



Submitted to
Physics Letters B

CERN-EP-PHYS 78-45 MW
C.

CERN - EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

CERN/EP/PHYS 78-45 Rev.
22 December 1978

13 FEB. 1979

OBSERVATION OF CHARMED D MESON PRODUCTION IN pp COLLISIONS

CERN-Collège de France-Heidelberg-Karlsruhe Collaboration

D. DRIJARD, H.G. FISCHER, W. GEIST, R. GOKIELI¹⁾, P.G. INNOCENTI, V. KORBEL²⁾,
A. MINTEN, A. NORTON, R. SOSNOWSKI¹⁾, S. STEIN³⁾, O. ULLALAND and H.D. WAHL
CERN, Geneva, Switzerland.

P. BURLAUD, M. DELLA NEGRA⁴⁾, G. FONTAINE, P. FRENKIEL, C. GHESQUIERE,
D. LINGLIN⁴⁾ and G. SAJOT
Collège de France, Paris, France.

H. FREHSE⁵⁾, E.E. KLUGE, M. HEIDEN, A. PUTZER and J. STIEWE
Institut für Hochenergiephysik der Universität Heidelberg, Germany.

P. HANKE⁵⁾, W. HOFMANN⁶⁾, M. PANTER⁶⁾, K. RAUSCHNABEL⁶⁾, J. SPENGLER⁶⁾ and
D. WEGENER⁶⁾
Institut für Experimentelle Kernphysik der Universität Karlsruhe, Germany.

-
- 1) Now at Institute for Nuclear Research, Warsaw, Poland
 - 2) Now at DESY, Hamburg, Germany
 - 3) Now at SLAC, Stanford, USA
 - 4) Now at LAPP, Annecy, France
 - 5) Now at CERN, Geneva, Switzerland
 - 6) Now at Institut für Physik der Universität Dortmund, Germany

HDW/jm

ABSTRACT

Evidence for the forward production of D^+ mesons has been found in pp collisions at a c.m. energy of 52.5 GeV. A signal of 92 ± 18 events above background has been observed in the decay channel $D^+ \rightarrow K^{*0} \pi^+ \rightarrow K^- \pi^+ \pi^+$.

The cross section in the phase space region of our experiment is given and implications for the total cross section for charmed D meson production are discussed.

1. INTRODUCTION

Ever since the prediction of charmed D-mesons [1] and their observation in e^+e^- annihilation [2], considerable experimental effort has been devoted to their detection in hadron-hadron interactions. So far, no direct evidence for charm production in hadron interactions has been found [3]. In an experiment done at higher energy than most prior searches, we have studied the mass distributions of $K^-\pi$ and $K^-\pi\pi$ systems, produced in the forward direction in pp collisions.

2. EXPERIMENT, APPARATUS, TRIGGER

The experiment was performed at the CERN Intersecting Storage Rings (ISR), using the Split Field Magnet (SFM) detector, at a c.m. energy $\sqrt{s} = 52.5$ GeV. The experimental set-up and procedures of data acquisition and analysis have been described in previous publications [4,5]. The results presented here are based on events obtained by triggering on a negative particle emitted at a c.m.s. polar angle $\theta \sim 8^\circ$ with a transverse momentum larger than 0.5 GeV/c (see table 1 for a summary of the phase space region covered by the trigger). A threshold Cerenkov counter covering the acceptance of the trigger allowed a separation into "pion" and "non-pion" samples (see [refs 5 and 6] for details). In the non-pion sample, about 80% of the triggering particles are K^- , 20% are \bar{p} , and the pion contamination is negligible. In the following presentation, all triggering particles in the non-pion sample are treated as kaons, and all other particles are assumed to be pions.

3. CUTS APPLIED TO THE DATA

The $K^-\pi$ and $K^-\pi\pi$ mass spectra presented below are uncorrected for acceptance and trigger efficiency and subject to some or all of the following selection criteria:

(A) Multiplicity: events are accepted when $n_{\text{obs}} \leq 10$ (n_{obs} = number of observed particles). This selection reduces the combinatorial background in the $K\pi$ and $K\pi\pi$ spectra.

(B) Leading system in the opposite direction: events are accepted only if $|x_{\text{opp}}| > 0.5$ or $|x_{\text{opp}}| < 0.1$, where x_{opp} is the total reduced longitudinal momentum seen in the hemisphere opposite to the longitudinal direction of the trigger. This cut selects those events in which a leading system is observed opposite in x to the trigger ($|x_{\text{opp}}| > 0.5$) or escapes detection ($|x_{\text{opp}}| < 0.1$).

(C) Longitudinal momentum of K^- : the triggering K^- is required to have reduced longitudinal momentum $x_K > 0.3$. This is similar to a P_T cut since the polar angle of the trigger has a small range (see table 1).

(D) Recoiling system in the same longitudinal direction: the "recoiling system" is defined as the set of particles emitted in the same longitudinal direction as the triggering K^- , but with a transverse component opposite to its transverse momentum \vec{q}_{TK} . We define a quantity p_R by

$$p_R = - \sum k_{Ti} = - \sum p_i \cdot (\vec{q}_{TK} / |\vec{q}_{TK}|),$$

where the sum is extended over the particles in the recoiling system, i.e. particles with $k_{Ti} < 0$ and $p_{Li} > 0$ (note that $x_K > 0$). Events are accepted only when $p_R > 0.2$ GeV/c. This selection favours events where additional particles are produced close in phase space to the D , which (at least partly) compensate its transverse momentum.

Table 2 contains a summary of the numbers of events fulfilling the various conditions.

4. MASS SPECTRA

The $K^- \pi^+$ mass spectrum is shown in fig. 1. It shows a clear $K^*(890)$ signal, but no evidence for an enhancement in the D^0 region. Also in the $K^- \pi^+ \pi^+$ mass spectrum no obvious signal corresponding to $D^+ \rightarrow K^- \pi^+ \pi^+$ is observed. Since decay of the D via $K^*(890)$ is reported to be an important fraction of its semi-leptonic decay (50 to 70%) [7,8], it is conceivable that the same

is also true for the hadronic decay into $K\pi\pi$. Furthermore, the Dalitz plot for the decay $D^+ \rightarrow K^- \pi^+ \pi^+$, as observed in e^+e^- annihilation [9], does not exclude the presence of a $\overline{K}^{*0} \pi^+$ contribution.

Therefore, we show in fig. 2 the $K^- \pi^+ \pi^+$ mass distribution under the condition that the mass of at least one of the two $(K^- \pi^+)$ pairs is inside the $K^*(890)$ region (of full width 60 MeV around the K^* mass). Indeed, an enhancement is seen in the D mass region, which becomes more pronounced if conditions (C) and (D) are imposed in addition to (A) and (B). Since the quantum numbers of the $\overline{K}^{*0} \pi^+$ system are exotic, the signal is not compatible with any known non-charmed resonance, but it is consistent with the decay $D^+ \rightarrow K^{*0} \pi^+$ (i.e. $\Delta C = \Delta S$). We have verified that the signal is not the reflection of a resonance whose decay products have been misidentified.

In order to estimate the significance of the signal and to verify that the signal is not due to kinematics and/or acceptance, three different background spectra have been considered, all obtained from the data with the same selection criteria as used for the $K^- \pi^+ \pi^+$ spectrum of fig. 3(a):

- (a) the $K^- \pi^- \pi^-$ and $K^- \pi^+ \pi^-$ mass distribution (shown in fig. 3(b));
- (b) the $K^- \pi^+ \pi^+$ distribution obtained by combining the triggering K^- with pions from different events;
- (c) the " K^- " $\pi^+ \pi^+$ spectrum obtained by assigning the K^- mass to the triggering π^- in the pion sample.

Since none of these background spectra (which all have very similar shapes) show an enhancement in the D-region, we conclude that the signal is genuine and is not generated by the experimental procedure.

The background shown in fig. 3 is obtained by adding the background spectra (a) and (b) mentioned above and normalizing to the number of $K\pi\pi$ combinations outside the D-region (i.e. the three bins from 1.8 to 2.04 GeV, in which the signal is seen). Subtraction of this background yields a signal of

$$92 \pm 18 \text{ events.}$$

We have estimated the mass of the D to be $M = 1.91 \text{ GeV}/c^2$. The difference of $40 \text{ MeV}/c^2$ between our estimation and published results is compatible

with local systematic errors added to the statistical uncertainty $\sigma = 20 \text{ MeV}/c^2$. The width of the signal is compatible with our experimental mass resolution (FWHM = $150 \text{ MeV}/c^2$).

5. CORRECTIONS AND CROSS SECTIONS

The inclusive cross section for D production has been determined by Monte-Carlo methods, taking account of the experimental acceptance and the cuts on kaon momentum and $K^- \pi^+$ mass. Generating a sample of K^- events into the same acceptance region as for the D, we obtain the absolute normalization by comparison with published K^- data [10] (after correction for the \bar{p} content of the trigger sample).

The losses in signal due to the cuts on event topology ((A), (B) and (D) in sect. 3) have been estimated by inspection of the complementary sample of rejected events; by the absence of a D signal in this sample we obtain an additional contribution which is at most 28 events.

The acceptance region of our trigger for D mesons (Ω_D) covers a sizeable fraction of phase space ($\phi \in (90, 270)$, $y \in (1.7, 3.2)$ and full p_T range), considerably larger than the corresponding region for the K^- trigger particle. Consequently our global acceptance for D mesons in Ω_D , and the resulting estimation of the cross-section $(B \cdot \sigma)_{\Omega_D}$, are sensitive to the momentum dependence of D production. We have therefore performed model calculations with various assumptions about the longitudinal and transverse momentum dependence of the D cross section.

We assume scaling and approximate factorization in the form

$$E \frac{d^3\sigma}{d^3p} \sim f(y/y_{\max}) e^{-Ap_T},$$

where $y_{\max}(p_T)$ is the maximum rapidity allowed by kinematics.

Two different p_T slope parameters, $A = 2$ and $A = 2.8 \text{ (GeV}/c)^{-1}$ are chosen. The function $f(y/y_{\max})$ is adjusted to generate three different production mechanisms: proportional to $(1-x)^3$, uniform in rapidity and uniform in x . Isotropic decay angle distributions are assumed for both the decays $D^+ \rightarrow K^{*0} \pi^+$ and $K^{*0} \rightarrow K^- \pi^+$. The corresponding inclusive cross sections in our acceptance region Ω_D , $(B \cdot \sigma)_{\Omega_D}$, are presented in table 3. They are rather sensitive to the slope parameter A and can vary by a factor of 6 under extreme assumptions.

Table 3 also gives estimates of the corresponding total cross section $(\sigma \cdot B)_{\text{tot}}$, where further uncertainties are introduced due to the extrapolation to the full rapidity range. In this case, the values change by up to a factor of 15 for the different production mechanisms. Apart from these uncertainties, the individual values for the different models have an estimated error of $\sim 40\%$, due to the following contributions:

- Uncertainty in the size of the signal (including the correction for the cuts on event topology): $\sim 30\%$;
- Uncertainty in the acceptance calculation: $\sim 20\%$;
- Uncertainty in the absolute normalisation: $\sim 30\%$.

If we assume [8] that $\sim 2/3$ of the measured decay $D^+ \rightarrow K^- \pi^+ \pi^+$ [11] proceeds via K^{*0} formation, we obtain a branching ratio of $(2.6 \pm 1.0)\%$ for our decay channel. This would correspond to a total cross section for D^+ production between 150 and 2000 μb (with errors of $\sim 60\%$), depending on the production mechanism. The higher value is incompatible with measured upper limits for D production at ISR energies [12]; hence our data indicate that the D is produced with a rather flat distribution in p_T and rapidity or Feynman x . This mechanism would also help in understanding the results of the CERN beam dump experiments [13].

Acknowledgement

This experiment was greatly helped by contributions from the SFM detector group, in particular by W. Bell, E. Chesi, B. Heck, L. Naumann and F. Piuz. We are indebted to H.F. Hoffmann and the ISR experimental support group. For on-line programming, the contributions of R. Cooper and M. Sciré were essential. We wish to thank U. Schlüpmann and the SFM software group for assistance and for development of the off-line programs. We are grateful to J. Moritz (Karlsruhe) for his important contribution to design and construction of the Cerenkov-counters. The Heidelberg and the Karlsruhe group have been supported by a grant from the Bundesministerium für Forschung und Technologie of the Federal Republic of Germany, the Collège de France group was funded by IN2P3 and the CEA.

REFERENCES

- [1] M.K. Gaillard, B.W. Lee and J.L. Rosner, Rev. Mod. Phys. 47 (1975) 277, and references quoted therein.
- [2] G. Goldhaber et al., Phys. Rev. Lett. 37 (1976) 255;
J. Peruzzi et al., Phys. Rev. Lett. 37 (1976) 569.
- [3] See, e.g. R. Diebold, Rapporteur's talk at the 19 Int. Conf. on High Energy Physics, Tokyo (1978).
- [4] CCHK Collaboration, M. Della-Negra, et al., Nucl. Phys. B127 (1977) 1;
CCHK Collaboration, D. Drijard et al., CERN/EP/PHYS 78-14, submitted to Nucl. Phys. B.
- [5] CCHK Collaboration, M. Della Negra et al., Phys. Lett. 59B (1975) 481.
- [6] P. Hanke, Ph. D. thesis, University of Karlsruhe, report KFK 2412 (1977).
- [7] DASP Collaboration, R. Brandelik et al., Phys. Lett. 70B (1977) 387;
J.M. Feller et al., Phys. Rev. Lett. 40 (1978) 274.
- [8] DELCO Collaboration, results presented by J. Kirkby at the Summer Institute on Particle Physics, Stanford, July (1978), SLAC-PUB-2231.
- [9] J.E. Wiss et al., Phys. Rev. Lett. 37 (1976) 1531.
- [10] British-Scandinavian Collaboration, B. Alper et al., Nucl. Phys. B100 (1975) 237;
Bologna-CERN Collaboration, P. Capiluppi et al., Nucl. Phys. B70 (1974) 1.
- [11] I. Peruzzi et al., Phys. Rev. Lett. 39 (1977) 1301.
- [12] A.G. Clark et al., Phys. Lett. 77B (1978) 339;
L. Baum et al., Phys. Lett. 68B (1977) 279;
J. Alder et al., Phys. Lett. 66B (1977) 401;
M. Block et al., Nucl. Phys. B140 (1978) 525.
- [13] Gargamelle Collaboration, P. Alibrant et al., Phys. Lett. 74B (1978) 134;
CDHS Collaboration, T. Hansl et al., Phys. Lett. 74B (1978) 139;
ABCLOS Collaboration, P.C. Bosetti et al., Phys. Lett. 74B (1978) 143.

TABLE CAPTIONS

Table 1 Parameters characterizing the kinematical region covered by the acceptance of the K^- trigger.

Table 2 Number of events fulfilling the various conditions imposed on the data sample.

Table 3 Summary of cross sections for D production obtained by different model calculations.

Table 1

Variable	Range	Mean	FWHM	r.m.s.
x	0.2 - 0.6	0.31	0.09	0.04
y	2 - 3	2.65	0.38	0.15
P_T (GeV/c)	> 0.5	1.15	0.55	0.27
θ	$5^\circ - 13^\circ$	8.0°	3.1°	1.28°
ϕ	$144^\circ - 168^\circ$	158°	10.4°	3.9°
	$174^\circ - 216^\circ$	193°	21.2°	7.4°

Table 2

Condition	Nb. of events fulfilling condition	
	before K* selection	after
none	98 884	8910
(C)	49 707	4055
(A) (B)	37 376	3457
(A) (B) (C)	18 561	1580
(A) (B) (C) (D)	12 124	977

Table 3

Longitudinal dependence	p_T slope A, $(\text{GeV}/c)^{-1}$	$(B \cdot \sigma)_{\Omega_D}$ μb	$(B \cdot \sigma)_{\text{tot}}$ μb
$d\sigma/dx \propto (1 - x)^3$	2	2.81	21.6
	2.8	8.40	62.0
$d\sigma/dy = \text{const}$	2	1.67	7.5
	2.8	4.38	19.3
$d\sigma/dx = \text{const}$	2	1.44	4.0
	2.8	3.80	10.5

FIGURE CAPTIONS

Fig. 1 The $K^- \pi^+$ mass distribution obtained without cuts (full line) and with conditions (A), (B) (broken line).

Fig. 2 The $K^- \pi^+ \pi^+$ mass distribution under the condition that at least one of the two $K^- \pi^+$ combination has a mass in the K^* region (± 30 MeV around the K^* mass), for events fulfilling various conditions: (A) and (B) (full line), (A), (B) and (C) (broken line), and (A), (B), (C) and (D) (shaded histogram).

Fig. 3 $K^- \pi \pi$ mass distributions under the condition that at least one of the two $K^- \pi$ combinations (independent of their charge) has a mass in the K^* region, for events fulfilling conditions (A), (B), (C) and (D):

- (a) $K^- \pi^+ \pi^+$,
- (b) Sum of $K^- \pi^+ \pi^-$ and $K^- \pi^- \pi^-$.

The broken line represents the background obtained by the procedure outlined in sect. 4.

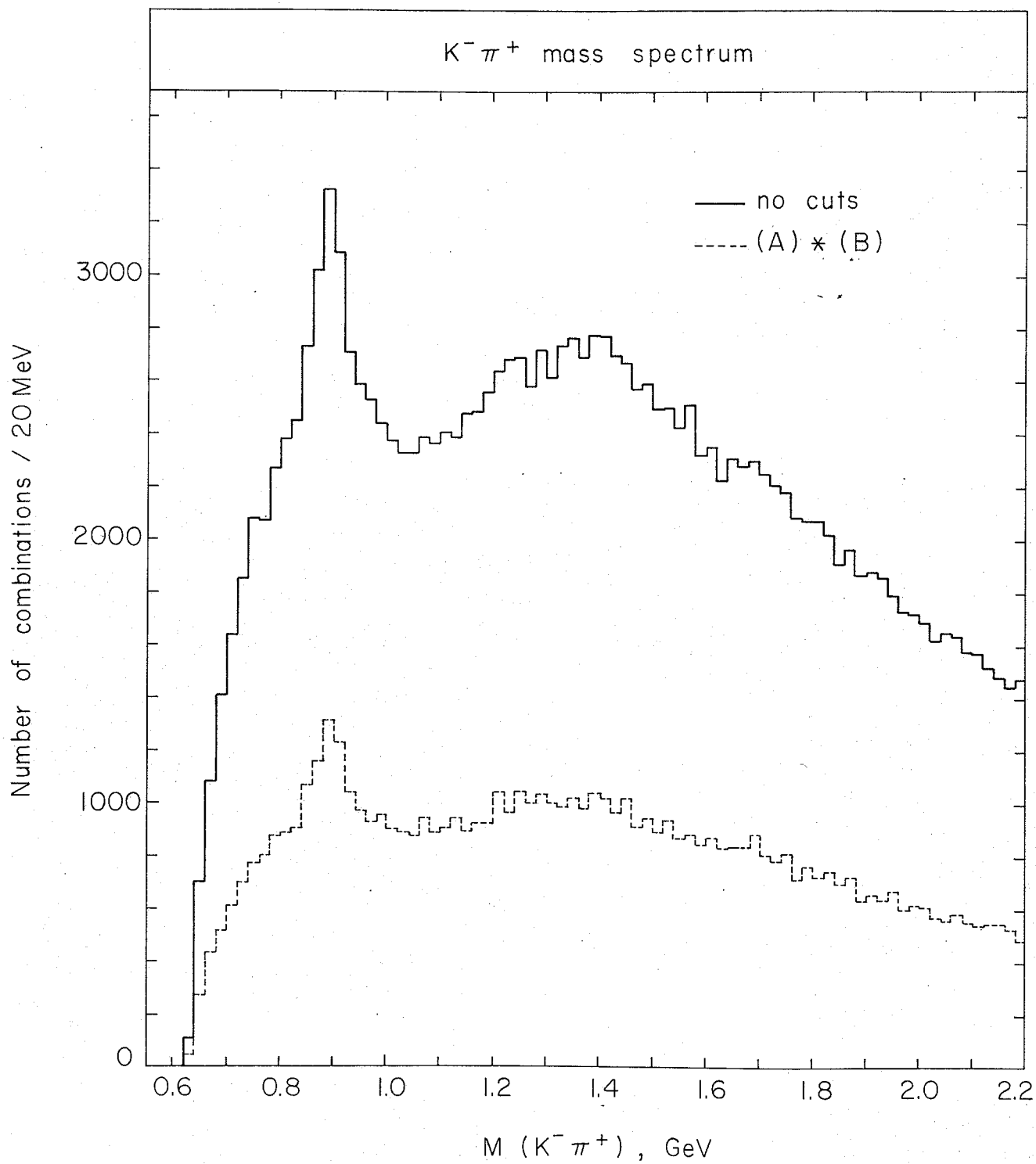


Fig. 1

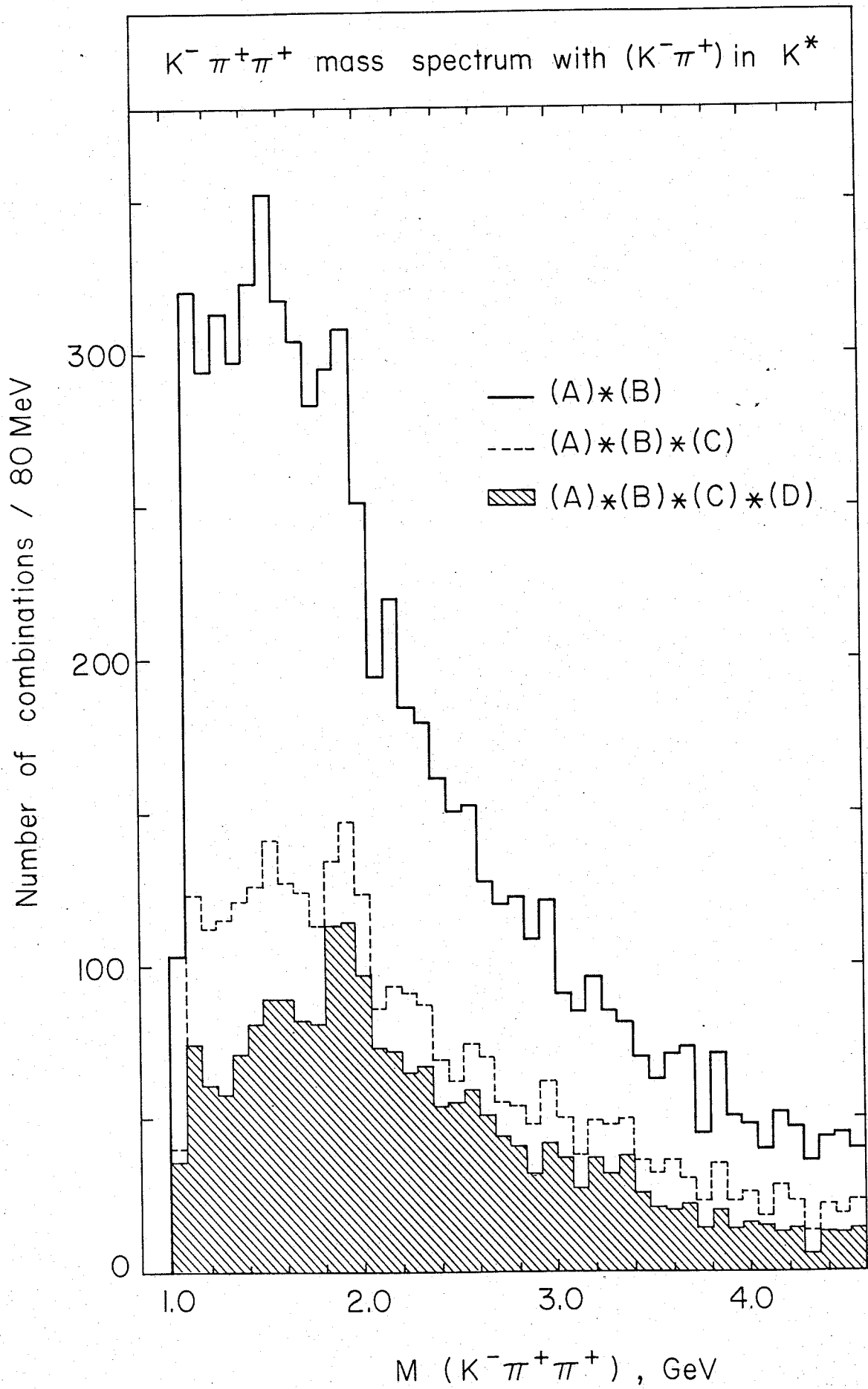


Fig. 2

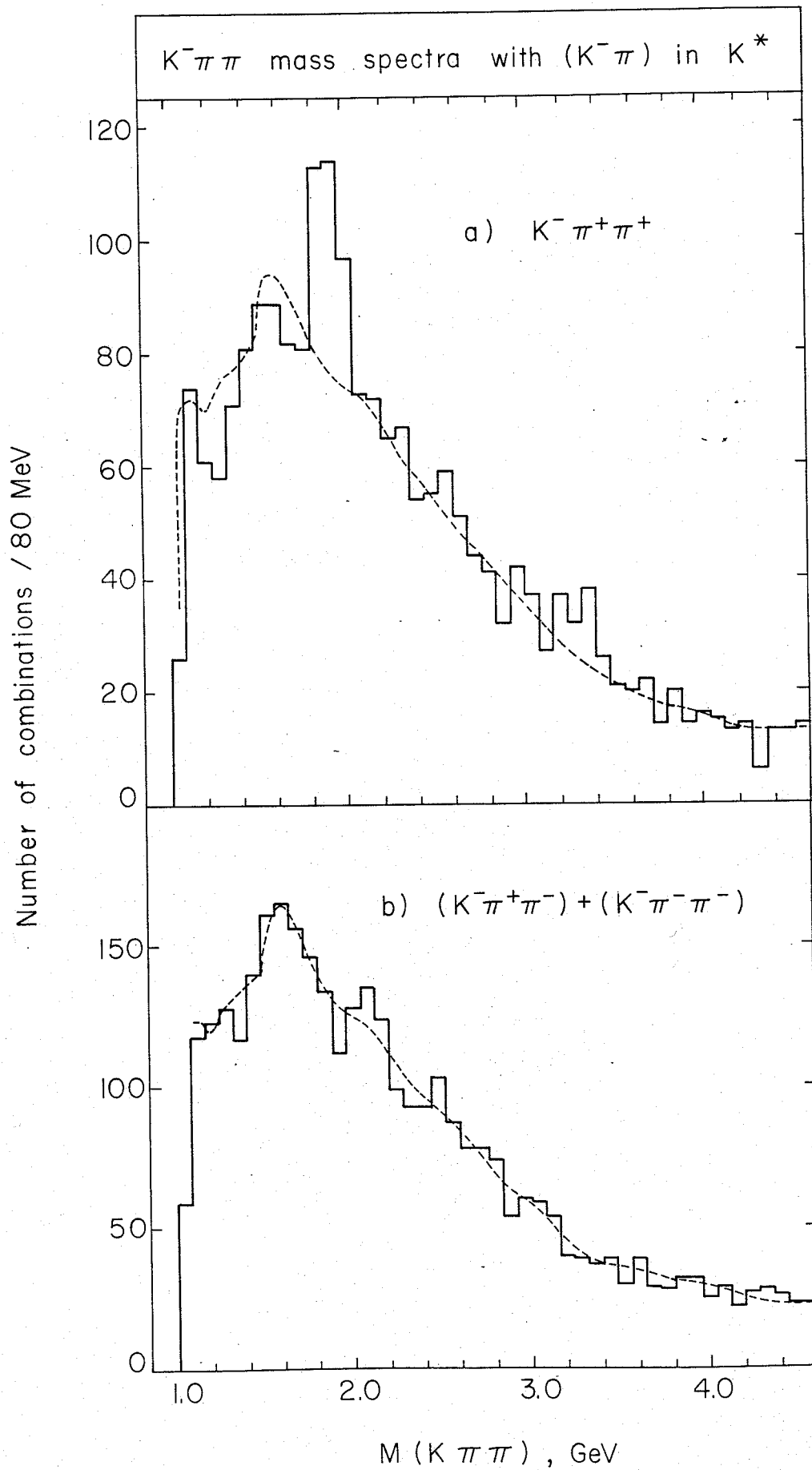


Fig. 3