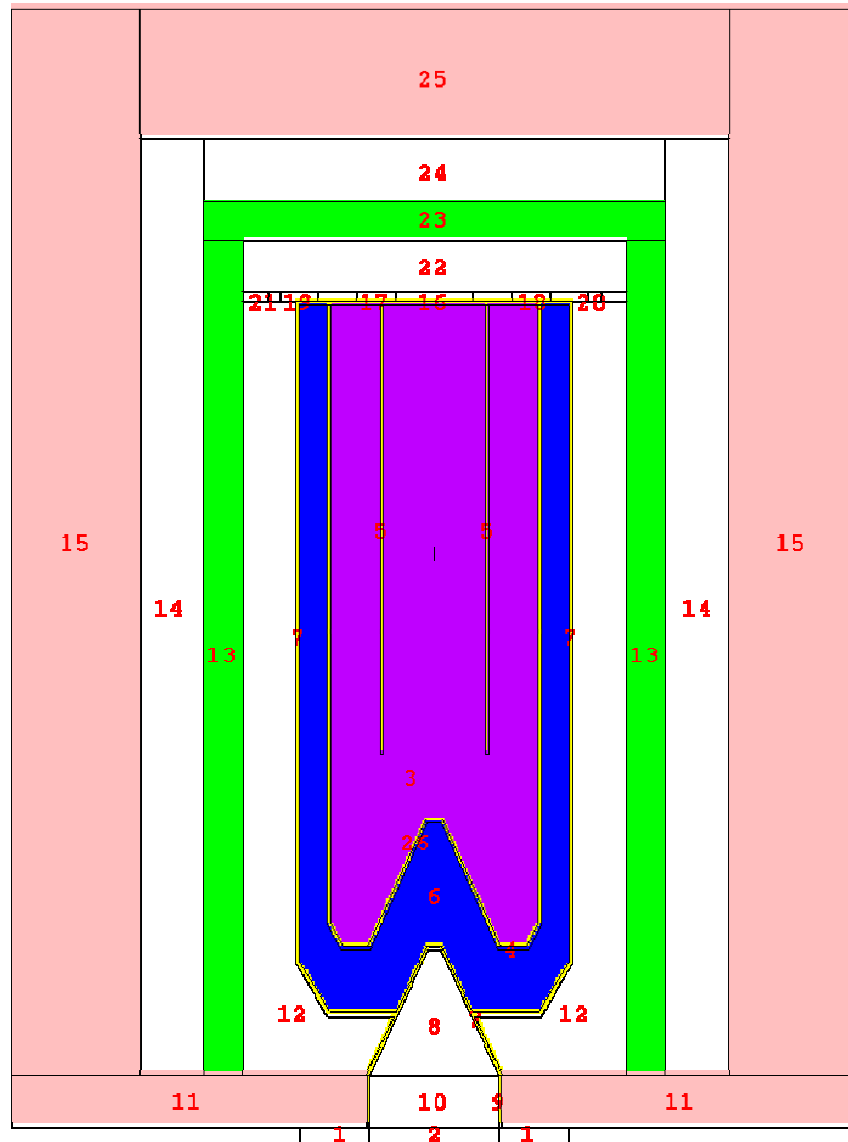


Figure 13: A possible design of the high-power converter target for EURISOL. The drawing shows a cut through the cylindrical assembly. The violet area in the centre denotes the mercury converter, which is hit by the 1 GeV proton beam, entering from the lower side. The green areas mark the uranium production target [33].

The neutron flux in the production target is dominated by evaporative neutrons with a mean energy of about 2 MeV as shown in figure 14.

Due to these low neutron energies, the fission process is induced at energies rather close to the fission barrier, leading to fission-fragment distributions, which are strongly influenced by shell effects and which are dominated by asymmetric fission as demonstrated in figure 9. The production concentrates on very neutron-rich isotopes of a few elements around Xe ($Z = 54$) and Sr ($Z = 38$). The elements between $Z = 42$ and $Z = 48$ are only weakly produced, and the distribution dies out for $Z < 32$ and $Z > 60$.

Due to the high neutron-capture cross section of ^{235}U for low neutron energies, more than half of the nuclide production in a production target of natural uranium is due to fission of the compound nucleus ^{236}U [33]. As can be seen in figure 15, the nuclide distribution is slightly shifted to the neutron-deficient side if compared to low-energy neutron-induced fission of ^{238}U , see figure 9. The difference of 3 neutrons in the fissioning nucleus leads to an average loss of about 1.5 neutrons in the fission fragments. Thus, the neutron-induced fission

of ^{235}U is less favourable for producing the most neutron-rich fission fragments. The total production in the high-power target is roughly given by the superposition of the nuclide distributions of figure 9 and figure 15.

As the total fission rate is limited from the point of view of the target handling, it would be advantageous to use depleted ^{238}U for the target in order to enhance the production of more neutron-rich nuclides.

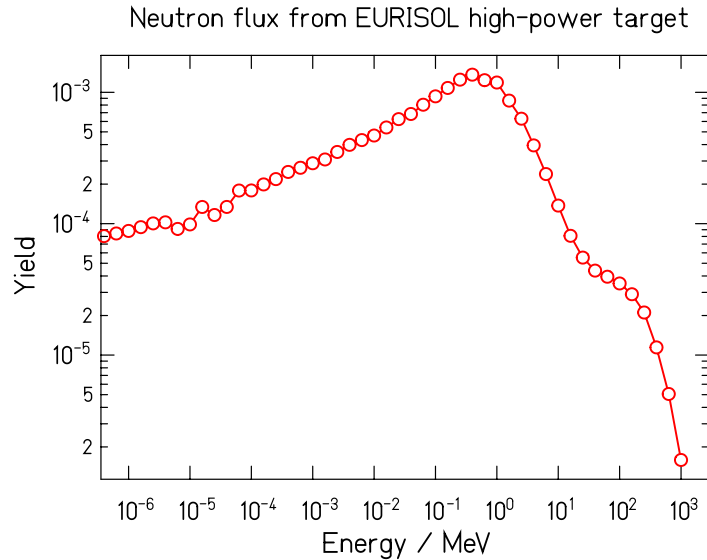


Figure 14: Calculated energy distribution of the neutron flux around the converter of the EURISOL high-power target [33].

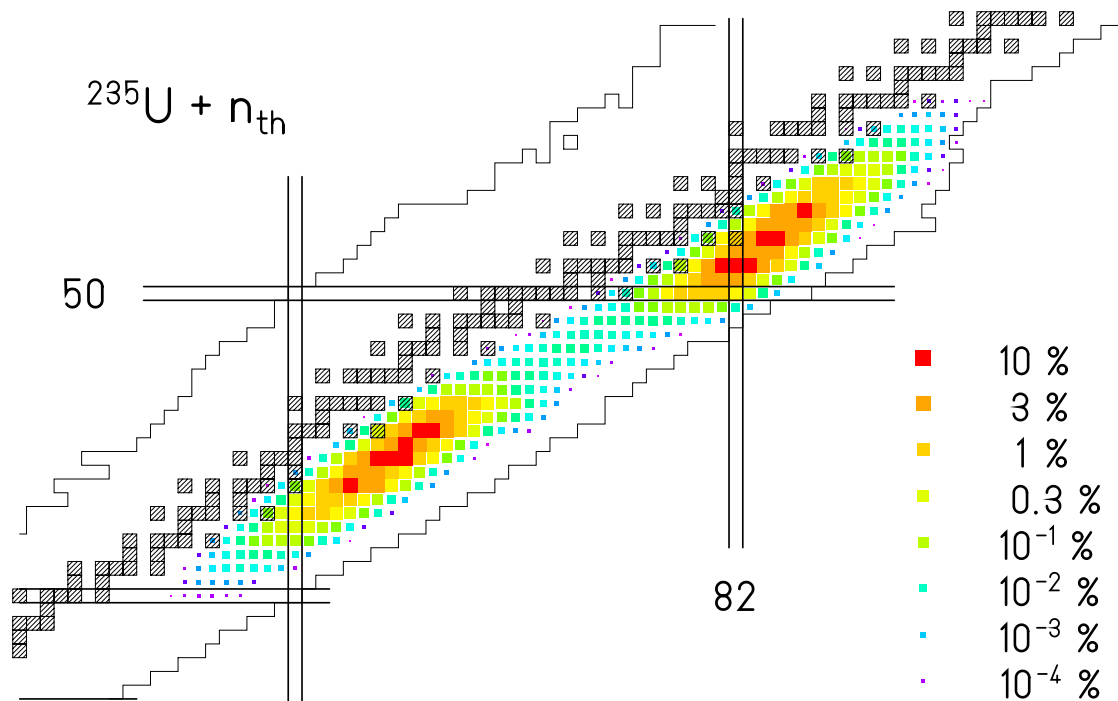


Figure 15: Nuclide distribution of fission products from thermal-neutron-induced fission of ^{235}U , calculated with the nuclear-reaction code ABRABLA. The colour code gives the yields in percent.

It has been discussed in the planning for RIA and it has become the standard option for SPIRAL 2 to apply a different converter scheme using deuterons as a primary beam, leading to a neutron flux of appreciably higher energies. The resulting energy spectrum of the

secondary neutrons peaks a bit below $\frac{1}{2}$ of the initial deuteron energy. That means that the neutrons essentially keep the velocity of the primary beam at the place in the target where they are produced. In addition, they are mostly forward directed. This is demonstrated in figure 16 for a primary beam energy of 50 MeV. Due to the condition that the deuterons should be stopped in the converter target and, at the same time, most of the neutrons should leave the converter target without being absorbed, the beam energy of this approach is limited to about 200 MeV. At SPIRAL 2, a primary energy of only 40 MeV will be used in order to limit the influence of neutron evaporation from the fragments, which tends to reduce the neutron excess of the fission products.

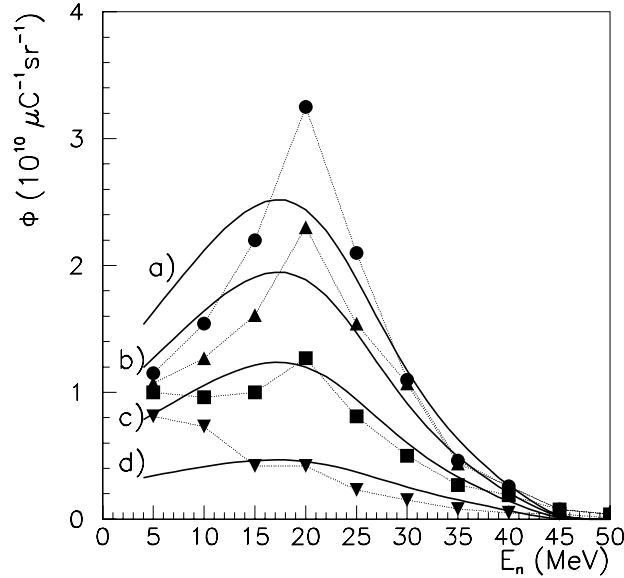


Figure 16: Experimental energetic distributions of neutrons for 50 MeV ^2H incident energy on a thick Be target: filled circles at 0° , filled triangles with up-point at 5° , filled squares at 10° , and filled triangles with down-point at 25° . These distributions are compared with theoretical ones: (a) 0° , (b) 5° , (c) 10° , and (d) 25° . The dotted lines are to guide the eye. The figure is taken from ref. [34].

The influence of the higher neutron energy on the nuclide distribution for the neutron-induced fission of ^{238}U can be seen by comparing figure 17, which shows the production in $^{238}\text{U} + 20$ MeV neutrons, with figure 9. Three tendencies can be observed: Symmetric fission is more abundant, populating the elements between $Z = 42$ and $Z = 48$ much more strongly, the distribution is slightly shifted to less neutron-rich isotopes, and the fission-fragment distribution extends slightly further to the heavier and to the lighter masses. As a consequence, the production of nuclei in the vicinity of ^{78}Ni seems to be slightly enhanced. The loss in the production of very neutron-rich isotopes is not strongly pronounced.

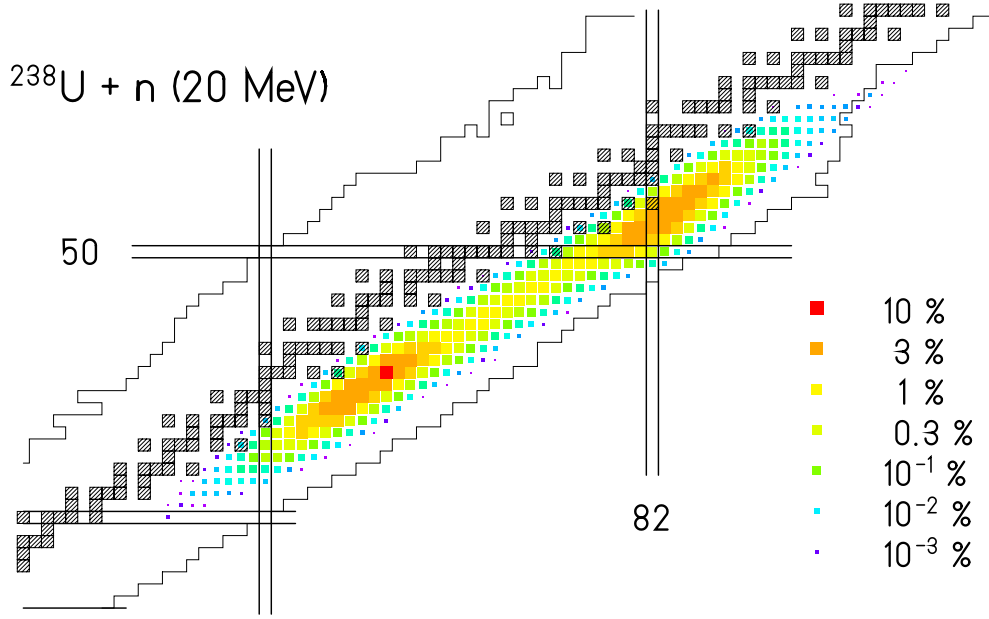


Figure 17: Nuclide distribution of fission products from the neutron-induced fission of ^{238}U for a fixed neutron energy of $E_n = 20$ MeV, calculated with the nuclear-reaction code ABRABLA. The colour code gives the yields in percent.

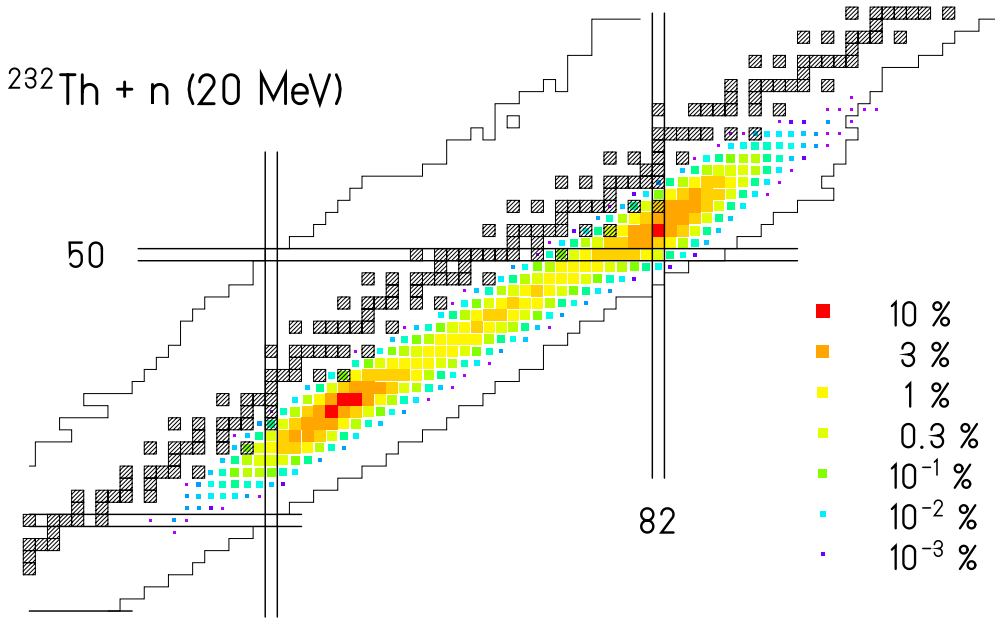


Figure 18: Nuclide distribution of fission products from the neutron-induced fission of ^{232}Th for a fixed neutron energy of $E_n = 20$ MeV, calculated with the nuclear-reaction code ABRABLA. The colour code gives the yields in percent.

If ^{232}Th is used as target material in the deuteron-converter option, the neutron excess of the fission fragments is expected to be reduced by about 0.4 units only. However, the light peak of the fission fragments moves to lighter masses by about 6 units. The calculated nuclide production for this option is shown in figure 18. This may be another efficient way to enhance the production of nuclides in the vicinity of ^{78}Ni . These conclusions are consistent with experimental results on fission-fragment mass distributions in proton-induced reactions with ^{232}Th and ^{238}U targets [35].

Quantitative considerations on these different options will be presented in section V.C.

D. Exploiting general properties of heavy-ion fragmentation reactions

It has been observed that the use of heavy ions instead of proton projectiles in irradiation of a uranium-carbide target gives enhanced yields and broader isotopic distributions for light ions. This finding is demonstrated in figure 19, which is taken from refs. [36, 37]. It seems that the production of long-lived isotopes close to beta stability is enhanced by a factor of 2 and 5 by ^3H and ^{12}C projectiles, respectively. As already discussed by Ravn et al. [3] the enhancement factor increases for the more neutron-rich isotopes. However, one should not forget that at the same time they are characterized by shorter life times. Thus, one may be cautious about concluding on the origin of the enhanced yields.

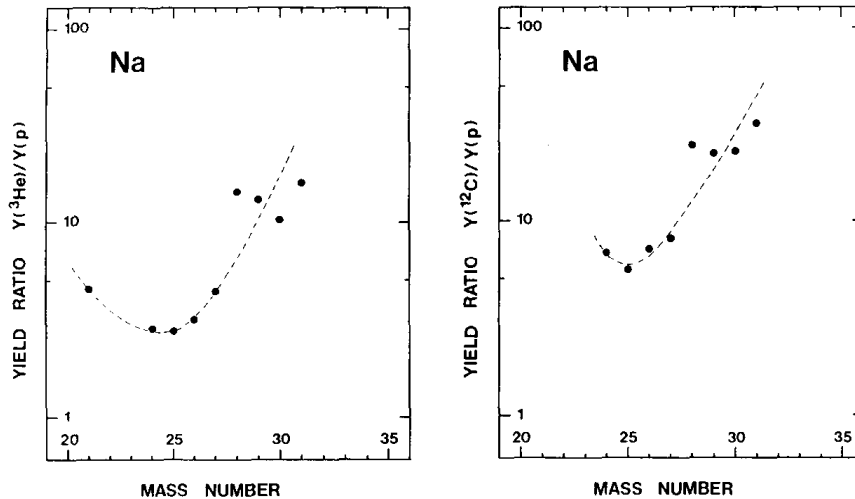


Figure 19: Ratio of sodium beam yields obtained by fragmentation of uranium by means of 300 A MeV ^3He and 78 A MeV ^{12}C over fragmentation of uranium by means of 600 MeV protons (figure taken from ref. [3]).

It is interesting that the pattern in the enhancement factor observed for the isotopes ^{28}Na ($t_{1/2} = 30.5$ ms), ^{29}Na ($t_{1/2} = 44.9$ ms), ^{30}Na ($t_{1/2} = 48$ ms), and ^{31}Na ($t_{1/2} = 17$ ms) does not scale with the mass number but rather with the inverse of the half life. Indeed, figure 20 illustrates, that the production ratios correlates with the half life rather well, in particular for the short-lived isotopes.

In this context, we would like to mention that from the operation of ISOLDE, a decay loss of more than a factor of hundred was deduced for ^{28}Na [11]. Thus, we are inclined to suppose that at least *part of the enhancement* and most of its variation with mass number is caused by *smaller decay losses in the heavy-ion-induced reactions*. The overall efficiencies for the long-lived isotopes seem also to be correlated with the different ranges of 600 MeV protons (321 g/cm²), 297 A MeV ^3He (82.2 g/cm²), and 78 A MeV ^{12}C (3.8 g/cm²) in uranium, which can be compared to the target thickness of 12 g/cm². (The range in carbon, which also is a constituent of the production target, scales about in the same way.) A similar correlation can be found with the specific energy-loss values in 1 mg/cm² uranium: 1.3 keV for 600 MeV protons, 6.8 keV for 297 A MeV ^3He , and 136 keV for 78 A MeV ^{12}C .

The actual physical reason for the presumed variation of the decay losses remains unclear to us. We may speculate that small differences in the parameters of the experiment, e. g. in the target temperature or in the beam position, might have caused slightly different overall-efficiency values and in particular different decay losses. In any case, it seems to us that such kind of differential measurements is a very demanding task.

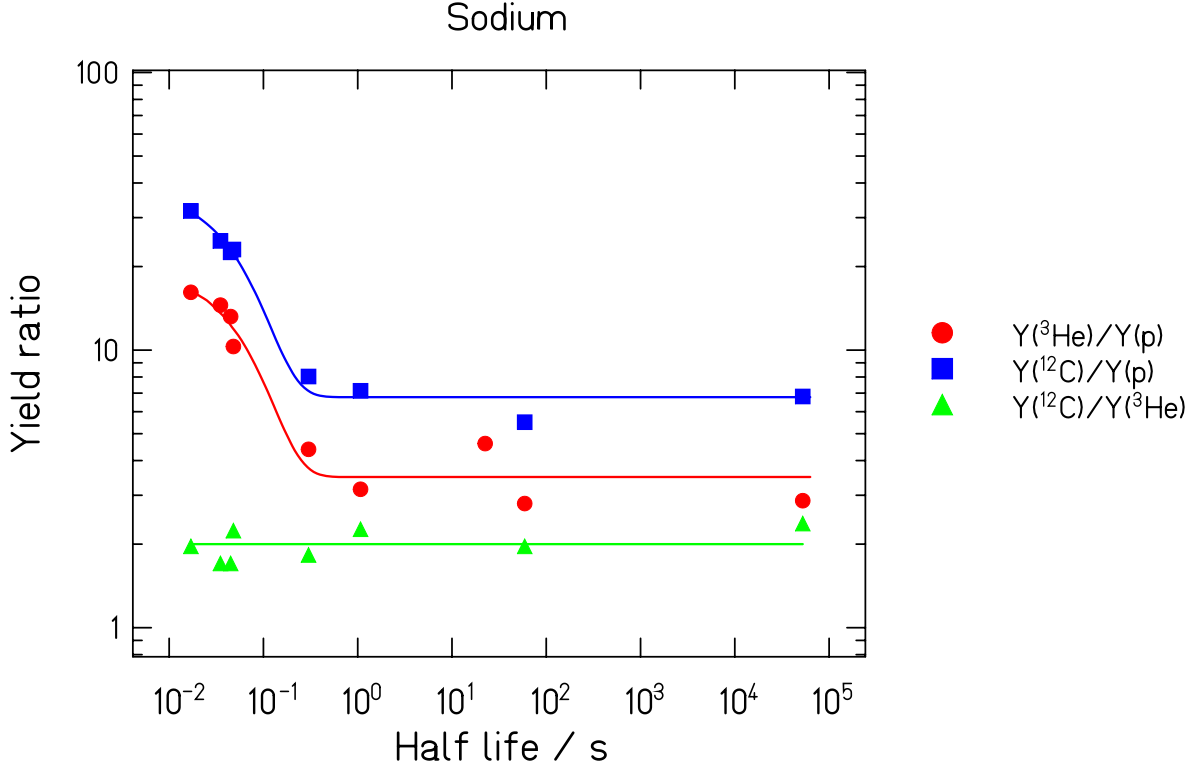


Figure 20: Measured yield ratios of different sodium isotopes obtained by fragmentation of uranium by 300 A MeV ^3He (dots) and 78 A MeV ^{12}C (squares) over spallation of uranium by 600 MeV protons as a function of the isotope half life. The ratio between ^{12}C - and ^3He -induced fragmentation (triangles) is also shown.

From the given dataset, it is difficult to draw a definite conclusion on the magnitude of the cross-section increase for the short-lived isotopes. The higher yields for long-lived isotopes, however, might partly be related to an enhanced emission of intermediate-mass fragments in the case of heavier projectiles, which will be discussed in section F.

E. Converter-catcher - ISOL scenario using heavy-ion beams

In the contribution of the Target-and-Ion-Source group (App C, Sect. 2.2.3) to the EURISOL RTD report [4], the use of an intense heavy-ion beam impinging on a beam-power-dissipating target of lithium or carbon, followed by a fragment catcher-diffuser, was proposed as an alternative option for the production of light exotic beams. This approach should offer several important advantages.

The first advantage of using heavy ions stems from the possibility of using different beams – including isotopically enriched projectiles – that can be optimally chosen. In this way, one could overcome the problems due to restrictions on the target material discussed in section III.A. The choice of a light heavy-ion beam can provide the opportunity to use nuclei which are in principle not available as target material, e.g. gaseous elements such as e.g. oxygen or noble gases. An example of such approach is the use of an ^{18}O beam for the production of ^{14}Be . In case of oxygen, there is a principal possibility to use an oxide target; however beryllium is not released from such target material. A similar possibility is the use of a neon beam for the production of the nuclide ^{18}Ne which is of interest for the beta-beam capability of EURISOL. A quantitative comparison for both cases will be carried out in section V.E.

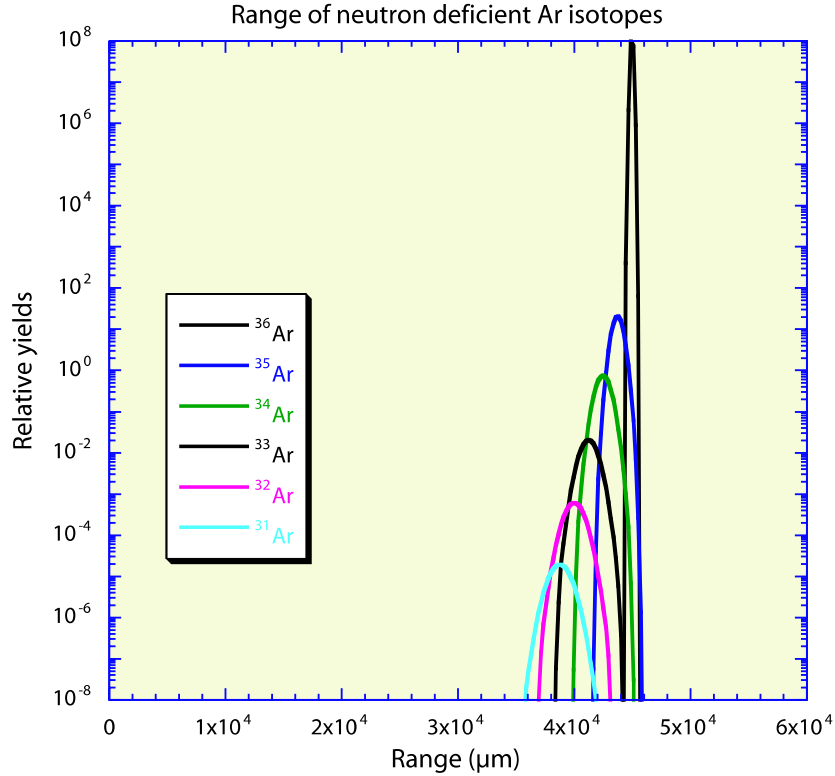


Figure 21: Range distribution of the projectile and neutron-deficient argon isotopes produced by fragmentation of a ^{36}Ar projectile at 333 A MeV. The decoupling of production target, diffusion catcher and beam stopper is possible even in the case of neutron-deficient reaction products, so that about 75% of the beam energy is deposited in the . The figure is taken from ref. [4].

Another important aspect is that for the production of non-fission fragments, this option also allows for splitting the target into two parts, with distinct functions, i.e. firstly a primary-to-secondary beam ‘converter’ target and secondly, a ‘catcher’ for the effusion-diffusion of the stopped secondary ions. Thus, the functions of the target for production and diffusion are separated, allowing one to choose the most resilient diffusion-catcher in most cases. Owing to the narrow forward emission cone of the fragments, the beam-power-dissipater/target may more easily be dissociated from the diffusion-catcher, and (because of their shorter range) this allows one to use a more compact catcher structure that enhances the effusion speed. The technique of a two-step target is already in use at the present SPIRAL facility for the production of $^{6,8}\text{He}$ isotopes in the reaction $^{13}\text{C} + \text{C}$. In this case, the first part stops the primary ^{13}C beam, and simultaneously produces $^{6,8}\text{He}$ by target fragmentation. The second part stops the projectile fragments of $^{6,8}\text{He}$. But also for neutron-deficient reaction products, the decoupling of production target, diffusion catcher and beam stop is possible, as demonstrated in figure 21. In this way, the respective optimum materials can be used for these different layers. A more quantitative discussion is given in section E.

F. Production of neutron-rich nuclides in heavy-ion reactions in the Fermi-energy regime related to nucleon exchange between target and projectile

Nucleus-nucleus collisions in the Fermi-energy domain exhibit a large variety of contributing reaction mechanisms and reaction products (see e.g. [38, 39, 40, 41]) and in principle offer the possibility to produce mid-heavy to heavy neutron-rich nuclei in very peripheral collisions. An enhancement of yields of neutron-rich nuclei was observed in the reaction $^{86}\text{Kr}+^{64}\text{Ni}$ [39] in very peripheral collisions (see figure 22), thus pointing to the possible importance of neutron and proton density profiles at the projectile and target surfaces. In the reactions of massive heavy ions such as $^{86}\text{Kr}+^{124}\text{Sn}$ [42] and $^{124}\text{Sn}+^{124}\text{Sn}$ [43] (see figure 23), an enhancement was observed over the yields expected in cold fragmentation

(as represented by EPAX [44] estimates), which is at present, besides fission, the method of choice to produce neutron-rich nuclei.

Peripheral nucleus-nucleus collisions can be described theoretically using the model of deep-inelastic transfer in combination with an appropriate model of de-excitation. Very good description of experimental observables was obtained [39, 42, 43, 45] using the DIT model implementation of Tassan-Got [46] and the de-excitation codes SMM [47] and GEMINI [48]. SMM represents the statistical model of multifragmentation (SMM) and GEMINI invokes the model of sequential binary decay (SBD). An enhancement of the yields of neutron-rich nuclei over the prediction of such calculations was observed in the reaction $^{86}\text{Kr}+^{64}\text{Ni}$ at 25 A MeV [39]. The shapes of the velocity spectra suggested a process with a short timescale such as very peripheral collisions where the details of neutron and proton density profiles at the projectile and target surfaces can play a significant role. In [49], the DIT model of Tassan-Got [46] was supplemented with a phenomenological correction introducing the effect of shell structure on the nuclear periphery. A consistent agreement with experimental data is achieved using a modified DIT calculation combined with SMM in the reactions of a 25 A MeV ^{86}Kr beam with three different target nuclei, specifically allowing to describe the deviation of the nucleon exchange from the path toward isospin equilibration.

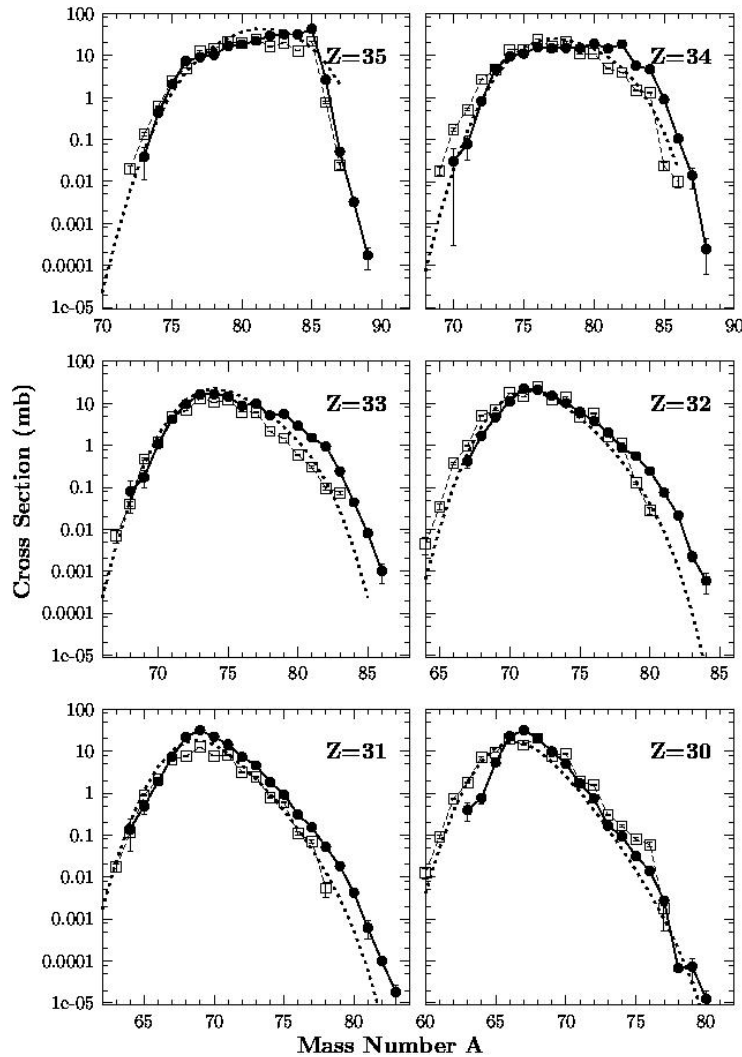


Figure 22: Observed excess of neutron-rich nuclei in the reaction $^{86}\text{Kr}+^{64}\text{Ni}$ at 25 A MeV [39]. Solid symbols - experimental data, open symbols - DIT+Gemini, dashed line - EPAX.

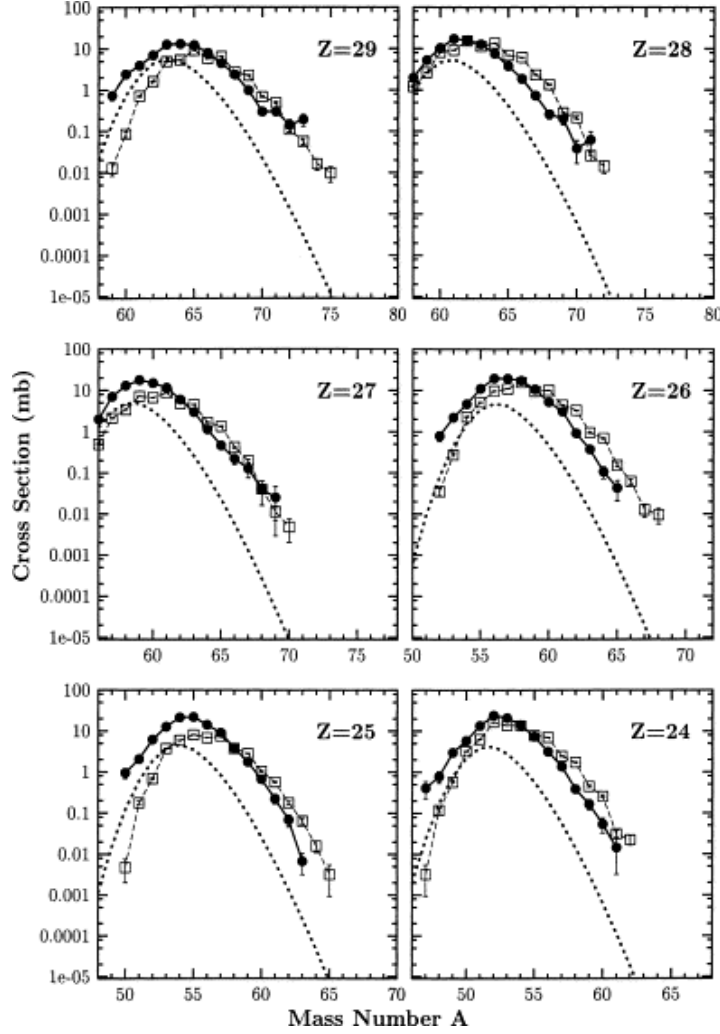


Figure 23: Observed excess of neutron-rich nuclei in reaction $^{124}\text{Sn}+^{124}\text{Sn}$ at 20 A MeV [43]. Solid symbols - experimental data, open symbols - DIT+Gemini, dashed line - EPAX.

The success of the modified DIT calculation can be explained as a correction reflecting the modification of neutron and proton transfer probabilities in peripheral collisions due to the effect of shell structure on isospin asymmetry at nuclear periphery. Such an effect is of interest for better prediction of the production rates of exotic nuclei with a wide range of N/Z -ratios at the new generation of rare-isotope beam facilities.

The limitation of the method is mainly in the use of relatively thin targets, and thus the method can only be effective under conditions where the losses of the in-target yield can be compensated by a corresponding increase of production cross sections. Such situation can be typically anticipated for the most neutron-rich products, where the excess of production cross sections over the relativistic fragmentation appears to increase. Within the context of EURISOL, the reactions in the Fermi-energy domain correspond to a driver-accelerator scenario with beam energy 28 A MeV for $A/q = 3$. A complementary scenario for use of reactions in the Fermi-energy domain at EURISOL would be the acceleration of stable heavy-ion beams at the planned EURISOL post-accelerator, which is a matter of discussion.

V. QUANTITATIVE INVESTIGATIONS

A. Extending the mass range of spallation-evaporation residues by 2 GeV primary beam

In the following, the benefit of a 2 GeV primary beam in comparison with a 1 GeV proton beam will be investigated quantitatively. The option considered here is a 2 GeV $^3\text{He}^{2+}$

beam, mostly because of its low $A/q = 1.5$ value. The quantitative answer will be based on two criteria. The first one considers the mass range populated in a spallation-evaporation reaction, while the second one focuses on the population of the most neutron-deficient isotopes. The first criterion exploits the fact that the mass range populated in a spallation-evaporation reaction grows with increasing beam energy. Thus, the beam energy should be chosen in a way that the production bridges the gaps between the different usable target materials. The second criterion is based on the observation that the nuclide production by spallation evaporation on the average shifts to the neutron-deficient side with increasing mass loss in the reaction. Thus, it might be advantageous to use a heavier target material for the production of the most neutron-deficient isotopes of a certain element.

We would like to mention that the choice of the specific target materials and the secondary products in this section is made independently of the technical conditions and the scientific interest. The case study should just show the principal tendencies which one has to face due to the relatively large gaps between different usable target materials, which typically extend up to 20 elements.

1. Filling gaps in mass

When the desired product is far below the target, the production cross section depends strongly on the projectile energy. That means that higher incident energy enhances the production significantly.

Figure 24 investigates the benefit of a 2 GeV primary beam on the production cross sections on the basis of experimental data. It shows a comparison of the measured production cross sections of hafnium ($Z = 72$) and samarium isotopes ($Z = 62$) in a lead target ($Z = 82$) by 1 GeV protons and 2 GeV deuterons. These elements are 10, respectively 20, units below the target. Actually, the measurement has been made in inverse kinematics. Model calculations suggest that the nuclide production by 2 GeV deuterons is rather similar to the one expected for 2 GeV ^3He projectiles.

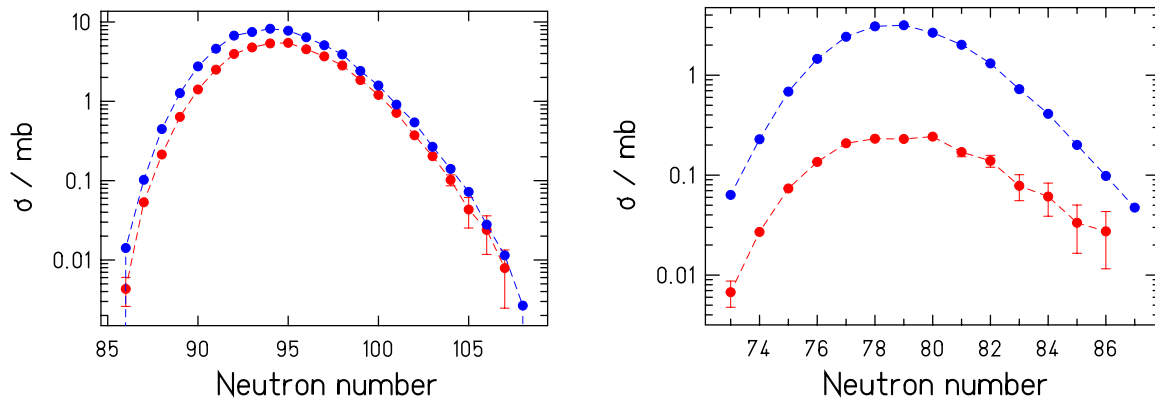


Figure 24: Production cross sections of hafnium isotopes ($Z = 72$), left, and samarium isotopes ($Z = 62$), right, in the spallation-evaporation reaction of ^{208}Pb with 1 GeV protons (red symbols) and 2 GeV deuterons (blue symbols) measured in inverse kinematics at GSI, Darmstadt.

Obviously, there is no benefit for the production cross sections of hafnium but a clear advantage for the production cross sections of samarium in favour of the 2-GeV option by more than an order of magnitude. The advantage decreases for the most neutron-rich isotopes.

For a realistic estimate of the problem, the yields in a thick target with appropriate beam intensities should be considered. The variations of the beam energy due to electronic energy loss and of the intensity due to beam attenuation by nuclear reactions are depicted in figure 25.

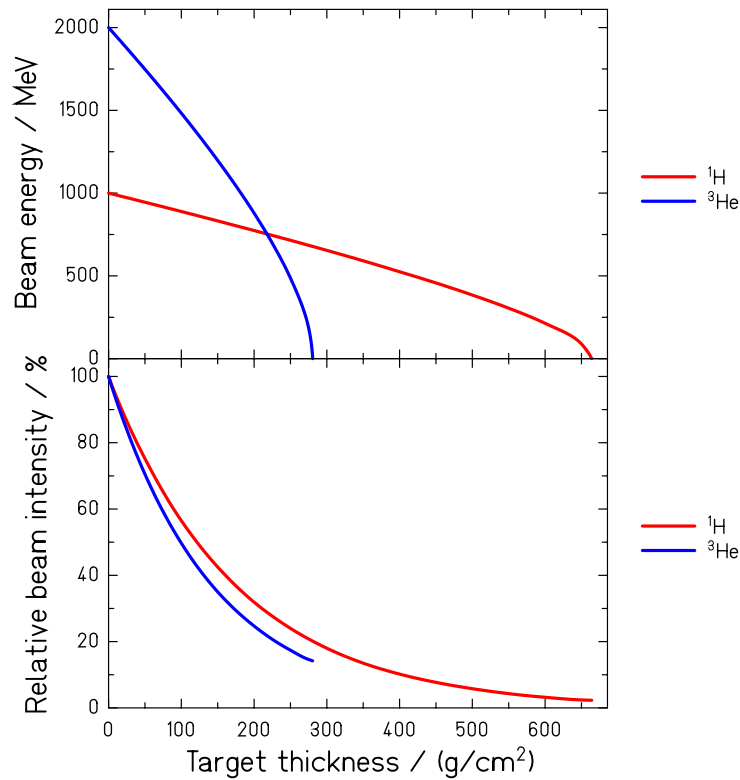


Figure 25: Evolution of beam energy and relative intensity along a thick uranium target for a 1 GeV proton and a 2 GeV ^3He beam, from calculations performed with the code AMADEUS [55].

In a target of 100 g/cm^2 , about half the projectiles interact in both cases, in 200 g/cm^2 , the fraction is about 75%. This defines a rough limit on the target thickness which is about the same for these two beams. After a 100 g/cm^2 target, the protons have a final energy of about 890 MeV, ^3He projectiles leave the target with about 1500 MeV. The energy per projectile deposited in the same target is about a factor of four higher for ^3He . This might impose a limitation on the beam intensity of ^3He (counted in projectile per second), which should be a factor of 4 weaker than the hydrogen beam, if the same deposited energy is considered. Considering the 100 kW (100 μA) proton beam according to the baseline option of EURISOL, this would imply a 50 kW (25 μA) ^3He beam.

If the production of a specific nuclide is considered, the situation complicates further. With the energy loss after a certain distance in the target, the production of the desired element may drop severely, and any further target length becomes useless. However, the slope of the production cross-section dependence on energy is higher at lower energies, and enters into saturation at higher energies, depending on the difference in Z between the target and the product.

The example of rhenium produced in a uranium compound target, which is a critical case for EURISOL on the basis of the presently available targets, is investigated in the following in some detail. Rhenium is 17 units below uranium. To take the energy loss and the attenuation of the beam into account in a simplified way, one can divide the target in three equal parts of 80 g/cm^2 each, and calculate the production rate in each segment. Using production cross sections calculated at the energy of the beam in the middle of each segment, and taking into account the reaction rate in each segment, one can obtain production rates for all the isotopes. The results in the case of the 1 GeV proton and 2 GeV ^3He beams are shown in the figure 26. The decrease in the production rates along the target is visibly sharper in the case of the proton beam.

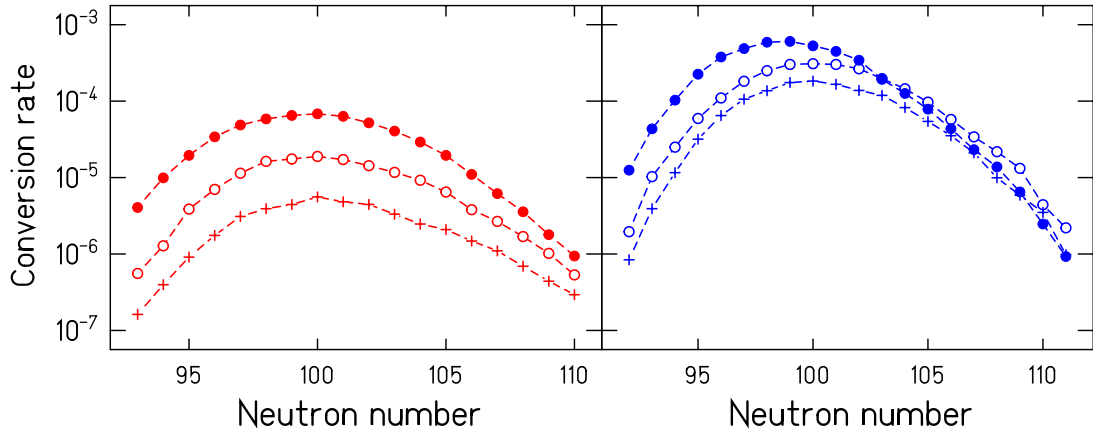


Figure 26: Production rates per beam particle for rhenium isotopes in three consecutive uranium-carbide target segments of equal thickness (80 g/cm^2). Left – the case of the proton beam; right – the case of ^3He beam. Solid dots - first segment, open dots - second segment, crosses – third segment.

Figure 27 compares the maximum obtainable in-target production rates for the isotopic chains of this element with a 1 GeV proton beam or a 2 GeV ^3He beam. The beam power has been normalized to 100 kW for protons and 50 kW for ^3He .

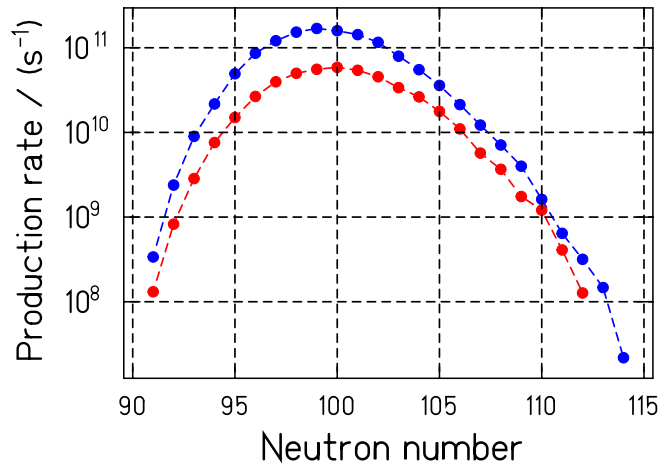


Figure 27: In-target production rates of rhenium isotopes ($Z = 75$), in the spallation-evaporation reaction induced by 1 GeV protons (red symbols) and 2 GeV ^3He (blue symbols) in a uranium-carbide target. The beam power is 100 kW in the case of protons, and 50 kW in the case of ^3He . A target thickness of 240 g/cm^2 was assumed, beyond which the production drops significantly in both cases.

Obviously, there is a clear advantage for the production of rhenium in favour of the 2-GeV option by up to a factor of 4, if the same thermal power density in the production target is considered. The gain factor decreases slightly towards the neutron-rich wing of the isotopic distribution.

The final judgement on the benefit of a 2 GeV ^3He beam depends on the size of the gaps in the choice of suitable target material for EURISOL. If suitable targets can be provided by future research and development all over the element range with a maximum distance of 10 units, there is no gain by the 2 GeV option with respect to the 1 GeV option. If one expects that gaps of about 20 units remain, there is a clear advantage up to a factor of 4 for the 2-GeV primary beam for certain elements due to the extended mass range produced.

2. Production of neutron-deficient isotopes

To investigate the population of the most neutron-deficient isotopes as the second criterion for the benefit of a 2 GeV primary beam, calculations on a hypothetical production of isotopes of the element number 60 in targets with element numbers 62, 72, and 82 were performed, as an example which does not necessarily obey the operating conditions for ISOL targets. Natural isotopic composition of the target materials was assumed. Again, the answer on the benefit of the 2 GeV beam depends on the size of the gap between available target materials. A target element slightly above the desired product gives the highest cross section for all isotopes, including the most neutron-deficient ones. In figure 28, this is the case of the target with $Z = 62$ for the production of the isotopes of the element number 60.

If the target element is 22 units above the desired product, the use of the 2 GeV beam has a clear advantage, as already seen in the previous figure. However, a combination of a target element 12 units above the desired fragment with the 1 GeV proton beam gives higher yields, if such target is available. Thus, if a reasonably close-by target is available, there is no specific gain for the production of neutron-deficient isotopes from increasing the beam energy and using a heavier target.

Our calculations with targets of different isotopic composition have shown that one should prefer to use an even element as a target material if exotic nuclei are to be produced, because it provides a larger number of primordial isotopes. This way, one obtains higher yields on both wings of the isotopic distributions, on the neutron-deficient and on the neutron-rich side.

The results obtained here for the production of neodymium isotopes in targets with Z by 2, 12, and 22 units larger can be generalized to other cases. The essential parameter is the difference in atomic number of the target material and the fragment to be produced.

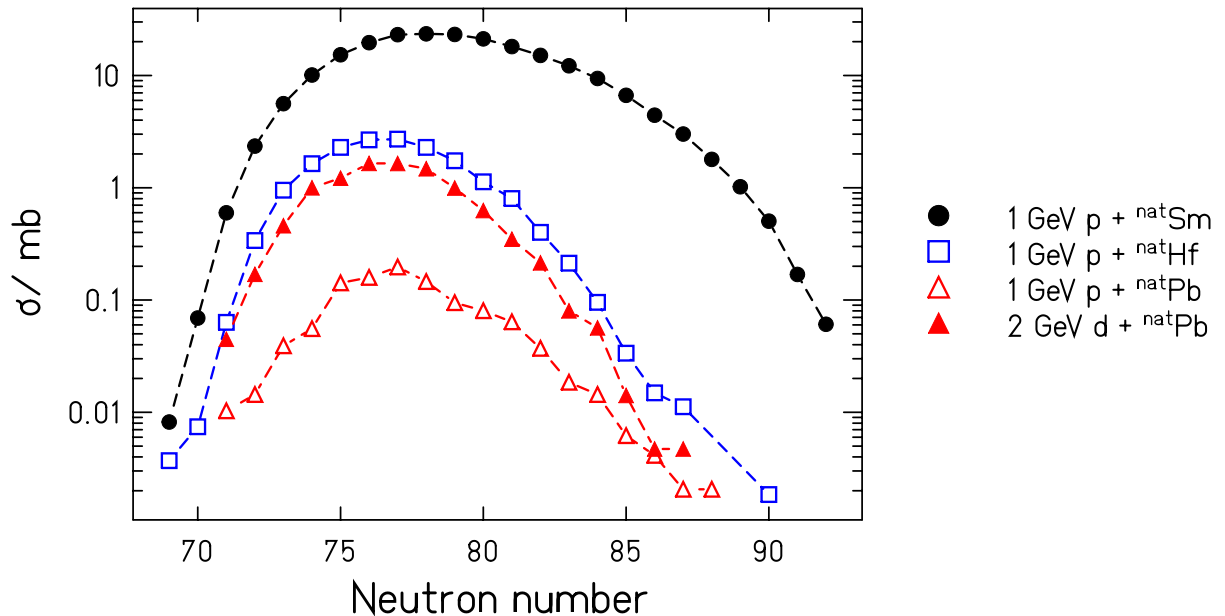


Figure 28: Calculated production cross sections of neodymium ($Z=60$) isotopes by 1 GeV protons and 2 GeV deuterons in different targets.

B. Enhancing IMF emission to produce light neutron-rich nuclides by a 2 GeV primary beam in a fissile target

The production of intermediate-mass fragments ($Z < 30$) in a fissile target is sensitive to the energy of the incoming beam. To investigate to what extent their production cross sections can be enhanced we have performed calculations on nuclide production in a ^{238}U

target using 1-GeV proton, 2-GeV proton and 2-GeV ${}^3\text{He}^{2+}$ beams. The calculations were performed using the intranuclear-cascade code INCL4 [50] coupled to ABLA.

In figure 29 we compare production cross sections of sodium calculated using the above mentioned options. In the same figure, we also show the sodium cross sections from the reaction ${}^1\text{H}(1\text{ GeV}) + {}^{238}\text{U}$ measured at GSI [18]. According to these calculations, with 2-GeV protons we can expect an increase in the production cross sections of IMFs by a factor of ~ 4 as compared to the 1-GeV-proton option. Moreover, if one would use ${}^3\text{He}$ beam one could expect ~ 8 times stronger production of IMF as compared to the standard 1-GeV option. The reason for the additional increase of the IMF production cross-sections with the use of ${}^3\text{He}$ projectiles is related to the stronger population of the high-energy tail of the calculated prefragment excitation-energy distribution due to the larger size of the projectile. However, there are no analyzed experimental data yet to confirm this effect.

We can conclude that an increase in the primary-beam energy and, additionally, the use of a ${}^3\text{He}^{2+}$ beam could lead to an important enhancement of light, $Z < 30$, neutron-rich nuclides. If normalized to the same deposited thermal energy, this enhancement amounts to a factor of about 2.

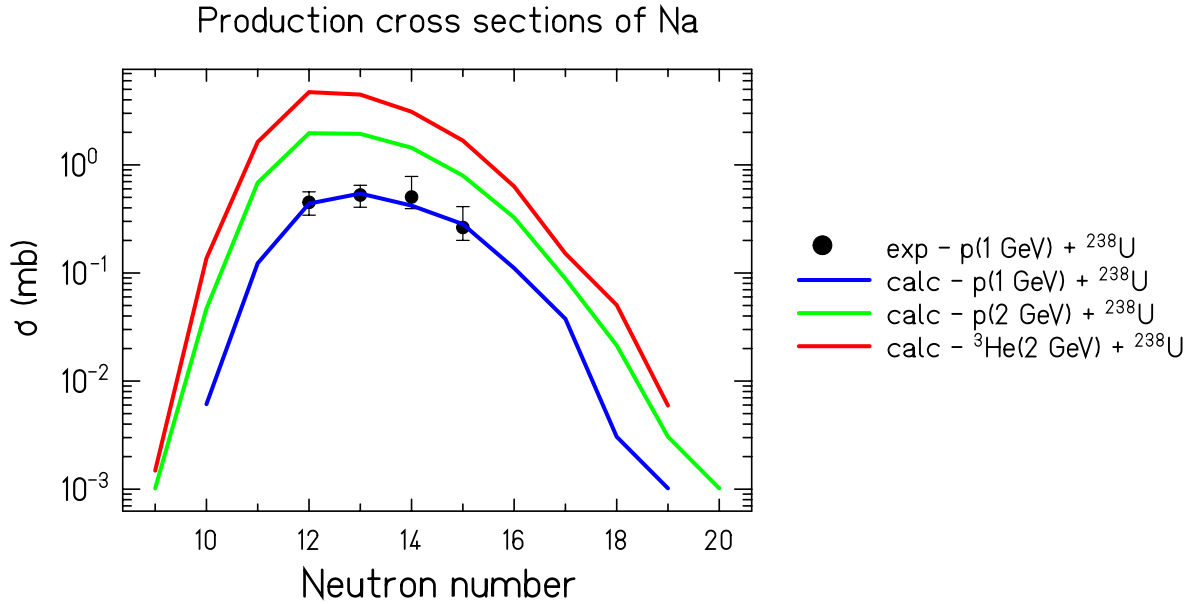


Figure 29: Production cross sections of Na in a ${}^{238}\text{U}$ target interacting with different beams – 1-GeV protons (blue line), 2-GeV protons (green line) and 2-GeV ${}^3\text{He}$ (red line). Calculations were made with INCL4+ABLA. For comparison, also sodium production cross sections (points) measured at GSI in the reaction ${}^1\text{H}(1\text{ GeV})+{}^{238}\text{U}$ [18] are shown.

C. Extended fission-fragment distribution in neutron-induced fission with $E_n = 20$ to 80 MeV by the deuteron-converter option

The converter-target option using deuterons below 100 MeV differs essentially in one aspect from the 1 GeV proton indirect-target option: The deuteron option provides neutrons with an energy spectrum peaking at about 40% of the primary-deuteron energy, while the proton option is dominated by evaporated neutrons in the few-MeV range. The different neutron-energy distributions lead to different nuclide distributions. In the deuteron case the gap close to symmetric fission around $Z = 46$ is filled up, and the elements in both tails below $Z = 30$ and above $Z = 62$ are populated more strongly. All these aspects were already qualitatively discussed in section IV.C. The small component of ${}^{235}\text{U}$ in the natural or depleted uranium does not affect the fission rate or the fission fragment distribution in this case, as there are few thermal neutrons in the spectrum.

In figure 30, we present the results of model calculations for the nuclide yields in the deuteron-converter option with a 50 MeV ^2H primary beam and a target geometry similar to the one used in SPIRAL 2, which mostly exploits the neutrons produced close to the projectile velocity in forward direction.

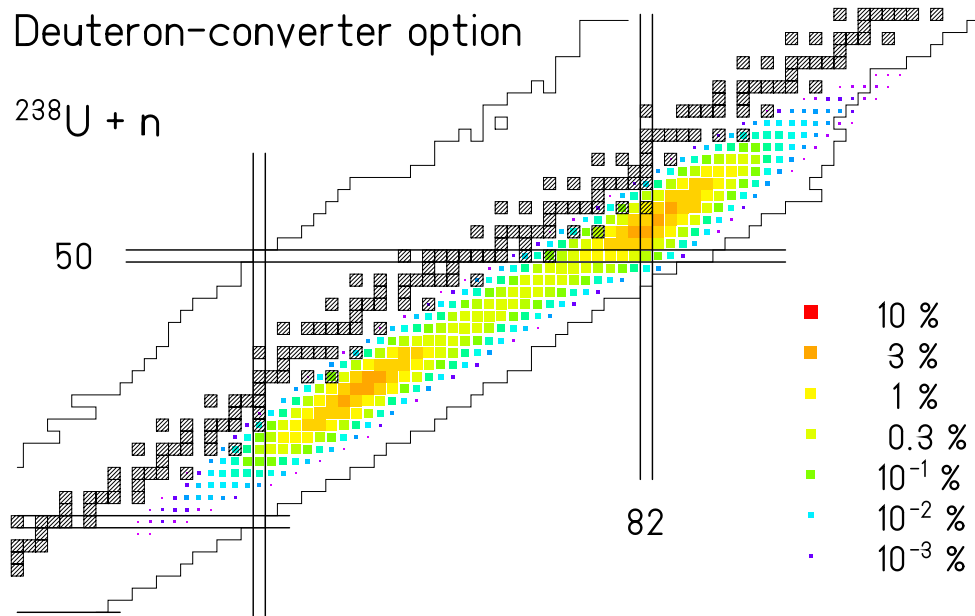


Figure 30: In-target nuclide-production yields for a 50-MeV deuteron-converter option, calculated with the ABRABLA code, using the experimental neutron-energy spectrum of Figure 16 within 10° emission angle. The yields are given on a colour scale in percent. The in-target production rates are obtained by multiplying these yields with the total fission rate.

The EURISOL high-power target is presently designed to support a primary beam of 1 GeV protons with a beam current of 3 to 4 mA, corresponding to a beam power of 3 to 4 MW [51]. Fission rates induced by secondary neutrons of about 10^{15} s^{-1} are expected.

In-target nuclide-production rates estimated with the ABRABLA code for the EURISOL high-power target are shown in figure 31.

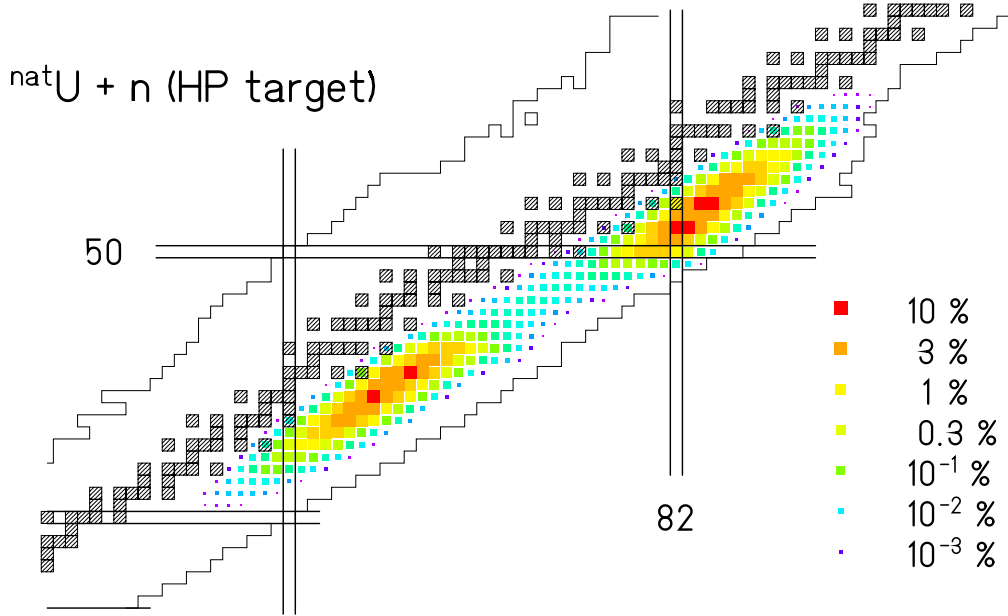


Figure 31: In-target nuclide-production yields for the EURISOL 2-step high-power option calculated with the ABRABLA code, using the neutron-energy spectrum of figure 14. The yields are given on a colour scale in percent. The in-target production rates are obtained by multiplying these yields with the total fission rate.

For orientation, we mention two other cases, where deuterons are considered as primary beam:

SPIRAL 2 is designed to provide a fission rate of 10^{13} s^{-1} with a standard low-density UC_x target with a 40 MeV deuteron beam of 5 mA or $3 \cdot 10^{16}$ projectiles per second [52]. This corresponds to a beam power of 200 kW. The fission rate could be increased to 10^{14} s^{-1} by using a high-density UC_x target.

In the latest design of RIA, a deuteron beam up to 800 MeV and a beam power of 100 kW were considered as one of the different options, however due to the target geometry used, the production mostly relies on fission by low-energy neutrons in the range from 0.5 to 3 MeV [53], thus providing a nuclide distribution similar to the one obtained in the EURISOL proton-converter option.

D. Exploiting general properties of fragmentation induced by heavy-ions

As discussed in section IV.D, it is not clear to which extent the broader isotopic distributions in fragmentation reactions induced by heavy-ions are caused by nuclear-reaction aspects, as supposed by H. Ravn et al. [9], and to which extent by lower decay losses.

A detailed study of the decay-loss aspect of the increase in the exotic-nuclide yields is beyond the scope of this work. Such a detailed study should carefully consider the different limitations e.g. due to available primary-beam intensities and the affordable heat load in the target.

E. Converter-catcher - ISOL scenario using heavy-ion beams

Estimates of attainable secondary-beam intensities by fragmentation of heavy ions were performed by the Target-and-Ion-Source group in the EURISOL RTD report (App C, Sect. 2.2.3) [4]. The calculations were performed assuming a multi-beam driver delivering any heavy-ion beam with $A/q = 6$ or 3 at 166.5 or 333 A MeV, respectively, impinging on a carbon converter target in which most of the primary beam energy is lost. The reaction products are stopped in a diffusion-catcher, which is placed just behind the converter and can in principle consist of any material, providing best conditions for extraction. The production

yields were only calculated in the target/converter, because only projectile fragmentation was taken into account. The projectile fragment yields were estimated via EPAX. Losses by secondary reactions of the radioactive fragments during their stopping and transfer were taken into account. The best suited projectile was chosen, to give the highest yields and for which ECRIS technology may be able to deliver the beam intensities assumed for the calculation. Results of these calculations showed that for light nuclei heavy ions present an advantage of one order of magnitude or more in the in-target yields (essentially due to the use of the 2-step method for light heavy ions) as compared to the proton option. For heavier ions, this advantage disappears and the tendency is reversed. Thus, a multi-beam driver, with both protons and light heavy ions could extend the 2-step concept over the whole chart of the nuclides.

The gain factors obtained in this elaborate study for the in-target production, compared to a 100 μ A beam of 1 GeV protons, are shown in figure 32. These gain factors should be adapted to the primary-beam intensities and energies obtainable with the different scenarios for the driver accelerator listed in chapter 2. Disregarding changes in the nuclear-reaction characteristics that may be important below 100 A MeV and in primary-beam intensities, the nuclide production rates scale with the usable target thickness, which strongly reduces with decreasing beam energy. Table I lists the energies and the intensities of the heavy-ion beams used for the comparison shown in figure 32. It also provides the relative energy-dependent scaling factors for the total nuclear-reaction rates in the usable target thickness for a few beam-energy values. E.g. the in-target production yields with a ^{40}Ar beam decrease by a factor of about $0.5/40 = 1/80$, when the energy is reduced from 333 A MeV, as assumed for the estimations in figure 32, to 28 A MeV, if the beam intensity remains the same.

Finally, it should be stressed that the investigations of the Target-and-Ion-Source group reported in the EURISOL RTD report (App C, Sect. 2.2.3) [4] assumed that the overall ISOL efficiencies obtainable for the heavy-ion fragmentation option are equal to the ones of the 1-GeV proton direct-target option of EURISOL. For a general judgment on this assumption, appropriate targets, which are now available at SPIRAL mostly for noble gases, need to be developed for the heavy-ion fragmentation option.

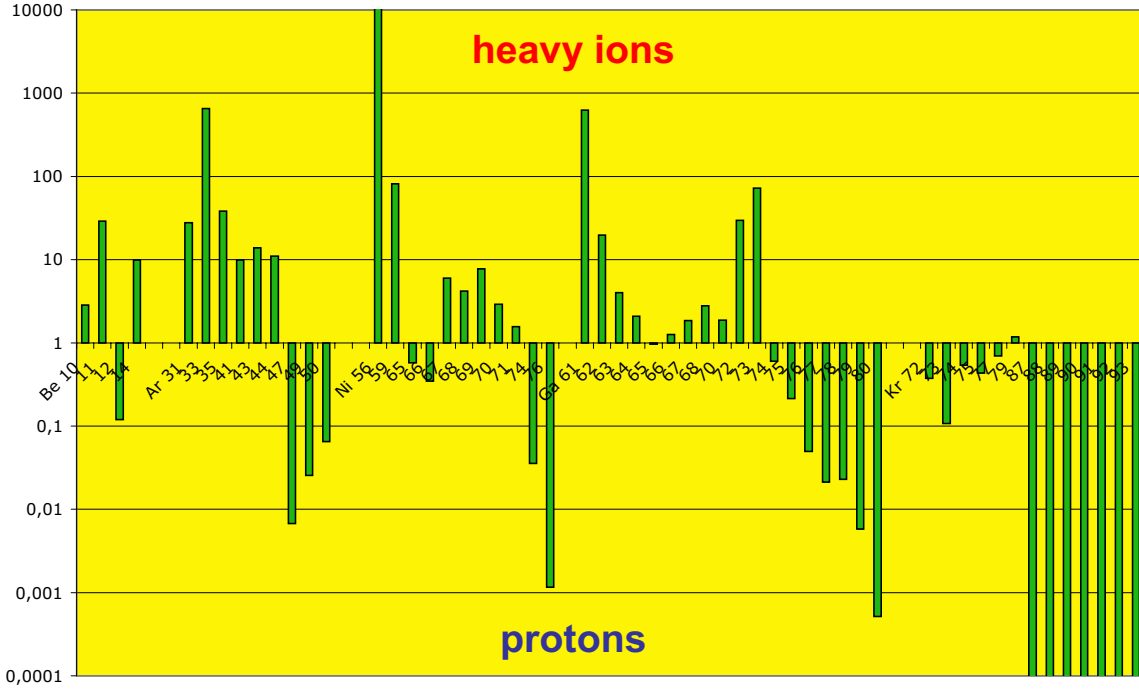


Figure 32: Comparison of the yields obtained with a beam of 1 GeV protons (100 μ A) and with an optimum heavy-ion beam. See text for details. The figure is taken from ref. [54].

Table I: Possible secondary-beam intensities using different primary heavy-ion-beam energies. Intensities used for the comparison shown in figure 32 are printed in bold. In addition, the fraction R of the projectiles, which react in a carbon target with a thickness of 80% of the range of the corresponding projectile, are given for different beam energies. The beam energies correspond to the different EURISOL driver-accelerator options listed in chapter 2.

Projectile	E/A [MeV]	I [mA]	R [%]						
			28 A MeV	43 A MeV	100 A MeV	150 A MeV	166.5 A MeV	333 A MeV	500 A MeV
B, C, O	333	1	~ 1	~ 2	8 – 13	16 – 24	20 – 28	50 – 65	73 – 86
Ar	333	0.5	~0.5	~ 1	~ 7	~ 13	~ 15	~ 40	~ 63
Ca	166.5	0.5	~ 0.5	~ 1	~ 5	~ 11	~ 13	~ 35	~ 55
Ni	166.5	0.4	~ 0.5	~ 1	~ 5	~ 10	~ 12	~ 32	~ 53
Zn, Ge, Se, Kr, Ag	166.5	0.3	~ 0.5	~ 1	~ 5	9 – 10	~ 11	~ 30	~ 50
Sn, Te, Xe, Ce	166.5	0.2	~ 0.5	~ 1	~ 4	~ 8	~ 10	~ 27	~ 45

As a specific case, we have also considered more closely the production of ^{18}Ne in three different scenarios as presented in Table II. The option using a ^{20}Ne beam is compared to two scenarios using a 1 GeV proton beam with MgO and Ta targets. The initial energy of 500 A MeV and $A/q=2$, as envisioned in one of the options for the EURISOL driver accelerator, was considered. In the case of light targets, the production cross sections are estimated using the EPAX parameterization. The validity of EPAX for the reaction channels close to the projectile was verified using data from the reaction $^{56}\text{Fe}+p$, obtained at GSI [13,

14]. EPAX reproduces the data well for channels with removal of up to six protons independent of the beam energy which ranged from 300 to 1500 A MeV. For heavier targets (tantalum and uranium carbide), the ABRABLA code was used to calculate the cross sections. The reaction rate in the target was estimated using the code AMADEUS [55], taking into account that, in the case of a heavy target, only a limited target thickness is usable due to the rapidly decreasing cross sections for complex-fragment emission. An initial beam intensity of 100 kW was considered in all cases. The overall efficiency for the reaction $^{20}\text{Ne}+\text{C}$ was estimated using the experimental value obtained at SPIRAL, as reported in [4] (App C, Sect. 2.2.3), for the reaction $\text{p}+\text{MgO}$ the same value was taken, assuming an ECRIS ion source, and for the reaction $\text{p}+\text{Ta}$ a rough estimate was obtained using the formula of Lukić et al. [11] with global parameters $\varepsilon_s=0.3$, $\alpha=1.5$ and $\tau=2\text{s}$. The estimated RIB intensity for the reaction $^{20}\text{Ne}+\text{C}$ thus appears lower by a factor of 4 than for the reaction $\text{p}+\text{MgO}$ and more than two orders of magnitude higher than in the reaction $\text{p}+\text{Ta}$.

Table II: Comparison of several scenarios to produce a ^{18}Ne secondary beam. See text for details.

<i>Reaction</i>	$\sigma(^{18}\text{Ne})$ [mb]	<i>Reaction rate in target</i> [%]	<i>In-target yield</i> [s^{-1}] ($I=100\mu\text{A}$)	<i>Overall efficiency</i> [%]	<i>RIB intensity</i> [s^{-1}]
$^{20}\text{Ne}+\text{C}$	2.3	73	3.4×10^{10}	3	1×10^9
$\text{p}+\text{MgO}$	0.7	97	1.4×10^{11}	3	4.2×10^9
$\text{p}+\text{Ta}$	4.9×10^{-4}	30	3.0×10^7	13	4.0×10^6

In Table III we show the results of a similar quantitative comparison of the production of the nuclide ^{14}Be for the reaction $^{18}\text{O}+\text{C}$ and several proton-beam options that are currently used at ISOLDE (Ta, UC_x), or could be used in the future (SiC_x). The production cross sections and the reaction rates in the target were estimated in the same way as in the previous example. As no data on the overall efficiency of beryllium are available, only in-target yields are considered. The energy of the ^{18}O beam was again fixed to 500 A MeV, even if the A/q ratio in this case slightly exceeds two. The RIB intensity in the heavy-ion option is two orders of magnitude higher than with the silicon-carbide or tantalum targets, but of the same order of magnitude as with the uranium-carbide target.

Table III: Comparison of several scenarios to produce a ^{14}Be secondary beam. See text for details.

<i>Reaction</i>	$\sigma(^{14}\text{Be})$ [mb]	<i>Reaction rate in target</i> [%]	<i>In-target yield</i> [s^{-1}] ($I=100\mu\text{A}$)
$^{18}\text{O}+\text{C}$	2×10^{-3}	85	3.4×10^7
$\text{p}+\text{SiC}_x$	2.6×10^{-6}	97	4.8×10^5
$\text{p}+\text{Ta}$	3.3×10^{-6}	30	2.2×10^5
$\text{p}+\text{UC}_x$	1.7×10^{-4}	30	1.2×10^7

Thus, in both cases, the heavy-ion option with 500 A MeV seems to lead to a final RIB intensity of the same order of magnitude as in the baseline scenario, with a carefully chosen target, and does not offer any significant advantage over it.

F. Production of neutron-rich nuclides in heavy-ion reactions in the Fermi-energy regime related to nucleon exchange between target and projectile

1. Production cross sections of nuclei around ^{78}Ni

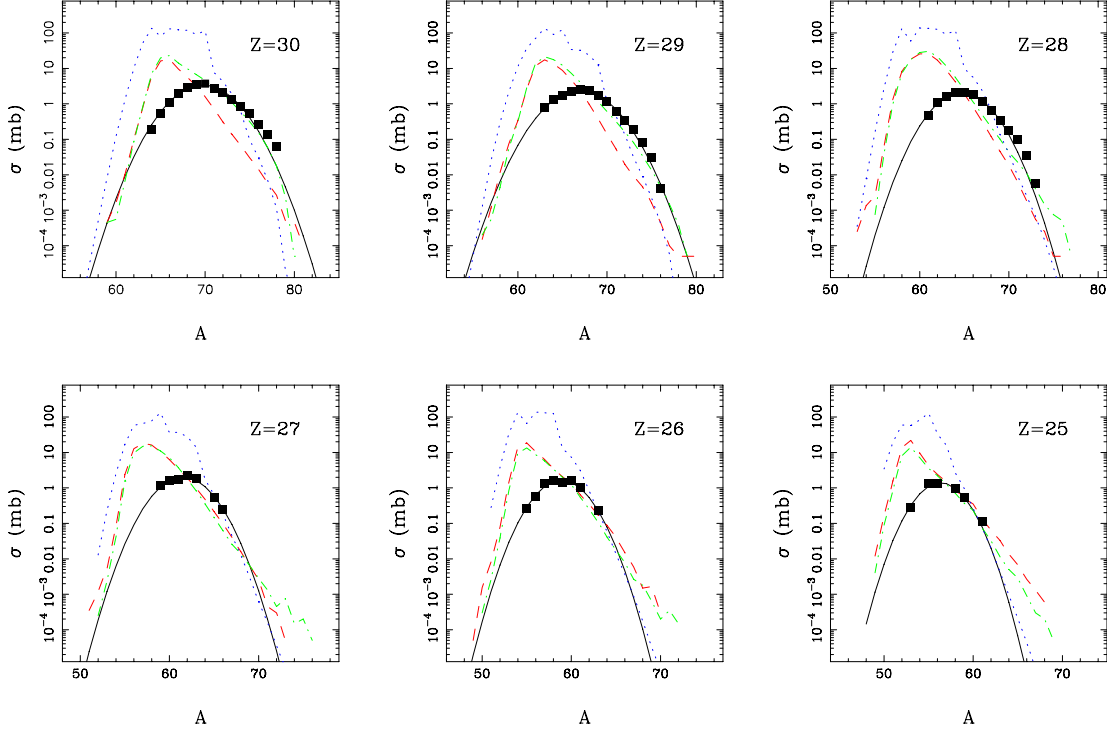


Figure 33: Production cross sections of isotopes of Zr and elements below. Symbols - experimental data for reaction of 1 GeV protons with ^{238}U [17], solid lines - Gaussian fits of experimental mass distributions, blue lines - estimates of production cross sections in fragmentation of stable nuclei, obtained using EPAX by optimization for each nucleus separately, red and green lines - production cross sections for reactions of ^{86}Kr , $^{82}\text{Se}+^{64}\text{Ni}$ at 25 A MeV, calculated using the modified version [49] of the deep-inelastic transfer code.

Figure 33 shows a comparison of the production of element Zr and below. The experimental spallation cross sections for the reaction of 1 GeV protons with ^{238}U [17] is represented by symbols and solid lines, representing Gaussian fits of the experimental mass distributions. The experimental cross sections for the reaction of 1 GeV protons with ^{238}U [17] were also compared to values calculated using the ABRABLA code. It appears that the calculation for elements below Zr (in the light fission tail) overestimates the production cross sections for neutron-rich nuclei, and thus a representation by Gaussian fits of experimental mass distributions was chosen.

The estimates of production cross sections in fragmentation of stable nuclei, obtained using EPAX by optimization for each nucleus separately are represented by blue lines. The production cross sections for the reactions of ^{86}Kr , $^{82}\text{Se}+^{64}\text{Ni}$ at 25 A MeV, was calculated using the modified version [49] of the deep-inelastic transfer code of Tassan-Got [46] (red and green lines). The production of cold fragmentation-like residues was modified by introducing the N/Z-dependent component of excitation energy and the cross section data from reactions $^{86}\text{Kr}+^{112,124}\text{Sn}$ at 25 A MeV, measured at 4 deg, were reproduced satisfactorily. One can see that the estimated optimum fragmentation cross sections essentially follow the experimental trend as described by the Gaussian fits. On the other hand, the production cross sections in the reactions of ^{86}Kr , $^{82}\text{Se}+^{64}\text{Ni}$ at 25 A MeV appear to dramatically overshoot the experimental trend for spallation of ^{238}U . In this aspect it is worthwhile to mention that the calculation for the reactions of $^{86}\text{Kr}+^{64}\text{Ni}$ at 25 A MeV, used in

figure 33, correctly reproduces the experimental cross sections, measured at angles below 3 degrees, as documented in figure 34 (the discrepancies at the proton-rich side are caused by restricted Bp-coverage for such products in the experiment, which was focused primarily on neutron-rich products [49, 39]).

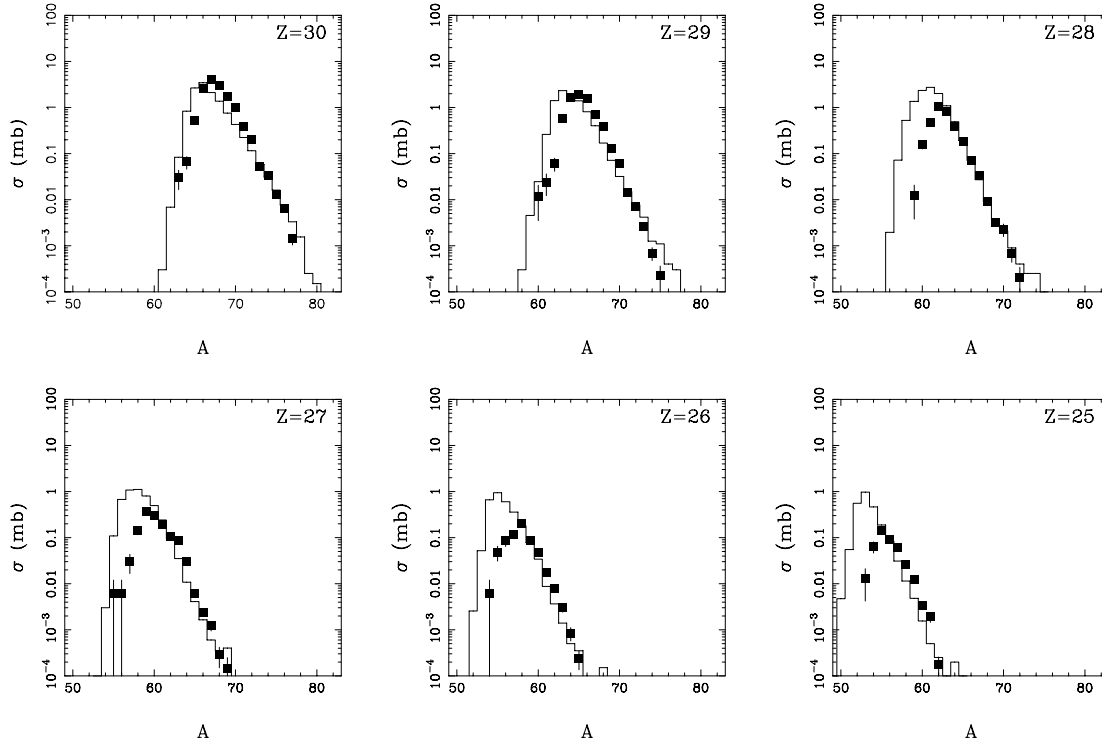


Figure 34: Experimental production cross sections of isotopes of Zr and elements below measured in the reaction $^{86}\text{Kr}+^{64}\text{Ni}$ at 25 A MeV for angles below 3 degrees [39] (symbols), compared to the results of calculation [49] (lines).

2. Estimate of the intensity of secondary beams

Figure 35 shows estimated RIB intensities of nuclides around ^{78}Ni using the production cross sections presented above. The solid lines represent estimated RIB intensities based on experimental spallation cross sections (Gaussian fits of measured mass distributions) for the reaction of 1 GeV protons with ^{238}U [17]. A primary proton-beam intensity of 100 μA impinging directly on the ^{238}U target was assumed. The target thickness corresponding to one interaction length is assumed. The efficiencies are estimated based on the work of Lukić et al. [11], where the RIB intensities for alkali elements, halogenides and noble gases measured at ISOLDE using thick UC_x target are compared to in-target yields calculated for a uranium target using ABRABLA. We adopt the parameters $\varepsilon_s=0.3$, $\alpha=1$ and $\tau=10$ s (see Equ. 1), which represent a global upper limit on extraction efficiencies for various chemical groups treated in work [11].

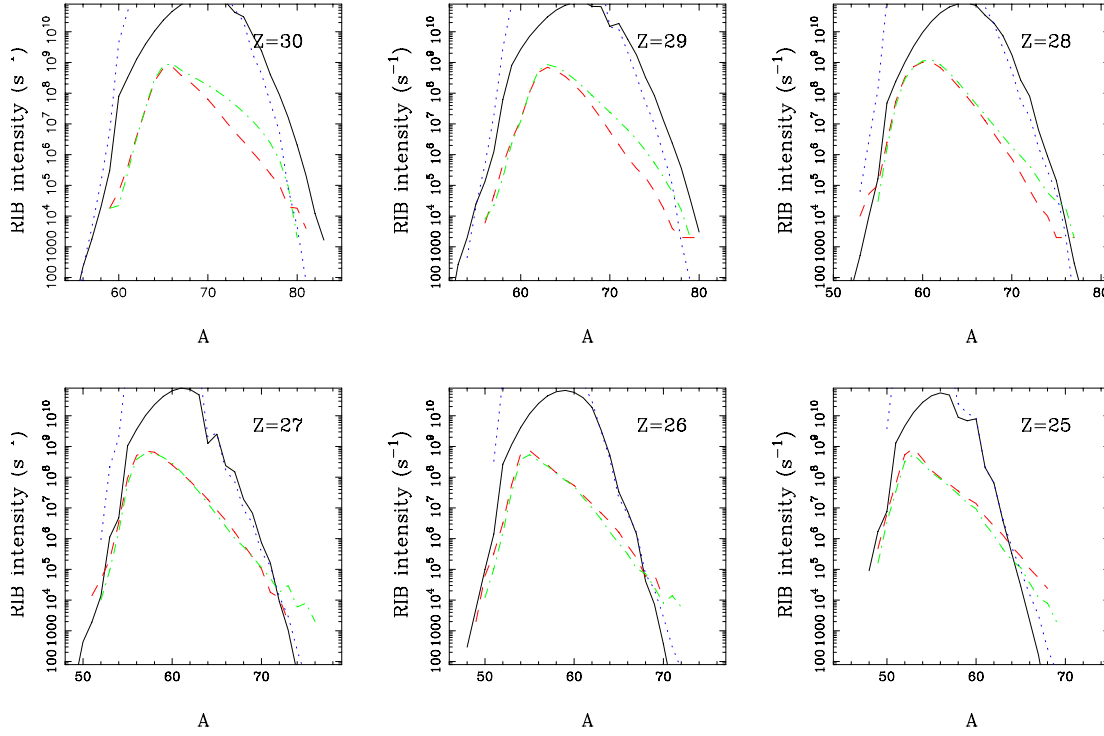


Figure 35: Estimated RIB intensities of nuclides around ^{78}Ni for primary beam intensity 100 μA . Solid lines - reaction of 1 GeV proton beam with ^{238}U target, blue lines - fragmentation of stable nuclei, red and green lines - reactions of ^{86}Kr , $^{82}\text{Se} + ^{64}\text{Ni}$ at 25 A MeV.

The blue dotted line represents estimated RIB intensities for the process of fragmentation of stable beams (the beam species optimized for each nuclide). The primary-beam intensity 100 μA was assumed. As discussed in the EURISOL Report resulting from the Feasibility Study RTD, App. C, section 2.2.3, the intensity of several hundreds of μA for beams up to krypton can be considered as achievable for the foreseeable future. The beam energy 500 A MeV was considered, high enough to allow the use of a thick target leading to reaction rates close to the proton-beam case. In particular, for the ^{86}Kr beam with the energy of 500 A MeV the reaction rate in the target will approach 60%, on the other hand the beam power will exceed 4 MW which is probably at the limit of what the converter may sustain. The beam energy 500 A MeV was not considered in the EURISOL driver heavy-ion options for beams up to krypton, where for $A/q=3$ nuclei the maximum considered energy is 100 A MeV, which however would lead to a loss of the in-target yield by one order of magnitude. The extraction efficiencies were determined as in the previous case.

Red and green lines represent estimated RIB intensities for reactions of ^{86}Kr , $^{82}\text{Se} + ^{64}\text{Ni}$ at 25 A MeV. As in the work [39], a rather moderate target thickness of 20 mg/cm^2 was assumed. The primary-beam intensity of 100 μA was chosen, and a transport efficiency of 30 % was assumed for the large-acceptance separator, thus representing the in-flight option. The angular coverage of the large-acceptance separator up to 10 degrees would be required in order to cover the most populated part of the angular distribution. The results for the 1 GeV proton beam and heavy-ion beam fragmentation are similar and dominate at the neutron-rich side of elements above Ni. However, for the extremely neutron-rich nuclides of nickel and below, the reactions with a relatively thin target and beam energy of 25 A MeV appear to lead to higher RIB intensities than in other scenarios, due to the high in-target yields. Nevertheless, the quality of such secondary beams would be lower than in other cases. The use of a gas cell may be an alternative, leading to better beam quality, provided that losses will not increase dramatically, which is an open question needing further R&D.

VI. CONCLUSION

A number of approaches for the production of radioactive beams have been examined with regard to their potential to be used as possible additional options for the EURISOL project in order to enhance the production of specific nuclear species with respect to the baseline 1-GeV proton case. We found that the spallation of suitable target material by 1 GeV protons and the fission induced by secondary neutrons in a uranium target provide overall high intensities for secondary beams almost all over the chart of the nuclides.

Still, there are cases where the 1 GeV proton beam together with the restrictions of the ISOL method in view of suitable target material do not allow optimizing the reaction parameters. In addition, low overall efficiencies for certain elements let gaps in available beams. Some of these problems can be overcome by providing heavier projectiles and/or higher beam energies:

- A 2 GeV $^3\text{He}^{2+}$ beam would fill the gaps in the nuclide production given by the limited choice of ISOL target materials. This would lead to gain factors up to a factor of 4 in particular for the production of neutron-deficient isotopes of many elements.
- A 2 GeV $^3\text{He}^{2+}$ beam would increase the production of neutron-rich isotopes of light to medium-heavy elements ($Z < 30$) by about a factor of 2.
- The deuteron-converter option with a primary-beam energy above 40 MeV provides fission-fragment nuclide distributions with appreciably higher fission yields (normalized to the total number of fission events in the target) for elements between technetium ($Z = 43$) and indium ($Z = 49$), below germanium ($Z = 32$), and above neodymium ($Z = 60$) compared to the standard EURISOL high-power-target option. However, only part of this advantage can be used, since many of the enhanced elements are poorly released from ISOL targets.
- Fragmentation of heavy-ion projectiles provides higher in-target yields for some neutron-deficient isotopes of light elements and presumably higher overall ISOL efficiencies for short-lived isotopes. It can also be useful to overcome limitations in the choice of the target material in the standard proton option and to divide production target and catcher. The gain factors depend strongly on the beam energy.
- Due to nucleon exchange between projectile and target, heavy-ion reactions in the Fermi-energy regime (around 20 to 30 A MeV) provide a substantial benefit for the production of neutron-rich isotopes of elements outside the main fission region.

The listed approaches require extended capabilities of the driver accelerator. Thus, the benefit of the options discussed in this report should be carefully considered in view of the extra cost involved.

The present work should be considered as preliminary. First, not all options have been investigated in full quantitative detail. Further on, we could not consider in a conclusive way the possible advance of technological limits, on e.g. currents provided by the ion source or heat load in the target, until the time EURISOL is constructed. Nevertheless, we hope that this report gives first guidance to the problem and stimulates a fruitful discussion on the questions that remain open.

Acknowledgement

We acknowledge the financial support of the European Community under the FP6 “Research Infrastructure Action – Structuring the European Research Area”, EURISOL DS Project, contract no. 515768 RIDS. The EC is not liable for any use that may be made of the information contained herein. We would like to thank Ulli Köster for bringing to our attention many of the ideas developed in the ISOLDE community, as well as for many stimulating

remarks and careful reading of the manuscript. We also acknowledge valuable discussions with Wolfgang Mittig on the converter-catcher - ISOL scenario using heavy-ion beams.

References:

- [1] <http://www.eurisol.org>
- [2] <http://isolde.cern.ch>
- [3] H. L. Ravn, P. Bricault, G. Ciavola, P. Drumm, B. Fogelberg, E. Hagebo, M. Huyse, R. Kirchner, W. Mittig, A. Mueller, H. Nifenecker and E. Roeckl, Nucl. Instr. Meth. B 88, 441 (1994)
- [4] http://www.ganil.fr/eurisol/Final_Report.html
- [5] <http://www.orau.org/ria>
- [6] J. Nolen, in *Proceedings of the Third international conference on radioactive nuclear beams*, East Lansing, May 24-27, 1993, edited by D.J. Morissey,
- [7] E. Cottureau, Nucl. Phys. A 701, 78 (2002)
- [8] U. Köster, Radiochim. Acta 89, 749 (2001)
- [9] H. L. Ravn, Nucl. Instrum. Methods B 26, 72 (1987)
- [10] K. Helariutta, J. Benlliure, M. V. Ricciardi and K.-H. Schmidt, European Physical Journal A 17, 181 (2003)
- [11] S. Lukic, F. Gevaert, A. Kelic, M. V. Ricciardi, K.-H. Schmidt and O. Yordanov, arXiv nucl-ex/0601031, Nucl. Instr. Meth. A, accepted
- [12] C. L. Jiang, B. B. Back, I. Gomes, A. M. Heinz, J. Nolen, K. E. Rehm, G. Savard and J. P. Schiffer, Nucl. Instr. Meth. A 492 (2002) 57
- [13] C. Villagrasa-Canton, PhD thesis, Université de Paris XI, 2003
- [14] P. Napolitani, K.-H. Schmidt, A. S. Botvina, F. Rejmund, L. Tassan-Got and C. Villagrasa, Phys. Rev. C 70, 054607 (2004)
- [15] P. Napolitani, PhD thesis, Université Paris XI, 2004
- [16] J. Taieb, K.-H. Schmidt, L. Tassan-Got, P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, E. Casarejos, S. Czajkowski, T. Enqvist, R. Legrain, S. Leray, B. Mustapha, M. Pravikoff, F. Rejmund, C. Stephan, C. Volant and W. Wlazlo, Nucl. Phys. A 724, 413 (2003)
- [17] M. Bernas, P. Armbruster, J. Benlliure, A. Boudard, E. Casarejos, S. Czajkowski, T. Enqvist, R. Legrain, S. Leray, B. Mustapha, P. Napolitani, J. Pereira, F. Rejmund, M. V. Ricciardi, K.-H. Schmidt, C. Stéphan, J. Taieb, L. Tassan-Got and C. Volant, Nucl. Phys. A 725, 213 (2003)
- [18] M. V. Ricciardi, P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, S. Czajkowski, T. Enqvist, A. Kelic, S. Leray, R. Legrain, B. Mustapha, J. Pereira, F. Rejmund, K.-H. Schmidt, C. Stephan, L. Tassan-Got, C. Volant and O. Yordanov, Phys. Rev. C 73, 014607 (2006)
- [19] P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, E. Casarejos, S. Czajkowski, T. Enqvist, S. Leray, P. Napolitani, J. Pereira, F. Rejmund, M.-V. Ricciardi, K.-H. Schmidt, C. Stephan, J. Taieb, L. Tassan-Got and C. Volant, Phys. Rev. Lett. 93, 212701 (2004)
- [20] J.-J. Gaimard and K.-H. Schmidt, Nucl. Phys. A 531, 709 (1991)
- [21] A. R. Junghans, M. de Jong, H.-G. Clerc, A. V. Ignatyuk, G. A. Kudyaev and K.-H. Schmidt, Nucl. Phys. A 629, 635 (1998)
- [22] J. Benlliure, A. Grewe, M. de Jong, K.-H. Schmidt and S. Zhdanov, Nucl. Phys. A 628, 458 (1998)
- [23] B. Jurado, K.-H. Schmidt and J. Benlliure, Phys. Lett. B 533, 186 (2003)
- [24] J. Benlliure, K.-H. Schmidt, D. Cortina-Gil, T. Enqvist, F. Farget, A. Heinz, A. R. Junghans, J. Pereira and J. Taieb, Nucl. Phys. A 660 (1999) 87
- [25] B. Fernandez-Dominguez, P. Armbruster, L. Audouin, J. Benlliure, M. Bernas, A. Boudard, E. Casarejos, S. Czajkowski, J. E. Ducret, T. Enqvist, B. Jurado, R. Legrain, S. Leray, B. Mustapha, J. Pereira, M. Pravikoff, F. Rejmund, M. V. Ricciardi, K.-H. Schmidt, C. Stephan, J. Taieb, L. Tassan-Got, C. Volant and W. Wlazlo, Nucl. Phys. A 747, 227 (2005)
- [26] L. Audouin, L. Tassan-Got, P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, E. Casarejos, S. Czajkowski, T. Enqvist, B. Fernandez-Dominguez, B. Jurado, R. Legrain, S. Leray, B. Mustapha, J. Pereira, M. Pravikoff, F. Rejmund, M.-V. Ricciardi, K.-H. Schmidt, C. Stephan, J. Taieb, C. Volant and W. Wlazlo, arXiv nucl-ex/0503021
- [27] T. Enqvist, W. Wlazlo, P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, S. Czajkowski, R. Legrain, S. Leray, B. Mustapha, M. Pravikoff, F. Rejmund, K.-H. Schmidt, C. Stephan, J. Taieb, L. Tassan-Got and C. Volant, Nucl. Phys. A 686, 481 (2001)
- [28] T. Enqvist, P. Armbruster, J. Benlliure, M. Bernas, A. Boudard, S. Czajkowski, R. Legrain, S. Leray, B. Mustapha, M. Pravikoff, F. Rejmund, K.-H. Schmidt, C. Stéphan, J. Taieb, L. Tassan-Got, F. Vivès, C. Volant and W. Wlazlo, Nucl. Phys. A 703, 435 (2002)
- [29] S. B. Kaufmann and E. P. Steinberg, Phys. Rev. C 22, 167 (1980)
- [30] J. Lettry, R. Catherall, P. Drumm, P. Van Duppen, A. H. M. Evensen, G. J. Focker, A. Jokinen, O. C. Jonsson, E. Kugler, H. Ravn and ISOLDE Collaboration, Nucl. Instr. Meth. B 126, 130 (1997)
- [31] S. Lukić, C. Clausen, T. Haarju and U. Köster, unpublished
- [32] H.-J. Kluge, *Isolde user's guide*, CERN, Geneva, 1986,
- [33] D. Ridikas, private communication

-
- [34] S. Ménard, M. Mirea, F. Clapier, N. Pauwels, J. Proust, C. Donzaud, D. Guillemaud-Mueller, I. Lhenry, A. C. Mueller, J. A. Scarpaci and O. Sorlin, *Phys. Rev. Spec. Topics* 2, 033501 (1999)
- [35] H. Penttillä, contribution to EURISOL-Task11 meeting, Caen, November 2005
- [36] T. Bjornstadt, H.-A. Gustafsson, B. Jonson, O. C. Jonsson, V. Lindfors, S. Mattsson, A. M. Poskanzer, H. L. Ravn and D. Schardt, *Z. Phys. A* 303, 227 (1981)
- [37] T. Bjornstadt, L. C. Carraz, H.-A. Gustafsson, J. Heinemeyer, B. Jonson, O. C. Jonsson, V. Lindfors, S. Mattsson, A. M. Poskanzer and H. L. Ravn, *Nucl. Instrum. Methods* 186, 391 (1981)
- [38] M. Veselsky, *Nucl. Phys. A* 705, 193 (2002)
- [39] G.A. Souliotis, M. Veselsky, G. Chubarian, L. Trache, A. Keksis, E. Martin, A. Ruangma, E. Winchester and S. J. Yennello, *Phys. Lett. B* 543, 163 (2002)
- [40] G.A. Souliotis, D.V. Shetty, M. Veselsky, G. Chubarian, L. Trache, A. Keksis, E. Martin, and S.J. Yennello, *Phys. Rev. C* 68, 24605 (2003)
- [41] M. Veselsky, G. A. Souliotis, G. Chubarian, L. Trache, A. Keksis, E. Martin, A. Ruangma, E. Winchester and S. J. Yennello, *Nucl. Phys. A* 724, 431 (2003)
- [42] G.A. Souliotis, M. Veselsky, G. Chubarian, L. Trache, A. Keksis, E. Martin, D.V. Shetty, and S.J. Yennello, *Phys. Rev. Lett.* 91, 022701 (2003)
- [43] G.A. Souliotis, M. Veselsky, G. Chubarian and S. J. Yennello, *Nucl. Instr. Meth. B* 204 (2003) 166
- [44] K. Sümmerer and B. Blank, *Phys. Rev. C* 61, 34607 (2000)
- [45] M. Veselsky, R.W. Ibbotson, R. Laforest, E. Ramakrishnan, D.J. Rowland, A. Ruangma, E.M. Winchester, E. Martin, and S.J. Yennello, *Phys. Rev. C* 62, 064613 (2000)
- [46] L. Tassan-Got, PhD Thesis, Orsay, France, 1988, IPNO-T-89-02, 1989; L. Tassan-Got and C. Stefan, *Nucl. Phys. A* 524, 121 (1991)
- [47] J. P. Bondorf, A. S. Botvina, A. S. Iljinov, I. N. Mishustin and K. Sneppen, *Phys. Rep.* 257, 133 (1995)
- [48] R. Charity, M.A. McMahan, G.J. Wozniak, R.J. McDonald, L.G. Moretto, D.G. Sarantites, L.G. Sobotka, G. Guarino, A. Pantaleo, L. Fiore, A. Gobbi and K.D. Hildenbrand, *Nucl. Phys. A* 483, 371 (1988)
- [49] M. Veselsky and G.A. Souliotis, *Nucl. Phys. A* 765, 252 (2006)
- [50] A. Boudard, J. Cugnon, S. Leray and C. Volant, *Phys. Rev. C* 66, 044615 (2002)
- [51] EURISOL town meeting, Caen, 2005
- [52] <http://www.ganil.fr/research/developments/spiral2/index.html>
- [53] C.L. Jiang, B.B. Back, I. Gomes, A. M. Heinz, J. Nolen, K. E. Rehm, G. Savard and J.P. Schiffer, *Nucl. Instrum. Methods A* 492, 57 (2002)
- [54] A. Villari, Contribution to the EURISOL town meeting Orsay, France, December 2003
- [55] <http://www-w2k.gsi.de/charms/amadeus.htm>